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## A 24 HR GLOBAL CAMPAIGN TO ASSESS PRECISION TIMING OF THE MILLISECOND PULSAR J1713+0747

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### ABSTRACT

The radio millisecond pulsar J1713+0747 is regarded as one of the highest-precision clocks in the sky and is regularly timed for the purpose of detecting gravitational waves. The International Pulsar Timing Array Collaboration undertook a 24 hr global observation of PSR J1713+0747 in an effort to better quantify sources of timing noise in this pulsar, particularly on intermediate (1–24 hr) timescales. We observed the pulsar continuously over 24 hr with the Arecibo, Effelsberg, GMRT, Green Bank, LOFAR, Lovell, Nançay, Parkes, and WSRT radio telescopes. The combined pulse times-of-arrival presented here provide an estimate of what sources of timing noise, excluding DM variations, would be present as compared to an idealized  $\sqrt{N}$  improvement in timing precision, where  $N$  is the number of pulses analyzed. In the case of this particular pulsar, we find that intrinsic pulse phase jitter dominates arrival time precision when the signal-to-noise ratio of single pulses exceeds unity, as measured using the eight telescopes that observed at  $L$  band/1.4 GHz. We present first results of specific phenomena probed on the unusually long timescale (for a single continuous observing session) of tens of hours, in particular interstellar scintillation, and discuss the degree to which scintillation and profile evolution affect precision timing. This paper presents the data set as a basis for future, deeper studies.

**Key words:** gravitational waves – ISM: structure – pulsars: individual (PSR J1713+0747)

**Online-only material:** color figures

### 1. INTRODUCTION

The International Pulsar Timing Array<sup>26</sup> (IPTA; Hobbs et al. 2010; Manchester & IPTA 2013) is a gravitational wave (GW) detector currently consisting of  $\sim 50$  pulsars distributed across the sky, monitored regularly by up to seven telescopes around the world: the Arecibo Observatory in the US, the Effelsberg radio telescope in Germany, the NRAO Green Bank Telescope (GBT) in the US, the Lovell radio telescope at Jodrell Bank Observatory in the UK, the Nançay radio telescope in France, the Parkes

telescope in Australia, and the Westerbork Synthesis Radio Telescope (WSRT) in the Netherlands. Some of the pulsars in the IPTA have been precision-timed for a decade or more. These observations are performed by the European Pulsar Timing Array (EPTA; Kramer & Champion 2013), the North American Nanohertz Observatory for Gravitational Waves (NANOGrav; McLaughlin 2013), and the Parkes Pulsar Timing Array (PPTA; Hobbs 2013; Manchester et al. 2013). The three collaborations combine their data as the IPTA.

Pulsar timing compares times-of-arrival (TOAs) to those predicted from a model that describes the pulsar's rotation, its binary motion, the interstellar medium (ISM) between us and the

<sup>26</sup> <http://ipta.phys.wvu.edu>

pulsar, and the Earth’s motion in the solar system. The measured TOAs are typically derived from pulsar profiles that have been averaged over the observation duration; and referenced against a high-precision frequency standard at the observatories (typically hydrogen masers); which in turn is referenced to an international timing standard (Lorimer et al. 2004). If the resulting differences between measured and modeled TOAs (the so-called timing residuals) deviate significantly from zero, this indicates astrophysical processes that are either not (or not completely) accounted for by the timing model. One possibility for such a process is long-period GWs perturbing the spacing between pulses as they propagate from a pulsar to the Earth. Obtaining accurate enough timing residuals to detect these GWs requires repeated measurements over many years. Sensitivity to GWs increases as observation duration grows, and the longest observation spans, as well as the red spectrum of the expected GWs, mean that the array is most sensitive at a frequency of about 10 yr (Sesana 2013). Individual TOAs are obtained by measuring the offset of emission beamed across the line-of-sight (LOS) at a given time from a template profile shape. Pulses can be averaged over a subintegration time, also known as “folding” according to a best-known pulse period. The template profile is high S/N and often averaged from long-term observations. From the radiometer equation relevant for pulsars in Lorimer et al. (2004) we have

$$S/N \propto \frac{G\sqrt{t_{\text{int}}\Delta f}}{T_{\text{sys}}} \quad (1)$$

in which S/N represents the integrated pulse S/N,  $G$  the telescope gain,  $t_{\text{int}}$  the pulse subintegration time,  $\Delta f$  the bandwidth, and  $T_{\text{sys}}$  the telescope’s system temperature. Thus, subintegration time, bandwidth, and gain are all important observational parameters, with  $G/T_{\text{sys}}$  most significantly impacting the reduction of radiometer noise, assuming we are comparing sensitivities for the same slice in frequency, and given that the telescopes are all equipped with receivers having state-of-the-art  $T_{\text{sys}}$  levels. (Throughout this paper, pulse S/N will refer to the ratio of the peak pulse amplitude to the standard deviation of the mean-subtracted off-pulse amplitudes.)

Pulsar timing arrays (PTAs) aim to detect perturbations due to GWs (Sazhin 1978; Foster & Backer 1990) in TOAs from millisecond pulsars (MSPs) on the order of 100 ns (Jenet et al. 2004) after the TOAs are corrected for many other effects. These include terrestrial clock calibration, solar system ephemeris, variations in dispersion measure (DM; proportional to the integrated LOS electron column density), proper motion, and position errors, all in the presence of noise due to other GWs at the source pulsars themselves. While the modeling uncertainties due to all the effects just listed are significant, GWs should perturb TOAs in a correlated way across the sky as a function of angle of separation between pulsars (Hellings & Downs 1983). This correlation makes the detection criterion less sensitive to any systematic errors in the TOAs or in the timing model for any one pulsar. Detectable strains (spatial strains due to GWs,  $h$ ) are expected to be on the order of  $h \sim 10^{-15}$  at nHz frequencies (Sesana 2013). Plausible sources producing GW strains in the PTA frequency range include: a stochastic background of GWs (Detweiler 1979, Hellings & Downs 1983) due to merging supermassive black hole binaries (SMBHBs), continuous wave sources from individual SMBHBs in  $z < 1.5$  galaxies (Sesana et al. 2009), bursts on timescales of months from SMBHBs in highly elliptical orbits (Finn & Lommen 2010), cosmic strings

(Starobinskiĭ 1979; Sanidas et al. 2013), phase transitions in the early universe (Caprini et al. 2010), and relic GWs from the era of inflation (Grishchuk 2005). Additionally, PTAs make possible the detection of GW bursts with memory, signals that are anticipated from events such as the final merger of SMBHBs and potentially from exotic phenomena at extremely high redshift (van Haasteren & Levin 2010, Cordes & Jenet 2012, Madison et al. 2014).

Through the IPTA consortium, all three PTAs (NANOGrav, the EPTA, and the PPTA) share timing data from their seven different observatories. The seven telescopes have different receivers, backend instruments, sensitivities, and radio frequency interference (RFI) environments, and have been observing their selected sets of pulsars for a range of epochs. Each telescope also has a history of regularly improving instrumentation, and thus TOAs obtained at later times are often of a much higher quality than those from earlier times. This trend is helpful for timing precision, but a wider bandwidth may require a more complicated frequency-dependent pulse profile model, due to frequency-dependent pulse shapes (Liu et al. 2014; Pennucci et al. 2014). Differences among PTAs include the number of standard observing frequencies and the methods for modeling DM variations. Fortunately, many of these difficulties in data combination are not insurmountable, and tremendous progress has already been made (see Manchester & IPTA 2013). The benefits of such a combination are many, and include an improved cadence, cross checks, better frequency coverage, and more pulsars correlated across the sky.

Apart from the need to combine data from many telescopes, there is also the need to better understand what might intrinsically limit timing quality. PSR J1713+0747 (Foster et al. 1993) is regularly observed by all IPTA telescopes, and provides much of the sensitivity for GW upper limit calculation (Arzoumanian et al. 2014) with a timing stability of  $\sim 100$  ns on timescales of five years or more (Verbiest et al. 2009). In contrast, the first MSP discovered, PSR B1937+21, is well known to be extremely stable on the order of weeks to months, but its residuals show a significant red noise power spectrum visible on timescales of years (Kaspi et al. 1994). As we design larger telescopes and observing programs, it is imperative that we know the fundamental limits of timing precision, i.e., at what point additional gain, observing time, or bandwidth will not increase our timing precision.

Upper limit papers such as those by Shannon et al. (2013), van Haasteren et al. (2011), and Demorest et al. (2013) have all calculated GW limits based on TOAs over 5+ yr. If the observation duration at a single epoch were increased from the typical subintegration time by a factor of  $X$ , the timing precision (in the absence of other limiting effects) would be naively expected to improve by  $\sqrt{X}$  as in Equation (1). This is simply due to the fact that the number of pulses collected would increase, if TOAs (in the absence of GWs) can be fitted to standard timing models such that the residuals are white noise, assuming there are no significant pulse shape changes between observation epochs.

Many properties of a pulsar along its particular LOS are not precisely predictable: DM variations, interstellar scintillation (ISS), scattering variations, and low-level glitches, to name a few. Glitches are not observed in PSR J1713+0747 or in MSPs in general (Espinoza et al. 2011), though they may be present at low amplitudes in many pulsars (see, however, Espinoza et al. 2014) and therefore may act as a limiting factor in searches for GWs; they may be especially problematic sources of noise

**Table 1**  
Observing Schedule and Parameters

Telescope	Obs. Mode	Start UT	End UT	Min. Subint. Length (s)	Cent. Freq. (MHz)	Bandwidth (MHz)	Min. Channel Width Possible (MHz)	Pulse Profile S/N <sup>a</sup> for Obs. Duration <sup>b</sup>	Time to Stated S/N (hr) <sup>c</sup>
Arecibo	Int <sup>d</sup>	Jun 23 03:17	Jun 23 04:44	...	1382	800	6.25	3138	1.45
Effelsberg	Fold <sup>e</sup>	Jun 22 18:26	Jun 23 03:40	10	1348	200	1.56	473	7.20
GBT	Int	Jun 22 00:56	Jun 23 10:15	...	1497	800	6.25	2200	8.95
GMRT	CF <sup>f</sup>	Jun 22 13:38	Jun 22 21:58	60	1387	33.3	1.04	80	7.63
Lovell	Fold	Jun 22 18:04	Jun 23 05:34	10	1532	400	0.25	404	9.85
Nançay	Fold	Jun 22 22:33	Jun 22 23:30	60	1524	512	16	125	0.95
Parkes	Fold <sup>g</sup>	Jun 22 10:20	Jun 22 16:20	60	1369	256	0.25	344	6.00
LOFAR	CF	Jun 22 18:11	Jun 23 03:00	5	148.9	78.1	0.195	8	8.82
Westerbork	CF	Jun 22 21:46	Jun 23 04:39	10	345 <sup>h</sup>	80	8.75	27	4.95

#### Notes.

<sup>a</sup> All S/N values are scaled to 512 phase bins. These S/N values are affected by both scintillation and observation length as well as the telescope parameters.

<sup>b</sup> The duration here refers to the duration of the folding portions of a telescope’s observation run only.

<sup>c</sup> In addition to observation time, scintillation also significantly influenced these S/N values.

<sup>d</sup> Intensity recording mode (nonfolding). Single pulses are the minimum subintegration time.

<sup>e</sup> Normal pulse folding mode, using coherent dedispersion.

<sup>f</sup> Offline coherent filterbank mode. For these multi-antenna telescopes acting as a phased array, folding and coherent dedispersion is applied offline.

<sup>g</sup> The Parkes DFB3/4 backend, one of three backends used in parallel, does not apply coherent dedispersion online.

<sup>h</sup> Observations alternated between 345 MHz and 1398 MHz as center frequencies. Refer to Figure 1 and Table 2 for details.

in searches for GW bursts with memory. M28A is another exception (Cognard & Backer 2004), though admittedly this may be because it is particularly young for an MSP.

Pulse phase jitter, which is independent of radio frequency, is also a limiting factor for pulsar timing. Jitter, also known as pulse-to-pulse modulation, was first described in Cordes & Downs (1985) for canonical pulsars, established for PSR B1937+21 by Cordes et al. (1990), and more recently measured in PSR J1713+0747 in Shannon et al. (2014) and Shannon & Cordes (2012). The term refers to the distribution of arrival times of single pulses about the peak of the averaged template pulse, which have slight offsets in pulse phase. Neither increasing telescope size nor bandwidth eliminates the presence of TOA jitter errors, jitter being both broadband and independent of pulse profile S/N. Generally, the only way to reduce jitter-induced white noise in most pulsars is to increase the observation duration (Cordes et al. 2004), though there are exceptions; see Osłowski et al. (2013) and Osłowski et al. (2011) for an example of a jitter mitigation technique on PSR J0437–4715.

These considerations motivate observations of PSR J1713+0747 for 24 continuous hours, using nine radio telescopes: the seven IPTA telescopes along with LOFAR (LOW Frequency ARray; van Haarlem et al. 2013, Stappers et al. 2011) in the Netherlands and the GMRT (Giant Meterwave Radio Telescope) in India. See Table 1 for details of the allotted frequencies for each telescope. The duration of 24 hr was chosen because MSP timing has been explored at the hour and week timescales, but not in the intermediate regime. The inclusion of LOFAR provided an ultra-low frequency (110–190 MHz) component to the observation, sampling a frequency range that features prominent effects from the interstellar medium. Observing at *L*-band/1.4 GHz is ideal for studying the timing properties of this particular pulsar, which is reasonably bright given its flux spectrum with a power-law of slope of  $-1.7$  (Kuzmin & Losovsky 2001), but not significantly affected by red noise in the timing residuals due to the ISM (Keith et al. 2013). The GMRT filled in the time coverage gap between Parkes and the European telescopes, enabling a continuous 24 hr of observing.

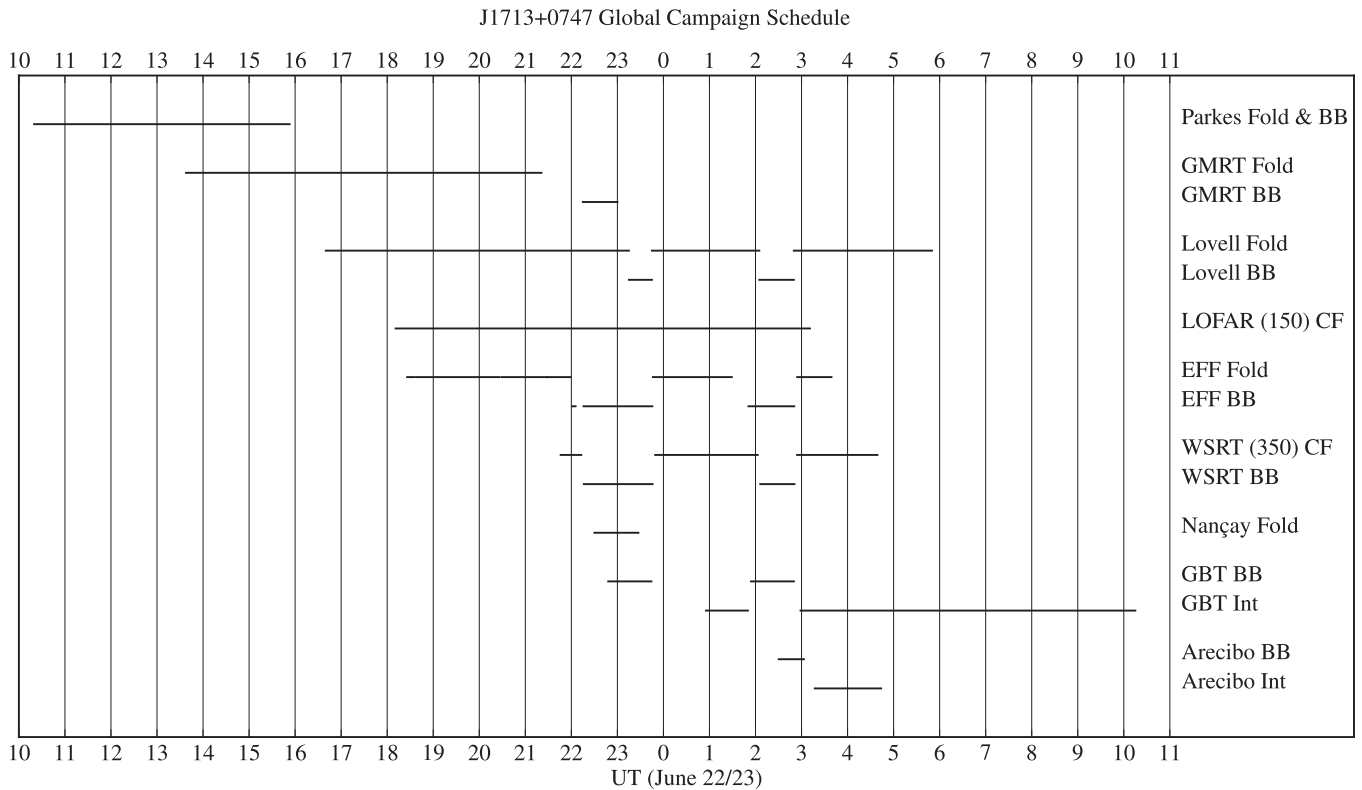
Another major science goal of the global PSR J1713+0747 observation relates to the LEAP project (Large European Array for Pulsars; C. Bassa et al. 2014, in preparation; Kramer & Champion 2013), which uses the Effelsberg, Nançay, Lovell, WSRT, and Sardinia radio telescopes, together as sensitive as Arecibo, as a phased array, with a comparable total collecting area ( $3 \times 10^4 \text{ m}^2$ ), but with a much greater observable declination range than Arecibo. As PTAs advance, this configuration may prove to be critical for the detection of GWs. The 24 hr global observation of PSR J1713+0747 helps the LEAP effort by adding three more telescopes—Arecibo itself, as well as the Green Bank Telescope and the GMRT. The present data set therefore opens up the possibility of experimenting with a telescope having over twice the collecting area of Arecibo alone. The combined effort will be referred to in the present paper as GiantLEAP.

This data set also represents a unique opportunity to measure clock offsets between telescopes and how they vary across overlapping time intervals. When combining TOAs from different telescopes (or when backends change on a single telescope) an offset or “jump” is needed. Such a jump can be due to delays in the backends themselves, such as cable delays, conspiring with factors that are difficult to quantify individually (see Lommen & Demorest 2013 and Kramer & Champion 2013 for further explanation). The jumps can then be quantified by fitting for one arbitrary timing offset per telescope/backend pair per frequency only in the overlapping region, such that the rms of the combined data set is minimized. Simultaneous observations—the longer the better—provide an opportunity to measure such offsets and their drifts with high accuracy.

Timing stability on the  $\sim 24$  hr timescale can also be quantified using the Allan variance of the residuals. The Allan variance was originally used to quantify the stability of atomic clocks (see Matsakis et al. 1997 for details). The present data set allows us to evaluate the Allan variance for clock frequencies of  $10^{-2}$  Hz all the way into frequencies corresponding to the 5 or even 20 year data sets (W. W. Zhu et al. 2014, in preparation) that exist for PSR J1713+0747.

Yet another goal of the PSR J1713+0747 24 hr global campaign is to assess a noise floor. How does rms precision increase





**Figure 1.** Timeline of the global observation of PSR J1713+0747 showing the various telescopes and observing modes. Here, “Fold” refers to ordinary pulse folding with coherent dedispersion. “BB” refers to baseband mode at  $L$  band/1.4 GHz, recording complex voltages without folding. “Int” refers to intensity integrations, also known as “coherent search mode,” which is similar to baseband mode except recording an intensity rather than full voltage information. “CF” refers to the offline coherent filterbank mode used at LOFAR and WSRT, in which the telescope recorded a coherent sum (tied-array beam) of all the antennas, which was written out as complex voltages and then coherently dedispersed and folded offline.

as longer timespans of data are analyzed? Does the improvement “bottom out” or continue indefinitely with the number of collected pulses on these timescales?

We intend the present paper to be a description of the data set itself as well as being an introduction to a series of papers, given the size of the data set and the large number of science goals. Here we present some first science results, and intend to expand upon them and address other topics in later papers. In Section 2, we describe the observation and the resulting data set in detail, supplemented by the Appendix. In Section 3, we explore a number of first results emerging from the all-telescope analysis. Section 4 shows some general first results on the timing error budget, and finally Section 5 mentions some future paper topics based on the data.

## 2. DESCRIPTION OF THE DATA

The observations were conducted on 2013 June 22 (MJD 56465–56466) starting with the Parkes Telescope and progressing through the other eight telescopes for as much time as possible between rise and set. The observation timeline can be seen in Figure 1. The time of year was such that local midnight roughly corresponded to the middle of each telescope’s observation. PSR J1713+0747 was approximately  $146^\circ$  from the Sun, minimizing the possibility of any solar effects on the data (You et al. 2007).

As Figure 1 shows, there were different modes used among the telescopes—ordinary pulse folding, baseband mode, coherent filterbank mode (formed offline), and nonfolded intensity integrations (also known as “coherent search mode”). For ordinary pulse folding, coherent dedispersion is applied in real

time (Hankins & Rickett 1975), correcting for pulse delays due to dispersion in the ISM across few-MHz channels (with the exception of the DFB backends at Parkes in which dedispersion is applied after the fact—see Appendix). All the nonbaseband data presented here have coherent dedispersion applied. Observing in baseband mode affords a number of advantages. For the science goal of expanding LEAP with Arecibo and GBT, baseband recording is a requirement of any phased array formed offline. It also allows us to evaluate how cyclic spectroscopy (Demorest 2011; Walker et al. 2013; Stinebring 2013) might improve the quality of some of the data sets as a method of obtaining a significantly higher frequency resolution (Archibald et al. 2014; G. Jones et al. 2014, in preparation). The baseband sessions were driven by when the transit telescopes (Arecibo and Nançay) could observe the pulsar. Hence, the first baseband session was when the source transited at Nançay, when baseband data were obtained at GMRT, Lovell, Effelsberg, Westerbork, GBT and Nançay, while the second session was when the source transited at Arecibo, and baseband data were obtained with the Lovell, Effelsberg, Westerbork, GBT, and Arecibo telescopes. Nonfolded intensity integrating is similar to baseband mode, except that intensities rather than signed voltages are recorded, without phase information, resulting in a more manageable data size. Both baseband data and nonfolded intensity recording yield single pulse information. Otherwise, single pulses are not recoverable due to the folding process.

Baseband was not, in general, the default observing mode. Taking baseband data for the entire rise-to-set time at each telescope would be cumbersome in terms of data volume (which as it is approached 60 TB total, largely due to the baseband

**Table 2**  
Baseband Observing Parameters

Telescope	Obs. Mode	Start UT	End UT	Cent. Freq. (MHz)	Bandwidth (MHz)
Arecibo	BB <sup>a</sup>	Jun 23 02:30	Jun 23 03:00	1378	200
Effelsberg	BB	Jun 22 22:15	Jun 22 23:45	1396	128
	BB	Jun 23 02:05	Jun 23 02:50	1396	128
GBT	BB	Jun 22 22:15	Jun 22 23:15	1378	200
	BB	Jun 23 01:50	Jun 23 02:52	1378	200
GMRT	BB	Jun 22 22:24	Jun 23 00:00	1387	33.3
Lovell	BB	Jun 22 22:14	Jun 22 23:44	1396	128
	BB	Jun 23 01:47	Jun 23 02:56	1396	128
Parkes	BB	Jun 22 10:20	Jun 22 16:20	1369	256
Westerbork	BB	Jun 22 22:16	Jun 22 23:45	1398	128
	BB	Jun 23 02:05	Jun 23 02:51	1398	128

**Note.** <sup>a</sup> Baseband voltage recording mode.

component) and in some cases would also limit bandwidth. For example, in the case of the GUPPI/PUPPI backends (Green Bank/Puerto Rican Ultimate Pulsar Processing Instrument; DuPlain et al. 2008), baseband recording can only be conducted over 200 MHz of bandwidth, as opposed to the 800 MHz available for the folding and intensity integration modes. The overlapping baseband portion of the observation for GiantLEAP consisted of two baseband allotments (see Figure 1). Arecibo could only participate in one of these sessions due to the limited range in zenith angle. Parkes observed with baseband data taking in parallel for its entire observation duration.

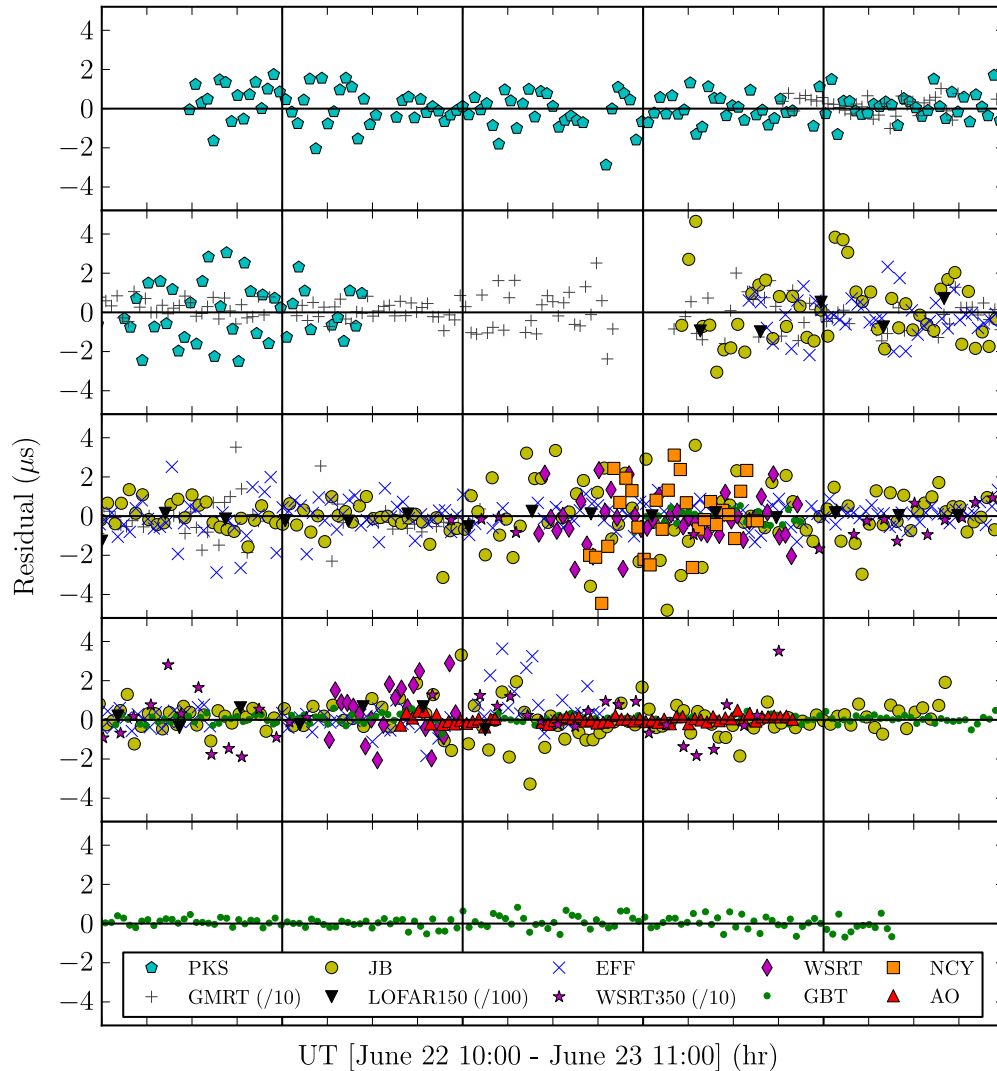
Reduced profiles are in the PSRFITS format processed with the PSRCHIVE and DSPSR software packages (Hotan et al. 2004; van Straten et al. 2012), with final timing residuals determined by the TEMPO2 software package (Hobbs et al. 2006). Residuals are generated with a common TEMPO2 parameter file, after having produced TOAs with PSRCHIVE based on standard observatory-specific template pulse profiles, noise-reduced using PSRCHIVE when necessary. Different profiles are used because of local bandpass, frequency range, and calibration differences. We create TOAs from 120 s fiducial subintegrations in common between *L*-band/1.4 GHz telescopes. This choice of subintegration time was a compromise between, on the one hand, having a sufficient number of TOAs in order to probe timing precision on long timescales (see Section 3.1), and on the other, having a minimum TOA S/N of approximately 1 across all *L*-band/1.4 GHz telescopes (see Section 3.2 for an application to measuring jitter). Further details about specific telescopes can be found in the Appendix and in Table 1, in which the S/N of each telescope’s 120 s TOAs can be found, with baseband details in Table 2. As can be seen in Table 1, the possible minimum subintegration length is less than 120 s for most telescopes, providing flexibility for future studies. Slots in which the table is blank indicate intensity integrations, which can be integrated to any time greater than  $2/\Delta\nu$ , where  $\Delta\nu$  is the channel width. WSRT, GMRT, and LOFAR, being multi-antenna telescopes, observed in a tied-array mode (formed online), which dumped complex voltage data to be coherently dedispersed and folded using DSPSR. We refer to this as the “coherent filterbank” mode because a filterbank is formed offline, except for GMRT, for which we still refer to the modes used as folding and baseband, given the additional optimization employed (see the Appendix for details).

We remove RFI automatically in post-processing using PSRCHIVE. The program defines an off-pulse window by

iteratively smoothing the profile and finding the minimum. Using the maximum minus minimum intensity values defined in this window, we apply a median filter to identify RFI spikes and flag spikes at greater than  $4\sigma$  from half the intensity difference. Any remaining particularly noisy frequency channels or time integrations are manually removed after visual inspection. At Lovell, RFI is excised in real time. For the other telescopes, post-facto algorithms are used—one that excises “rows” and “columns” of RFI in frequency and time, and another that excises RFI within the pulse profile itself according to bad phase bins.

The residuals shown in Figure 2 were generated by folding the respective data sets from each individual telescope with a parameter file that was created from 21 yr of PSR J1713+0747 data (W. W. Zhu et al. 2014, in preparation), all with the same Bureau International des Poids et Mesures (BIPM) correction table to the global atomic timescale, TT (BIPM2012), with the 2013 extrapolation, and the Jet Propulsion Laboratory DE421 planetary ephemeris. The DM ( $15.99113 \pm 0.00001$  pc cm<sup>−3</sup>) is the only fitted astrophysical parameter determined from the present observation across all *L*-band/1.4 GHz telescopes simultaneously. Relative offsets were also fitted between telescopes. These are simply free parameters that align the residuals and do not represent absolute clock offsets, and are not shown here for this reason. The remainder of the parameters were held fixed at the W. W. Zhu et al. (2014, in preparation) values (see Table 3 in Appendix for further details). The timing residuals at other, higher subintegration times are computed from this base set of residuals. To test that residuals can be averaged down without a loss in modeled timing precision, we also generated a set of 10 s subintegration time TOAs with PSRCHIVE from the GBT, and found that simply averaging the 10 s residuals produced new residuals with rms values different from the 120 s TEMPO2 residuals at  $\ll 1\sigma$ . Therefore, we can create 10 s residuals, average them, and obtain nearly identical 120 s residuals to those resulting from 120 s TOAs.

The eight telescopes at *L*-band/1.4 GHz saw a changing spectrum due to interstellar scintillation, as shown in the dynamic spectrum, or plot of pulsar intensity versus time and frequency, in Figure 3. Scintillation is due to the scattering and refraction of pulsed emission in the ionized ISM, and can significantly change the pulse profile S/N as can be seen in Table 1. Figure 3 shows that bright scintles contribute to the high S/N of telescopes after about June 23 00:00. The image was created by fitting the telescopes’ templates via matched filtering for a given telescope with a given profile  $P(\nu, t)$  and reporting the amplitude with the off-pulse mean subtracted. Amplitudes are converted to corresponding intensities depending on each telescope’s calibration data and then scaled to an identical color scaling for Figure 3. Dynamic spectra from different telescopes were scaled empirically to match them in Figure 3, lacking absolute calibration for some telescopes, but using the noise diode calibrations employed (see Appendix for details). A cross-hatched pattern is clear. Such a pattern is usually only observable at low frequencies due to the small scintle size in frequency and time even for modest bandwidths (Rickett 1970). Note the narrowing of the scintillation bandwidth ( $\Delta\nu$ , the typical scintle width in frequency) as a function of frequency, along with a decrease in scintillation timescale ( $\Delta t$ , the typical timescale for scintillation) with decrease in frequency. Scintillation patterns are shown to overlap well between telescopes, thus establishing a typical spatial scale of a waveform—see Section 3.3. In some cases there appear to be deviations between the scintillation



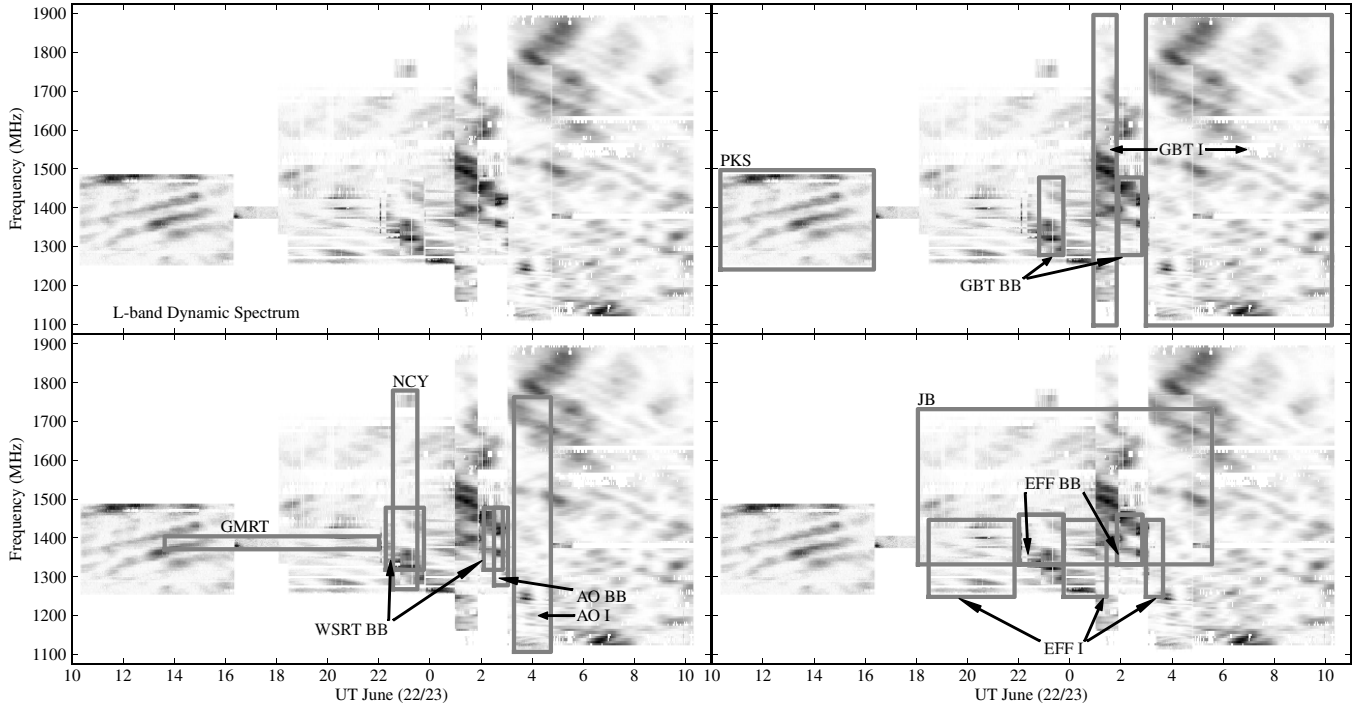
**Figure 2.** 24 hr timing residuals: timing residuals for the nine telescopes as a function of time. Values shown are for  $L$ -band/1.4 GHz observations unless otherwise noted. Uncertainties on each residual are not shown here in order to maintain clarity. In all cases, the fitting error on individual residuals is on the order of the scatter of all residuals shown for a particular telescope. All residual values are for 120 s integrations, except LOFAR, which is for 20 minutes. The increase in residual values in the third and fourth rows from the top is due to most telescopes having switched to the lower bandwidth baseband mode. Residual values from the GMRT, WSRT-350 MHz, and LOFAR-150 MHz are scaled down by factors of 10, 10, and 100 respectively in order to show all residuals in a single panel. (A color version of this figure is available in the online journal.)

patterns seen at different telescopes, but this always corresponds to times when the source was close to rising or setting at one of the telescopes involved.

Narrowband template fitting (Taylor 1992) assumes a relatively constant profile with frequency. In addition to distorting the pulse phase, merely averaging across frequency would result in drifting residuals with a nonwhite appearance (Craft 1970), due to the intrinsic profile evolution acting in combination with ISS. Profile shape changes with frequency are present in all canonical pulsars—see Hankins & Rickett (1986) for multifrequency observations on many pulsars, and Hassall et al. (2013), who uses observations from LOFAR and other telescopes. Similar shape changes have also been found in MSPs (Kramer et al. 1999), including PSR J1713+0747.

Figure 4 shows the presence of profile evolution with frequency in this observation’s GBT data. Starting with 8 hr of GBT data, we use the fiducial subintegration length of 120 s and a subband size of 50 MHz. We sum profiles in time to get 16 profiles as a function of observing frequency and phase,

$P(\nu, \phi)$ . These profiles are dedispersed using the best fit,  $L$ -band/1.4 GHz DM. For each  $P(\nu, \phi)$ , we fit a NANOGrav standard template  $T(\phi)$  to the data profile  $P(\nu, \phi)$  to find the best-fit phase offset  $\delta\phi(\nu)$  and amplitude  $A(\nu)$ . We create difference profiles by shifting and scaling  $T(\phi)$  using the best fit phase offset  $\delta\phi(\nu)$  and amplitude  $A(\nu)$  for each frequency and then subtracting as  $D(\nu, \phi) = P(\nu, \phi) - A(\nu)T(\phi - \delta\phi(\nu))$ , where  $D(\nu, \phi)$  are the difference profiles. These are plotted in the main panel of Figure 4. The right panel shows the timing offsets as a function of frequency,  $\delta\phi(\nu)$ , with the 1422 MHz profile set to zero offset because the template most closely resembles these data in the center of the band. These timing offsets will be dependent on the value of our measured DM, which in turn is dependent on the frequency-dependent (FD) model parameters (Z. Arzoumanian et al. 2014) used in TEMPO2. The FD parameters correspond to the coefficients of polynomials of the logarithm of radio frequency that show the TOA shift due to profile evolution. If the profile evolution within a subband is small and if the FD model parameters quantify the offsets well



**Figure 3.** Dynamic spectrum: L-band/1.4 GHz dynamic spectrum for eight of the nine telescopes (excluding LOFAR at 150 MHz and WSRT at 350 MHz). Note the narrowing of the scintillation bandwidth as a function of frequency, along with a decrease in scintillation timescale with decrease in frequency. Scintillation patterns are shown to overlap well between telescopes, thus establishing (in Section 3.3) a typical spatial scale for the scintillation pattern. A cross-hatched pattern is clear, which is usually only observable at low frequencies, but observable here because of the large bandwidth and 24 hr observation time. When observations overlap, the regions of the dynamic spectrum shown are from telescopes with the highest sensitivity per patch  $dA = dvdt$ .

**Table 3**  
Timing Model Parameters

Parameter	Value	Held Fixed in Fit? <sup>a</sup>	Parameter Uncertainty
Right Ascension, $\alpha$ (J2000)	17:13:49.5331497	Y	$5 \times 10^{-7}$
Declination, $\delta$ (J2000)	07:47:37.492844	Y	$1.4 \times 10^{-5}$
Proper motion in $\alpha$ , $\mu_\alpha$ (mas yr <sup>-1</sup> )	4.922	Y	0.002
Proper motion in $\delta$ , $\mu_\delta$ (mas yr <sup>-1</sup> )	-3.909	Y	0.004
Parallax, $\pi$ (mas)	0.88	Y	0.03
Spin Frequency (Hz)	218.81184381090227	Y	$7 \times 10^{-14}$
Spin down rate (Hz <sup>2</sup> )	$-4.083907 \times 10^{-16}$	Y	$8 \times 10^{-22}$
Reference epoch (MJD)	54971	Y	
Dispersion Measure (pc cm <sup>-3</sup> )	15.99113	N	$1 \times 10^{-5}$
Profile frequency dependency parameter, FD1	$1.328 \times 10^{-5}$	Y	$4 \times 10^{-8}$
Profile frequency dependency parameter, FD2	$-3.73 \times 10^{-5}$	Y	$2 \times 10^{-7}$
Profile frequency dependency parameter, FD3	$3.24 \times 10^{-5}$	Y	$7 \times 10^{-7}$
Profile frequency dependency parameter, FD4	$-1.07 \times 10^{-5}$	Y	$5 \times 10^{-7}$
Solar system ephemeris	DE421	Y	
Reference clock	TT(BIPM)	Y	
Binary Type	T2 <sup>b</sup>	Y	
Projected semi-major axis, $x$ (lt-s)	32.34242245	Y	$1.2 \times 10^{-7}$
Eccentricity, $e$	$7.49414 \times 10^{-5}$	Y	$6 \times 10^{-10}$
Time of periastron passage, $T_0$ (MJD)	54914.0602	Y	0.0003
Orbital Period, $P_b$ (day)	67.825147	Y	$5 \times 10^{-6}$
Angle of periastron, $\omega$ (deg)	176.1978	Y	0.0015
Derivative of periastron angle, $\dot{\omega}$ (deg)	0.00049	Y	0.00014
Companion Mass, $M_c$ ( $M_\odot$ )	0.29	Y	0.01

**Notes.**

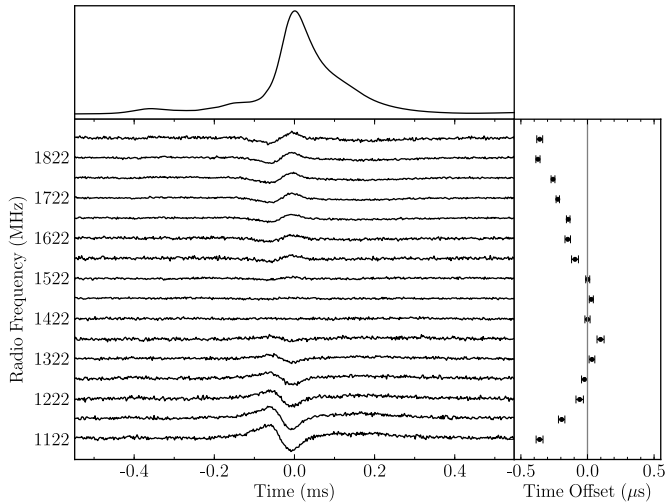
<sup>a</sup> We also fit for arbitrary jumps between telescopes, which are not astrophysical and not shown here.

<sup>b</sup> Damour & Deruelle (1986).

when each subband is independently used to create a set of timing residuals, then the weighted broadband residuals should be consistent with the white noise expected from scintillation. As Figure 2 shows for all the telescopes, the residuals are qualita-

tively white-noise-like in character. The broadband weighting, then, appears to correctly take the profile evolution/scintillation interaction into account. The profile evolution shown with frequency in Figure 4 is likely to be intrinsic and not an instrumental



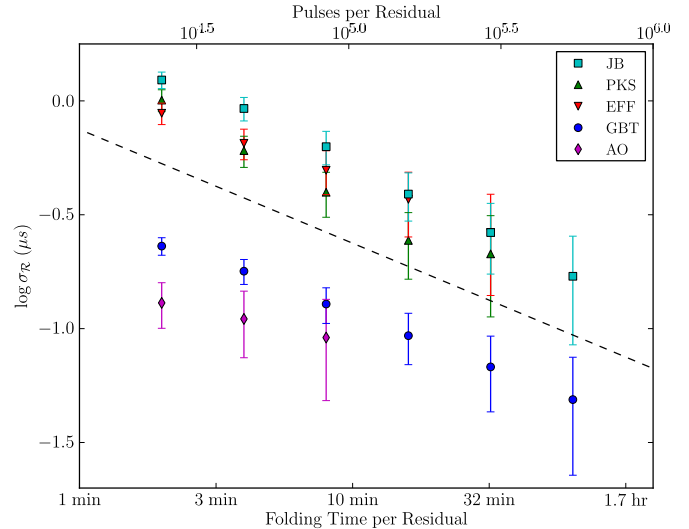


**Figure 4.** Differential offset times vs. frequency for the GBT. The top panel shows the standard NANOGrav template used for GUPPI at  $L$  band/1.4 GHz. The main panel shows the difference profiles as a function of frequency, calculated by subtracting a best-fit template from the data profiles. The right panel shows the mean-subtracted, best-fit phase offsets for each data profile vs. frequency. These offsets are a function of the FD (frequency dependent) polynomial parameters in TEMPO2 that model timing offsets due to pulse profile frequency evolution. The shape of the time offsets vs. frequency curve is covariant with any residual dispersion delay across the band.

artifact because the equivalent Arecibo plot (i.e., using a different receiver at a different telescope) is nearly identical across the same bandwidth. A more detailed analysis of the observed profile frequency evolution is a subject of future work.

We address this profile-evolution problem for all telescopes with bandwidths of 100 MHz or more (that is, all telescopes except the GMRT) by computing TOAs for multiple narrow frequency channels, using PSRCHIVE as described above. For the GBT and for Arecibo, the band is divided into 16 bins of 50 MHz each. The data from the other telescopes are split into subbands in similar fashion. We then obtain sets of narrowband timing residuals using the FD parameters in TEMPO2. The four best-fit FD parameters can be found in Table 3 in the Appendix, representing third-order polynomial coefficients starting with the lowest order first. We then perform a weighted mean of these values in order to obtain the broadband residuals. For all telescopes we use the FD parameters with TEMPO2 independently computed from W. W. Zhu et al. (in preparation).

Intrinsic pulse profile evolution is thought to arise from varying offsets between the emission region and the surface of the neutron star, with higher frequency emission being produced closer to the surface (see Cordes 2013 for a more detailed discussion). If a profile at a high narrowband frequency differs significantly from a profile at a low narrowband frequency, then any frequency-dependent pulse shape changes will be highly covariant with the DM measurements at each epoch. Multi-frequency timing minimizes such covariances. The timing offsets due to intrinsic pulse shape changes with radio frequency are constant in time. Effects due to interstellar scintillation and scattering will depend on time, however. The former produces a varying S/N across the band due to scintillation that changes the relative weighting of each subband as part of the final TOA; the latter broadens the pulse, resulting in a scattering delay. (The  $L$ -band/1.4 GHz frequency is chosen for the present data set and for most standard timing observations, so that scatter broadening is minimal.)



**Figure 5.** Improvement of  $L$ -band/1.4 GHz timing rms with number of collected pulses  $N$ . Both subintegration time per residual and number of pulses are shown in the abscissa for reference. Symbols, from top to bottom: teal squares, Lovell/JB; green triangles, Parkes; red inverted triangles, Effelsberg; blue circles, GBT; magenta diamonds, Arecibo. Error bars are the standard error of the scattered rms values. The dashed lines show a  $1/\sqrt{N}$  law for reference. The residuals are derived from telescopes with different bandwidths, and the resulting timing rms values are dependent both on collecting area and bandwidth.

(A color version of this figure is available in the online journal.)

### 3. ANALYSIS OF THE MULTI-TELESCOPE DATA

The following assumptions and terminology will be used throughout this section.

The rms of the timing residuals over the total time span is consistent with the errors expected from a finite S/N ratio and from single-pulse stochasticity that is intrinsic to the pulsar. From Cordes & Shannon (2010), finite S/Ns yield an approximate template fitting error of

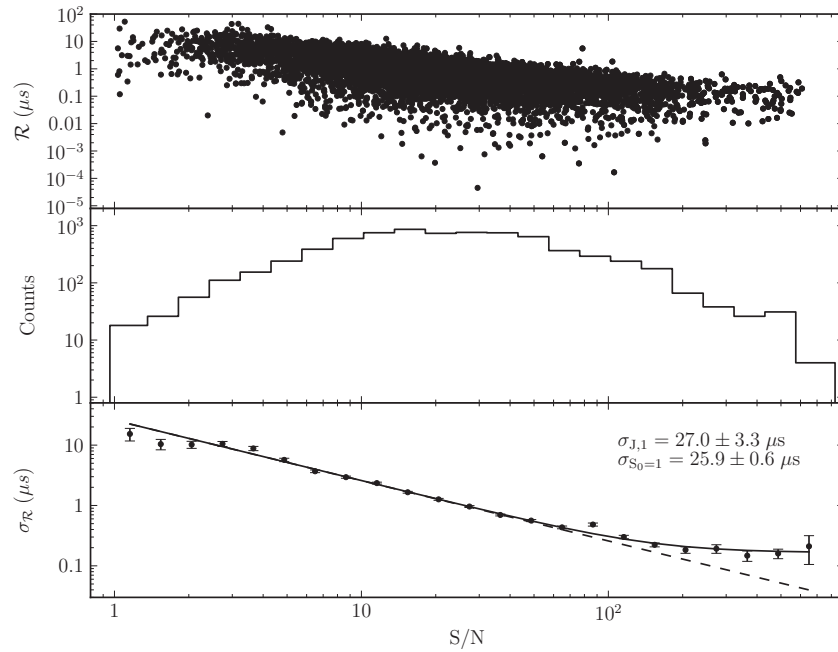
$$\sigma_{S/N} = \frac{W_{\text{eff}}}{S(N_\phi)\sqrt{N_\phi}}, \quad (2)$$

where  $S(N_\phi)$  is the S/N of the pulse profile (peak to rms off-pulse) that has  $N_\phi$  phase bins and  $W_{\text{eff}}$  is the effective pulse width. Cordes & Shannon (2010) give an expression for  $W_{\text{eff}}$  that we use for PSR J1713+0747, yielding 0.54 ms.

Here,  $N_\phi$  is included because for the GMRT,  $N_\phi$  was 64, while for the other telescopes,  $N_\phi$  was 512, and this difference has been noted in all relevant calculations. These values for  $N_\phi$  are chosen such that  $N_\phi$  is small enough to afford submicrosecond timing precision at some telescopes, while at the same time, producing a  $S(N_\phi)$  small enough such that the pulse peak measurement is reasonably accurate. When we use  $\sigma_R$ , it will refer to the total residual rms, whether template fitting error, jitter, radiometer white noise, or white noise due to the ISM.

#### 3.1. Timing Residual Precision versus Integration Time

In Figure 5 we show the logarithmic change of the  $L$ -band/1.4 GHz TOA residual rms, a proxy for timing precision, as both a function of subintegration time  $T$  and of the corresponding number of pulses  $N$ . The time per TOA is plotted on the abscissa, and the corresponding number of pulses for each residual subintegration time is also shown. The ordinate shows the rms of the residuals within the entire observation time of the telescope. For this reason, the data points for Arecibo do not extend



**Figure 6.** Improvement of residual rms with S/N, for  $\sim 50$  MHz timing residuals. Shown in the top plot are the timing residual values as a function of S/N, using data from all telescopes. The middle plot shows the number of residuals in each bin. In the bottom plot, rms values on the residuals are shown using logarithmic bins with eight bins/decade. We fit Equation (3), shown as the solid line fit, which yields a white noise in the timing residuals due to pulse phase jitter of  $\sigma_{J,1} = 27.0 \pm 3.3 \mu\text{s}$ . Scaled to 120 s integrations,  $\sigma_J = 0.17 \pm 0.02 \mu\text{s}$ . All residuals shown in the top panel of Figure 6 are for an integration/folding time of 120 s, removing residuals below a S/N of 1. The single pulse jitter timing rms is  $\sigma_{J,1}$ , and  $\sigma_{S_0=1}$  is the timing rms in the absence of jitter for a S/N of 1. The dashed line represents the expected timing uncertainties in the absence of pulse phase jitter.

to as long of a timescale as the other telescopes despite the high sensitivity. We start from TOAs from the base subintegration time of 120 s and integrate down (i.e., to larger subintegrations) for each successive step. We show this function for five of the *L*-band/1.4 GHz telescopes, choosing the maximum integration time at each telescope which corresponds to at least eight subintegrations in order to ensure that small-number statistics (due to having only a few long-subintegration residuals) are not important. Nançay and WSRT are not shown because of their short, noncontiguous observing times. GMRT’s residuals did not probe small values of  $N$  and are not shown. Error bars are  $1\sigma$  and are simply the standard error of the scattered subintegration rms values in a block of TOAs. For reference, the expected  $1/\sqrt{N}$  slope is plotted. Each successive data point is not independent of the data points in Figure 5 for small values of  $N$ . In all telescopes, there are no significant deviations from this simple improvement in timing rms with number of pulses collected. This is expected behavior for the backends used in this observation (see Stairs et al. 2002, which uses data from PSR B1534+12 as a demonstration that this kind of integrating-down behavior works efficiently for coherent dedispersion machines in contrast to filterbank machines). For PSR J1713+0747, this means that on timescales of  $\sim 1$  hr (the largest time which on this timescale we can make multiple samples with a minimum of eight TOAs), there is no significant evidence of an absolute noise floor.

The comparative sensitivity of the telescopes can be seen along with the fact that longer tracks produce smaller uncertainties in the residual rms  $\sigma_{\mathcal{R}}$ . Arecibo and GBT have up to eight times the bandwidth as some of the other telescopes and, considering also the sensitivities, fall significantly beneath the others in terms of timing rms. Previous studies of PSR J1713+0747 (Shannon & Cordes 2012) have shown a tracking of this  $1/\sqrt{N}$

for an  $N$  of  $1\text{--}10^5$ , corresponding to subintegration times of 4.57 ms to 457 s.

### 3.2. Timing Residual Precision From Template Fitting and Pulse Jitter

Radiometer noise is always reduced by additional bandwidth but jitter noise, measurable in high S/N timing observations, is not improved because it is identical and correlated across frequency. Thus, the minimum expected rms in the broadband timing residuals is somewhat higher than what we would predict from merely reducing the narrowband timing rms by the increased bandwidth factor. The error from jitter, given in Cordes & Shannon (2010), is  $\sigma_J \propto W_{\text{eff}}/\sqrt{N}$ . The noise value  $\sigma_J$  has no dependence on bandwidth or telescope sensitivity. Knowing how much the rms for an individual TOA would be composed of  $\sigma_J$  typically requires either directly measuring the jitter via single pulses, or measuring the correlated TOAs across frequency. Shannon & Cordes (2012) report a value of  $26 \mu\text{s}$  for a single-pulse  $\sigma_J$  from PSR J1713+0747 also based on Arecibo observations, using both measurement methods. Shannon et al. (2014) have also measured PSR J1713+0747’s single pulse phase jitter rms as  $31.1 \pm 0.7 \mu\text{s}$ . Here, we show the presence of jitter in PSR J1713+0747 by demonstrating the nondependence of a noise component on telescope sensitivity.

In order to directly measure the presence of jitter in Figure 6, we binned residuals from *all telescopes* into approximately eight bins/decade in S/N and then took the scatter of the arrival times within each bin,  $\sigma_{\mathcal{R}}$ . The subintegrations used were 120 s within frequency bins of  $\sim 50$  MHz, and only residuals for which  $S/N > 1$  were included.

We fit a curve given using the following equation, which comes from the assumption that the white noise timing residuals,  $\sigma_{\mathcal{R}}$ , are composed of two other white noise components added

in quadrature:

$$\sigma_{\mathcal{R}} = \sqrt{\sigma_J^2 + (\sigma_{S_0}(S_0/S))^2} \quad (3)$$

where  $S$  simply represents  $S/N$ ,  $S_0$  refers to a particular fiducial  $S/N$ , and  $\sigma_{S_0}$  is the timing rms due to template fitting. Using the fiducial 120 s TOAs, we find that  $\sigma_J = 0.17 \pm 0.02 \mu\text{s}$ , which implies a single pulse jitter of  $27.0 \pm 3.3 \mu\text{s}$ , consistent with the measurement in Shannon & Cordes (2012) of  $26 \mu\text{s}$ . This value is also consistent within  $<2 \sigma$  of the more recent measurement of Shannon et al. (2014). The other fitted parameter,  $\sigma_{S_0=1}$ , signifying the white noise value in the absence of jitter, was  $25.9 \pm 0.6 \mu\text{s}$ .

The presence of jitter in PSR J1713+0747 does mean that for a telescope as sensitive as Arecibo, LEAP, or GiantLEAP, including the future Square Kilometer Array (SKA) or the Five-hundred-meter Aperture Spherical Telescope (FAST) telescopes, both  $\sigma_J$  and the ordinary timing rms scale as  $1/\sqrt{N}$  (Cordes et al. 2004; Cordes & Shannon 2010; Lazio 2013). The dominance of pulse phase jitter seen in the Arecibo portion of the present study may necessitate the use of such long tracks for all future highly sensitive telescopes to further reduce  $\sigma_{\mathcal{R}}$ . This is seen in a particularly dramatic fashion over eight of the nine telescopes here. Even when using one telescope alone, GBT or Arecibo for instance, the fit to  $\sigma_{\mathcal{R}}$  yields  $\sigma_{J,1}$  values of  $21.6 \pm 4.1 \mu\text{s}$  (GBT) and  $27.9 \pm 5.3 \mu\text{s}$  (Arecibo). The  $\sim 1\sigma$  consistency of each single-telescope value with the all-telescope value implies that the jitter numbers reported are not telescope-dependent, and are intrinsic to the pulsar as expected.

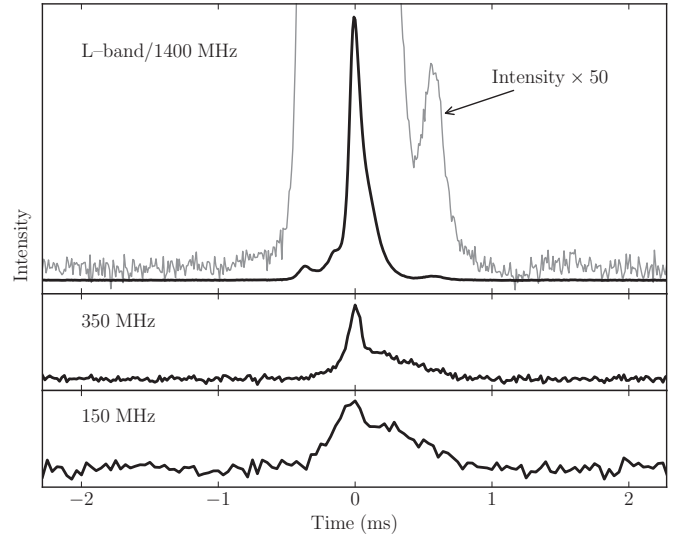
Single pulse phase jitter causes a timing error  $\propto 1/\sqrt{N}$  that is independent of  $S/N$ . The two contributions are equal for single-pulse  $(S/N)_1 \sim 1$  (Shannon & Cordes 2012).  $S/N > 1$  single pulses should be present, given that a telescope is sensitive to jitter noise (Cordes & Shannon 2010).

A more detailed analysis is deferred to a separate publication. Single pulses can be extracted from any of the observations with baseband data or intensity integrations, which were taken at various times during the global campaign in all nine telescopes. However, it is interesting that fitting Equation (3) across the eight telescopes in Figure 6 allows a jitter measurement without probing into residuals with a subintegration time of  $<120$  s, much less with single pulses. The residuals in the brightest bin, rescaled to subintegrations of a single period, correspond to a single pulse  $S/N$  of  $\sim 3$ .

### 3.3. Strong Correlation of Diffractive Scintillation Between Telescopes

Figure 3 shows that the frequency–time structure in the dynamic spectrum is qualitatively identical between simultaneous measurements from different telescopes, apart from low-elevation-angle observations and from masked episodes of RFI. This high correlation includes telescope pairs with the largest separations (up to 9000 km), Parkes and GMRT; GMRT with the European telescopes (Jodrell Bank, Effelsberg, and WSRT), between the GBT and the European telescopes, and between the GBT and Arecibo.

The observations are consistent with the expectation that the dynamic spectra for PSR J1713+0747 should be highly correlated between all terrestrial telescopes because of the low level of scattering along the line of sight. We estimate the spatial scale  $\ell_d$  of the diffraction pattern from the parallax distance  $d = 1.05 \pm 0.06$  kpc (Chatterjee et al. 2009) and the scintillation



**Figure 7.** Grand average profile for all telescopes. The top panel was created across all bandwidths from those telescopes that observed at  $L$  band/1.4 GHz. While it has been shown in Section 2 that there is some profile smearing occurring due to the pulse profile evolution with frequency, we sum in weighted fashion from individual subbands’ residual values. The resulting profile has a  $S/N$  of  $\sim 4000$ , where the signal value is taken as the amplitude at the maximum value of the summed pulse and the noise is taken in the first 100 phase bins of the off-pulse part of the combined profile. The bottom two panels show the grand average profiles for low frequencies, manifesting the significant profile evolution with frequency. The DM for all telescopes is set by the fitted  $L$ -band/1.4 GHz DM.

bandwidth  $\Delta\nu_d \approx 0.6 \pm 0.2$  MHz at 0.43 GHz (Bogdanov et al. 2002) using Equation (9) of Cordes & Rickett (1998),

$$\ell_d = \frac{1}{\nu} \left( \frac{cd\Delta\nu_d}{4\pi C_1} \right)^{1/2}, \quad (4)$$

where  $C_1 = 1.16$  using a default, uniform Kolmogorov scattering medium. This yields  $\ell_d \approx 5 \times 10^4$  km at 0.43 GHz and scaling by  $\nu^{1.2}$ , the diffraction scale at 1.4 GHz is  $\ell_d \approx 2 \times 10^4$  km, much larger than the Earth.

## 4. INITIAL RESULTS ON THE NOISE BUDGET OF THE TIMING RESIDUALS

In this section we briefly consider some aspects of PSR J1713+0747’s noise budget, in other words, whether the  $S/N$  across telescopes corresponds to expectations from general considerations. A more detailed consideration of the noise budget will be found in a forthcoming paper.

Figure 7 (top panel) shows the grand average profile for the eight telescopes that observed at  $L$  band/1.4 GHz. Low frequencies are shown in the lower panels. Profiles are summed across the full band from individual subbands’ residual values, weighted by the off-pulse noise values, and folded according to the measured  $L$  band/1.4 GHz value of DM. The resulting  $L$ -band/1.4 GHz profile has a  $S/N$  of  $\sim 4000$ , where the signal value is taken as the amplitude at the maximum value of the summed pulse and the noise is taken from the off-pulse part of the combined profile.  $S/N$  values were calculated using the first 100 bins of the profile for the noise region. The profile was centered at maximum, with  $N_\phi = 512$ , before summation. See Table 1 for estimates of the degree to which each telescope contributes to the total pulse profile  $S/N$ . Low frequency profiles are also shown for reference, showing significant profile evolution with frequency. The pulsar is weak at low frequencies



because it appears to turn over somewhere above the LOFAR band (T. E. Hassall et al. 2014, in preparation).

The minimum rms on timing residuals from the eight-telescope *L*-band/1.4 GHz profile can be estimated using Equation (2). Given an effective pulse width of 0.54 ms, that we have 512 phase bins, and that the 24 hr pulse profile S/N in the grand average profile was about 4000, this yields a template fitting error,  $\sigma_{S/N}$ , of about 3 ns. Given a 24 hr implied template fitting error of 3 ns and an implied 24 hr jitter timing error of  $170/\sqrt{30}/\sqrt{24} = 6.3$  ns (rescaling from the 120 s jitter value given in Section 3.2 to a 24 hr value), we add these values in quadrature to arrive at an approximate timing uncertainty of 7 ns. Jitter and template fitting alone would then yield a timing residual error of 7 ns on a 24 hr TOA.

The uncertainty from our *L*-band/1.4 GHz DM measurement of  $0.00001 \text{ pc cm}^{-3}$  from Section 2 corresponds to 24 ns of smearing across the band. A better measurement of the DM on MJD 56465–56466 would require incorporation of the 150 MHz LOFAR data and the 350 MHz WSRT data, but taking into account a model of the pulse profile evolution with frequency that extends to these lowest two frequencies. Such a model is a topic of exploration for a future paper on interstellar electron density variations. Any discovered DM variations, along with an improved DM smearing value, would need to inform the noise floor assessment.

## 5. FURTHER WORK

The initial results presented will be important for the three PTAs and for the IPTA as a whole. For some telescopes, PSR J1713+0747 is timed (or is under consideration to be timed) for longer observation durations at each epoch, or at a higher observing cadence. Being among a small set of pulsars with the lowest timing residual rms values, it strongly influences the sensitivity of the entire IPTA, despite the necessity of calculating angular correlations in order to populate the Hellings & Downs diagram and detect a stochastic background of GWs (Hellings & Downs 1983). Increasing the observation duration for these pulsars helps the sensitivity of the IPTA to other types of GW source populations, such as burst, continuous wave, and memory bursts (see Arzoumanian et al. 2014 for current limits on continuous wave sources). The first results presented in this work provide a starting point on the subtleties that may emerge with an increasing dependence on PSR J1713+0747 and similar pulsars such as PSR J0437–4715 and PSR J1909–3744.

We plan to release papers on the following subjects, among others, based on this data set:

*The Noise Budget of the 24 hr Global Observation of PSR J1713+0747:* the question to be explored here is the degree to which one can dissect the noise present on the different timescales relevant in this observation. From single pulses at the  $\mu\text{s}$  resolution all the way to the full 24 hr, the statistical structure of noise in timing residuals can be probed using various diagnostics. Structure on different timescales can be probed by looking at the pulsar with the Allan variance function. Single pulses can also be exploited in order to search for smaller timescale structure such as giant pulses, mode changes, and drifting subpulses (see Shannon & Cordes 2010). Shape changes can be probed and possibly mitigated using various methods (Cordes 1993; Demorest 2007; Osłowski et al. 2011).

*Interstellar Electron Density Variations and Pulse Profile Frequency Evolution:* the all-telescope dynamic spectrum can yield interesting information in a further analysis. Given the data ob-

tained at low frequencies with LOFAR and WSRT, it will be informative to search for correlations between events occurring in the *L*-band/1.4 GHz dynamic spectrum and the highly scattered structure at 150 MHz and 350 MHz, respectively. Analysis can be done using the LOFAR and WSRT data to obtain more accurate DM measurements while taking into account the significant profile evolution between the two low-frequency observations and the *L*-band/1.4 GHz observations.

*GiantLEAP:* one of the signature objectives of this observation is to use the European telescopes, Arecibo, the GMRT, and the GBT as a single phased array, or at least to expand LEAP with some subset thereof. In particular, RFI excision using simultaneous data from a subset of telescopes might significantly improve the quality of the phased array over one more locally situated. Once the proper correlations are performed, in principle the timing rms of PSR J1713+0747 from the largest simultaneous collecting area ever used will be obtained. However, what practical limitations will come into play at realizing this ideal would be the subject of future studies. Undoubtedly, whatever timing results will be obtained will be highly affected by the presence of pulse phase jitter.

*Polarization Studies.* Most telescopes in this study took polarimetric data (see Appendix) and studying the timescales of PSR J1713+0747's polarization over the 24 hr could provide new insights, particularly at the single pulse level.

## 6. CONCLUSIONS

We have presented an overview of the goals and data products of the 24 hr global campaign on PSR J1713+0747. This  $\sim 60$  TB data set is useful for many goals, which will be explored in future papers, including but not limited to: better determination of the overall noise budget for PTAs, a wide-bandwidth, long-timespan examination of the effects of the ISM on pulsar timing, combining baseband data from simultaneous observations for the GiantLEAP experiment, an examination of single pulses and their phenomenology over the 24 hr, and many others.

In the first results presented here, some interesting conclusions can already be drawn. PSR J1713+0747's intrinsic pulse phase jitter ( $\sim 27.0 \mu\text{s}$  for single pulses) can be measured by fitting a noise model across all telescopes, even when TOA integration times are as long as 120 s. The improvement of timing residual rms is not found to depart significantly from a factor of  $1/\sqrt{N}$ , where  $N$  is the number of integrated pulses. Finally, the diffraction scale at 1.4 GHz was seen to be  $\ell_d \approx 2 \times 10^4 \text{ km}$ , much larger than the Earth, from the overlapping scintillation pattern seen in the dynamic spectrum in Figure 3.

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## APPENDIX

### ADDITIONAL OBSERVING DETAILS

Most information about the observation can be found in Table 1. TEMPO2 parameters are in this Appendix in Table 3, most of which derive from the parameters calculated in W. W. Zhu et al. (2014, in preparation).

#### A.1. EFFELSBERG 100 m RADIO TELESCOPE

Effelsberg’s data taking began with 9.3 hr of contiguous observing at 1380 MHz. The PSRIX instrument (R. Karuppuswamy et al. 2014, in preparation) was used for both baseband and folding modes (after real-time coherent dedispersion). There were two baseband sessions, one 30 minutes and the other 1 hr. In folding mode, PSRIX was configured to coherently dedisperse and fold  $8 \times 25$  MHz bands. Each resulting file has 200 MHz of bandwidth (though some channels are removed due to RFI), 128 channels, 1024 phase bins of  $4.47 \mu\text{s}$  each, 10 s subintegrations, and full polarization information. In total there

were  $\sim 6$  hr (2.1 GB) of folding mode data. In baseband mode, data was recorded as  $8 \times 16$  MHz (128 MHz) subbands in order to be compatible with other LEAP telescopes. The data were flux calibrated using the noise diode, which in turn was calibrated using a North-On-South triplet of observations of 3C 218 following the 24 hr campaign. The data were recorded in “Timer Archive” format. They were converted to PSRFITS format as part of Effelsberg’s standard data reduction pipeline.

#### A.2. GIANT METERWAVE RADIO TELESCOPE

The GMRT used 22 antennas, employing two observing modes—a total offline coherent filterbank mode with 65.1 kHz spectral and  $61.44 \mu\text{s}$  time resolution, and a coherent array voltage mode with a single subband for the baseband portion of the observation. Both these modes are described in Roy et al. (2010). The frequency range was from 1371 MHz to 1404 MHz. There were  $7 \times 1$  hr recording scans interleaving with phasing scans for the array, as well as a 50 minutes coherent array baseband voltage recording scan. This resulted in 436 GB of raw filterbank data and 460 GB of voltage data. The GMRT filterbank data ( $61.44 \mu\text{s}$  time resolution) are 16 bit and in a format compatible with the PRESTO<sup>27</sup> searching suite. The GMRT coherent array voltage data (15 ns time resolution) are 8-bit and in a DSPSR friendly format. The GMRT coherent array provides some built-in immunity to RFI as the processing pipeline adjusts the antenna phases to correct for the effect of rotation of the sky signals, which in turn de-correlates the terrestrial signals. Interleaved calibrator observations (QSO J1822–096 in this case) every 2 hr were required to optimize the coherent array sensitivity at the observing frequencies. The antenna-based gain offsets are also corrected using this calibrator before making the coherent beam.

#### A.3. LOVELL TELESCOPE AT JODRELL BANK OBSERVATORY

L-band/1.4 GHz observations of PSR J1713+0747 were obtained with the Lovell telescope at Jodrell Bank over an 11.5 hr timespan. The data are continuous except for a few brief gaps due to the telescope being parked for wind constraints. Two instruments were used, (1) the DFB performed real-time folding with incoherent dedispersion, producing folded 10 s subintegrations of 1024 pulse phase bins of  $4.47 \mu\text{s}$  in size, 0.5 MHz channels over a 384 MHz wide band centered at 1532 MHz, and (2) the ROACH, using a Reconfigurable Open Architecture Computing Hardware FPGA board performing real-time folding with coherent dedispersion using the PSRDADA<sup>28</sup> and DSPSR software packages. This provided 10 s subintegrations with 2048 pulse phase bins of  $2.23 \mu\text{s}$  in size and covered 400 MHz wide band centered at 1532 MHz, split into 25 subbands of 16 MHz, each channelized to provide 0.25 MHz channels. The ROACH was also used to record baseband data during the times when Nançay and Arecibo observed the pulsar. Dual polarization, Nyquist sampled baseband data, 8 bit digitized, was recorded for the lower eight subbands of 16 MHz (1332–1460 MHz), while the remaining seventeen 16 MHz subbands performed real-time folding with coherent dedispersion as before. At the end of the observations the baseband data was folded and coherently dedispersed with the same parameters as for the real-time folding to give one continuous observing run. The spectral

<sup>27</sup> <http://www.cv.nrao.edu/~sransom/presto/>

<sup>28</sup> <http://psrdada.sourceforge.net>

kurtosis method by Nita et al. (2007) for identifying and flagging RFI, as implemented in DSPSR, was used to excise RFI in real time. After the observations a combination of manual and automatic RFI excision was performed to clean the data further. The folded profiles were polarization calibrated using the Single Axis model (van Straten 2004) using observations of the noise diode and observations of pulsars with known polarization properties.

#### A.4. LOW FREQUENCY ARRAY

LOFAR observed from 110 to 190 MHz using the BG/P beam-former and correlator (see van Haarlem et al. 2013). The sampling time was  $5.12 \mu\text{s}$  and 400 subbands of 0.195 MHz each were recorded. Full polarization information was taken in complex voltage mode; see Stappers et al. (2011) for more information on pulsar observing modes with LOFAR. The raw data volume (32 bit) was  $4.5 \text{ TB hr}^{-1}$ , yielding 40 TB of complex-voltage raw data (of which only 1 hr, 4.5 TB, of raw data has been archived long-term; the rest is only available as folded archives, as summarized in Table 1) The 40 TB value did not include in the total the 60 TB value for the all-telescope data. These data were coherently dedispersed and folded offline using DSPSR. The following 23 LOFAR Core stations were combined for the 9 hr observation: CS001, CS002, CS003, CS004, CS005, CS006, CS007, CS011, CS017, CS021, CS024, CS026, CS028, CS030, CS031, CS032, CS101, CS103, CS201, CS301, CS302, CS401, and CS501. See van Haarlem et al. (2013) for more specific location information; by default the phase center of the tied-array beam is placed at the position of CS002.

#### A.5. NANÇAY DECIMETRIC TELESCOPE

Nançay observed using the NUPPI (Nançay Ultimate Pulsar Processing Instrument) backend at  $L$  band/1.4 GHz with a total bandwidth of 512 MHz, split into  $32 \times 16$  MHz channels and 8 bits digitized. The profiles were folded and integrated over 1 minute and finally stored in a 29 MB PSRFITS file. All data were coherently dedispersed and the total Nançay observation lasted for  $\sim 1$  hr. The PSRCHIVE program PAC was used to do the polarization calibration with the Single Axis model and automatic zapping was then applied with the PAZ program. The TOAs were produced with PAT using a high S/N template.

#### A.6. PARKES 64 m TELESCOPE

The Parkes 64 m radio telescope observed PSR J1713+0747 at 1362 MHz using both fold-mode and baseband mode in parallel for  $\sim 6$  hr. This time included four 3 minutes noise diode calibration scans, between 64 minute blocks of folding. The backends DFB3/4 (incoherent filterbank, 60 s foldings, 1024 frequency channels over 256 MHz bandwidth, 1024 phase bins of  $4.47 \mu\text{s}$  each), APSR (real-time coherent filterbank, 30 s foldings, 512 frequency channels over 256 MHz bandwidth, 1024 phase bins of  $4.47 \mu\text{s}$  each), and CASPSR (30 s foldings, 400 MHz bandwidth with  $\sim 10$  MHz band edges, 1024 phase bins of  $4.47 \mu\text{s}$  each) observed in parallel, allowing for simultaneous baseband and folding mode observations. RFI was removed, and consistent results were obtained with all backends using a median filter in the frequency domain. Polarization and flux calibration for the DFB data used the standard monthly flux calibrations on the Hydra A radio galaxy (Manchester et al. 2013). CASPSR was calibrated for differential gain and phase but nothing else.

RFI was mitigated using the radio-frequency domain filter implemented in the PSRCHIVE command PAZ (Hotan et al. 2004). Additionally, the CASPSR instrument mitigates RFI in real time by rejecting portions of the data that show distribution inconsistent with receiver noise by using a spectral kurtosis filter. Digital filterbank data were calibrated for cross-coupling using a model for the feed derived from long-track observations of the bright source PSR J0437–4715. Observations over these wide parallactic angles enable the measurement of the feed cross coupling and ellipticity to be measured (van Straten 2004). The model used was the average of many long-track observations because the feed parameters were not found to change significantly with time. CASPSR observations were not calibrated for this cross coupling. PSR J1713+0747 is only a modestly polarized pulsar and the effects of correction for these effects were found to be negligible improvement on its long term timing precision (Manchester et al. 2013). Details of this calibration are further described in Manchester et al. (2013). All the data are digitized with 8 bit digitizers. The baseband data set is 30 TB, with 7.5 GB for the fold-mode data set. The fold mode used the PSRFITS data format, and the baseband mode used the PSRDADA format.

#### A.7. WESTERBORK SYNTHESIS RADIO TELESCOPE

WSRT is a  $14 \times 25$  m dish East–West array, which for pulsar observations is used in tied-array mode and is phased up before the observations for all observing bands used. There were no absolute flux calibrations done for these observations. Full polarization information was stored. Due to the array-nature of the telescope there is usually very little RFI and therefore any remaining leftover narrow-channel RFI zapping is done offline using PSRCHIVE tools. For these observations we had 11 out of the total 14 dishes available. The WSRT observations made use of PuMa-II instrument (Karuppusamy et al. 2008), observing at both  $L$  band/1.4 GHz and 350 MHz each with eight separate bands, which are each either 10 MHz wide (for the 350 MHz observations) or 20 MHz wide (for the 1380 MHz observations). The individual bands overlapped by 2.5 MHz of the 10 MHz bandwidth for the 350 MHz observations. At  $L$  band/1.4 GHz the bands overlapped by 4 MHz out of the 20 MHz bandwidth to ensure full overlap with the 16 MHz bands at the Effelsberg, Lovell, and Nançay telescopes and allow for coherent addition of the data. For each there are 10 s subintegrations stored with 256 bins of  $17.9 \mu\text{s}$  each across the profile and 64 channels across each band (which is 10 or 20 MHz depending on the frequency). The low frequency TOAs were made using PSRCHIVE templates, created using the PAAS (analytic template) routine based on a high-S/N summation of many other observations.

#### A.8. ARECIBO OBSERVATORY AND THE NRAO GREEN BANK TELESCOPE

Both Arecibo and the GBT, as mentioned earlier, used non-folded intensity recording for their nonbaseband portions of the observation so that single-pulse data would be available over the entire span of the observation, without the cost in bandwidth. Unlike baseband data, no voltage, and thus no electromagnetic phase information, is present. This observation mode is essentially a pulsar search mode with coherent dedispersion, according to the source’s DM. Calibrations at the start of the observations at both GBT and Arecibo were performed with a noise diode switched at 25 Hz, including the polarization calibration. Absolute flux measurements, also including the polarization calibration, were performed on QSO B1442+101 at the GBT and

QSO J1413+1509 at Arecibo. We apply these calibrations via the Single Axis model using the PSRCHIVE program PAC.

#### A.8.1. Green Bank Telescope

GBT began its observation starting with baseband mode, switching to an hour of intensity integration observing mode, switching back to baseband mode for 30 minutes, and then returning to intensity integration mode for the remaining 7 hr. The two baseband sessions would be simultaneous with other telescopes. The intensity integration mode had  $256 \times 3.125$  MHz frequency channels and had a time resolution of  $5.12 \mu\text{s}$ . The same was planned for the PUPPI backend at Arecibo. Full polarization information was recorded. Intensity integrations are effectively search mode but coherently dedispersed with the known DM. During the switch from the first baseband time block to the first intensity integration time block, some time was lost due to a problem with the observing mode on GUPPI, and so to ensure a safer data rate, the observing mode was switched to  $128 \times 6.25$  MHz channels with a time resolution of  $2.56 \mu\text{s}$ . For this reason, the settings on GUPPI were different from PUPPI. The tradeoff is that with slightly wider channels in PUPPI, although yielding a better set of phase bins, cannot in principle excise RFI as efficiently. This interruption and restart caused a small gap in the data, visible in Figure 3. GBT's backend employing real-time cyclic spectroscopy (G. Jones et al. 2014, in preparation) was also used in parallel during the folding mode observations, with 65 MHz of bandwidth centered at 1398 MHz.

#### A.8.2. Arecibo Observatory

Finally, Arecibo joined the observation. Beginning with 30 minutes of baseband observing in order to contribute to GiantLEAP, it then switched over to PUPPI intensity integration mode, a different mode used than that at GBT as just described.

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