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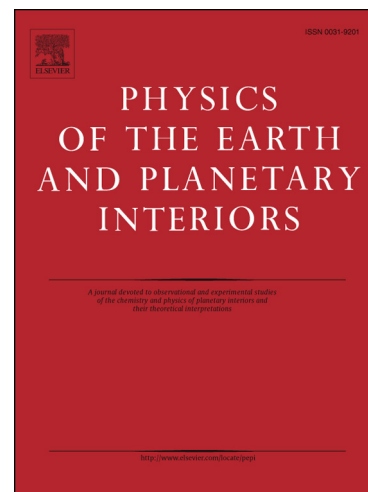
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Laboratory micro-seismic signature of shear faulting and fault slip in shale

J. Sarout^a, Y. Le Gonidec^b, A. Ougier-Simonin^c, A. Schubnel^d,
Y. Guéguen^d, and D.N. Dewhurst^a

^a*CSIRO Energy, Perth, Australia*

^b*Géosciences Rennes - CNRS/INSU UMR6118, Rennes, France*

^c*British Geological Survey, Engineering Geology, Keyworth, UK*

^d*Ecole Normale Supérieure, CNRS-UMR 8538, Laboratoire de Géologie, Paris, France*

Abstract

This article reports the results of a triaxial deformation experiment conducted on a transversely isotropic shale specimen. This specimen was instrumented with ultrasonic transducers to monitor the evolution of the micro-seismic activity induced by shear faulting (triaxial failure) and subsequent fault slip at two different rates. The strain data demonstrate the anisotropy of the mechanical (quasi-static) compliance of the shale; the P-wave velocity data demonstrate the anisotropy of the elastic (dynamic) compliance of the shale. The spatio-temporal evolution of the micro-seismic activity suggests the development of two distinct but overlapping shear faults, a feature similar to relay ramps observed in large-scale structural geology. The shear faulting of the shale specimen appears quasi-aseismic, at least in the 0.5 MHz range of sensitivity of the ultrasonic transducers used in the experiment. Concomitantly, the rate of micro-seismic activity is strongly correlated with the imposed slip rate and the evolution of the axial stress. The moment

Email address: joel.sarout@csiro.au (J. Sarout)

tensor inversion of the focal mechanism of the high quality micro-seismic events recorded suggests a transition from a non-shear dominated to shear-dominated micro-seismic activity when the rock evolves from initial failure to larger and faster slip along the fault. The frictional behaviour of the shear faults highlights the possible interactions between small asperities and slow slip of a velocity-strengthening fault, which could be considered as a realistic experimental analogue of natural observations of non-volcanic tremors and (very) low-frequency earthquakes triggered by slow slip events.

Keywords: Shale, P-wave velocity, Anisotropy, Micro-seismicity, Focal mechanism, Shear faulting, Fault slip, Friction

1. Introduction

Changes in the stress state can induce brittle damage and fracturing in rocks that can radiate mechanical energy in the form of elastic waves. At the field scale, the radiated energy is often referred to as Micro-Seismic (MS) activity; in the laboratory, it is often called acoustic emissions [28, 31, 32]. The phenomena of micro-seismicity and acoustic emission are similar in nature, although the frequency content of the radiated elastic perturbation might be different due to the scale of the fracturing. Therefore, in this manuscript we will use the term micro-seismicity (and its derivatives such as micro-seismic activity or micro-seismic events) to name the events recorded in the laboratory. In general, the accumulation of damage can ultimately lead to the mechanical failure of the rock. Among the various rock failure mechanisms listed in the literature, we focus here on brittle faulting pertaining to numerous geological settings observable during the deformation of rocks in the Earth's upper crust.

It is generally accepted that for a given material, MS activity is prominently observed during deformation under the following conditions: (i) relatively low normal stresses; (ii) relatively high shear stresses; and/or (iii) relatively high stress loading rates, e.g., [1, 51]. In the past, most research efforts published in the literature involving micro-seismic monitoring of deformation processes in the laboratory have focused either on:

- crystalline rocks in relation to earthquake/fault mechanics, geotechnical or geothermal applications, e.g., [6, 26, 29, 33]; or
- conventional reservoir rocks in relation to oil and gas exploration, production and monitoring (reservoir integrity, compartmentalisation, injection-

26 induced fracture/fault reactivation...), e.g., in sandstones [7, 8, 9, 15, 17, 45];
 27 or to a lesser extent in porous carbonate rocks [16].

28 At the field scale, several studies on the monitoring of MS activity in
 29 granites and carbonates have been published. These include the monitoring
 30 of: thermally-induced MS activity potentially associated with radioactive
 31 waste disposal in boreholes drilled in a tunnel's floor at Äspö's Hard Rock
 32 Laboratory in Sweden [37], in the Excavation Damage Zone (EDZ) in the Un-
 33 derground Research Laboratory in a granitic rock mass in Canada [53, 54]
 34 and injection-induced MS activity in a limestone formation in the Laboratoire
 35 Souterrain à Bas Bruit in France [18]. Fewer field-scale studies on the MS
 36 activity induced by faulting or fault slip in shale formations have been pub-
 37 lished. A recent study demonstrated the feasibility of monitoring the time
 38 evolution of MS activity associated with the EDZ in the Opalinus Clay for-
 39 mation at the Mont-Terri Underground Research Laboratory in Switzerland
 40 [27]. The MS activity associated with fluid injection in the Colorado Shale
 41 formation was successfully monitored by [46]. In contrast, the monitoring of
 42 the spatial extent of anthropogenic hydraulic fractures in stimulated oil/gas
 43 reservoirs have been an active field of research since the 1980's, strongly sup-
 44 ported by industry funding, especially in the recent years with the advent
 45 and development of commercially-viable unconventional reservoirs such as
 46 gas shales, e.g., [49].

47 At the laboratory scale, experiments have been reported on shales uni-
 48 axially deformed at room conditions under large loading rates (see [2] and
 49 references therein). However, no determination of spatial locations or focal
 50 mechanisms of the recorded MS events (MSEs) was carried out. MS activity

51 and location in shale samples containing quartz veins have been reported by
52 [30]. In this particular case, and as expected, the MS activity seemed to
53 coincide with the location of quartz veins favorably oriented with respect to
54 the maximum principal compressive stress.

55 To our knowledge, no data on spatio-temporal localisation and focal
56 mechanism estimation of MS activity have been reported on deforming clay-
57 rich rocks such as conventional reservoir-sealing shales. Under triaxial defor-
58 mation at realistic subsurface stress conditions, the shale specimens fail in
59 shear, leading to the formation of a shear fracture. The first questions that
60 arise then for these rocks are the following: (i) can we expect precursory
61 micro-seismic activity prior to the macroscopic faulting? (ii) Would the slip
62 on the newly generated fault induce any micro-seismic activity? (iii) How
63 would the signature of the MS activity be affected by the deformation rate?

64 Due to their fine-grained nature, it is generally thought that clays act as a
65 lubricant in frictional geological environments, e.g., [36]. Also, the brittleness
66 of clay behaviour is known to be controlled by their degree of hydration (the
67 more hydrated, the less brittle), their mineral composition, and the imposed
68 deformation rate (higher rates induce a more brittle response). The lack
69 of published experimental studies on the MS activity of shales subjected to
70 stress conditions typical of the upper crust can probably be explained by the
71 inherent complexity of shales and the associated difficulty in conducting lab-
72 oratory deformation experiments on them under well-controlled conditions.
73 In addition, there is considerable technical complexity in conducting and pro-
74 cessing laboratory experiments aimed at monitoring and locating with high
75 accuracy the MS activity induced by deforming relatively small specimens.

In this regard, the difficulty in locating the MS activity is exacerbated by the directional dependency (anisotropy) of wave propagation in shales, e.g., [11, 13, 24, 39, 40, 48].

In this paper, the results and analysis of a laboratory deformation experiment in a shale specimen are reported. The specimen was triaxially deformed to beyond the failure point under subsurface stress conditions while associated MS activity was recorded. The aim was to analyse the contrast in the MS signature of shear faulting and subsequent fault slip as well as the effect of the deformation rate on the fault's micro-seismic and frictional response.

In the following pages, the experimental conditions are detailed (section 2) along with the main results in terms of stress-strain data, ultrasonic P-wave velocity data, and micro-seismic activity (section 3). The fourth section is dedicated to an analysis and discussion of these results in terms of the MS signatures of shear faulting and fault slip (slow/fast slip), frictional behaviour of the shear fault in relation to the associated MS activity, and a comparison to other rock lithologies.

2. Description of the experiment

2.1. Shale material

A large core was recovered from the North Sea at a depth of 1643 m below sea bed in a clay-rich shale formation (Campanian, upper Cretaceous). The core was preserved since recovery from depth in several layers of plastic and aluminium wrap with an additional external wax coating. After unpacking, this shale appeared relatively homogeneous, dark grey in colour, with bedding visible inclined at 45° to the core axis. Twin cylindrical specimens

40 mm in diameter have been cored along the axis of the original core so that the bedding was also inclined at 45° to their axis. Their end faces were trimmed and ground to be parallel to each other to within 0.02 mm. The final length of the specimens was 81 mm (long specimen) and 40 mm (short specimen), respectively. For the coring, trimming and grinding operations, compressed air was used as the cooling fluid. After preparation, the specimens were equilibrated for several days at room conditions (20°C , relative humidity of 50%) until stabilisation of their mass at these conditions. After this initial treatment the specimens turned to a light grey colour. The mass evolution of the samples during this initial treatment and their change in color suggest that they lost water (dehydration) by exchange with the atmosphere. The porosity of the shale was estimated to be of the order of 19% (density: 2370 kg/m^3) based on mass measurements conducted on a separate block cut from the original core in its preserved state (immediately after unpacking the core) and its state after mass stabilisation at a room conditions (20°C , relative humidity of 50%). Note that this porosity is only a lower bound estimate of the actual porosity of the shale assuming that the core was fully water-saturated in its preserved state and is fully dry in its final equilibrated state (20°C and relative humidity of 50%). It is expected that only the so-called "free" water could have evaporated during this treatment, so that the shale specimens are likely in a partially saturated state.

The shorter specimen was used to conduct permeability measurements with nitrogen gas under increasing effective pressure using a steady state method, i.e., constant gas flow imposed at one end of the specimen, and monitoring of the differential pressure build-up and stabilisation across its

two ends [25]. The permeability results are summarised in TABLE 1. The permeability of this shale to nitrogen decreases by almost two orders of magnitude from 2.1×10^{-5} mD down to 6.9×10^{-7} mD when the effective confining pressure increases from 4 MPa up to 65 MPa. This seems to indicate that stress-sensitive pre-existing micro-cracks (damage) are closed by the increasing effective pressure. Such micro-cracks might have been induced by stress release following the recovery of the shale core from depth and/or the dehydration of the specimen at room conditions during initial treatment.

The longer specimen was used to conduct the triaxial deformation experiment with MS monitoring detailed in the remainder of this article.

2.2. Experimental equipment

In order to characterise the MS response of the shale to changes in the triaxial stress state, a specific laboratory setup is required to monitor both the deformation of the specimen and the induced MS activity. The experimental setup consists mainly of: (i) a *Sanchez Technologies* axisymmetric triaxial stress vessel in which a radial and an axial stress can be independently applied to a cylindrical rock specimen; (ii) an *Applied Seismology Consultants* multi-channel ultrasonic/micro-seismic monitoring system (Fig. 1). This apparatus allows the simultaneous acquisition of various types of data on a single rock specimen: (i) radial and (ii) axial deformations, (iii) active ultrasonic monitoring, i.e., ultrasonic P-wave velocities along numerous propagation paths at selected stages of the deformation (called velocity surveys); and (iv) passive monitoring, i.e., induced micro-seismicity (also called acoustic emissions). Note that both active and passive monitoring are conducted using the same array of ultrasonic transducers as described below.

After the initial drying treatment of the long shale specimen at a temperature of 20°C and a relative humidity of 50%, four strain gauges are glued onto its lateral surface so that four independent directions of deformation are measured (see Figs. 1 and 2): Gauge 1 measures the axial strain along the specimen's axis, at 45° to the bedding orientation. Gauge 2 measures the circumferential strain orthogonal to the specimen's axis, at 45° to the bedding; this strain also corresponds to the radial strain, and for sake of simplicity, it will be referred to as radial strain in the remaining of the article. Gauge 3 measures the strain orthogonal to the bedding, at 45° to the specimen's axis. Gauge 4 measures the strain along the bedding, at 45° to the specimen's axis. In addition, the average axial displacement between the two ends of the specimen was monitored using three contactless Eddy current displacement transducers located outside the pressure vessel.

2.3. Experimental protocol

The shale specimen is inserted into a flexible *Viton* sleeve and placed inside the pressure chamber of the triaxial stress vessel, which is then closed and filled with oil. The purpose of the flexible sleeve is to isolate the specimen from the hydraulic oil used to apply the radial stress [40]. This specimen is instrumented with: (i) four strain gauges glued directly to its lateral surface, at mid-height; (ii) an array of 16 miniature ultrasonic transducers (6 mm in diameter) made of piezo-ceramic material with a central resonant frequency of about 0.5 MHz. These transducers can be used as ultrasonic sources or receivers attached directly to the lateral surface of the specimen, through sealable holes in the flexible *Viton* sleeve (Fig. 2).

The experimental deformation protocol consists of: (i) an isotropic stress

loading to subject the specimen to a simulated *in situ* condition with a confining pressure of 10 MPa; (ii) a deviatoric stress loading at a constant axial displacement rate of 1 mm/h ($3.5 \times 10^{-6} \text{ s}^{-1}$) up to a point beyond the specimen's failure, which is indicated by a peak in the recorded deviatoric stress; then (iii) a sudden increase of the displacement rate to 10 mm/h ($3.5 \times 10^{-5} \text{ s}^{-1}$) until stabilisation of the recorded deviatoric stress (Fig. 3). The deformation experiment is conducted without injecting water and without controlling the pore pressure at the two ends of the specimen.

The aim of the deviatoric stress loading is two-fold: (i) assess the effect of shear faulting and fault slip on the MS response of a shale; and (ii) assess the effect of fault slip rate on the MS activity. The active and passive monitoring equipment is controlled with the *Xtream* software, while the data management and processing is conducted with the *Insite Seismic Processor* software.

As part of the active ultrasonic monitoring, at selected stages of the experiment, a P-wave velocity survey is conducted. Each survey consists of 16 consecutive shots, one from each transducer acting as a source. For each source transducer shot, the transmitted waveforms are recorded on the 15 remaining transducers which act as receivers. The waveform recorded at each receiver corresponds to the mechanical vibration transmitted through the rock specimen from the source transducer to that particular receiver. In order to improve the signal-to-noise ratio (SNR), each waveform is in fact the result of the stack of several tens of shots from a given source transducer. The waveforms are recorded with a sampling rate of 10 MHz and an amplitude resolution of 12 bits. Each source-receiver pair defines a particular ray path

within the specimen, i.e. different directions of wave propagation relative to the specimen's axis and therefore relative to the shale bedding. Each velocity survey typically lasts 30 seconds and consists of 240 waveforms (recorded over 82 microseconds), half of which corresponds to different ray paths within the volume of the specimen. The ultrasonic survey data set acquired during the experiment consists of 10 surveys recorded during the isotropic stress loading after every one or two MPa of confining pressure, and 11 surveys recorded during the deviatoric stress loading.

Between two consecutive velocity surveys, the ultrasonic/micro-seismic system is switched to the passive monitoring mode in order to record any MS activity induced by the stress loading. In this mode, the voltages generated by the ultrasonic transducers sensing a given Micro-seismic events (MSE) are recorded according to a pre-defined trigger logic. Typically, if five transducers exceed a voltage threshold of 15 mV within a time window of 100 microseconds, the waveforms from all 16 transducers are recorded for a time window of 82 microseconds. These waveforms are also recorded with a sampling rate of 10 MHz and an amplitude resolution of 12 bits. At the end of the experiment, nearly 500 events have been detected according to this protocol.

3. Shear faulting and post-failure slip

3.1. Identification of the faulting dynamics

The shale deformation experiment can be divided into an isotropic stress loading (Phase 0), followed by a deviatoric stress loading. The deviatoric loading stage itself is composed of three phases as discussed below (Figs. 4

224 and 5).

225 During Phase 0, the specimen reaches the simulated *in situ* stress con-
 226 dition with a confining pressure of 10 MPa (point A in Figs. 4 and 5).
 227 This phase consists of a step-wise increase of the confining pressure and an
 228 equilibration of the specimen at the target condition over several days.

229 Phase 1 corresponds to the *shear faulting* (yellow area in Figs. 4 and 5).
 230 Axial loading is applied to the specimen at a controlled vertical displacement
 231 rate of 1 mm/h until the peak axial stress is slightly passed and a first
 232 moderate stress drop of about 1 MPa is observed, most probably concomitant
 233 with a first slip of the newly formed shear fault (point B in Figs. 4 and 5).
 234 The dip angle of the slip surface with respect to a horizontal plane have been
 235 estimated post mortem to be about 45°, coinciding approximately with the
 236 orientation of the shale bedding. Such an orientation is probably due to the
 237 presence of weak bedding planes (see failed sample in Fig. 3).

238 Phase 2 corresponds to the *slow fault slip* (blue area in Figs. 4 and 5). The
 239 vertical displacement rate is maintained constant so that the newly formed
 240 shear fault is slipping at constant rate, while the axial stress drop of about
 241 7 MPa is more pronounced than in Phase 1.

242 Phase 3 corresponds to the *fast fault slip* (pink area in Figs. 4 and 5) The
 243 vertical displacement rate is suddenly increased to 10 mm/h, which leads to
 244 a sudden, moderate and temporary increase of the axial stress of less than 1
 245 MPa (point C in Figs. 4 and 5). While the axial displacement is maintained
 246 constant at that higher rate, after a temporary stabilisation, the axial stress
 247 starts to slowly increase to reach a plateau by the end of the experiment
 248 (point D in Figs. 4 and 5).

249 In addition to the evolution with time of the axial stress and displacement,
 250 Figs. 4 and 5 also display the evolution of the micro-seismic activity in terms
 251 of cumulated number of MSEs and rate of occurrence, respectively. Overall,
 252 the cumulated number of events is linearly related to the axial displacement
 253 rate, except temporarily after the increase in the imposed displacement rate
 254 from 1 to 10 mm/hour and until the axial stress reaches a plateau. Con-
 255 sistently, the rate of micro-seismic activity is strongly correlated with the
 256 imposed displacement rate and the evolution of the axial stress. More de-
 257 tails about this part of the dataset are provided in Section 4.3.

258 *3.2. Analysis of the stress-strain data*

259 At the end of the isotropic stress loading (Phase 0 aimed at reaching a
 260 confining pressure of 10 MPa), Gages 1, 2, and 3 display a similar amount
 261 of strain (0.123%), whereas Gage 4 (along the bedding and at 45° to the
 262 specimen's axis) displays about half that amount of strain (0.072%). This
 263 suggests a significant stress-induced anisotropy of the shale in which the bed-
 264 ding direction is significantly less compliant than the three other measured
 265 directions. However, the difference in the magnitude of the recorded strain
 266 between Gages 1, 2 and 3 does not clearly reflect a larger compliance in a
 267 direction orthogonal to the bedding compared to the two other intermediate
 268 orientations (at 45° to the bedding). Over all, the amount of deformation
 269 experienced by the specimen during this isotropic stress loading is relatively
 270 small, which may explain the lack of sensitivity of the strain gauge recordings
 271 and therefore the lack of discrimination between the three directions probed
 272 by Gages 1, 2 and 3.

273 During the deviatoric stress loading, the four gauges record a significantly

274 larger amount of strain (Fig. 6). The whole dataset recorded during Phases
 275 1, 2 and 3 is displayed in this figure. Note however that past the point
 276 of strain localisation (shear faulting, slightly beyond the peak stress corre-
 277 sponding to Point B in Fig. 4-6), the local strain measurement provided by
 278 the strain gauges is no longer representative of the average strain field over
 279 the volume of the specimen because most of the imposed axial displacement
 280 is then accommodated by the slipping shear fault. The largest deformation is
 281 expectedly recorded along the specimen's axis (about 1% at the peak stress,
 282 along the maximum principal compressive stress), while the radial strain
 283 along the minimum principal stress is negative due to Poisson's effect (about
 284 -0.1% at the peak stress). Gages 3 and 4 record an intermediate amount of
 285 strain, consistent with their orientation with respect to the principal stress
 286 axes. The difference in magnitude of strain recorded by these two gauges
 287 highlights again the existence of a significant anisotropy in the mechanical
 288 compliance of the shale. Indeed, in view of their similar orientation with
 289 respect to the principal compressive stress axis (45°), they should record a
 290 similar deformation if the shale was isotropic. However, it turns out that
 291 Gage 3 oriented normal to the bedding records a larger strain than gauge 4
 292 oriented along the bedding due to the mechanical anisotropy of the shale.

293 These observations suggest that the quasi-static mechanical compliance
 294 of this shale exhibits a significant directional dependency (anisotropy), that
 295 is, the compliance across the bedding plane is measurably larger than that
 296 along the bedding. This phenomenon has been extensively reported in the
 297 literature for many shales of different origin and geological history (e.g., [11,
 298 14, 40, 41, 42] and references therein). It has also been reported for other

sedimentary rocks (e.g., [10] and references therein). It is therefore reasonable to assume that while subjected only to a confining pressure, this shale is transversely isotropic (TI) in terms of mechanical properties with a symmetry axis orthogonal to the bedding plane. This symmetry might not hold during deviatoric stress loading because the applied axial stress does not coincide with the shale's original axis of transverse isotropy.

4. Micro-seismic signature

4.1. Analysis of the P-wave velocity data

The 21 P-wave velocity surveys recorded during the experiment were processed with the *Insite* software. The flight time of the P-wave recorded in each waveform is picked manually rather than by using an automatic algorithm because of the reasonable number of acoustic surveys. This allows systematic quality control of the results with a high degree of confidence. For each source-receiver pair, the P-wave velocity V_p is calculated using the shortest straight path between the transducers, that is from the closest edge of each transducer to the other (known from the spatial location and dimension of the transducers).

At a given stage of the experiment, the P-wave velocity along five directions of propagation are estimated, which are referred to as $V_p(90^\circ)$, $V_p(60^\circ)$, $V_p(45^\circ)$, $V_p(30^\circ)$, and $V_p(0^\circ)$, where the angles in degrees indicate the propagation direction with respect to the bedding plane. Note that for each nominal ray path orientation θ with respect to the shale bedding, V_p is averaged over all source-receiver pairs yielding a ray path orientation comprised in the interval $[\theta-5^\circ, \theta+5^\circ]$.

323 The uncertainty in the estimation of the relative variation of V_p along a
 324 given direction during the experiment is of the order of 1%. This estimate is
 325 based on: (i) a waveform sampling period of $0.1 \mu s$ for a propagation time
 326 within the specimen comprised between 10 and $15 \mu s$, and (ii) an uncertainty
 327 in the determination of the propagation distance of about 0.1 mm (caliper) for
 328 an average travel distance of about 30 mm. The uncertainty in the estimation
 329 of the absolute value of V_p along a given direction is expected to be higher,
 330 of the order of 10%, mainly due to the inherently higher uncertainty of about
 331 $1 \mu s$ with which a human operator can decide for the P-wave arrival time
 332 from an experimentally recorded waveform.

333 During the isotropic loading (Phase 0), and for all propagation direc-
 334 tions, a significant increase in V_p with a confining pressure increase from
 335 0 to 3 MPa is observed, with only slight increase between 3 and 10 MPa
 336 (Fig. 7a). Despite the uncertainty in the estimation of the absolute value
 337 of V_p (the worst case scenario is represented by the error bars in Fig. 7a),
 338 the relative magnitudes of V_p along the different propagation directions can
 339 be considered as reliable. The elastic anisotropy of the shale is clearly high-
 340 lighted, with a slow $V_p(90^\circ)$ and a fast $V_p(0^\circ)$ velocity across and along the
 341 bedding, respectively. We also observe that $V_p(60^\circ)$, $V_p(45^\circ)$ and $V_p(30^\circ)$
 342 exhibit intermediate values, inversely proportional to their angular inclina-
 343 tion with respect to the bedding plane. This suggests that the shale specimen
 344 can reasonably be assumed to be transversely isotropic (TI) in terms of its
 345 dynamic elastic response. This phenomenon has also been extensively re-
 346 ported in the literature for many shales of different origin and geological
 347 history (e.g., [40, 41] and references therein)

348 In view of the size of the ultrasonic transducers and the propagation
349 distances within the specimen, the estimated P-wave velocities are assumed
350 to be group (ray) velocities ([12]). However, along the symmetry axis and
351 the symmetry plane of the TI shale, group (ray) and phase velocity coincide.
352 Therefore, Thomsen's parameter [47] $\varepsilon = (V_p(0^\circ)^2 - V_p(90^\circ)^2)/2V_p(90^\circ)^2$
353 quantifying the P-wave anisotropy in a TI medium can be estimated using
354 the measured group velocities (Fig. 7b, d).

355 The P-wave velocity and the corresponding P-wave anisotropy as mea-
356 sured by Thomsen's ε parameter exhibit a significant dependency to the
357 confining pressure (Fig. 7a, b): ε drops from 1.8 to 0.8 between 0 and 3 MPa
358 and remains almost constant from 3 to 10 MPa. This suggests a closure
359 of pre-existing micro-cracks (damage) sub-parallel to the bedding with the
360 increase in effective pressure, which is consistent with the dependency of the
361 gas permeability to effective pressure reported in Section 2.1. In contrast,
362 during the deviatoric stress loading (Phases 1 to 3), P-wave velocities appear
363 nearly constant or rise slightly (Fig. 7c), and Thomsen's parameter ε exhibits
364 a moderate dependency to deviatoric stress (Fig. 7d), decreasing to 0.6 as
365 differential stress increases from 0 to 35 MPa.

366 4.2. P-wave velocity model of the shale sample

367 In order to spatially locate the MSEs recorded during the experiment,
368 a P-wave velocity model is required. Based on the analysis of the P-wave
369 velocity data, the velocity model should in principle account for the TI nature
370 of the elastic properties of the shale and the variation of the P-wave velocities
371 with stress. However, as the aim is only to locate MSEs recorded during the
372 deviatoric stress loading (Phases 1 to 3), and accounting for the fact that the

P-wave velocities are not significantly affected by the deviatoric stress during these phases (Fig. 7c and d), the velocities recorded at the start of Phase 1 are used to build the required velocity model of the shale, that is when the confining pressure is 10 MPa and the axial stress is zero. Note that this model is only a pragmatic approximation assuming that the shale specimen is homogeneous.

In addition, because their spatial location is known, the ultrasonic sources shot during the velocity surveys can first be used to assess the validity of both the location (inversion) algorithm and the selected TI velocity model. A Simplex algorithm implemented in the *Insite* software, and a velocity model based on a slow velocity $V_p(90^\circ) = 2000$ m/s and an $\varepsilon = 0.78$ are used. The orientation of the symmetry axis of this model is inferred from the known orientation of the bedding in the specimen, that is at 45° to the specimen's axis.

Although this velocity model accounts for the experimentally estimated velocity and anisotropy, at the scale of the specimen used in this experiment, this combination of values produced a distorted pattern of location of the source shots. In an attempt to improve the results and optimise the procedure, several values of the slow velocity $V_p(90^\circ)$ and the value of ε are tested. The combination that produces the best source shots locations is found to be $V_p(90^\circ) = 1900$ m/s and $\varepsilon = 0.625$. Because the velocity field is not strongly affected by the axial load and related displacement during Phases 1 to 3 (see Fig. 7c), we use this velocity field for all source shots location and MSEs recorded during these phases. The inversion using these values and applied to 176 ultrasonic shots (for all the velocity surveys conducted during Phases

1 to 3) reflects reasonably well the known position of the ultrasonic array, i.e. the sources clearly locate in the vicinity of the transducers (Fig.8). The uncertainty associated specifically with the location of these source shots for all surveys conducted during Phases 1 to 3 amounts to 0.9 mm in average for an average number of triggered sensors of 12. This residual mismatch between the recovered and the actual sensor positions can reasonably be attributed to: (i) the progressive loss of transverse isotropy of the shale during the application of the deviatoric stress (not aligned with the original symmetry axis); (ii) the natural and stress induced heterogeneity of the velocity field in the shale specimen (including the effects of shear fracturing/faulting); and (iii) the relative motion of the sensors during the shear fracturing and faulting in Phases 2 and 3. The pragmatic selection of the adequate velocity structure ($V_p(90^\circ) = 1900$ m/s and $\varepsilon = 0.625$) partially compensates for these uncertainties in view of the results of the source shots location (Fig.8).

4.3. Analysis of the induced micro-seismicity

4.3.1. Spatio-temporal evolution

According to the passive monitoring protocol described in Section 2.3, nearly 500 events are detected during the whole experiment, although not all of them are identified as MSEs. Due to the reasonable number of events recorded, a manual check of the acquired data set was possible. A number of events are identified as electronic noise while others are discarded due to the low SNR of the recorded waveforms. Finally, only the events that could be reliably located within the volume of the specimen are selected for further analysis (Fig. 9). This procedure finally leads to the selection of a total of 280 MSEs: 34 during Phase 1 (yellow spheres in Fig. 9), 14 during Phase 2 (blue

spheres), and 232 during Phase 3 (pink spheres). The average location error for the whole dataset is 2.3 mm for an average number of triggered sensors of 12. This location error is computed for each MSE during the inversion based on the residuals and knowing the velocity structure of the medium. The causes of this uncertainty are related to the experimental uncertainties in the determination of the exact position of the sensors, the heterogeneity of the rock sample, the uncertainty associated with the determination of the (homogeneous but anisotropic) velocity structure to match the actual velocity along the various ray paths. For an imposed axial displacement of 1 mm/hour (Phase 2), the average rate of MSEs is 0.07 MSE/second. This value reaches an average of 0.19 MSE/second over the whole Phase 3 of imposed axial displacement at 10 mm/hour.

Only 64 events are detected during Phase 0 of confining pressure loading applied to reach the simulated *in situ* stress. Out of these events, 15 MSEs with sufficient SNR have been identified and spatially localised. They were randomly located in the volume of the shale specimen. For sake of clarity and because they are not induced by the triaxial loading, these events have been discarded and are not represented in Fig. 9.

The spatial distribution of the MSEs is clearly not random: they appear distributed along two main planar structures, sub-parallel to the shale bedding (Fig. 10). A first structure, highlighted as a yellow plane, is initiated during Phase 1: few yellow MSEs seem to be distributed over the volume of the specimen, but most of them appear to cluster along the highlighted yellow plane. This reflects an initial diffuse damage, then a first pattern of strain localisation in the vicinity of the yellow plane. The second structure,

highlighted as a pink plane, is initiated during Phase 2 (slow slip, blue MSEs) and largely develops during Phase 3 (fast slip, pink MSEs). Note however that the MSEs occurring during Phase 3 do not locate only in the vicinity of the pink plane, but also in the overlap volume between the yellow and pink planes, and on the yellow plane to a lesser extent. In addition, there are few yellow MSEs located on the pink plane, which suggests that shear faulting could have been initiated simultaneously on both planes, then the upper shear plane takes over the lower one and accommodates most of the rock shortening at the end of the experiment.

The above results are derived from the combined use of active ultrasonic and passive MS monitoring of the deformation process. Both monitoring techniques are based only on the picking of the time of arrival of the first phase in the recorded waveforms.

4.3.2. *Moment tensor analysis*

The first motion polarities and relative amplitudes of the waveforms recorded for a given MSE can be used to estimate its source mechanism, similar to the approach widely used in seismology to define the source mechanism of earthquakes. This method, generally known as the Moment Tensor Inversion (MTI), is implemented in the *Insite* software and is used here to characterise the focal mechanism of the recorded MSEs [37, 52, 53, 54]. However, in order to obtain reliable MTI results, the analysis must be restricted to MSEs of sufficiently high quality, which represent a relatively small subset of all the spatially located MSEs. The MTI has been carried out on all the MSEs located spatially. The results reported Figure 10 fulfil the additional criteria: (i) a spatial location error strictly lower than 5 mm; (ii) a mean

473 error factor lower than 17; (iii) an inversion quality index lower than 4.4;
 474 and (iii) a T-k error norm lower than 0.3. The mean error factor measures
 475 the difference between the amplitude residual and the estimated uncertainty
 476 in the original amplitude measurement. The inversion quality factor is based
 477 on the 6x6 covariance matrix and depends on the Green's functions used,
 478 rather than the amplitudes. It is computed from the sum of the squares of
 479 the elements of the covariance matrix. The T-k error norm is the RMSE of
 480 the errors on the deviatoric (T) and isotropic (k) parameters representing the
 481 source [23]. The threshold values of the mean error factor, inversion quality
 482 index and T-k error norm have been selected as the mean values obtained for
 483 the whole set of spatially located MSEs to which the MTI has been carried
 484 out. With such criteria, 42 MSEs have been selected: 11 MSEs in Phase
 485 1, 6 MSEs in Phase 2 and 25 in Phase 3. The average amplitude residual
 486 parameter for these 42 MSEs is 0.21, and the standard deviation is 0.08.

487 Figure 11 shows the spatial distribution of the 42 MSEs within the shale
 488 specimen. For each MSE, the detecting ultrasonic sensors covered a reason-
 489 able portion of the solid angle around it, which allowed for a reliable MTI. In
 490 this figure, MSEs #79 in Phase 1, #128 in Phase 2 and #259 in Phase 3 have
 491 been highlighted because they exhibit the largest location magnitude for each
 492 phase. For each of these three MSEs, the focal mechanism is represented by a
 493 focal sphere plot, i.e., the so-called *beachballs* widely used in seismology. The
 494 sensors that detected the MSE are represented by small discs in the *beach-*
 495 *balls*, with the convention that black and white discs represent compressional
 496 and dilatational first motion, respectively. The fault plane is calculated using
 497 the first-motion polarity of the P-wave picked in the waveform recorded by

each sensor that detected this MSE and is represented by red circles in each focal sphere plot. The orientation of the fault plane is consistent with that of the fault planes identified statistically by the spatial distribution of the MSEs and by the post-mortem observation of the sample.

The MTI procedure yields the focal mechanism of each MSE as a combination of three basic modes, with usually a dominant mode: ISO, stands for isotropic dilatation, DC for double-couple (shear), and CLVD for compensated linear vector dipole [44, 52]. Hudson's so-called T-k source-type plot ([23]) is well-suited to display such decomposition in an equal-area graph (Fig. 11) where the T stands for the deviatoric component of the mechanism (shear deformation, zero volume change) and k stands for the normal/isotropic component (volumetric deformation, either positive-explosive or negative-implosive).

Figure 12 reports graphically the results of the moment tensor decomposition of the selected high quality MSEs. Figure 12a shows the detailed decomposition of the focal mechanisms into Hudson's T-k source types; figure 12b shows the corresponding decomposition into DC, CLVD and ISO MSEs; figure 12c shows the corresponding decomposition into pure shear (DC) and non-shear (ISO+CLVD) MSEs; and figure 12d represents the fault plane orientation for the selected MSEs in terms of azimuth and dip angles. In each of these plots, the raw data as obtained through the moment tensor inversion are represented with coloured symbols; the corresponding coloured lines are obtained by a moving average procedure with a window size of 9 points. The purpose of these lines is only to identify possible trends during the three phases of the experiment. The T-k decomposition suggests that

at the early stages of Phase 1, damage is dominated by non-shear MSEs (k or ISO+CLVD MSEs). A transition toward more shear MSEs occurs during Phases 2 and 3 (Fig. 12a and c). During most of Phase 3, the most common mechanisms is double couple, and by the end of the experiment, all mechanisms (DC, CLVD and ISO) become equiprobable (Fig. 12b). Figure 12d suggests that the most probable dip angle of the fault plane is comprised between 30 and 60; the most probable azimuth angle is comprised between 120 and 180. These angles are qualitatively consistent with the macroscopic shear fractures observed on the rock specimen post-mortem. The variability of the dip and azimuth angles around their nominal values of 45 and 180 can be attributed to the experimental uncertainties associated with the measurements and the inversion procedure. They could also be related to the shear failure of small asperities on the fault plane that are not well aligned with the macroscopic failure plane (see discussion below).

5. Discussion

5.1. *Shear faulting in the laboratory and relay ramp structures in the field*

The post-mortem picture of the failed specimen and the location of the recorded MSEs are in good agreement (see Fig. 10). Although the picture of the specimen cannot show the internal structure of the shear faults, their emergence at the external boundary of the specimen is in agreement with the location of the MSEs at this boundary. Two different planar structures are identified from the spatio-temporal location of the 280 MSEs (Figs. 9 and 10).

These results suggest that the lower shear fault (yellow plane) is most

547 active (accommodates most of the imposed axial displacement) at the early
 548 stages after strain localisation (Phase 1), although few yellow events are
 549 already located on the top part of the upper shear fault. However, during
 550 this phase no clustering of MSE is observed on this upper plane. During
 551 Phase 2, a transition of the micro-seismic activity is observed from the lower
 552 shear fault toward the upper shear fault (pink plane). During Phase 3 most
 553 of the imposed axial displacement is accommodated by the upper shear fault,
 554 although few events are still located on the lower shear fault, indicating that
 555 it is not entirely inactive. This is consistent with the sequence of events
 556 associated with a typical relay ramp structure formed during the growth of
 557 normal fault systems in large scale geology

558 This upward transition from the lower to the upper SF is particularly
 559 visible in Figure 10 where the MSEs in each phase have been colour-coded
 560 according to their time of occurrence within the phase. More precisely, once
 561 the yellow SF is formed and its activity slows down at the end of Phase
 562 1, the blue MSEs of Phase 2 first appear at the lower end of the pink SF
 563 then the MS activity migrates upward along this SF and approaches the
 564 boundary of the specimen. Once the pink SF is largely developed, part of
 565 the MS activity (pink MSEs) locates in the overlap volume between the two
 566 SF planes. In summary, it seems that the lower shear fault forms first (yellow
 567 plane), before the micro-seismic activity (blue spheres) migrates upward and
 568 the upper shear fault forms and accommodates most of the subsequently
 569 imposed axial displacement (pink plane). This is essentially similar to typical
 570 sequence of events associated with either (i) the formation of a relay ramp
 571 structure during the growth of normal fault systems in large scale geology;

572 or (ii) fractures growing towards one another and overlapping.

573 *5.2. Silent failure, slow slip and slip rate dependency*

574 Phase 1 is quasi-aseismic (only 34 MSEs), at least in the 0.5 MHz range
 575 of sensitivity of the ultrasonic transducers used in the experiment (about 0.1
 576 to 1 MHz). This is surprising because Phase 1 corresponds to the failure of
 577 the clay-rich rock and contrasts with other sedimentary or crystalline rocks
 578 (e.g., sandstones, granites) for which large amounts of precursory MSEs are
 579 usually recorded prior to the macroscopic failure, and failure itself has been
 580 reported to generate a much stronger MS activity (thousands of events).
 581 Phase 2 of slow slip on the yellow shear fault induces very small amount of
 582 MS activity: clays might be acting as a lubricant on the fault(s) at that slip
 583 rate. Silent or almost silent failures have already been reported in materials
 584 being deformed close to the brittle ductile transition, for instance Carrara
 585 marble [43], or Volterra gypsum at room temperature [5]. In all cases, silent
 586 failures are accompanied by slow slip and stress drop, i.e. the macroscopic
 587 fault releases the stress too slowly for the rupture and the slip to accelerate
 588 and start radiating elastic waves. As such, slow failures can be viewed as
 589 quasi-static failures in the Griffith sense, i.e., the entire energy release rate is
 590 dissipated at the rupture tip into fracture surface, damage and plastic strain.
 591 Note that slow failures are not always silent, because at the microscopic
 592 scale, damage at the crack tip can actually also radiate elastic waves and be
 593 associated to MSEs, as for instance during quasi static fault growth in granite
 594 [33], slow failure in porous basalt [4] or shear or compaction band formation
 595 in sandstones [15, 17]. Hence, both the growth of macroscopic fracture and
 596 the accumulation of microscopic damage are silent in the frequency range

investigated in these experiments. This suggests that shale and clays are indeed potential good candidate to host slow slip within shallow accretionary prism [19, 22], or in the shallow section of continental faults [50].

In contrast, Phase 3 of slip acceleration from 1mm/h to 10mm/h, i.e. slip velocities slightly larger than that observed during slow earthquakes which are typically of the order of several tens of cm per year only [20], generates a significant amount of MS activity. During that fast slip phase, the AE rate and the slip are proportional so there seems to be a significant rate dependency of the lubrication potential of clays. In Figure 13, we can see that the slip acceleration triggers an instantaneous increase in the friction coefficient, which is typical of the direct effect [35]. After that, the fault first weakens with increasing slip, then starts to re-strengthen after a few millimetres of slip, exhibiting thus the typical velocity strengthening behaviour observed for clay minerals [38]. It is interesting to note that during that phase, nevertheless, numerous MSEs are observed, probably linked to the dynamic shear failure of small asperities on the fault plane, as demonstrated by the inverted focal mechanisms (see Fig 13). These observations highlight the possible interactions between small asperities and slow slip of a velocity-strengthening fault [3], which could be considered as a realistic experimental analogue of natural observations of non-volcanic tremors and (very) low-frequency earthquakes triggered by slow slip events [19, 21].

6. Conclusion

We have demonstrated that it is possible to apply laboratory techniques usually employed for monitoring micro-seismicity on reservoir or crystalline

621 rocks to anisotropic shale specimens. The data acquired during this triaxial
 622 experiment allowed us (i) to quantify the P-wave (dynamic) anisotropy of
 623 the shale and its evolution with stress; (ii) monitor the micro-seismic activ-
 624 ity occurring during failure and subsequent fault slip at two different rates.
 625 The gas permeability as well as the P-wave velocity data and their respective
 626 sensitivity to pressure suggest the existence of micro-cracks in this partially
 627 dry shale specimen at room conditions. Although these micro-cracks tend to
 628 close with increasing effective confining pressure. The spatio-temporal loca-
 629 tion of the MSEs recorded during the three phases of the experiment (failure,
 630 slow fault slip, fast fault slip) indicates that two shear fault planes where in
 631 competition after the initial strain localisation that occurred near the peak
 632 axial stress. The evolution of these two shear fault planes as derived from the
 633 micro-seismic monitoring is consistent with the sequence of events associated
 634 with a typical relay ramp structure formed during the growth of normal fault
 635 systems in large scale geology. The moment tensor inversion carried out on
 636 the highest quality MSEs suggests a transition from non-shear dominated
 637 to shear-dominated micro-seismic activity when the rock evolves from initial
 638 failure to larger and faster slip along the fault. The spatial orientation of the
 639 fault plane obtained on the highest magnitude MSE for each phase is con-
 640 sistent with the macroscopic orientation of the shear faults. The frictional
 641 behaviour of the shear faults highlights the possible interactions between
 642 small asperities and slow slip of a velocity-strengthening fault, which could
 643 be considered as a realistic experimental analogue of natural observations of
 644 non-volcanic tremors and (very) low-frequency earthquakes triggered by slow
 645 slip events.

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References

- [1] Amitrano D., Brittle-ductile transition and associated seismicity: Experimental and numerical studies and relationship with the b value, *Journal of Geophysical Research-Solid Earth*, **108**(B1), doi:10.1029/2001JB000680, 2003.
- [2] Amann F., Button E.A., Evans K.F., Gischig V.S., and Blümel M., Experimental study of the brittle behavior of clay shale in rapid unconfined compression, *Rock Mechanics and Rock Engineering*, **44**, 415 - 430, 2011.
- [3] Ariyoshi K., Hori T., Ampuero J.P., Kaneda Y., Matsuzawa T., Hino R. and Hasegawa A., Influence of interaction between small asperities on various types of slow earthquakes in a 3-D simulation for a subduction plate boundary, *Gondwana Research*, **16**, 534 - 544, 2009.
- [4] Benson P.M., Thompson B.D., Meredith P.G., Vinciguerra S. and Young R.P., Imaging slow failure in triaxially deformed Etna basalt using 3D acoustic emission location and X-ray computed tomography, *Geophysical Research Letters*, **34**(3), 2007.
- [5] Brantut N., Schubnel A. and Guéguen Y., Damage and rupture dynamics at the brittle-ductile transition: The case of gypsum, *Journal of Geophysical Research-Solid Earth*, **116**(B1), 1978 - 2012, 2011.
- [6] Chang S.-H., and Lee C.-I., Estimation of cracking and damage mechanisms in rock under triaxial compression by moment tensor analysis of acoustic emission, *International Journal of Rock mechanics and Mining Sciences*, **41**, 1069 - 1086, 2004.

- [7] Charalampidou E.-M., Hall S.A., Stanchits S., Lewis H., and Viggiani G., Characterization of shear and compaction bands in a porous sandstone deformed under triaxial compression, *Tectonophysics*, **503**, 8 - 17, 2011.
- [8] Charalampidou E.-M., Stanchits S., Kwiatek G., and Dresen G., Brittle failure and fracture reactivation in sandstone by fluid injection, *European Journal of Environmental and Civil Engineering*, **19**, 564 - 579, 2015.
- [9] Dautriat J., Sarout J., David C., Bertauld D., and Macault R., Remote monitoring of the mechanical instability induced by fluid substitution and water weakening in the laboratory, *Physics of the Earth and Planetary Interiors*, *This issue*.
- [10] David C., Dautriat D., Sarout J., Delle Piane C., Menéndez B., Macault R., and Bertauld D., Mechanical instability induced by water weakening in laboratory fluid injection tests, *Journal of Geophysical Research-Solid Earth*, **120**, 4171 - 4188, doi:10.1002/ 2015JB011894.
- [11] Delle Piane C., Dewhurst D.N., Siggins A.F. and Raven M., Stress-induced anisotropy in brine saturated shale, *Geophysical Journal International*, **184**, 897 - 906, 2011.
- [12] Dellinger J., and Vernik L., Do travel times in pulse transmission experiments yield anisotropic group or phase velocities?, *Geophysics*, **41**, 1774 - 1779, 1994.
- [13] Dewhurst D.N. and Siggins A.F., Impact of fabric, microcracks and

- 694 stress field on shale anisotropy, *Geophysical Journal International*, **165**,
695 135 - 148, 2006.
- 696 [14] Dewhurst D.N., Siggins A.F., Sarout J. and Raven M., Geomechanical
697 and ultrasonic characterization of a Norwegian Sea shale, *Geophysics*,
698 **76**, WA101 - WA111, 2011.
- 699 [15] Fortin J., Stanchits S., Dresen G., and Gueguen Y., Acoustic emission
700 and velocities associated with the formation of compaction bands in
701 sandstone, *Journal of Geophysical Research-Solid Earth*, **111**, B10203,
702 doi:10.1029/2005JB003854, 2006.
- 703 [16] Fortin J., Stanchits S., Dresen G., and Gueguen Y., Damage evolution,
704 acoustic emissions and elastic wave velocities in porous carbonate rocks,
705 *AGU Fall Meeting Abstracts*, #T23D-0539, 2009.
- 706 [17] Fortin J., Stanchits S., Dresen G., and Gueguen Y., Acoustic emissions
707 monitoring during inelastic deformation of porous sandstone: Compar-
708 ison of three modes of deformation, *Pure and Applied Geophysics*, **166**,
709 823 - 841, 2009.
- 710 [18] Guglielmi Y., Cappa F., Avouac J.-P., Henry P., and Elsworth D., Seis-
711 micity triggered by fluid injection-induced aseismic slip, *Science*, **348**,
712 1224 - 1226, 2015.
- 713 [19] Hirose H., Asano Y., Obara K., Kimura T., Matsuzawa T., Tanaka S.
714 and Maeda T., Slow earthquakes linked along dip in the Nankai subduc-
715 tion zone, *Science*, **330**, 1502 - 1502, 2010.

- [20] Ikari M.J., Ito Y., Ujiie K. and Knopf A.J., Spectrum of slip behaviour in Tohoku fault zone samples at plate tectonic slip rates, *Nature Geosciences*, 2015.
- [21] Ito Y., Obara K., Shiomi K., Sekine S. and Hirose H., Slow earthquakes coincident with episodic tremors and slow slip events, *Science*, **315**, 503 - 506, 2007.
- [22] Ito Y., Hino R., Kido M., Fujimoto H., Osada Y., Inazu D. and Mishina M., Episodic slow slip events in the Japan subduction zone before the 2011 Tohoku-Oki earthquake, *Tectonophysics*, **600**, 14 - 26, 2013.
- [23] Hudson J.A., Pearce R.G. and Rogers R.M., Source type plot for inversion of the moment tensor, *Journal of Geophysical Research-Solid Earth*, **91**, 765 - 774, 1989.
- [24] Johnston J.E. and Christensen N.I., Seismic anisotropy of shales, *Journal of Geophysical Research-Solid Earth*, **100**(B4, 5991 - 6003, 1995.
- [25] Josh M., Esteban L., Delle Piane C., Sarout J., Dewhurst D.N. and Clennell M.B., Laboratory characterisation of shale properties, *Journal of Petroleum Science and Engineering*, **88-89**, 107 - 124, 2012.
- [26] Kusunose K., Lei X., Nishizawa O. and Satoh T., Effect of grain size on fractal structure of acoustic emission hypocenter distribution in granitic rock, *Physics of the Earth and Planetary Interiors*, **67**, 194 - 199, 1991.
- [27] Le Gonidec Y., Sarout J., Wassermann J. and Nussbaum C., Damage initiation and propagation assessed from stress-induced microseismic

- 738 events during a mine-by test in the Opalinus Clay, *Geophysical Journal*
739 *International*, **198**, 126 - 139, 2014.
- 740 [28] Lei X., Kusunose K., Satoh T. and Nishizawa O., The hierarchical rup-
741 ture process of a fault: an experimental study, *Physics of the Earth and*
742 *Planetary Interiors*, **137**, 213 - 228, 2003.
- 743 [29] Lei X., Masuda K., Nishizawa O., Jouniaux L., Liu L., Ma W., Satoh T.
744 and Kusunose K., Detailed analysis of acoustic emission activity during
745 catastrophic fracture of faults in rock, *Journal of Structural Geology*,
746 **26**, 247 - 258, 2004.
- 747 [30] Lei X., Nishizawa O., Kusunose K., Cho A., Satoh T. and Nishizawa O.,
748 Compressive failure of mudstone samples containing quartz veins using
749 rapid AE monitoring: the role of asperities, *Tectonophysics*, **328**, 329 -
750 340, 2000.
- 751 [31] Lei X. and Satoh T., Indicators of critical point behavior prior to rock
752 failure inferred from pre-failure damage, *Tectonophysics*, **431**, 97 - 111,
753 2007.
- 754 [32] Lockner D.A., The role of acoustic emission in the study of rock fracture,
755 *International Journal of Rock Mechanics and Mining Sciences*, **30**, 883
756 - 899, 1993.
- 757 [33] Lockner D.A., Byerlee J.D., Kuksenko V., Ponomarev A. and Sidorin A.,
758 Quasi-static fault growth and shear fracture energy in granite, *Nature*,
759 **350**, 39 - 42, 1991.

- [34] Lowry A.R., Larson K.M., Kostoglodov V. and Bilham R., Transient fault slip in Guerrero, southern Mexico, *Geophysical Research Letters*, **28**, 3753 - 3756, 2001.
- [35] Marone C., Laboratory-derived friction laws and their application to seismic faulting, *Annual Review of Earth and Planetary Sciences*, **26**, 643 - 696, 1998.
- [36] Niemeijer A.R. and Spiers C.J., Influence of phyllosilicates on fault strength in the brittle-ductile transition: insights from rock analogue experiments, *Special publication-Geological Society of London*, **245**, 303, 2005.
- [37] Pettitt W.S., Baker C., Young R.P., Dahlström L.-O. and Ramqvist G., The assessment of damage around critical engineering structures using induced seismicity and ultrasonic techniques, *Pure and Applied Geophysics*, **159**, 179 - 195, 2002.
- [38] Safer D.M. and Marone C., Comparison of smectite-and illite-rich gouge frictional properties: application to the updip limit of the seismogenic zone along subduction megathrusts, *Earth and Planetary Science Letters*, **215**, 219 - 235, 2003.
- [39] Sarout J., Delle Piane C., Nadri D., Esteban L. and Dewhurst D.N., A robust experimental determination of Thomsen's δ parameter, *Geophysics*, **80**, A19 - A24, 2015.
- [40] Sarout J., Esteban L., Delle Piane C., Maney B. and Dewhurst D.N.,

- 782 Elastic anisotropy of Opalinus Clay under variable saturation and tri-
783 axial stress, *Geophysical Journal International*, **198**, 1662 - 1682, 2014.
- 784 [41] Sarout J. and Guéguen Y., Anisotropy of elastic wave velocities in de-
785 formed shales: Part 1–Experimental results, *Geophysics*, **73**, D75 - D89,
786 2008.
- 787 [42] Sarout J., Molez L., Guéguen Y. and Hoteit N., Shale dynamic proper-
788 ties and anisotropy under triaxial loading: Experimental and theoretical
789 investigations, *Physics and Chemistry of the Earth*, **32**, 896 - 906, 2007.
- 790 [43] Schubnel A., Walker E., Thompson B.D., Fortin J., Guéguen Y. and
791 Young R.P., Transient creep, aseismic damage and slow failure in Car-
792 rara marble deformed across the brittle?ductile transition, *Geophysical*
793 *Research Letters*, **33**(17), 2006.
- 794 [44] Šílený J. and Milev A., Source mechanism of mining induced seismic
795 events: Resolution of double couple and non double couple models,
796 *Tectonophysics*, **456**, 3 - 15, 2008.
- 797 [45] Stanchits S., Mayr S., Shapiro S. and Dresen G., Fracturing of porous
798 rock induced by fluid injection, *Tectonophysics*, **503**, 129 - 145, 2011.
- 799 [46] Talebi S., Boone T.J. and Eastwood J.E., Injection-induced microseis-
800 micity in Colorado shales, *Pure and Applied Geophysics*, **153**, 95 - 111,
801 1998.
- 802 [47] Thomsen L., Weak elastic anisotropy, *Geophysics*, **51**, 1954 - 1966, 1986.

- [48] Vernik L. and Liu X., Velocity anisotropy in shales: A petrophysical study, *Geophysics*, **62**, 521 - 532, 1997.
- [49] Warpinski N.R., Du J. and Zimmer U., Measurements of hydraulic-fracture-induced seismicity in gas shales, *Society of Petroleum Engineers*, **27**, SPE-151597-PA, doi:10.2118/151597-PA, 2012.
- [50] Wei M. and Kaneko Y., Liu Y. and McGuire J.J., Episodic fault creep events in California controlled by shallow frictional heterogeneity, *Nature Geoscience*, **6**, 566-570, 2013.
- [51] Wong T.F. and Baud P., The brittle-ductile transition in porous rock: A review, *Journal of Structural Geology*, **44**, 25-53, 2012.
- [52] Young R.P., Hazzard J.F. and Pettitt W.S., Seismic and micromechanical studies of rock fracture, *Geophysical Research Letters*, **27**, 1767 - 1770, 2000.
- [53] Young R.P. and Collins D.S., Seismic studies of rock fracture at the underground research laboratory, Canada, *International Journal of Rock Mechanics and Mining Sciences*, **38**, 787 - 799, 2001.
- [54] Young R.P., Collins D.S., Reyes-Montes J.M. and Baker C., Quantification and interpretation of acoustic emission and microseismicity at the underground research laboratory, Canada, *International Journal of Rock Mechanics and Mining Sciences*, **41**, 1317 - 1327, 2004.

823 **Tables**

Table 1: Nitrogen gas permeability of the partially saturated shale measured at three effective pressure states using a steady state method

Confining pressure (MPa)	Pore pressure (MPa)	Effective pressure (MPa)	Permeability $\times 10^{-19} \text{ m}^2$	Permeability $\times 10^{-7} \text{ mD}$
10	6	4	211	213
50	15	35	28.3	28.7
80	15	65	6.8	6.9

824 **List of Figures**

- 825 1 Experimental setup including (clockwise from top left): the
826 triaxial stress vessel; the rock specimen enclosed in a flexible
827 *Viton* sleeve, instrumented with 16 ultrasonic P-wave trans-
828 ducers and connected to 16 pulser-amplifiers and 4 strain gages;
829 the strain monitoring computer; and the ultrasonic/micro-
830 seismic monitoring computer. 45
- 831 2 Spatial location of the ultrasonic sensors represented around
832 the cylindrical shale specimen (right panel), and in an an-
833 tipodal equal-angle projection (left panel). Two ultrasonic
834 transducers became inoperative at the early stages of the ex-
835 periment (represented in red). The four strain gages attached
836 to the shale sample are also represented. 46
- 837 3 Triaxial loading path: (i) confining pressure loading to reach
838 the simulated in situ stress state of 10 MPa (green line); (ii)
839 axial loading up to the peak stress (36.77 MPa) and stress
840 drop (30.36 MPa) at a constant axial displacement rate of
841 1 mm/hour (plain red line); (iii) axial loading at a constant
842 displacement rate of 10 mm/hour during which the axial stress
843 variation in non monotonic (sudden increase to 31.09 MPa,
844 slower decrease to 29.92 MPa, then even slower increase to
845 reach a plateau at 32.82 MPa. The failed specimen obtained
846 after the experiment is pictured on the right hand side. 47

847	4	Evolution with time of the total axial and radial stresses, axial	
848		displacement and cumulative number of micro-seismic events	
849		during Phases 1 (A to B in yellow), 2 (B to C in blue) and 3	
850		(C to D in pink) of the experiment. Over all, the cumulative	
851		number of events is linearly related to the axial displacement	
852		rate, except temporarily after the increase in the imposed dis-	
853		placement rate from 1 to 10 mm/hour and until the axial stress	
854		reaches a plateau.	48
855	5	Evolution with time of the total axial and radial stresses, axial	
856		displacement and rate of micro-seismic activity during Phase	
857		1 (A to B in yellow), 2 (B to C in blue) and 3 (C to D in	
858		pink) of the experiment. The rate of micro-seismic activity	
859		(amplitude of the green curve) is strongly correlated with the	
860		imposed displacement rate (slope of the blue cuve in the lower	
861		graph) and the evolution of the axial stress (amplitude of the	
862		blue curve in the upper graph).	49
863	6	Stress-strain data during Phase 1 (A to B), 2 (B to C) and 3 (C	
864		to D) of the experiment. The orientation of the strain gauges	
865		with respect to the shale bedding and the specimen's axis are	
866		also shown. The strain recorded by the gauges illustrates the	
867		mechanical anisotropy of the shale.	50

- 868 7 Evolution of P-wave group (ray) velocity and anisotropy (Thom-
869 sen's ε parameter) with confining pressure and deviatoric load-
870 ing. P-wave velocity data are indicated with a 10% error bar
871 (+/- 5%). The magnitude of the P-wave velocity as a func-
872 tion of the propagation direction with respect to the bedding
873 illustrates the elastic anisotropy of the shale. This anisotropy
874 decreases significantly with increasing confining pressure, and
875 is virtually not sensitive to the axial stress, at least until the
876 strain localises in a shear fault and the specimen fails. The
877 non-linear variation of the P-wave velocity with confining pres-
878 sure up to about 4 MPa suggests the existence of damage in
879 the shale specimen at room pressure; the linear variation of the
880 P-wave velocity with confining pressure above 4 MPa suggests
881 the existence of intrinsic anisotropy, most likely associated the
882 preferred alignment of clay platelets/particles. 51
- 883 8 Spatial and temporal location of the ultrasonic sources shot
884 during the velocity surveys. The squares represent the nominal
885 position of the centre of the ultrasonic sensors; the spheres
886 represent the location of the sources obtained by inversion
887 using the selected velocity model. 52

888	9	Spatial and temporal location of the micro-seismic events recorded	
889		during the three phases of the deviatoric loading: in yellow	
890		for Phase 1; in blue for Phase 2; and in pink for Phase 3.	
891		The micro-seismic activity suggests the existence of two over-	
892		lapping shear fault planes. Part of the micro-seismic activity	
893		locates in the overlap volume between these two planes. A fea-	
894		ture similar to relay ramps observed in large scale structural	
895		geology.	53
896	10	Spatial and temporal location of the recorded MSEs separated	
897		into the three phases of the experiment. For each phase, the	
898		color of each event is scaled to its time of occurrence, i.e., first	
899		events of the phase in green and last events in red. These	
900		results suggest that the lower shear fault (yellow plane) is	
901		most active (accommodates most of the imposed axial dis-	
902		placement) at the early stages after strain localisation (Phase	
903		1). During Phase 2, a transition of the micro-seismic activ-	
904		ity is observed from the lower shear fault toward the upper	
905		shear fault (pink plane). During Phase 3 most of the imposed	
906		axial displacement is accommodated by the upper shear fault	
907		although few events are still located on the lower shear fault,	
908		indicating that it is not entirely inactive. This is consistent	
909		with the sequence of events associated with a typical relay	
910		ramp structure formed during the growth of normal fault sys-	
911		tems in large scale geology.	54

912	11	Spatial location, T-k decomposition in Hudson's diagram [23],	
913		and moment tensor solution of the MSE with largest magni-	
914		tude in each of the three phases of the experiment: MSE #79	
915		in Phase 1, MSE #128 in Phase 2 and MSE #259 in Phase 3.	55
916	12	Results of the moment tensor decomposition of the selected	
917		high quality MSEs: (a) detailed decomposition of the focal	
918		mechanisms into Hudson's T-k source types; (b) correspond-	
919		ing decomposition into DC, CLVD and ISO MSEs; (c) corre-	
920		sponding decomposition into pure shear (DC) and non-shear	
921		(ISO+CLVD) MSEs; and (d) fault plane orientation for the	
922		selected MSEs in terms of azimuth and dip angles.	56
923	13	Fault frictional behaviour and MS activity during Phases 2	
924		and 3. The fault slip is calculated from the measured post-	
925		failure axial displacement and the orientation of the fault plane	
926		determined post mortem to be approximately at 45° to the	
927		specimen's axis.	57

928 **Figures**

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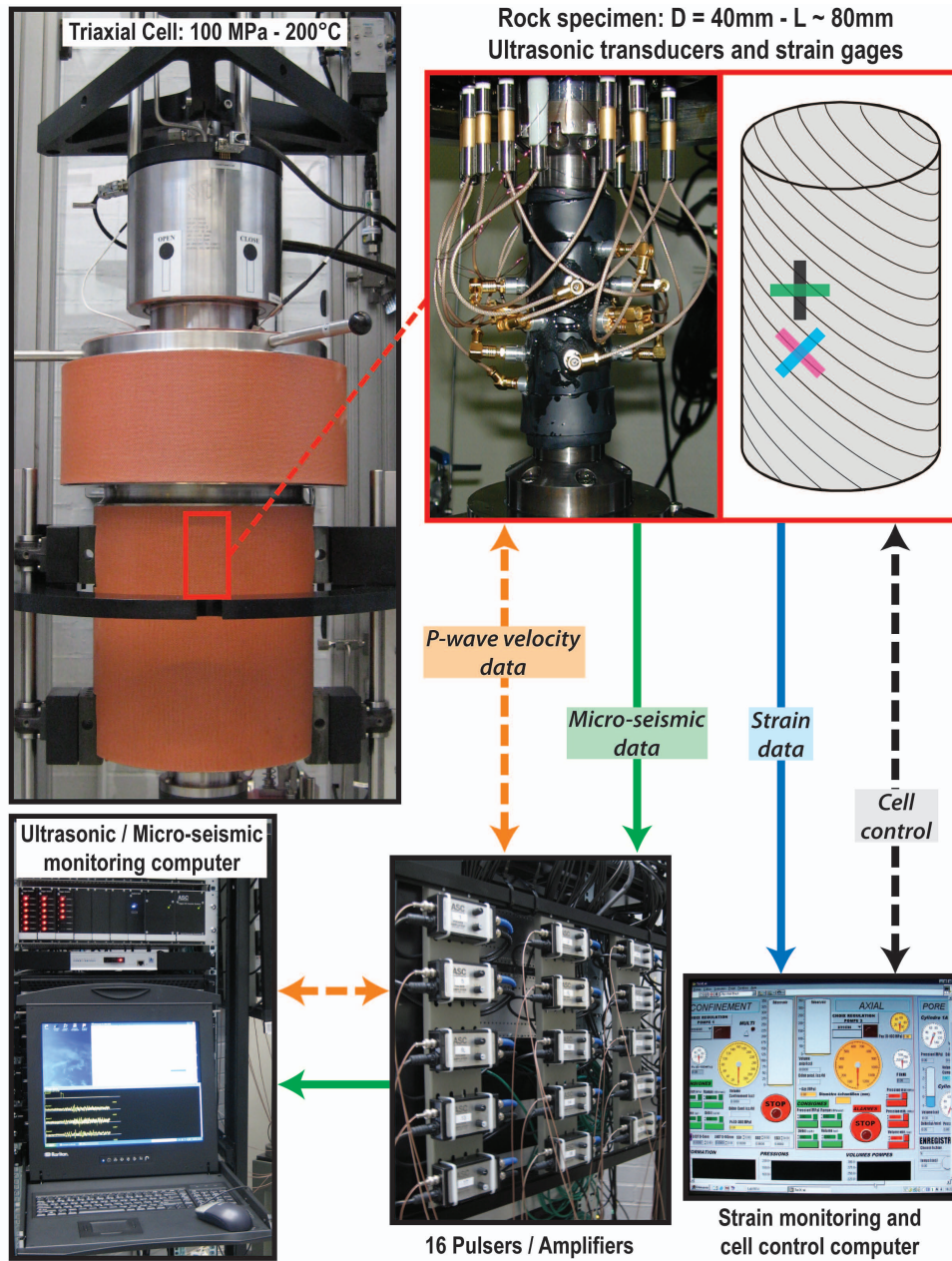


Figure 1: Experimental setup including (clockwise from top left): the triaxial stress vessel; the rock specimen enclosed in a flexible *Viton* sleeve, instrumented with 16 ultrasonic P-wave transducers and connected to 16 pulser-amplifiers and 4 strain gages; the strain monitoring computer; and the ultrasonic/micro-seismic monitoring computer.

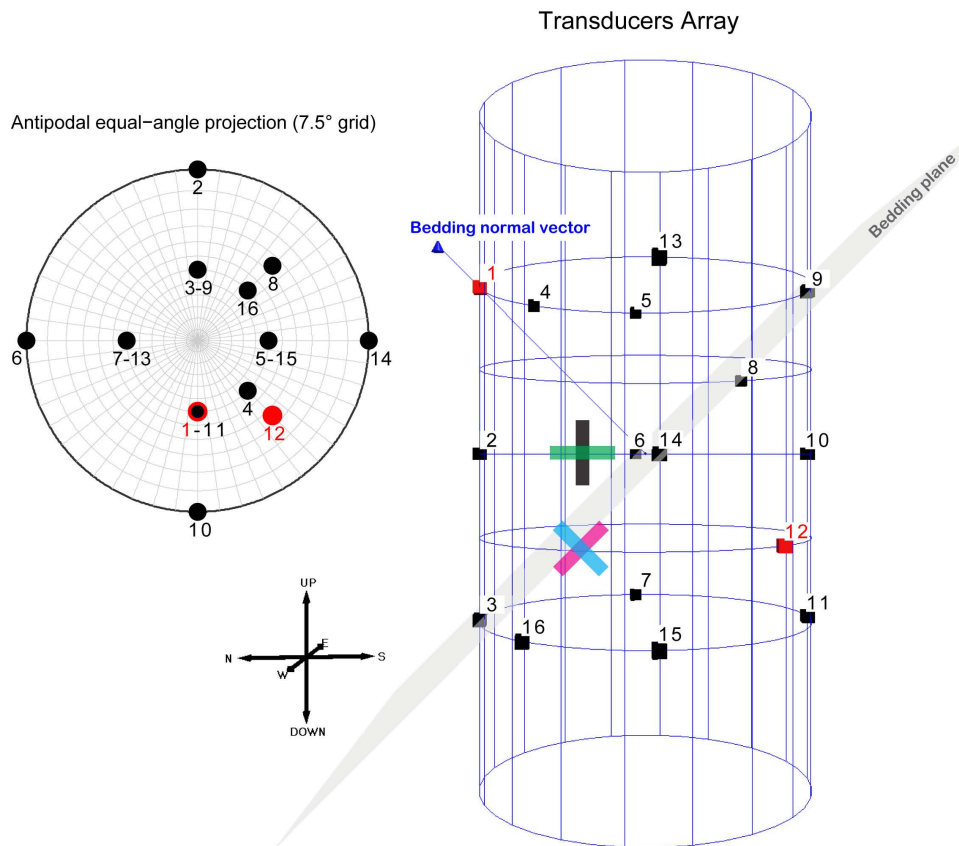


Figure 2: Spatial location of the ultrasonic sensors represented around the cylindrical shale specimen (right panel), and in an antipodal equal-angle projection (left panel). Two ultrasonic transducers became inoperative at the early stages of the experiment (represented in red). The four strain gages attached to the shale sample are also represented.

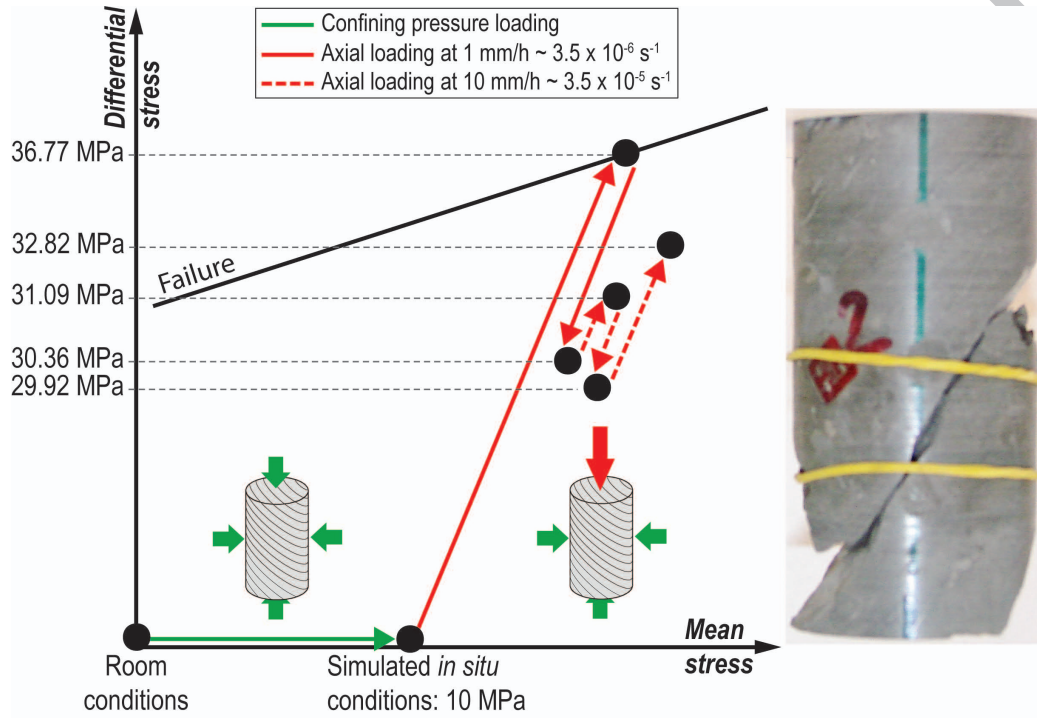


Figure 3: Triaxial loading path: (i) confining pressure loading to reach the simulated in situ stress state of 10 MPa (green line); (ii) axial loading up to the peak stress (36.77 MPa) and stress drop (30.36 MPa) at a constant axial displacement rate of 1 mm/hour (plain red line); (iii) axial loading at a constant displacement rate of 10 mm/hour during which the axial stress variation is non monotonic (sudden increase to 31.09 MPa, slower decrease to 29.92 MPa, then even slower increase to reach a plateau at 32.82 MPa). The failed specimen obtained after the experiment is pictured on the right hand side.

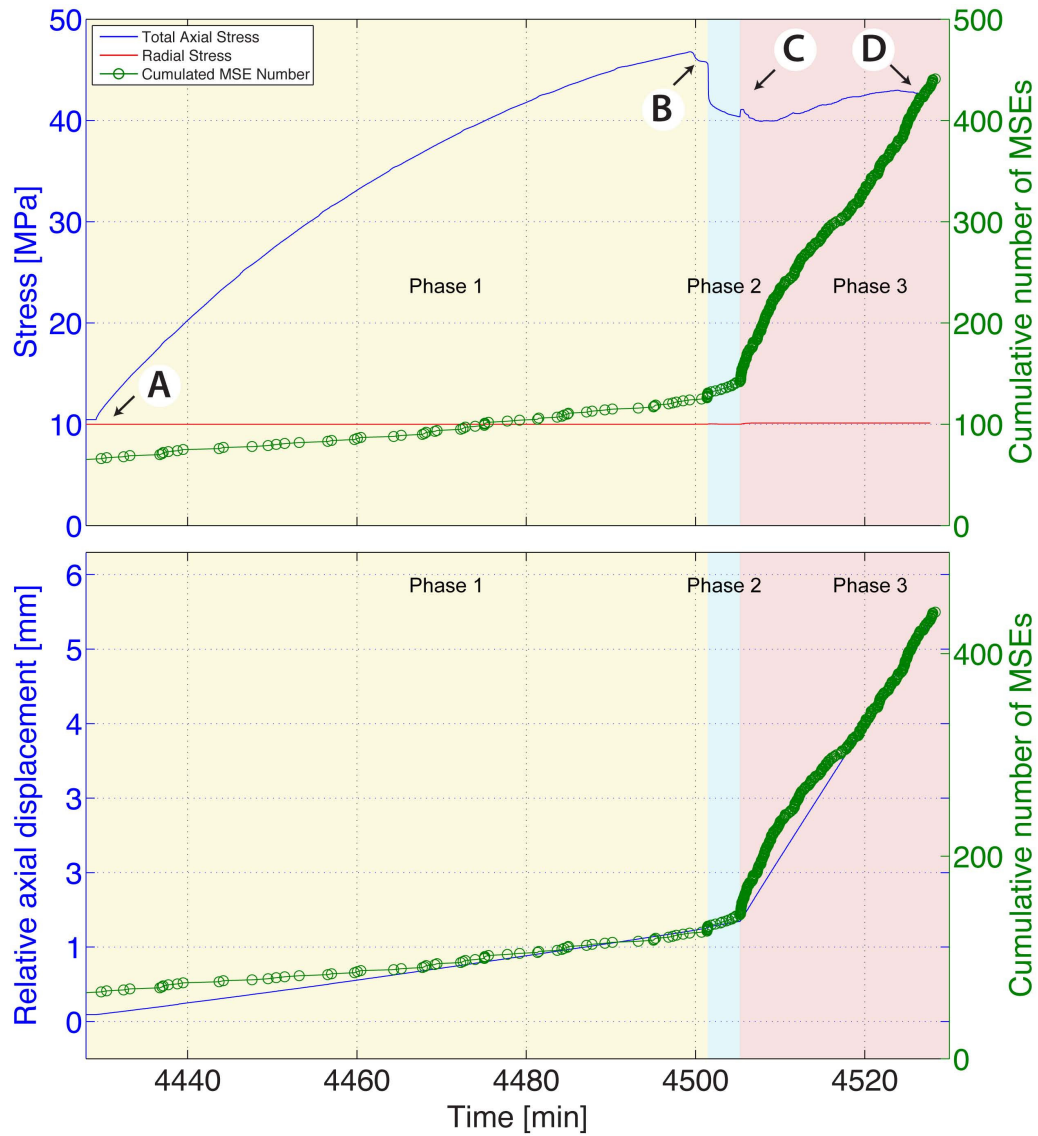


Figure 4: Evolution with time of the total axial and radial stresses, axial displacement and cumulative number of micro-seismic events during Phases 1 (A to B in yellow), 2 (B to C in blue) and 3 (C to D in pink) of the experiment. Over all, the cumulative number of events is linearly related to the axial displacement rate, except temporarily after the increase in the imposed displacement rate from 1 to 10 mm/hour and until the axial stress reaches a plateau.

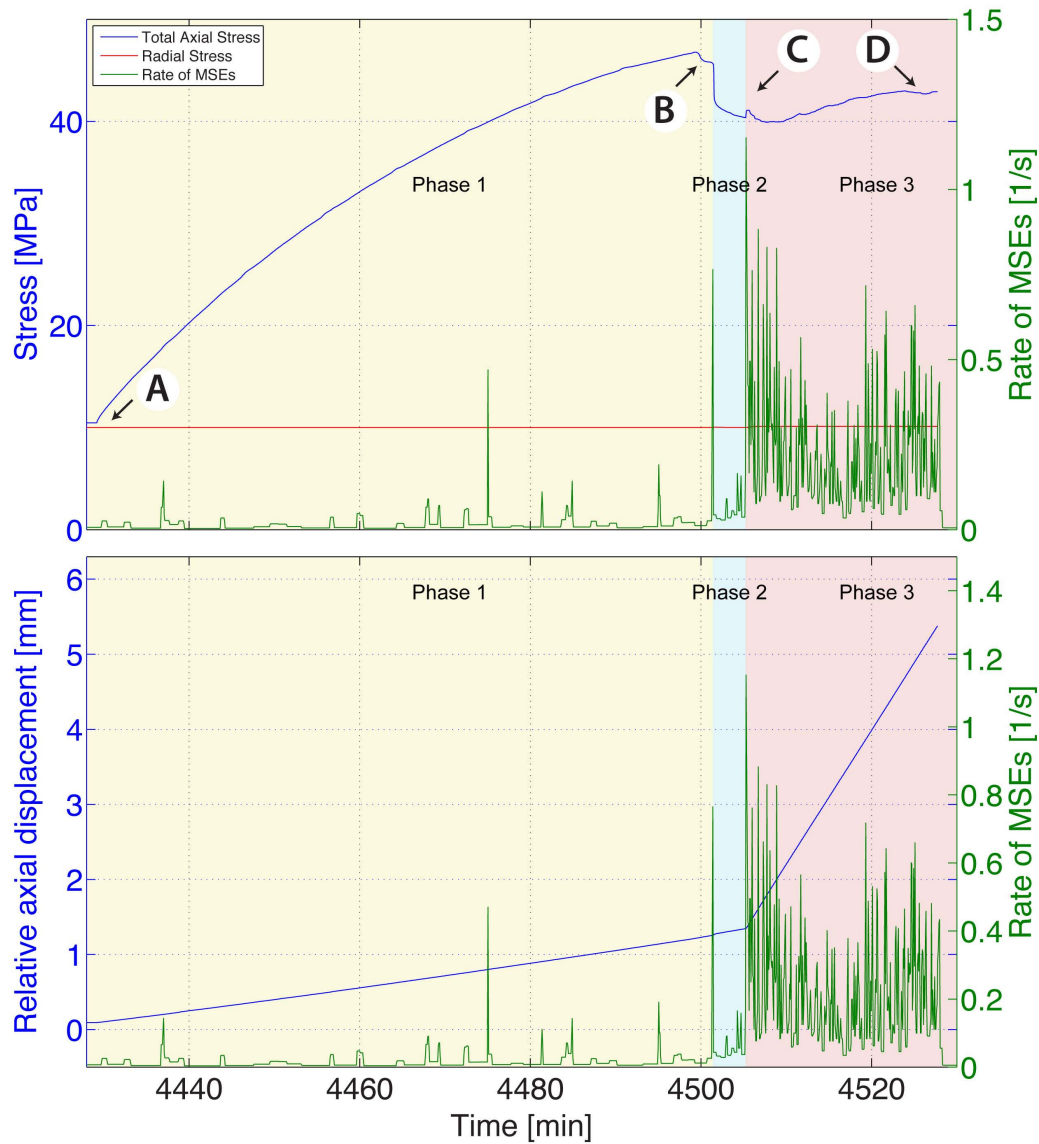


Figure 5: Evolution with time of the total axial and radial stresses, axial displacement and rate of micro-seismic activity during Phase 1 (A to B in yellow), 2 (B to C in blue) and 3 (C to D in pink) of the experiment. The rate of micro-seismic activity (amplitude of the green curve) is strongly correlated with the imposed displacement rate (slope of the blue curve in the lower graph) and the evolution of the axial stress (amplitude of the blue curve in the upper graph).

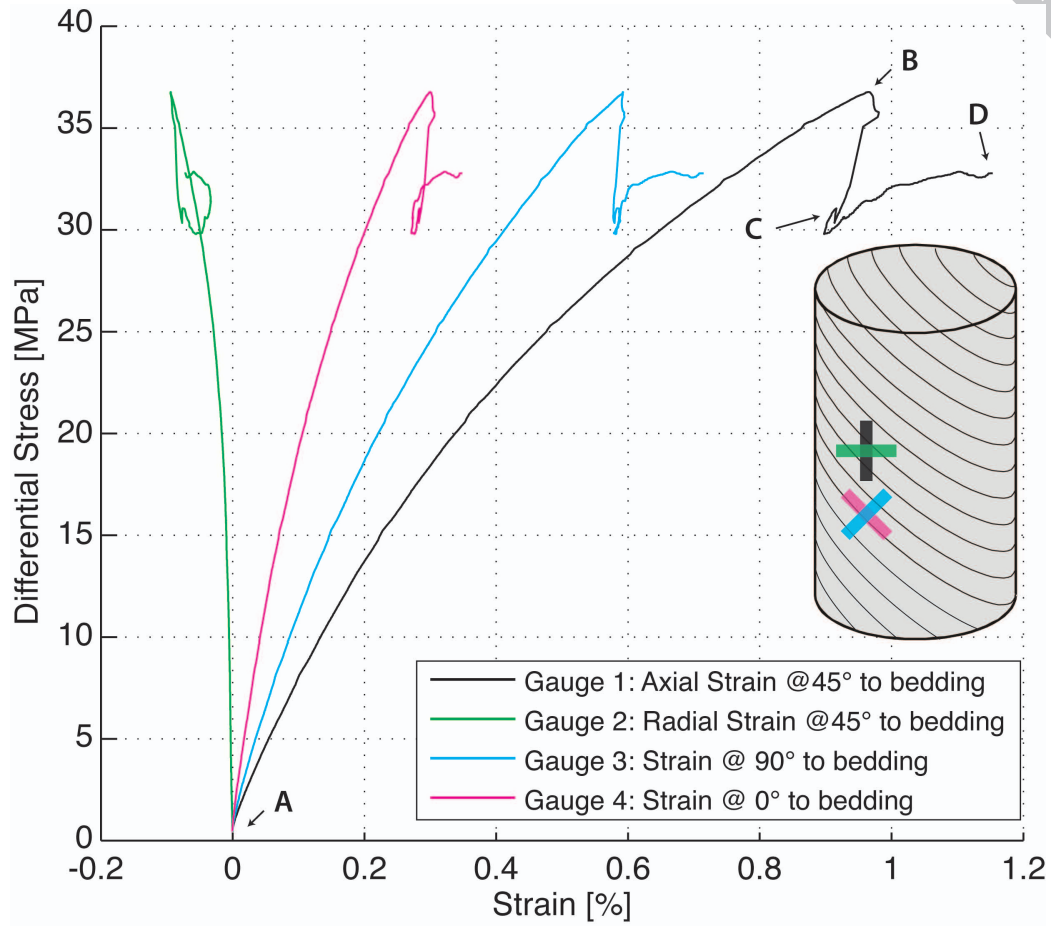


Figure 6: Stress-strain data during Phase 1 (A to B), 2 (B to C) and 3 (C to D) of the experiment. The orientation of the strain gauges with respect to the shale bedding and the specimen's axis are also shown. The strain recorded by the gauges illustrates the mechanical anisotropy of the shale.

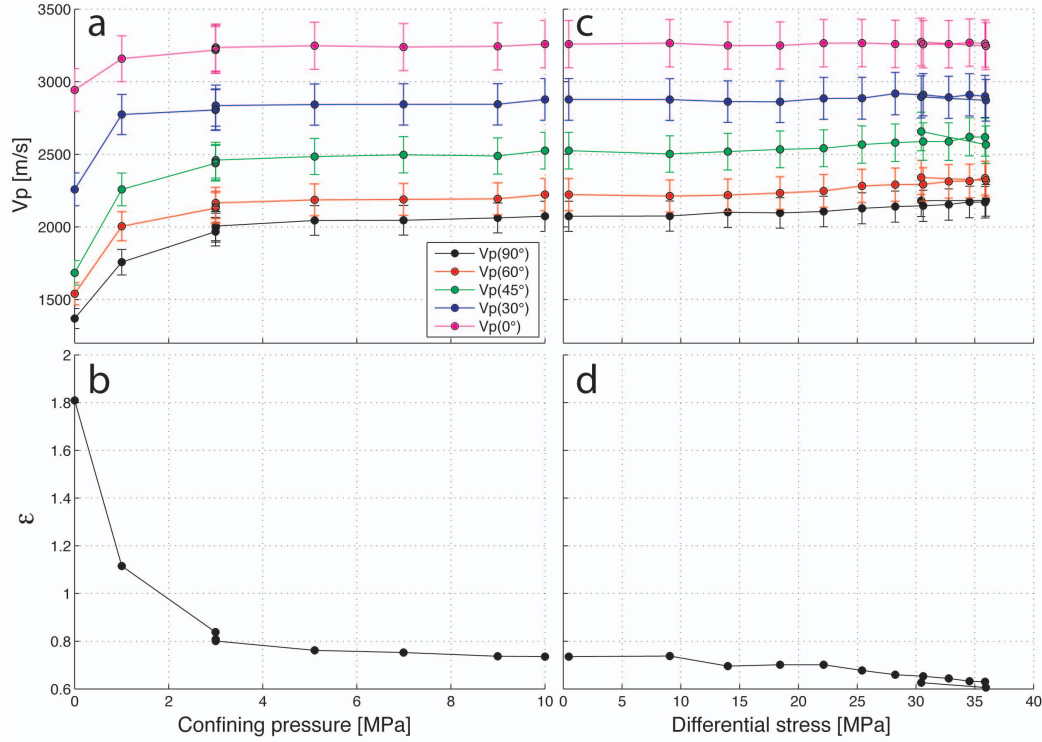


Figure 7: Evolution of P-wave group (ray) velocity and anisotropy (Thomsen's ε parameter) with confining pressure and deviatoric loading. P-wave velocity data are indicated with a 10% error bar (+/- 5%). The magnitude of the P-wave velocity as a function of the propagation direction with respect to the bedding illustrates the elastic anisotropy of the shale. This anisotropy decreases significantly with increasing confining pressure, and is virtually not sensitive to the axial stress, at least until the strain localises in a shear fault and the specimen fails. The non-linear variation of the P-wave velocity with confining pressure up to about 4 MPa suggests the existence of damage in the shale specimen at room pressure; the linear variation of the P-wave velocity with confining pressure above 4 MPa suggests the existence of intrinsic anisotropy, most likely associated the preferred alignment of clay platelets/particles.

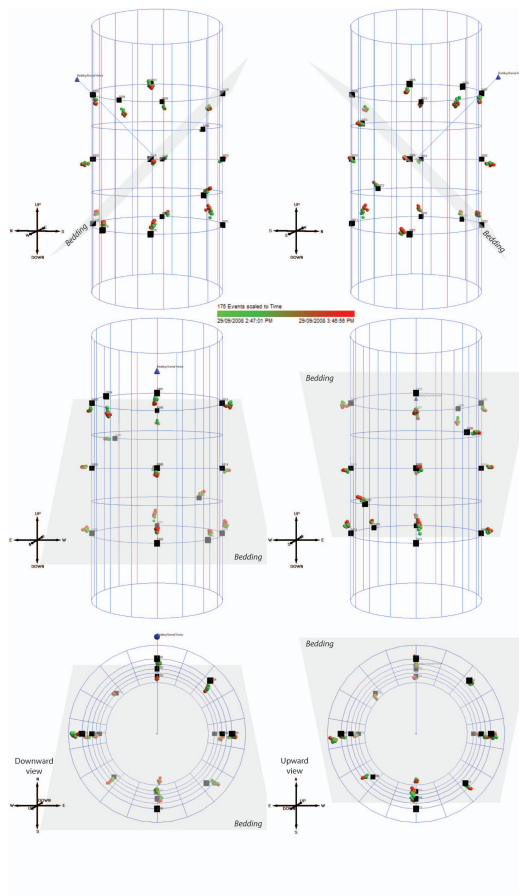


Figure 8: Spatial and temporal location of the ultrasonic sources shot during the velocity surveys. The squares represent the nominal position of the centre of the ultrasonic sensors; the spheres represent the location of the sources obtained by inversion using the selected velocity model.

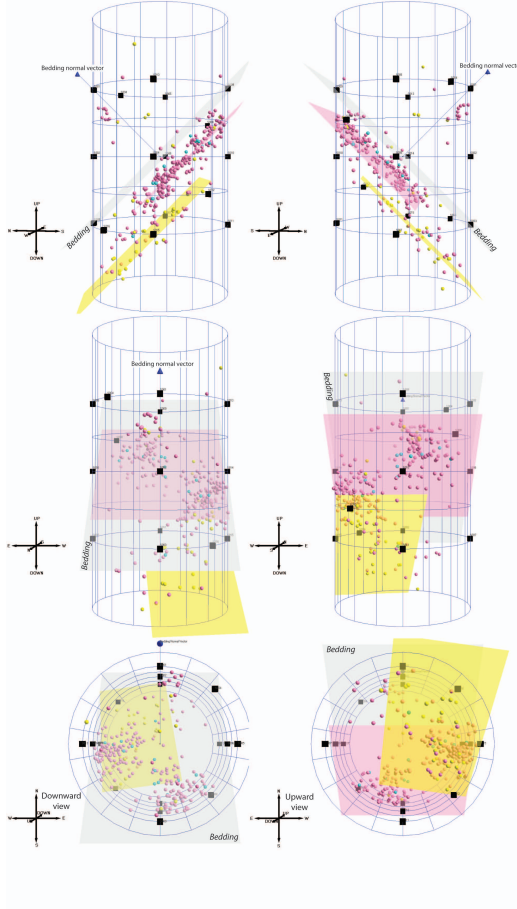


Figure 9: Spatial and temporal location of the micro-seismic events recorded during the three phases of the deviatoric loading: in yellow for Phase 1; in blue for Phase 2; and in pink for Phase 3. The micro-seismic activity suggests the existence of two overlapping shear fault planes. Part of the micro-seismic activity locates in the overlap volume between these two planes. A feature similar to relay ramps observed in large scale structural geology.

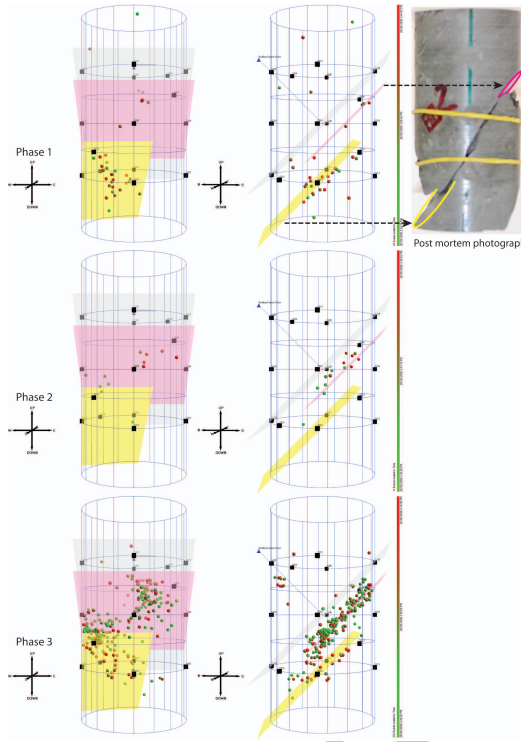


Figure 10: Spatial and temporal location of the recorded MSEs separated into the three phases of the experiment. For each phase, the color of each event is scaled to its time of occurrence, i.e., first events of the phase in green and last events in red. These results suggest that the lower shear fault (yellow plane) is most active (accommodates most of the imposed axial displacement) at the early stages after strain localisation (Phase 1). During Phase 2, a transition of the micro-seismic activity is observed from the lower shear fault toward the upper shear fault (pink plane). During Phase 3 most of the imposed axial displacement is accommodated by the upper shear fault although few events are still located on the lower shear fault, indicating that it is not entirely inactive. This is consistent with the sequence of events associated with a typical relay ramp structure formed during the growth of normal fault systems in large scale geology.

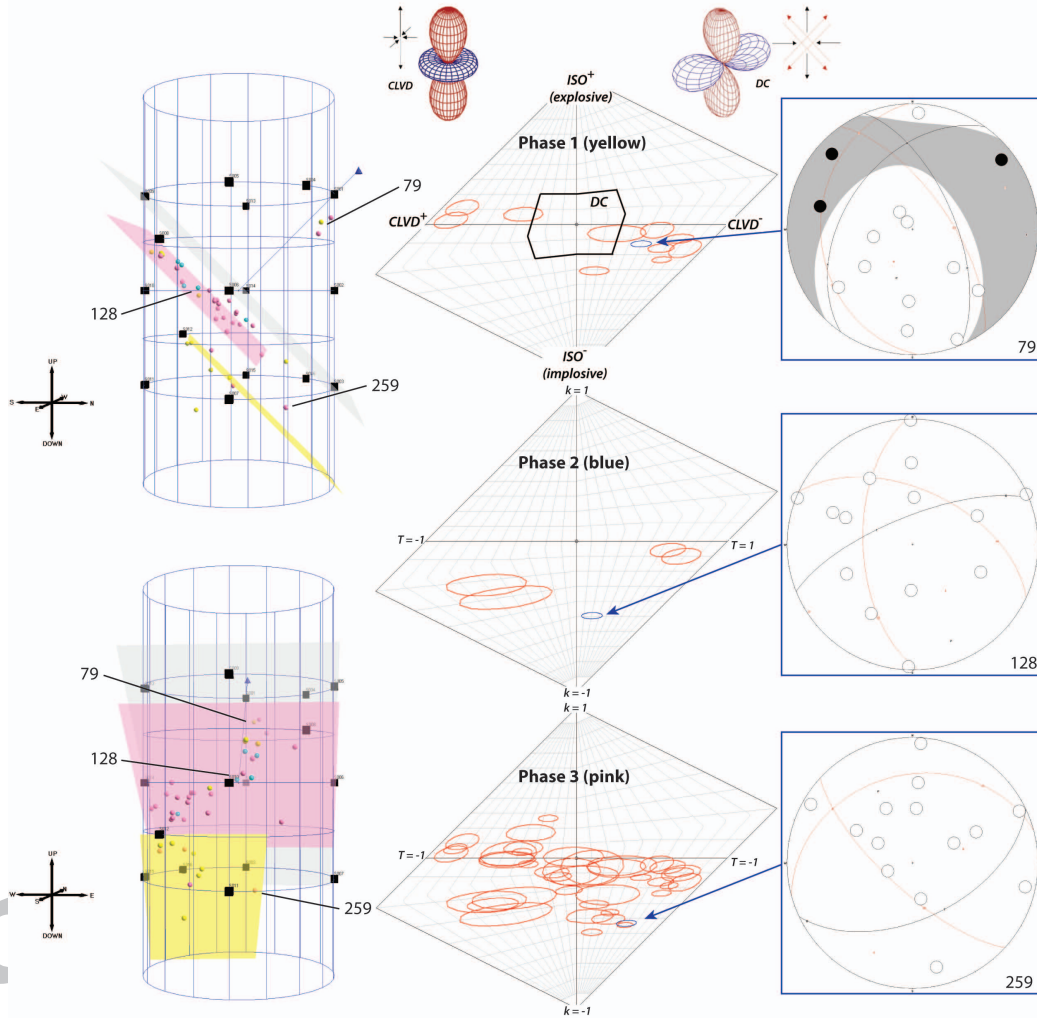


Figure 11: Spatial location, T-k decomposition in Hudson's diagram [23], and moment tensor solution of the MSE with largest magnitude in each of the three phases of the experiment: MSE #79 in Phase 1, MSE #128 in Phase 2 and MSE #259 in Phase 3.

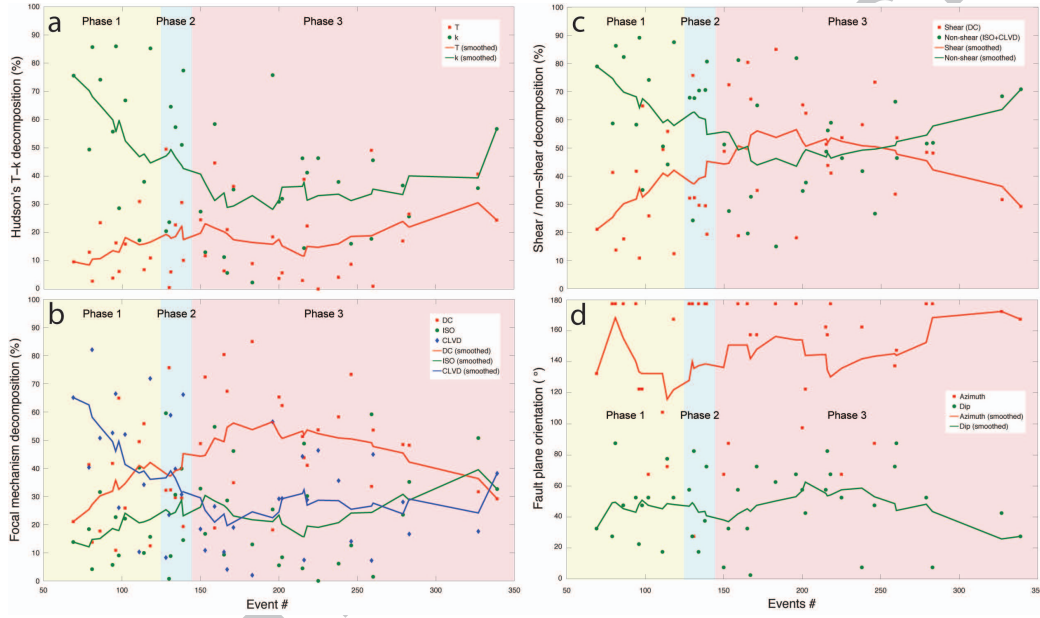


Figure 12: Results of the moment tensor decomposition of the selected high quality MSEs: (a) detailed decomposition of the focal mechanisms into Hudson's T-k source types; (b) corresponding decomposition into DC, CLVD and ISO MSEs; (c) corresponding decomposition into pure shear (DC) and non-shear (ISO+CLVD) MSEs; and (d) fault plane orientation for the selected MSEs in terms of azimuth and dip angles.

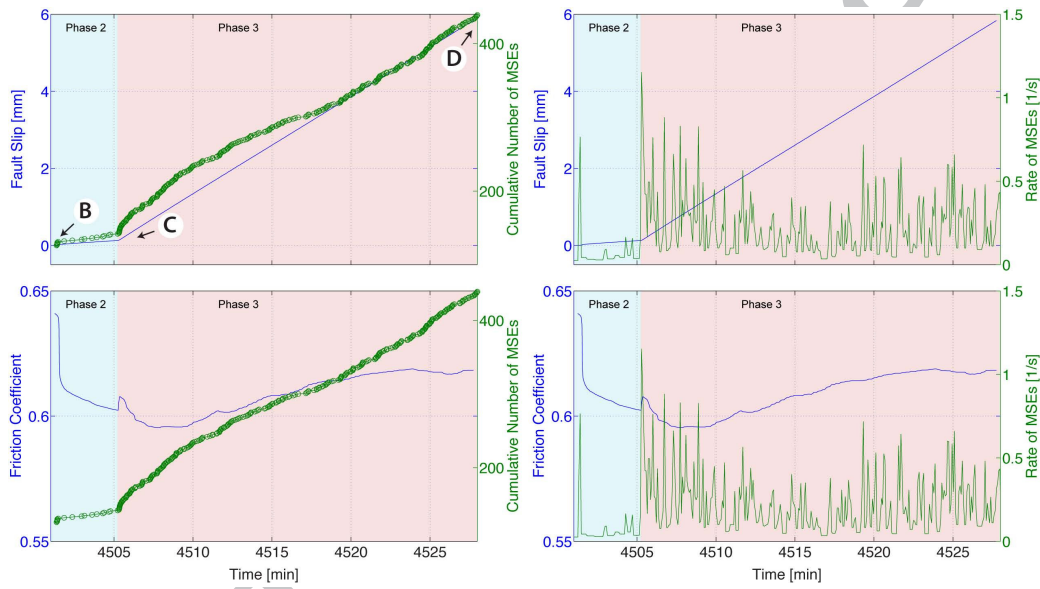


Figure 13: Fault frictional behaviour and MS activity during Phases 2 and 3. The fault slip is calculated from the measured post-failure axial displacement and the orientation of the fault plane determined post mortem to be approximately at 45° to the specimen's axis.