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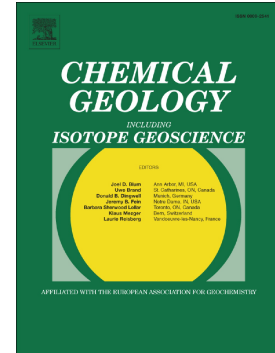
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X-RAY ABSORPTION SPECTROSCOPIC INVESTIGATION OF THE CA AND MG  
ENVIRONMENTS IN CO<sub>2</sub>-BEARING SILICATE GLASSES

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**Abstract**

Recent investigation showed that the CO<sub>2</sub> solubility in silicate melts is strongly affected by the #Mg (MgO/MgO+CaO). CO<sub>2</sub> solubility is decreasing with increasing #Mg implying that CO<sub>2</sub> molecules dissolves in the vicinity of Ca atoms to form CO<sub>3</sub><sup>2-</sup> molecular groups. We have investigated several CO<sub>2</sub>-bearing (up to 17.2 wt.%) silicate glasses using X-ray Absorption Spectroscopy (XAS) at the Mg and Ca K-edge in order to determine the structural environments of Mg<sup>2+</sup> and Ca<sup>2+</sup> cations. Analyses of the Mg XANES region show that the Mg environment is not affected by the presence of CO<sub>2</sub> dissolved as CO<sub>3</sub><sup>2-</sup> groups regardless of the CO<sub>2</sub> content. The Mg K-edge EXAFS simulations and XANES modelling show that the average Mg-O distance is close to 2 Å and the average Mg coordination is close to 6.

On the contrary, the position of the Ca XANES peak main resonance related to first-coordination shell is negatively correlated with the CO<sub>2</sub> content and shifts from 4051.4 to 4050.8 eV with CO<sub>2</sub> increasing from 0 to 17.2 wt.%. The Ca EXAFS simulations showed that Ca-O average distance is longer (~2.5Å) than the Mg-O distance. Furthermore, we observed that the Ca coordination number is positively correlated with the CO<sub>2</sub> content: Ca coordination number increases from 7 to 9 for CO<sub>2</sub> content changing from 0 to 17.2 wt.%.

Mg and Ca XAS results suggest that CO<sub>3</sub><sup>2-</sup> groups are dissolved in the vicinity of Ca<sup>2+</sup> cations and not in the surrounding of Mg<sup>2+</sup> cations which environment remains unaffected by CO<sub>2</sub> dissolution. The tighter Mg atomic environment as well as the lower coordination number for Mg cations as compared to Ca atomic environment could explain such behaviour. One implication of this result is the lower CO<sub>2</sub> transport capacity for Mg-rich silicate melts such as komatiites or kimberlites.

*Number of word: 293*

Keywords: CO<sub>2</sub> dissolution mechanism, XANES, EXAFS, Mg and Ca structural environment.

## 1. Introduction

Carbon dioxide (CO<sub>2</sub>) is an ubiquitous volatile species involved in many magmatic systems. CO<sub>2</sub> becomes dominant in some volcanic systems having very peculiar magmatic compositions which are low in silica such as melilitites or kimberlites. The study of these natural magmatic occurrences shows that CO<sub>2</sub> is known to dissolve in large quantities typically above 10 wt.% (Kjarsgaard et al., 2009; Russell et al., 2012; Bosshard-Stadlin et al., 2014). Previous experimental investigations (Brey and Ryabchikov, 1994; Brooker et al., 2001; 2011; Moussallam et al., 2015, 2016; Morizet et al., 2017a,b) emphasised that low silica melts such as melilitites and kimberlites are highly efficient in carrying CO<sub>2</sub> from the source at depth to the surface resulting in high CO<sub>2</sub> degassing.

Reconstruction of the CO<sub>2</sub> solubility (i.e. degassing) paths for these silicate melts has benefited from numerous experimental investigations since decades. For instance, the effect pressure and temperature induce on the CO<sub>2</sub> solubility is well-known (e.g. Blank and Brooker, 1994; Thibault and Holloway, 1994; Brooker et al., 1999; Morizet et al., 2002; Moussallam et al., 2015) and is well-reproduced by existing thermodynamic models (Newman and Lowenstern, 2002; Papale et al., 2006; Iacono-Marziano et al., 2012). The impact of redox conditions on CO<sub>2</sub> solubility is more complicated to apprehend (Pawley et al., 1992; Kadik et al., 2004; Morizet et al., 2010; Wetzel et al., 2013; Eguchi and Dasgupta, 2018) and further investigations are necessary to better constrain the CO<sub>2</sub> solubility change with  $fO_2$  variations. Many studies have focused on the effect of silicate melt compositions on CO<sub>2</sub> solubility and several layers of knowledge have to be distinguished. Whereas a clear effect of the SiO<sub>2</sub> content in silicate melt has been demonstrated on CO<sub>2</sub> solubility (Blank and Brooker, 1994;

Brooker et al., 1999; Massuyeau et al., 2015), Brooker et al. (2001) went further and showed that the CO<sub>2</sub> solubility is positively correlated to the increase in the Non-Bridging Oxygen per Tetrahedron ratio (NBO/T) calculated on the stoichiometric basis (Mysen, 1990). This parameter calculated from chemical composition links to the atomic structure and high NBO/T corresponds to a strongly depolymerized structure (i.e. more disrupted network). Therefore, the NBO/T also relates at some point to the CO<sub>2</sub> dissolution mechanisms: at NBO/T = 0, CO<sub>2</sub> will dissolve as CO<sub>2</sub> molecules in fully polymerized compositions (Brooker et al., 1999); at NBO/T < 1, CO<sub>2</sub> will dissolve as network carbonates (Brooker et al., 2001b); NBO/T > 1, CO<sub>2</sub> will form Free Ionic Carbonates (Morizet et al., 2017b,c). However, this is a first order approach as the calculation does not address the possible effect on CO<sub>2</sub> solubility of the cation nature and which was early suspected (Brey and Ryabchikov, 1994; Dixon, 1997); furthermore it only provides a partial picture of the structural molecular units within the glass and the fine molecular structure (i.e. different Al or Mg coordination numbers) is not depicted in the current NBO/T calculation.

The recent experimental work we have conducted aimed at investigating this second order of the effect of silicate melt chemical composition on CO<sub>2</sub> solubility. In two consecutive studies we showed that CO<sub>2</sub> solubility is affected by the #K = K<sub>2</sub>O/(K<sub>2</sub>O+Na<sub>2</sub>O) (Morizet et al., 2014) and #Mg = MgO/(MgO+CaO) (Morizet et al., 2017a). The CO<sub>2</sub> solubility in silicate melts increases with increasing #K and decreases with increasing #Mg. Such changes are related to the change in the cation field strength ( $z/r^2$  with  $z$  cationic charge and  $r$  ionic radius) which increases in the following order  $K^+ > Na^+ > Ca^{2+} > Mg^{2+}$ . Using a compilation of CO<sub>2</sub> solubility experimental data we have established a global parameter (Melt Ionic Field Strength dependent on the  $z/r^2$ ) which can predict the maximum CO<sub>2</sub> solubility in a continuum of silicate melt chemical composition from basalt to carbonatite. This result has major implications either on CO<sub>2</sub> degassing in volatile-rich magmatic systems or on the CO<sub>2</sub>

transport capability of silicate melts from the mantle source towards the surface. In particular, the impact of the high Mg concentration in natural kimberlite compositions onto CO<sub>2</sub> solubility and CO<sub>2</sub> degassing has never been addressed (Le Roex et al., 2003; Kjarsgaard et al., 2009; Russell et al., 2012; Kamenetsky et al., 2013), because most of the available experimental investigations have focused on Mg-poor silicate melt compositions.

One reason for the lower CO<sub>2</sub> solubility in Mg-rich silicate melt compositions is to be found in the CO<sub>2</sub> dissolution mechanisms into the melt structure. It has been demonstrated from NMR investigation that in low silica melts CO<sub>2</sub> dissolution uses available NBOs from the silicate network to form CO<sub>3</sub><sup>2-</sup> molecular groups which are totally disconnected from the silicate network and called Free Ionic Carbonates (FIC, Morizet et al., 2017b, c) in agreement with the early theory for CO<sub>2</sub> dissolution mechanisms (Eggler, 1978; Mysen and Virgo, 1980). Considering that CO<sub>2</sub> solubility was decreasing with increasing Mg content therefore it was suspected that CO<sub>2</sub> dissolves in the vicinity of Ca<sup>2+</sup> cations forming Ca<sup>2+</sup>..CO<sub>3</sub><sup>2-</sup> FIC. However, the preferred CO<sub>2</sub> dissolution in the surrounding of Ca<sup>2+</sup> cations instead of Mg<sup>2+</sup> cations remained to be explained. The present goal of this work is to determine if such CO<sub>2</sub> affinity is due to a difference in the structural environment of Ca<sup>2+</sup> and Mg<sup>2+</sup> or just a difference in the thermodynamic affinity.

Ca and Mg in silicate melts have similar behaviour and act principally as network modifying cations. However, in details, the structural environment of Mg<sup>2+</sup> is different from the Ca<sup>2+</sup> one reflecting the difference in the cation field strength. Most of the structural information on the atomic environments of these two elements was gathered from X-ray diffraction/absorption, neutron scattering and Solid-State Nuclear Magnetic Resonance (NMR); however, NMR spectroscopy reveals itself limited in providing structural information on these two particular elements. The extremely broad signal obtained from <sup>25</sup>Mg NMR in silicate glasses (Fiske and Stebbins, 1994; Kroeker and Stebbins, 2000; Shimoda et al., 2007) does not allow to fully



characterizing the Mg atomic environments and the clear discrimination of the Mg coordination in silicate glasses using  $^{25}\text{Mg}$  NMR spectroscopy is still to be improved. The expensive nature of the NMR active  $^{43}\text{Ca}$  isotopes has limited the number of structural investigations on that element (Angeli et al., 2007; Gambuzzi et al., 2015).

Mg and Ca X-ray Absorption Spectroscopy (XAS) have proven to be efficient in determining the atomic environment of these elements present either in minerals (e.g. Ildefonse et al., 1995; Paris et al., 1995; Andrault et al., 1998; Cabaret et al., 1998; Mottana et al., 1999; Quartieri et al., 2008; Fleet and Liu, 2009; Trcera et al., 2009; Yoshimura et al., 2013; Blanchard et al., 2016) or in silicate glasses (Taniguchi et al., 1997; Cormier and Neuville, 2004; Neuville et al., 2004; Guignard and Cormier, 2008; Trcera et al., 2009; 2011; Moulton et al., 2016; Ragoen et al., 2017). The application of Mg and Ca X-ray Absorption Near-Edge Spectroscopy (XANES) to a large variety of minerals has allowed establishing empirical curves for the Ca-O (Sowrey et al., 2004) and Mg-O (Li et al., 1999) coordination environments. Furthermore, the combined approach on both minerals and silicate glasses are often used to qualitatively determine the Ca and Mg structural environments in silicate glasses (Cormier and Neuville, 2004; Trcera et al., 2009). As for Extended X-ray Absorption Fine-Structure (EXAFS), it is currently applied to silicate glasses to determine 1) element coordination and 2) spatial relationship with nearest neighbours (Binsted et al., 1985; Jackson et al., 2005; Newville, 2014; Mastelaro and Zanotto, 2018). The recent study of Pohlenz et al. (2018) is the only one to show that the presence of  $\text{CO}_2$  molecules affects the average distance to oxygen for strontium and lanthanum. If the change in cationic structural environment as a function of the silicate glasses bulk chemical composition has been widely studied; up to now, no investigation was conducted on the effect of  $\text{CO}_2$  onto the Mg or Ca cationic environments in silicate glasses using XAS.

In the present study, we investigated the change in Ca and Mg structural environments in CO<sub>2</sub>-bearing silicate glasses using XAS (XANES and EXAFS). We studied silicate glasses with various compositions (from kimberlite to basalt) previously reported and having CO<sub>2</sub> concentrations up to 17 wt.%. As suspected, we observed that the Mg XANES in all investigated glasses is not affected by the presence of CO<sub>2</sub>. On the contrary, the Ca XANES spectra exhibit changes in peak positions with increasing CO<sub>2</sub> content which can be interpreted as a change in the Ca structural environments upon CO<sub>2</sub> dissolution. We conducted the analyses of the EXAFS spectra in order to extract the Ca and Mg structural parameters (bond length and coordination number to oxygens). A discussion is proposed in order to explain the CO<sub>2</sub> preferred dissolution in the surrounding of Ca<sup>2+</sup> and not Mg<sup>2+</sup> cations.

## 2. Experimental method

### 2.1 *CO<sub>2</sub>-bearing silicate glasses*

The elemental composition of the investigated glasses is reported in Table 1. The studied glasses are CO<sub>2</sub>-bearing silicate glasses with CO<sub>2</sub> content from 1.5 to 17.2 wt.%; an additional CO<sub>2</sub>-free wollastonite glass (CaSiO<sub>3</sub>) has been also studied for Ca environment, hence constraining the effect of CO<sub>2</sub> onto the Ca atomic environment. In all the investigated glasses, the CO<sub>2</sub> content has been determined using Raman spectroscopy and the method of Morizet et al. (2013). The accuracy on the determined CO<sub>2</sub> is 10% in relative to the value. The CO<sub>2</sub>-bearing silicate glasses were synthesised under high pressure conditions using piston-cylinder apparatus between 0.5 and 3.0 GPa for pressure and at temperature above 1400°C. The CO<sub>2</sub>-free wollastonite glass has been prepared in a Pt crucible placed in a 1 atm vertical furnace at 1500°C for several hours. Except for the TA series, the investigated glasses were studied

previously and the details of the synthesising method can be found in the subsequent publications. Most of the glass compositions are Fe-free in the Na<sub>2</sub>O-CaO-MgO-Al<sub>2</sub>O<sub>3</sub>-SiO<sub>2</sub> system. The RB8, HK, X and  $\beta$ MDE-5 were studied and reported in Morizet et al. (2017a). The TA samples are newly synthesised CO<sub>2</sub>-bearing glasses and were prepared from a natural peralkaline lamprophyre from Torre Alfina (Fe-bearing, Mitchell, 1989) to which carbonate (CaCO<sub>3</sub>, Na<sub>2</sub>CO<sub>3</sub>, CaMg(CO<sub>3</sub>)<sub>2</sub> and K<sub>2</sub>CO<sub>3</sub>; as a source of CO<sub>2</sub>) and oxides (SiO<sub>2</sub>, Al<sub>2</sub>O<sub>3</sub> and MgO) were added to attain the TA12 and TA10b major element compositions. TA12-C1, -C3 and TA10b-C1, -CH1 powders were loaded into Pt capsule and CO<sub>2</sub>-bearing glasses synthesised at pressure of 1.0 and 3.0 GPa using ½ inch piston-cylinder assemblies; temperature was set to 1400°C and monitored with type B thermocouple. The run duration was 30 min and the rapid quench (>100°C/s) was conducted by cutting-off the power at the end of the experiment.

The investigated compositions are depolymerized in nature. For instance, the NBO/T ranges from moderately depolymerized (NBO/T ~1) for  $\beta$ MDE-5 to strongly depolymerized (NBO/T > 3) for HKM and TA10b samples. One reason for studying depolymerizing silicate glass compositions is that these compositions dissolve a large amount of CO<sub>2</sub> in their structure (see Table 1) and the way CO<sub>2</sub> dissolves into the silicate melt structure is well-constrained (Brooker et al., 1999, 2001; Morizet et al., 2015, 2017a,b,c). We also reported the value of #Mg; considering that the concentration of Mg in glass has been inferred to be a major control of the CO<sub>2</sub> solubility (Morizet et al., 2017a).

The investigated database regroups samples synthesised at different pressure (up to 3.0 GPa), showing various CO<sub>2</sub> content (up to 17.2 wt.%), having different #Mg (up to 0.67); and different degrees of polymerization (i.e. 1 < NBO/T < 3.4). Hence, considering the large variety of glass samples, we expect to extract the effect of pressure, the effect of NBO/T and

#Mg on cation atomic environments, and by extension the CO<sub>2</sub> dissolution mechanisms and its effect on cation atomic environments.

## 2.2 XAS measurements

Mg and Ca K-edges XAS (EXAFS and XANES) spectra were collected on LUCIA beamline at SOLEIL synchrotron (Flank et al., 2006; Vantelon et al., 2016) using KTP and Si(111) monochromator, respectively. We used silicon drift diode detector and the spectra on the glass were acquired in fluorescence mode. Energy calibration was performed according to the Mg K-edge of periclase and Ca K-edge of diopside, respectively. The incoming beam size was 3x5 mm<sup>2</sup>. The samples were stick to a copper plate mounted onto an indium foil or with silver lacquer. The sample orientation had a 30° angle normal to the detector in order to increase the signal to noise ratio. Mg XANES spectra had to be corrected from self-absorption. We used the Athena Self-Absorption tool (Ravel and Newville, 2005) to do so. No self-absorption correction was needed for Ca XANES spectra.

The Mg XANES and EXAFS spectra were collected in the range 1250-1300 eV with a step of 2 eV, in the range 1300-1322 eV with a step of 0.1 eV, in the range 1322-1338 eV with a step of 0.5 eV and in the range 1338-1565 eV with a step of 1 eV. The Ca XANES and EXAFS spectra were collected in the range 3920-4035 eV with a step of 2 eV, in the range 4035-4070 eV with a step of 0.1 eV, in the range 4070-4125 eV with a step of 0.5 eV, in the range 4125-4500 eV with a step of 1 eV, in the range 4500-4800 eV with a step of 2 eV and in the range 4800-5000 eV with a step of 3 eV. At least two spectra were acquired on each sample in order to obtain a better signal to noise. The spectra were averaged using the Athena software (Ravel and Newville, 2005). The minimum counting time was set at 1s per point. In the present work the nomenclature for the spectra interpretation and peak assignments is the following: we used

italic capital letters for the peaks in crystalline material and small letters for the peaks in CO<sub>2</sub>-bearing silicate glasses.

In addition to the acquisition of the Mg and Ca K-edges XAS spectra on glass, several acquisitions were done on crystalline materials for comparison and fingerprinting. Mg K-edge spectra were collected for spinel (MgAl<sub>2</sub>O<sub>4</sub>), periclase (MgO), diopside (CaMgSi<sub>2</sub>O<sub>6</sub>) and San Carlo olivine (Mg<sub>2</sub>SiO<sub>4</sub>). In these crystalline compounds the Mg environment varies from 4-coordinated to 6-coordinated, covering the usually observed coordination number for Mg in silicate glasses (Taniguchi et al., 1997; Li et al., 1999; Kroeker and Stebbins, 2000; Kohara et al., 2004; Trcera et al., 2009; 2011; Wilding et al., 2012). Ca K-edge spectra for crystalline calcite (CaCO<sub>3</sub>), dolomite (CaMg(CO<sub>3</sub>)<sub>2</sub>), diopside (CaMgSi<sub>2</sub>O<sub>6</sub>) and anorthite (CaAl<sub>2</sub>Si<sub>2</sub>O<sub>8</sub>); therefore covering a large range of Ca coordination from 6 to 12, also in agreement with the usually considered Ca coordination number in silicate glasses (Taniguchi et al., 1997; Cormier and Neuville, 2004; Shimoda et al., 2008; Cormier and Cuello, 2013; Cicconi et al., 2016).

The acquired spectra were background subtracted using a polynomial function below the edge and then spectra were normalized in the EXAFS region using a cubic function. The EXAFS region could be turned into k space ( $\chi(k)$ ). For the EXAFS analyses, we used k<sup>2</sup>- and k<sup>3</sup>-weighted for Ca and Mg spectra, respectively. The simulation of the EXAFS region was done with the Artemis software and the DEMETER package (Ravel and Newville, 2005). The simulations were conducted in order to obtain the average distance to oxygen as well as the coordination number to oxygens; which are contained in the EXAFS signal for the first coordination sphere.

For the Mg EXAFS, the simulation is conducted with k ranging from 2 and up to 7 Å<sup>-1</sup> and R ranging from 1 and up to 2.2 Å. To simulate the Mg EXAFS of the CO<sub>2</sub>-bearing silicate

glasses, we used a single scattering path Mg to O from periclase standard crystalline structure: Coordination Number (CN) = 6,  $R_{\text{Mg-O}} = 2.11 \text{ \AA}$ . Because of the low number of independent points computed from the Fourier transform and fitting range, it was more appropriate from statistical point of view to only use a single path. Furthermore, simulations lead to relevant results and additional results were not necessary. Prior the simulation, the Mg EXAFS spectrum for periclase was simulated with theoretical single scattering factors obtained from a periclase crystallographic structure. We obtained the scattering amplitude  $S_0^2 = 0.85$  and a refined value of the E0 ( $\Delta E_0 = -2.06 \text{ eV}$ ). We kept these values for the simulation of the silicate glasses EXAFS spectra.

For the Ca EXAFS, the simulation is conducted with k ranging from 2 to  $8 \text{ \AA}^{-1}$  and R ranging from 1 to  $3 \text{ \AA}$ . We used several single scattering paths Ca to O from diopside to extract the first coordination sphere: O<sub>1</sub>, Coordination Number (CN) = 4,  $R_{\text{Ca-O}} = 2.36 \text{ \AA}$ ; O<sub>2</sub>, CN = 2,  $R_{\text{Ca-O}} = 2.56 \text{ \AA}$ ; O<sub>3</sub>, CN = 2,  $R_{\text{Ca-O}} = 2.72 \text{ \AA}$  and a single scattering path Ca to Si: CN = 2,  $R_{\text{Ca-Si}} = 3.09 \text{ \AA}$ . Prior to the simulation of the spectra for silicate glasses we conducted a simulation on the Ca EXAFS spectrum obtained on standard crystalline diopside and obtained the scattering amplitude  $S_0^2 = 0.977$  and the  $\Delta E_0 = 1.93 \text{ eV}$ . Due to the low number of independent points (~5 for 9 parameters), the optimization was accomplished in several steps: the average  $d_{\text{Ca-O}}$  for each path was optimized then followed by the optimization of the Debye-Waller attenuation factor for each path and finally an optimization of the coordination number (CN) was performed. The initial guesses were 0 for the deviation in the  $d_{\text{Ca-O}}$ , 0.003 for the Debye-Waller parameter and 4, 2 and 2 for the CN corresponding to the initial CN of the single scattering in the crystalline diopside. The three steps of the optimization were repeated several times in order to obtain a robust simulation of the EXAFS oscillation in k- and R-space. Considering that we keep the  $S_0^2$  and the  $\Delta E_0$  constant for both Mg and Ca EXAFS simulations therefore simulation results for silicate glasses can be compared in

relative to each other and the effect CO<sub>2</sub> could induce on the Mg and Ca atomic environments deciphered.

### 3. Results

Typical Mg and Ca K-edges spectra are shown in Figure 1 and 2, respectively. The entire set of spectra can be found in Supplementary material. Spectra for CO<sub>2</sub>-bearing silicate glasses are collated along with spectra for crystalline materials: periclase and spinel for Mg, dolomite and diopside for Ca. The spectra for CO<sub>2</sub>-bearing silicate glasses are categorized: A, for kimberlite or melilitite with NBO/T > 2; B, haplobasalt with NBO/T = 1, P = 1.5 GPa.

#### 3.1 *Mg XAS spectra*

The line shape of the Mg spectra in the XANES region for CO<sub>2</sub>-bearing glasses is different from the one in crystalline standards as seen in Figure 1. The Mg spectra in the XANES region are consistent with previous works for crystalline material (Ildefonse et al., 1995; Cabaret et al., 1998; Neuville et al., 2009; Trcera et al., 2009; 2011) and for Mg-bearing silicate glasses (Curti et al., 2009; Trcera et al., 2009; 2011; Moulton et al., 2016). Typically, the spectrum obtained for spinel is composed of three sharp peaks located at 1309.1, 1314.3, 1318.5 eV and two broad peaks at 1332 eV and 1350-1375 eV. The different peaks are denoted as *A*, *B*, *C*, *D* and *E* respectively (following the commonly used nomenclature, Trcera et al., 2009). For periclase, the same peaks can be retrieved but at different absorption energy: the *D* and *E* peaks are shifted to lower energy whereas the *A* peak is shifted to higher energy. Both spectra do not show any evidence of pre-edge peak such as the one observed for enstatite and pyrope (Figure 1B, Trcera et al., 2009). The lower energy for *D* and *E* peaks in periclase than for spinel would be related to the higher Mg coordination number for periclase than for spinel (Trcera et al., 2009).

For CO<sub>2</sub>-bearing silicate glasses in Figure 1, we observe that the XAS spectrum line shape is similar in between the investigated samples and comparable to previously reported spectra for silicate glasses (Ildefonse et al., 1995; Curti et al., 2009; Trcera et al., 2009; 2011; Moulton et al., 2016). The Mg spectrum in the XANES is a composition of several components and does not show pre-edge feature. The four peaks labelled *a* to *d* can be identified for silicate glasses (Li et al., 1996, Moulton et al., 2016). The comparison with crystalline spectra shows that the peak *C* and *D* are merged into *c* peak; and the latter is not always clearly observed for silicate glasses on the contrary to crystalline compounds. An assignment of these peaks has been proposed in previous work and attributed to different electronic transition: *a* peak is assigned to 1s to empty 3p; *b* and *d* peaks are related to the first coordination shell of Mg atoms, 1s → 3p first coordination shell electronic transition for *b* and first shell resonance of the Mg-O coordination for *d*; peak *c* is related to multiple scattering from coordination spheres around the absorbing Mg (Cabaret et al., 1998; Ildefonse et al., 1998; Li et al., 1999; Trcera et al., 2009; Moulton et al., 2016).

In between the Mg XAS spectra of the different CO<sub>2</sub>-bearing sample, the differences are not clearly visible. For instance, in Figure 1A, the spectra line shape is similar for the natural kimberlite TA10b-C1 and the synthetic kimberlite (HK). HK-3 spectrum appears slightly different from the other HK samples; this difference is not clearly explained: it could be due to an analytical artefact or it could be inherent to the sample itself and the presence of nanocrystals. Except for this particular sample, the HK sample spectra are similar in shape; the position of the labelled *a*, *b*, *c* and *d* appears comparable even though the samples have been synthesised at different pressure and have different CO<sub>2</sub> content.

For the haplobasalt compositions, the main difference observed is on the *a* and *b* peaks in between X1-2 and X3-2 spectra. There is a substantial increase in the *a* peak from X1-2 to X3-2 whereas the *b* peak appears to decrease in relative intensity. The main difference in



between these samples is #Mg which varies from 0.25 to 0.67 in between X1-2 and X3-2. Previous works (Trcera et al., 2009; Moulton et al., 2016) showed that the change in the *a* peak intensity could be related to the Mg content with an increase in the *a* peak intensity with decreasing Mg content. In Figure 1B, surprisingly it appears to be the other way around with an increase in the *a* peak with increasing Mg content. It can be suggested that there is a change in the Mg environment not related to the Mg content itself but to another factor. Considering that those samples have been synthesised at identical pressure (1.5 GPa) and that it is difficult to relate such a change with the Mg content, therefore the observed change is currently unexplained. Although, there is a difference in the CO<sub>2</sub> content (from 1.5 to 7.5 wt.%), it appears highly speculative to suggest that CO<sub>2</sub> is responsible for the change in the spectrum line shape as it has only been observed for these two particular samples. Notwithstanding, the change in the peak *a* intensity might correspond to a change in the Mg-O bond length and polyhedral distortion (Cabaret et al., 1998; Li et al., 1999; Trcera et al., 2009; 2011; Moulton et al., 2016).

Using the crystalline standards as a fingerprint to decipher the Mg environment in silicate glasses is currently difficult. The investigated silicate glasses exhibit Mg XANES spectra showing features which can be found in both the crystalline spinel and periclase spectra. For example, the *A* peak prominent in spinel spectrum is also visible in the periclase one but at slightly higher energy and the *a* peak in silicate glasses can be a combination of both. Similarly, peak *c* or *d* in silicate glasses are broad and could represent a combination of several Mg environments ranging from 4- to 6-coordinated as inferred by previous works (Wilding et al., 2012). However, one common aspect in between the Mg XAS spectra for CO<sub>2</sub>-bearing silicate glasses is that CO<sub>2</sub> content does not appear to induce a dramatic change on the Mg atomic environments. This is consistent with our current idea that CO<sub>2</sub> molecules

do not dissolve in the vicinity of  $\text{Mg}^{2+}$  cations or only to a limited extent and  $\text{CO}_2$  molecules will have stronger interactions with  $\text{Ca}^{2+}$  or  $\text{Na}^+$  cations (Morizet et al., 2017a).

### 3.2 Ca XAS spectra

The Ca XAS spectra are presented in Figure 2 showing the XANES and EXAFS regions. Ca XAS spectra for crystalline dolomite and diopside are comparable to previous work (Cormier and Neuville, 2004; Blanchard et al., 2016). In a similar way to the Mg XAS spectra, within each category of silicate glasses (Figure 2), the differences in between the Ca XAS spectra are weak and all spectra show almost identical line shape in the XANES and EXAFS regions. As a result, the possible effect  $\text{CO}_2$  will induce on the Ca environment in silicate glass will be difficult to establish. The method of fingerprinting with crystalline standard to investigate the effect  $\text{CO}_2$  will induce on the Ca atomic environment will also be limited as very subtle changes are suspected. The Ca XAS spectra for crystalline compounds and silicate glasses show the evidence of a pre-edge peak around 4040 eV (*p*) which corresponds to a  $1s \rightarrow 3d$  type transition (Combes et al., 1991; Cormier and Neuville, 2004). The pre-edge peak is present in all spectra of silicate glasses regardless of the  $\text{CO}_2$  content. The main resonance (*b*, assigned to  $1s \rightarrow n p$  transition) is observed at ~4051 eV for the silicate glasses. For crystalline compounds, the peak maximum position varies from 4048 to 4050 eV for dolomite and diopside, respectively. A first comparison in between Ca XAS for silicate glasses and minerals suggests that the Ca environments in silicate glasses is probably closer to the one in diopside than in dolomite; however, the presence of the shoulder (*a*) for Ca XAS spectra in silicate glasses before the main line could be explained by a possible combination of both Ca environments: 6- and 8-coordinated Ca (Skinner et al., 2012; Cormier and Cuello, 2013; Gambuzzi et al., 2015). Furthermore, the examination of the first EXAFS oscillation in

silicate glasses (*c*, broad peak above 4075 eV) also indicates that the Ca atomic environments in CO<sub>2</sub>-bearing silicate glasses is complex and a combination of multiple environments as inferred in previous works (Cormier and Cuello, 2013).

#### 4. Discussion

##### *4.1 Calcium and magnesium atomic environments from XANES in CO<sub>2</sub>-bearing silicate glasses*

We have investigated the Mg and the Ca XANES region with a particular attention to determine the position of the main line peak maximum (*b* in Figure 1 and 2) as a function of the CO<sub>2</sub> content. The results are reported in Table 2 and represented in Figure 3. The main line peak for Mg is related to the first coordination shell; the same applies to the Ca main line which is related to the 1s → n p transition associated to the first coordination shell. Therefore any sort of correlation with the CO<sub>2</sub> content could indicate to some extent a relationship between the CO<sub>2</sub> dissolution mechanism and the investigated atomic environment.

In Figure 3, the same scaling (1 eV) is considered for the Mg and Ca peak maximum and the error on the peak determination is ±0.1 eV. We observe that the main line peak maximum for Ca is negatively correlated to the CO<sub>2</sub> content. The peak maximum changes from 4051.4 eV for CO<sub>2</sub>-free Woll glass to 4050.8 eV for HKM sample with 13.2 wt.% CO<sub>2</sub>. The negative correlation appears to be non-linear. Although, there is a significant scatter at high CO<sub>2</sub> content; the observed changes are above the possible analytical error (±0.1 eV in this region). The change in the peak position appears not correlated to other parameters such as pressure or NBO/T. For instance, βMDE-5 and X3-2 samples synthesised at 2.0 and 1.5 GPa, respectively; show almost identical peak position at ~4051.2 eV for comparable CO<sub>2</sub> content.

For the HK series at constant NBO/T  $\sim$  2.5, there is a change in the peak position from 4050.9 to 4051.1 eV from HK-1 to HK-2.

Although we did not analyse a CO<sub>2</sub>-free Mg-bearing silicate glass, we believe that there is not a similar trend for Mg main resonance and within error, the Mg peak maximum of the Mg XANES spectra remains constant with increasing CO<sub>2</sub> content. For example, the change in the peak maximum does not exceed 0.3 eV from 1.5 (X3-2) to 13.2 wt.% CO<sub>2</sub> (HKM). Whereas a clear change is observed for the Ca XANES peak maximum with increasing CO<sub>2</sub> content, the Mg XANES peak maximum does not seem to be affected by the presence of CO<sub>2</sub> in the silicate glass structure. Therefore, the CO<sub>2</sub> dissolution influences the Ca atomic environment whereas the Mg atomic environment remains unaffected in the silicate glass structure. This observation appears consistent with our previous work showing that 1) increasing the Mg content in silicate melt composition induces a decrease in the CO<sub>2</sub> solubility and 2) CO<sub>2</sub> molecules dissolve preferentially in the vicinity of the Ca<sup>2+</sup> ionic environments (Morizet et al., 2017a).

We used the model of Li et al. (1999) using the Mg XANES spectra to constrain the Mg local structural environment in the investigated CO<sub>2</sub>-bearing silicate glasses and propose a possible coordination environment of the Mg atoms. The model of Li et al. (1999) uses the position of the *a* peak (Figure 1) to constrain the Mg-O distance in Å. The *a* peak line position was determined through the simulation of the Mg XANES spectra for both crystalline compounds and silicate glasses. Typical Mg XANES deconvolutions are provided in Supplementary material and results of the Li et al. (1999) model are shown in Figure 4. The position of the *a* line from the simulation is reported in Table 2 for silicate glasses and the corresponding Mg-O distance ( $d_{\text{Mg-O}}$ ) has been calculated with the equation given by Li et al. (1999). For the crystalline compounds, the *a* line was determined at 1309.2, 1310.09 and 1310.60 eV for spinel, periclase and diopside, respectively. The corresponding calculated  $d_{\text{Mg-O}}$  is 1.92, 1.99

and 2.02 Å. These values are in close agreement with previous works (Ildefonse et al., 1995; Li et al., 1999; Trcera et al., 2009). However, the energy of the *a* line is in contrast with the results obtained by Ragoen et al. (2018) for similar crystalline compounds. Ragoen et al. (2018) derived energy values for the *a* peak which is lower than the one determined in the present study: they determined an energy of ~1307.5 eV for diopside which is 3 eV below the one we obtained. This difference is possibly ascribed to an analytical measurement or simply difference in the studied crystalline materials.

The *a* line values derived for the silicate glasses are comprised in between the values determined for spinel and diopside and change from 1309.4 to 1310.4 eV for TA10b-CH1 and HKM, respectively (see Table 2). The  $d_{\text{Mg-O}}$  calculated ranges from 1.94 to 2.01 Å. The  $d_{\text{Mg-O}}$  values are in agreement with the determined values for silicate glasses based on different methods (Wilding et al., 2004; 2008; 2012; Shimoda et al., 2008; Kohara et al., 2011; Cormier and Cuello, 2013). In the investigated silicate glasses (Figure 4), the  $d_{\text{Mg-O}}$  values do not appear to be correlated to the CO<sub>2</sub> content. For example, the TA10b samples have high CO<sub>2</sub> content (>12 wt.%) and show a  $d_{\text{Mg-O}}$  at 1.94 Å; and the HKM sample has CO<sub>2</sub> content at 13.2 wt.% and have a  $d_{\text{Mg-O}}$  at 2.01 Å. Again, this observation tends to confirm that Mg environment changes are not due to the CO<sub>2</sub> dissolution considering that both  $d_{\text{Mg-O}}$  extrema show high CO<sub>2</sub> content. The origin of the change in  $d_{\text{Mg-O}}$  is not clearly established in the present work; however it can be due to several parameters such as synthesis conditions (i.e. pressure conditions; Wilding et al., 2012), the silicate glass chemical composition itself (Guignard and Cormier, 2008; Cormier and Cuello, 2013; Moulton et al., 2016) or the structural unit connectivity within the glass (corner- or edge-sharing; ring or non-ring configuration) and the nature of oxygens (BO and/or NBO) surrounding the Mg cations (Kohara et al., 2011). Notwithstanding, the observed range in the  $d_{\text{Mg-O}}$  is in favour of a multiple Mg coordination environment as inferred by previous works for silicate glasses

(Cormier and Cuello, 2013; Ragoen et al., 2018). Similarly, Kohara et al. (2011) showed that the  $d_{\text{Mg-O}}$  is related to the chemical strength which can be by extension ascribed to a change in the coordination number of Mg. According to Kohara et al. (2011), at  $d_{\text{Mg-O}} \sim 2 \text{ \AA}$ , the chemical strength relatively high indicative of a strong ionic bond to  $\text{Mg}^{2+}$  cation.

#### *4.2 Magnesium atomic environments from EXAFS in CO<sub>2</sub>-bearing silicate glasses*

The structural refinement of the Mg EXAFS reduced data for silicate glasses was conducted in order to obtain an estimate of the first coordination sphere of the Mg atoms: average coordination to oxygen atoms and average Mg-O bond length. Typical extracted Mg EXAFS obtained by Fourier transform of the k-weighted ( $k^3 \chi(k)$ ) oscillatory signal is shown in Figure 5 for HKM and X3-2. Because of the limited Mg EXAFS energy range (i.e. limited by Al absorption edge at  $\sim 1560 \text{ eV}$ ), the range in k-space to investigate the fine Mg environments is reduced. In Figure 5, the simulation of the signal has been limited in the range  $2.0$  to  $7.0 \text{ \AA}^{-1}$  in k-space and  $1.0$  to  $2.2 \text{ \AA}$  for radial distance R.

Results of the simulation for the CO<sub>2</sub>-bearing silicate glasses are reported in Table 2. The  $d_{\text{Mg-O}}$  from EXAFS varies from  $1.98 \pm 0.00$  to  $2.04 \pm 0.05 \text{ \AA}$  for X1-2 and HK1, respectively. The derived  $d_{\text{Mg-O}}$  values from EXAFS are consistent with the derived values from XANES and the model of Li et al. (1999); however, the value extracted from EXAFS data is longer than the one extracted from XANES data (see Table 2). Similar to XANES, the change in the  $d_{\text{Mg-O}}$  value appears uncorrelated to the CO<sub>2</sub> content confirming the initial hypothesis that CO<sub>2</sub> does not dissolve in the surrounding of Mg atoms and does not influence the atomic environment of Mg atoms. On average, the  $d_{\text{Mg-O}}$  is  $2.00 \pm 0.02 \text{ \AA}$  and this value is in agreement with previous works on silicate glasses. For example, Guignard and Cormier (2008) derived Mg-O bond length between  $1.98$  and  $2.05 \text{ \AA}$  for glasses synthesised in the MgO-Al<sub>2</sub>O<sub>3</sub>-SiO<sub>2</sub> system

at 1 bar. Based on Ab Initio Molecular Dynamic simulations, Vuilleumier et al. (2009) determined comparable value for  $d_{\text{Mg-O}}$ . The derived coordination number for Mg atoms for silicate glasses is also shown in Table 2. The Mg coordination varies from  $3.7 \pm 1.3$  to  $6.5 \pm 0.9$  for HK1 and TA10b-CH1. It should be pointed out that the derived Mg coordination for HK1 has a large error associated with a moderately satisfactory simulation (see Supplementary material). The coordination value of  $3.7 \pm 1.3$  appears anomalously low considering that in silicate glasses Mg atoms adopts a complex structural environment corresponding to a mixture of several structural environments ranging from 4 to 7 in coordination (Cormier and Cuello, 2013). Except for this outlier, all the derived coordination for Mg is consistent with previous investigations for silicate glasses. For instance, Wilding et al. (2012) determined an increase in the Mg coordination number from 4.6 to 5.7 for MgO-SiO<sub>2</sub> glasses with increasing pressure from 1 bar to 8.6 GPa. The varying coordination number for Mg atoms suggests that there is a mixture of Mg atom environments with different coordination ranging from 4 to 7 as inferred in previous works (Trcera et al., 2009; Shimoda et al., 2008b; Vuilleumier et al., 2009; Kohara et al., 2011; Wilding et al., 2012; Cormier and Cuello, 2013). We applied the bond valence theory equations from Brese and O'Keefe (1991) to estimate the Mg coordination number. The Mg valence ( $v_{\text{Mg-O}}$ ) is related to the bond length  $d_{\text{Mg-O}}$  with the following equation:

$$v_{\text{Mg-O}} = \exp\left(\frac{R_{\text{Mg-O}} - d_{\text{Mg-O}}}{b}\right) \quad \text{Eq. 1}$$

Where  $R_{\text{Mg-O}}$  is the bond valence parameter and is 1.693 for Mg and O;  $b$  is a constant 0.37 (Brese and O'Keefe, 1991). The coordination number for Mg can be extracted considering the Mg cation charge:  $\text{CN} = 2 / v_{\text{Mg-O}}$ . The CN for Mg using the  $d_{\text{Mg-O}}$  obtained from the XANES is slightly lower than the one obtained using the  $d_{\text{Mg-O}}$  obtained from the EXAFS: CN is comprised between 3.9 and 4.7 from XANES whereas CN is comprised between 4.3 and 5.1

from EXAFS. This latter range is in better agreement with the CN determined using the Mg EXAFS spectrum simulations, although it is slightly lower. Concluding if such small difference is due to the possible dissolution of CO<sub>2</sub> as CO<sub>3</sub><sup>2-</sup> in the surrounding of Mg atoms is currently not conceivable. We did not observe any correlation between the Mg coordination number and the CO<sub>2</sub> content. We obtained a coordination of 5.6±0.4 for βMDE-5 with 1.7 wt.% CO<sub>2</sub> and an identical coordination of 5.7±0.6 for HKM with 13.2 wt.% CO<sub>2</sub>. Hence, within error, all the investigated glasses have comparable Mg atomic environments (average Mg coordination = 5.7±0.7) regardless of the CO<sub>2</sub> content.

Recent atomistic investigations for molten carbonate structure (Vuilleumier et al., 2014; Du et al., 2018) showed that the pair-distribution function for C-O bond has distance below 2 Å; however, from geometrical consideration a CO<sub>3</sub><sup>2-</sup> molecular group will have a size exceeding the empty space surrounding the Mg atoms in silicate glasses as determined from EXAFS. Reverse calculation of the bond valence theory to derive a d<sub>Mg-O</sub> value from the Mg CN extracted from EXAFS simulations lead to d<sub>Mg-O</sub> values close to 2 Å. Furthermore, Kohara et al. (2011) suggest that the empty space in an Mg-rich glass will be close to 0 considering that MgO<sub>x</sub> polyhedra will organize with a ring corner-shared structure hence substantially decreasing the free volume within the glass structure. As a result, one possible explanation for the absence of CO<sub>2</sub> dissolution in the vicinity of Mg<sup>2+</sup> cations is likely to be related to the size of empty space or free volume for accommodating CO<sub>2</sub> molecules in the melt structure.

#### *4.3 Calcium atomic environments from EXAFS in CO<sub>2</sub>-bearing silicate glasses*

As shown earlier (Figure 3), the analysis of the Ca XANES spectra for silicate glasses points towards an effect of the CO<sub>2</sub> dissolution on the Ca environments. In a similar way to Mg, we treated the Ca EXAFS spectra for the investigated silicate glasses in order to determine the



average bond length in between the Ca atoms and the first nearest neighbour ( $d_{\text{Ca-O}}$ ) and provide an estimate of the Ca coordination number.

We show in Figure 6 the typical simulated Ca EXAFS spectra in  $k$ - and  $R$ -space. It can be observed that the EXAFS oscillation in  $k$ -space is well-reproduced up to  $8 \text{ \AA}^{-1}$ . In the range  $1.0$ - $3.0 \text{ \AA}$ , the radial distribution for the first Ca coordination shell is also adequately simulated. The derived average  $d_{\text{Ca-O}}$  (corresponding to the sum of  $d_{\text{Ca-O}}$  normalized to the coordination number) is  $2.51 \pm 0.41$  and  $2.53 \pm 0.43 \text{ \AA}$  for TA12-C3 and X3-2, respectively. The whole set of derived  $d_{\text{Ca-O}}$  is given in Table 2. The calculated range of  $d_{\text{Ca-O}}$  is  $2.48$  to  $2.55 \text{ \AA}$  for TA12-C1 glass and RB8-2 samples, respectively. The average error on the  $d_{\text{Ca-O}}$  determined from error propagation on the different single scattering  $d_{\text{Ca-O}}$  is  $\pm 0.32 \text{ \AA}$  which prevents from determining the real impact of the  $\text{CO}_2$  dissolution on the Ca-O bond length in the first coordination shell; hence because of this large error the  $d_{\text{Ca-O}}$  appears constant regardless of the  $\text{CO}_2$  content. According to Sowrey et al. (2004) such a range in  $d_{\text{Ca-O}}$  would correspond to an increase in the Ca CN of 2.

Studying the Ca coordination values derived from the simulation (see Table 2) reveals that there is an increase in the Ca CN with increasing  $\text{CO}_2$  content. This increase is reported in Figure 7. Although the simulation of the Ca EXAFS spectra is probably not unique and the associated error in the EXAFS simulation is rather large, we observe that the Ca coordination number for  $\text{CO}_2$ -free Woll glass is almost 2 units below the coordination number determined for  $\text{CO}_2$ -bearing silicate glasses:  $7.1 \pm 0.5$  for Woll glass and  $9.5 \pm 1.1$  for RB8-2. Nevertheless, the coordination value for Woll glass that we determined is probably overestimated; Taniguchi et al. (1997) reported a Ca coordination of 5.6 with a  $d_{\text{Ca-O}}$  of  $2.49 \text{ \AA}$  from EXAFS measurements; and more recently Skinner et al. (2012) reported a Ca coordination of 6 in wollastonite glass. Early work by Binsted et al. (1985) inferred that the change in Ca CN could be related to the presence of Mg which would increase the Ca CN. In the investigated

glasses, the increase in the Ca CN is independent of the Mg content and is attributed to the presence of CO<sub>2</sub> dissolved as CO<sub>3</sub><sup>2-</sup>. For instance, we determined an increase in Ca CN from 7.1 to 9.3 in two Mg-free silicate glasses: Woll glass and RB8E-3. The increase in the Ca CN appears not correlated to an increase in the d<sub>Ca-O</sub> distance as inferred from atomistic calculations from Katz et al. (1996) who showed that an increase in the Ca CN upon hydration is positively correlated to the increase in d<sub>Ca-O</sub>. The results of Sowrey et al. (2004) also showed that an increase of 2 units in the Ca coordination number would be accompanied by an increase in d<sub>Ca-O</sub>. In the present work, we do not obtain a sufficient accuracy in the determined d<sub>Ca-O</sub> to make the same conclusion (see Supplementary material).

Surprisingly, the increase in the Ca coordination value in Figure 7 appears non-linear with a strong increase in the Ca CN then followed by a plateau starting at ~5 wt.% CO<sub>2</sub>. This increase does not result from a change in the synthesis conditions (i.e. pressure) or the CaO content of the studied glasses as the higher Ca coordination values (~9) are observed for samples synthesised at various pressures (from 0.5 to 3.0 GPa) and with various CaO content (from ~15 to 40 wt.%). Furthermore, Mountjoy (2007) showed that the increase in Ca CN is only limited in the covered range of pressure. Consequently, the observed change in the Ca coordination number is ascribed to the dissolution of CO<sub>2</sub> molecules in the surrounding of the Ca atoms. However, the observed non-linearity in Figure 7 is currently not explained; whether this increase is due to the type of CO<sub>2</sub> dissolution mechanism is still to be clarified. Morizet et al. (2017b,c) showed that the CO<sub>2</sub> dissolution as Free Ionic Carbonates molecular groups is accompanied by a strong rearrangement of the silicate network structure (i.e. Q<sup>n</sup> distribution). One hypothesis is that the dissolution of the CO<sub>2</sub> molecules and the corresponding Q<sup>n</sup> change reach an optimal favoured energetic configuration to explain the plateau; however, further work such as First Principle Molecular Dynamic simulations (e.g. Vuilleumier et al., 2009) or Reverse Monte Carlo simulations (e.g. Cormier and Cuello, 2013) would be necessary to

confirm this point. The second hypothesis is that there are several stages in the CO<sub>2</sub> dissolution mechanism and the number of CO<sub>3</sub><sup>2-</sup> groups surrounding a Ca atom can be above 1. For instance, βMDE-5 glass with 1.7 wt.% CO<sub>2</sub> has a low Ca CN determined (8.0±0.6). For this particular sample investigated in Morizet et al. (2015), CO<sub>2</sub> molecules dissolve partially as Network CO<sub>3</sub><sup>2-</sup> molecular groups (a CO<sub>3</sub><sup>2-</sup> molecule intercalated between two tetrahedral units) whereas in RB8E-3 glass with 17.2 wt.% CO<sub>2</sub>, CO<sub>2</sub> is only dissolved as Free Ionic Carbonate group as demonstrated by Morizet et al (2017b). The proposed chemical reaction to form Free Ionic Carbonates in Morizet et al. (2017c) based on <sup>29</sup>Si NMR results also suggests that more than one CO<sub>2</sub> molecule can be present in the surrounding of Ca atom and could explain the non-linear increase in the Ca CN.

As compared to Mg, the Ca structural environments in silicate glasses have longer bond length: ~2 Å for d<sub>Mg-O</sub> and ~2.5 Å for d<sub>Ca-O</sub>; and higher coordination values: <6 for Mg and >7 for Ca. From geometrical point of view, such differences could explain the preference of CO<sub>2</sub> molecules for dissolving in the vicinity of the Ca atoms rather than Mg atoms. Such a difference in the bond length to oxygen could also reflect a difference in the nature of the surrounding oxygen, either bridging (BO) or non-bridging (NBO). It has been shown (Kelsey et al., 2008; Stebbins, 2016) that Ca and Mg in silicate glasses are surrounded by both types of oxygens. For instance, Cormier and Cuello (2013) showed that increasing the #Mg induces an increase in the BO and a decrease in the NBO. Considering that CO<sub>2</sub> molecules scavenge NBOs to form free CO<sub>3</sub><sup>2-</sup> groups in the investigated silicate glass compositions, therefore the lower concentration of NBO in the surrounding of the Mg atoms could prevent CO<sub>2</sub> molecules to dissolve close to Mg atoms.

## 5. Summary

It was recently suggested that CO<sub>2</sub> solubility in silicate glasses decreases with increasing #Mg and a possible explanation for that behaviour is either the lower affinity of CO<sub>2</sub> molecules for or a particular Mg atomic environment in silicate glasses. In the present study, we investigated CO<sub>2</sub>-bearing silicate glasses using Mg and Ca K-edge XAS spectroscopy (XANES and EXAFS) at a synchrotron beamline. The studied glasses cover a wide range of chemical composition (from basalt to kimberlite) and a large range of CO<sub>2</sub> content (up to 17 wt.%), various synthesis conditions (P up to 3.0 GPa). For the investigated samples, the CO<sub>2</sub> dissolves as CO<sub>3</sub><sup>2-</sup> molecular groups.

The XAS results show that the Mg atomic environments are not affected by the dissolution of CO<sub>2</sub> molecules whereas Ca atomic environments are changing upon CO<sub>2</sub> dissolution. These results corroborated the initial hypothesis which suggested that CO<sub>2</sub> molecules were dissolving preferentially in the Ca vicinity and not in the surrounding of the Mg atoms in silicate glasses.

The analysis of the fine structure with EXAFS revealed that the Ca coordination number is higher than the Mg coordination number in the structure and that it increases with increasing CO<sub>2</sub> content. The distance to the first neighbour is shorter for Mg than for Ca: ~2 Å for d<sub>Mg-O</sub> and ~2.5 Å for d<sub>Ca-O</sub>. Such geometrical differences could explain the CO<sub>2</sub> solubility behaviour as a function of #Mg. In other words, the Mg atomic environment in CO<sub>2</sub>-bearing silicate glasses is too tight to allow the accommodation of CO<sub>2</sub> molecules in its surrounding.

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**Figure caption**

Figure 1: XAS spectra obtained at the Mg K-edge for crystalline materials (spinel and periclase) and CO<sub>2</sub>-bearing silicate glasses: A, for depolymerized (NBO/T > 2) kimberlite glasses; B, for moderately depolymerized (NBO/T = 1) haplobasalt synthesized at 1.5 GPa. The CO<sub>2</sub> content, synthesis pressure, degree of polymerization (NBO/T) and #Mg are reported next to each spectrum. The spectra can be decomposed into two regions: XANES up to ~1340 eV and the EXAFS region above 1340 eV. Crystalline compounds exhibit five features denoted *A* to *E*, whereas silicate glasses only show four peak from *a* to *d*.

Figure 2: XAS spectra obtained at the Ca K-edge for crystalline materials (diopside and dolomite) and CO<sub>2</sub>-bearing silicate glasses: A, for depolymerized (NBO/T > 2) melilitite and kimberlite glasses, the spectrum for CO<sub>2</sub>-free wollastonite has also been added; B, for moderately depolymerized (NBO/T = 1) haplobasalt synthesized at 1.5 GPa. The CO<sub>2</sub> content, synthesis pressure, degree of polymerization (NBO/T) and #Mg are reported next to each spectrum. The spectra can be decomposed into two regions: XANES up to ~4075 eV and the EXAFS region above 4075 eV. As inferred in previous works (see references in the text), several features can be distinguished in the spectra for silicate glasses and denoted *p* corresponding to the pre-edge peak and *a* to *c*.

Figure 3: Change in the main line of the Ca and Mg XANES spectrum as a function of CO<sub>2</sub> content. A 10% in relative to the value for the wt.% CO<sub>2</sub> and ±0.1 eV for the absorption energy are considered.



Figure 4: Calculated  $d_{\text{Mg-O}}$  distance in Å using the model of Li et al. (1999) as a function of the experimentally determined Mg K-edge in eV. The  $d_{\text{Mg-O}}$  is reported for the analysed crystalline compounds and CO<sub>2</sub>-bearing silicate glasses.

Figure 5: A) Normalized  $k^3$ -weighted Mg K-edge EXAFS  $\chi(k)$  and fit for X3-2 silicate glass; B) Magnitude  $\chi(R)$  in Å<sup>-3</sup> in the R-space for X3-2 silicate glass; C) Normalized  $k^3$ -weighted Mg K-edge EXAFS  $\chi(k)$  and fit for HKM silicate glass; D) Magnitude  $\chi(R)$  in Å<sup>-3</sup> in the R-space for HKM silicate glass. The fit of the  $k^3$ -weighted  $\chi(k)$  and the magnitude  $\chi(R)$  in the R-space were done between 2 and 7.1 Å<sup>-1</sup> and 1 and 2.2 Å, respectively. We used a one Mg to O single scattering path from periclase crystallographic structure to fit the Mg EXAFS signal.

Figure 6: A) Normalized  $k^2$ -weighted Ca K-edge EXAFS  $\chi(k)$  and fit for TA12-C3 silicate glass; B) Magnitude  $\chi(R)$  in Å<sup>-4</sup> in the R-space for TA12-C3 silicate glass; C) Normalized  $k^2$ -weighted Ca K-edge EXAFS  $\chi(k)$  and fit for X3-2 silicate glass; D) Magnitude  $\chi(R)$  in Å<sup>-4</sup> in the R-space for X3-2 silicate glass. The fit of the  $k^2$ -weighted  $\chi(k)$  and the magnitude  $\chi(R)$  in the R-space were done between 2 and 8 Å<sup>-1</sup> and 1 and 3 Å, respectively. We used a three Ca to O single scattering paths and single Ca to Si scattering path from diopside crystallographic structure to fit the Ca EXAFS signal.

Figure 7: Change in the Ca coordination number (CN) as a function of CO<sub>2</sub> content (in wt.%) for silicate glasses. The Ca CN is determined from the subsequent Ca EXAFS simulations shown in Figure 6. The error on the CO<sub>2</sub> content is 10% in relative to the value and the error on the Ca CN is determined from error propagation obtained from the Ca EXAFS fit.

**Table caption:**

Table 1: Sample composition (major elements and volatiles), synthesis conditions and chemical parameters (NBO/T and XMg).

Table 2: Mg and Ca K-edge XAS results for CO<sub>2</sub>-bearing silicate glasses: XANES peak maximum, EXAFS simulation results (Mg and Ca coordination numbers, calculated distance to first neighbour).

Table 1

Sample	Temperature (°C)	Pressure (GPa)	SiO <sub>2</sub> (wt.%) <sup>a</sup>	Al <sub>2</sub> O <sub>3</sub>	FeO	MgO	CaO	Na <sub>2</sub> O	K <sub>2</sub> O	CO <sub>2</sub> <sup>b</sup>	NBO/T <sup>c</sup>	XMg <sup>d</sup>
Woll	1500	0.0001	51.7	0	0	0	48.3	0	0	0	2	0
HK-1	1550	1.5	40.2	4.3	0.0	21.9	22.0	4.5	0.0	7.1	2.6	0.58
HK-2	1550	0.5	40.1	4.4	0.0	21.3	22.3	4.5	0.0	5.6	2.5	0.57
HK-3	1550	1.0	37.2	4.2	0.0	21.8	21.4	4.5	0.0	2.9	2.7	0.59
HKM	1550	1.5	31.8	3.5	0.0	21.9	22.0	3.5	0.0	13.2	3.2	0.58
RB8E-3	1550	1.5	30.1	8.7	0.0	0.2	39.4	5.5	0.0	17.2	1.9	0
RB8-2	1550	0.5	33.2	9.8	0.3	0.2	40.4	4.9	0.0	6.0	2.1	0
X1-2	1525	1.5	39.7	14.2	0.1	6.0	25.1	4.8	0.0	7.5	1.1	0.25
BX2	1525	1.5	42.0	15.4	0.7	12.3	17.5	4.8	0.0	3.6	1.2	0.50
X3-2	1525	1.5	43.1	16.1	0.0	17.6	12.0	5.2	0.0	1.5	1.1	0.67
βMDE-5	1550	2.0	48.9	17.4	0.0	17.4	16.3	0.0	0.0	1.7	1.0	0.60
TA12-C1	1400	1.0	42.5	9.5	3.2	9.6	21.2	0.8	5.8	3.7	1.4	0.39
TA12-C3	1400	3.0	38.6	8.8	3.8	8.8	20.1	0.7	5.8	10.5	1.5	0.38
TA10b-C1	1400	1.0	26.9	6.3	1.4	19.7	26.2	0.5	3.6	12.9	3.4	0.51
TA10b-CH1	1400	1.0	26.6	6.1	1.6	19.5	26.0	0.5	3.6	12.7	3.4	0.51

<sup>a</sup> The glass composition was determined by EPMA analyses and has been reported in previous studies. The observed error based on replicated EPMA analyses associated to each oxide concentration is less 0.2 wt.% in relative.

<sup>b</sup> The CO<sub>2</sub> solubility has been determined using the method of Morizet et al. (2013) and using the derived CO<sub>3</sub>/HF value obtained from the deconvolution of the Raman spectra. The associated error on the CO<sub>2</sub> content using this Raman method is 10% in relative to the value.

<sup>c</sup> The NBO/T corresponds to the concentration of Non-Bridging Oxygen per Tetrahedron and is calculated from the chemical composition (Mysen, 1990).

<sup>d</sup> The XMg is defined as molar ratio MgO/(MgO+CaO).

Table 2

Sample	CO <sub>2</sub>	Mg XANES		Mg EXAFS <sup>c</sup>		Ca XANES	Ca EXAFS <sup>c</sup>	
		Peak maximum (eV) <sup>a</sup>	Calc. dMg-O (Å) <sup>b</sup>	Calc. dMg-O (Å)	Mg CN	Peak maximum (eV) <sup>a</sup>	Calc dCa-O (Å)	Ca CN
Woll	0					4051.4	2.54±0.24	7.1±0.5
HK-1	7.1	1309.8	1.96	2.04±0.05	3.7±1.3	4050.9	2.52±0.51	9.1±1.3
HK-2	5.6	1309.7	1.96	2.01±0.01	6.1±0.5	4051.1	2.49±0.32	8.8±0.9
HK-3	2.9	1309.6	1.95	2.00±0.01	5.6±0.5	4051.0	2.53±0.31	8.9±0.9
HKM	13.2	1310.4	2.01	2.00±0.01	5.7±0.6	4050.8	2.50±0.31	9.0±0.8
RB8E-3	17.2					4050.9	2.53±0.23	9.3±0.4
RB8-2	6.0					4050.9	2.55±0.39	9.5±1.1
X1-2	7.5	1309.8	1.96	1.98±0.00	6.0±0.2	4051.0	2.51±0.42	9.3±1.3
BX2	3.6	1309.9	1.97	2.01±0.02	5.9±0.6	4051.1	2.53±0.28	8.9±0.8
X3-2	1.5	1309.9	1.97	1.98±0.01	6.0±0.3	4051.2	2.53±0.43	8.6±1.1
βMDE-5	1.7	1309.8	1.96	2.00±0.01	5.6±0.4	4051.2	2.50±0.23	8.0±0.6
TA12-C1	3.7					4051.1	2.48±0.26	8.4±0.6
TA12-C3	10.5	1310.1	1.98	2.00±0.01	6.2±0.5	4050.8	2.51±0.41	9.3±1.2
TA10b-C1	12.9	1309.5	1.94	1.99±0.01	5.5±0.4	4051.0	2.51±0.24	9.1±0.7
TA10b-CH1	12.7	1309.4	1.94	1.99±0.02	6.5±0.9	4051.0	2.48±0.26	9.2±0.7

<sup>a</sup> The peak maximum corresponds to the peak maximum of the main line *b* and has been determined from the simulation of the XANES region. The error on the peak is 0.2 eV.

<sup>b</sup> The  $d_{\text{Mg-O}}$  in Å has been calculated from the model of Li et al. (1999) relating the position of the Mg K-edge to the distance between Mg and O atoms in the first coordination sphere.

<sup>c</sup> The Mg and Ca EXAFS results (distance to O and CN) were obtained from the simulation of the oscillatory EXAFS signal (see Figure 5 and 6) and using the Artemis software (Ravel and Newville, 2005). The error on each parameter is calculated from error propagation method.

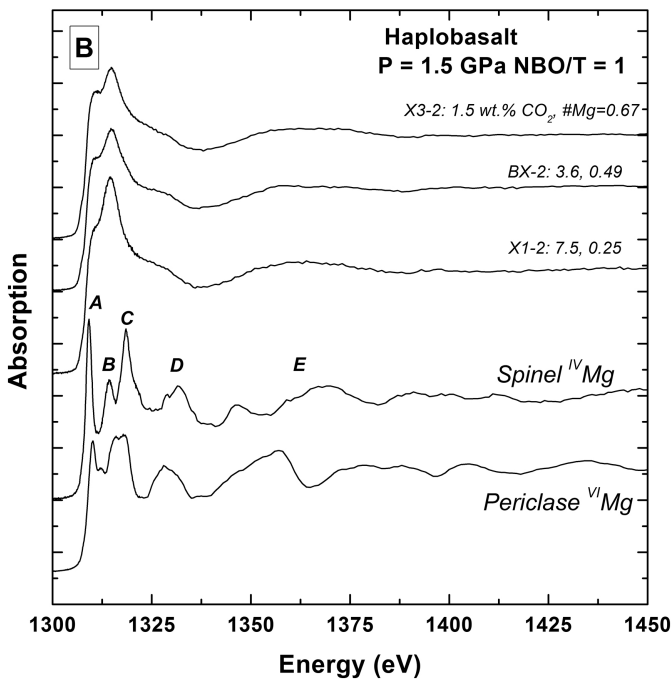
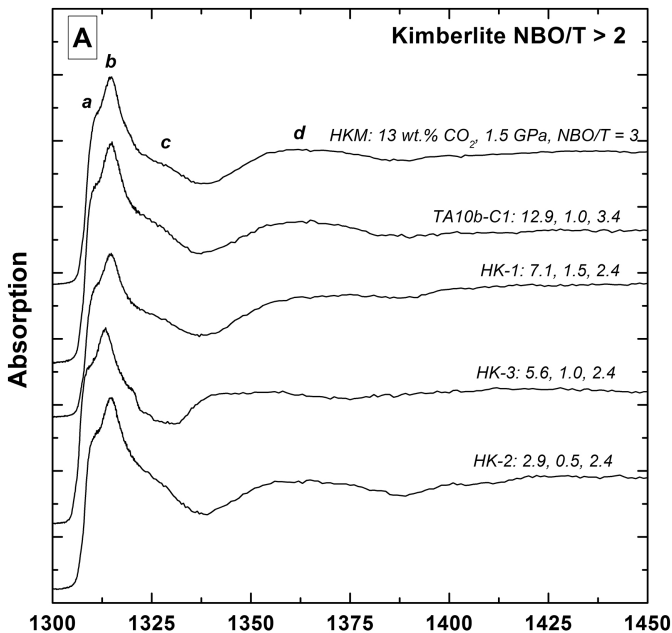


Figure 1

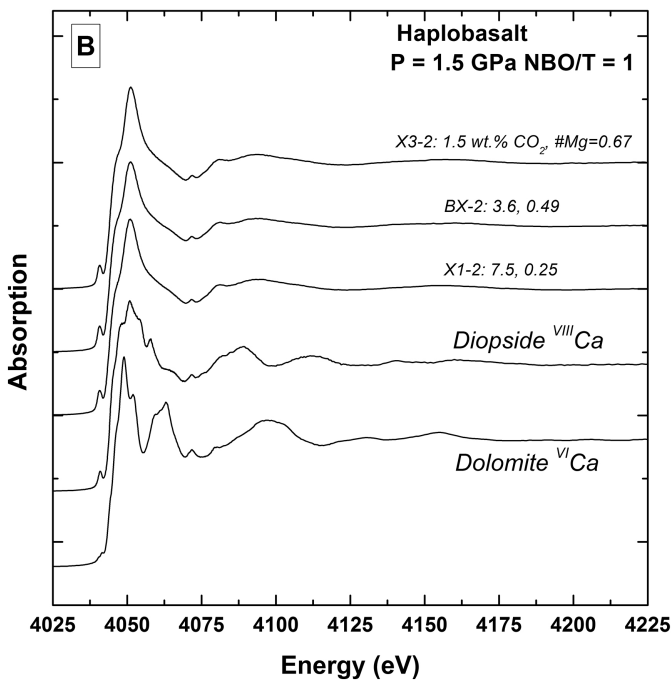
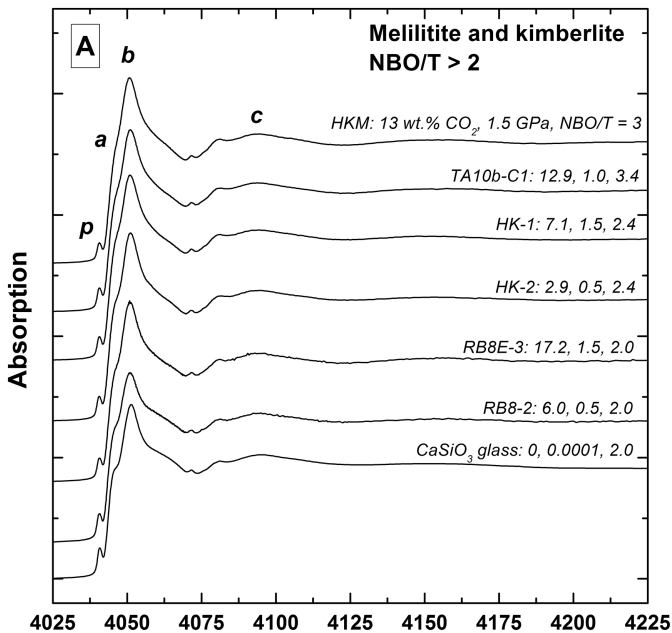


Figure 2

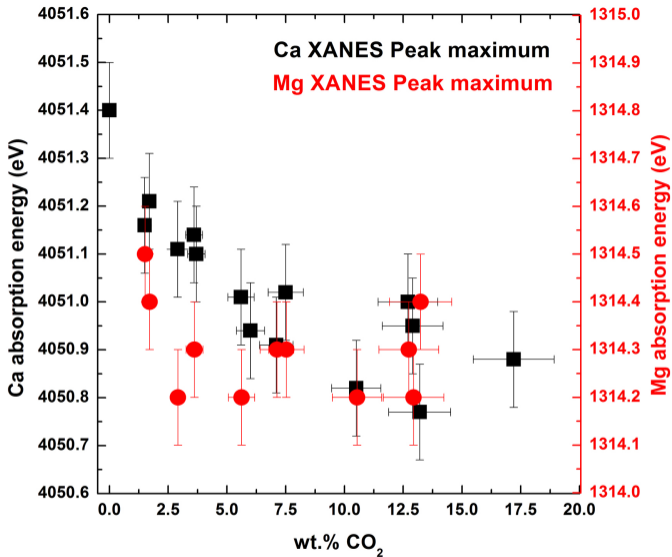


Figure 3

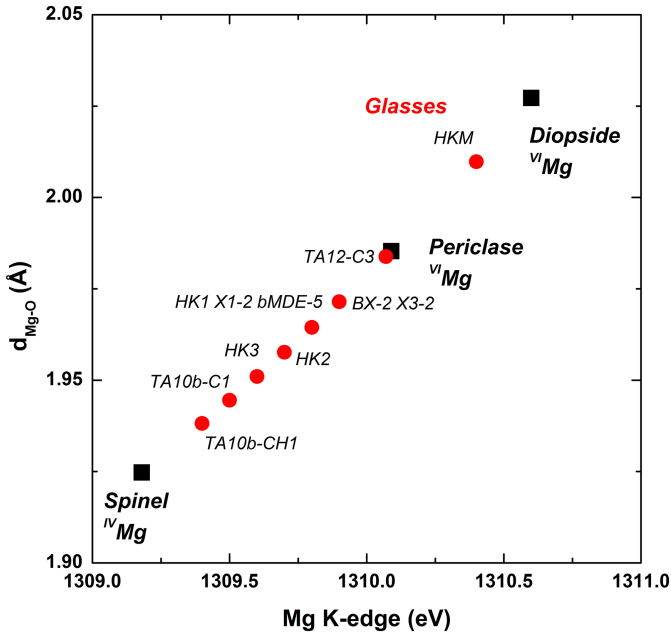


Figure 4



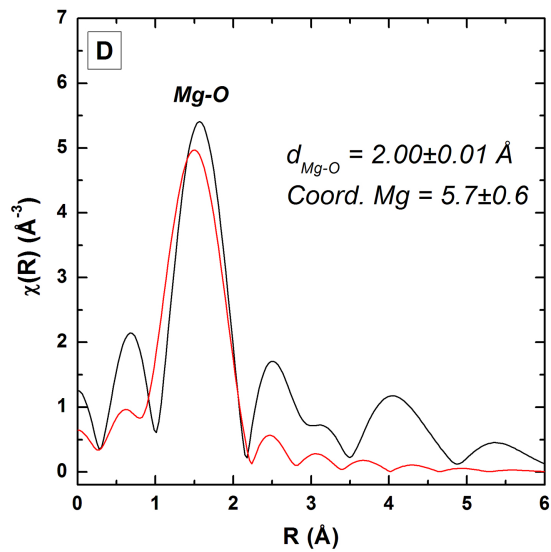
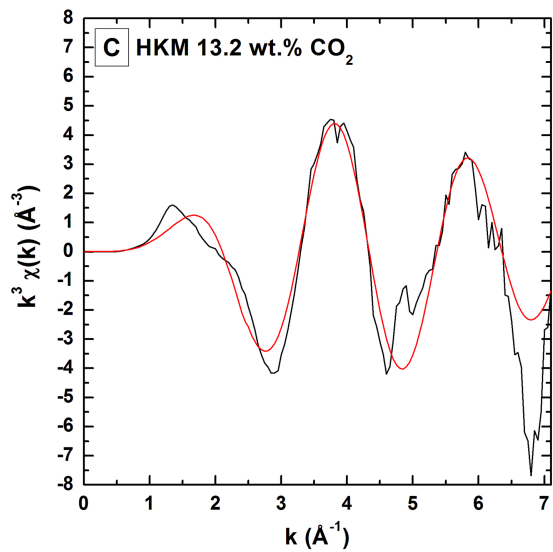
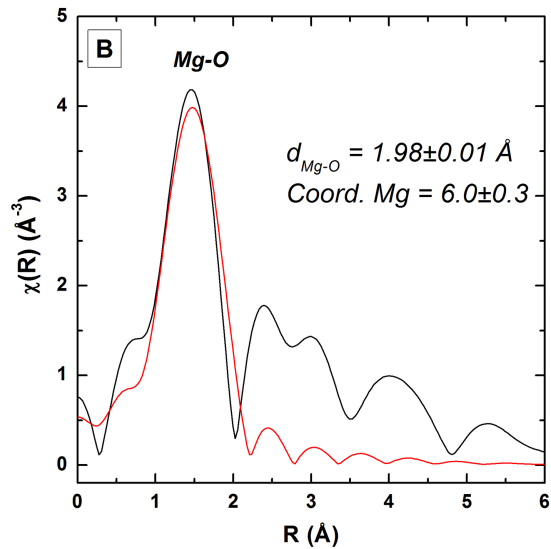
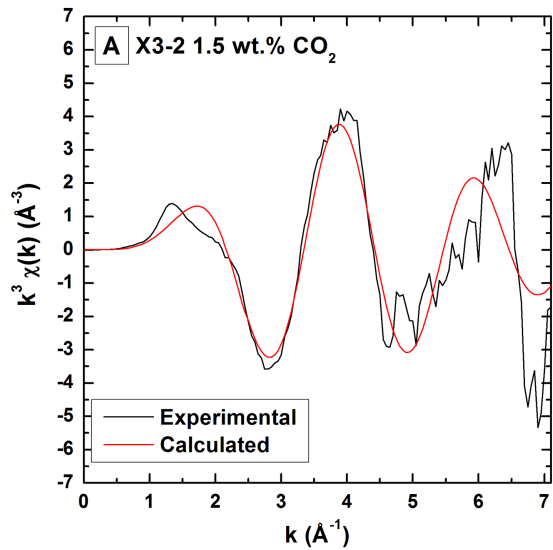


Figure 5

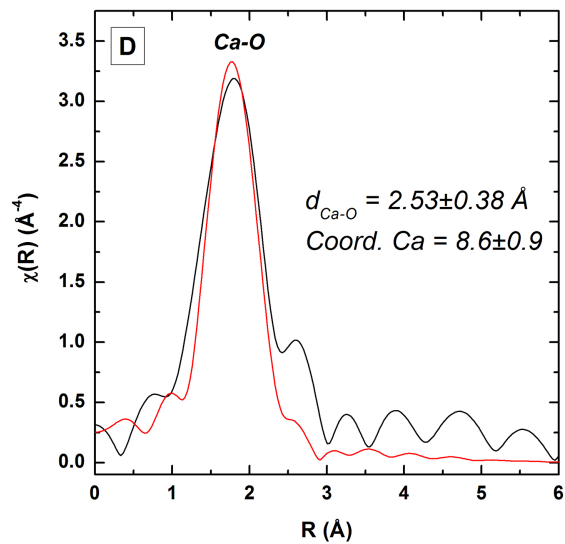
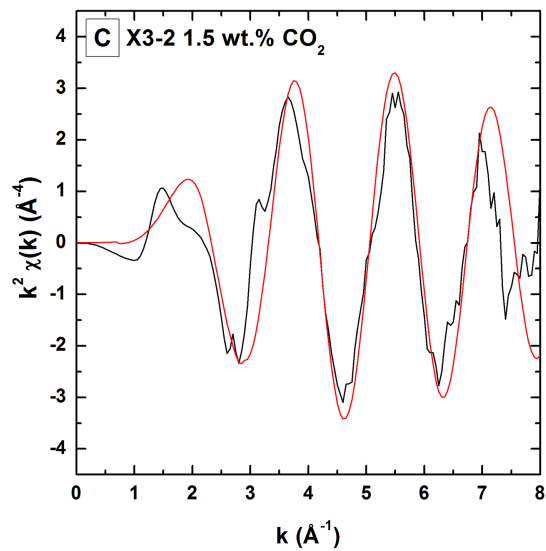
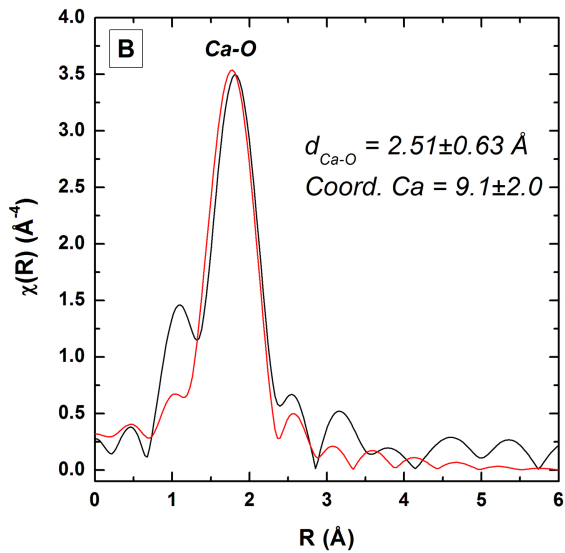
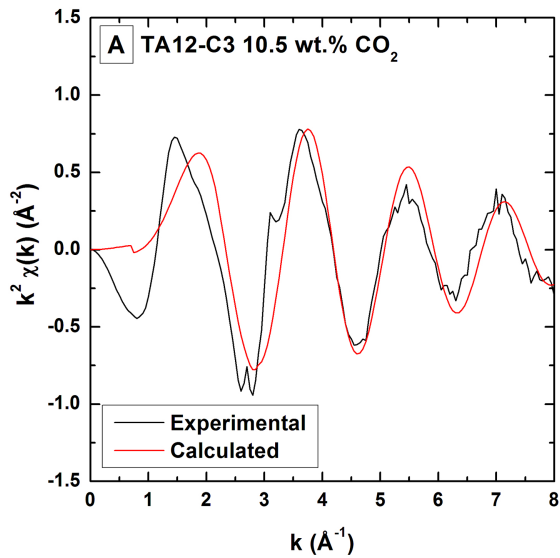


Figure 6

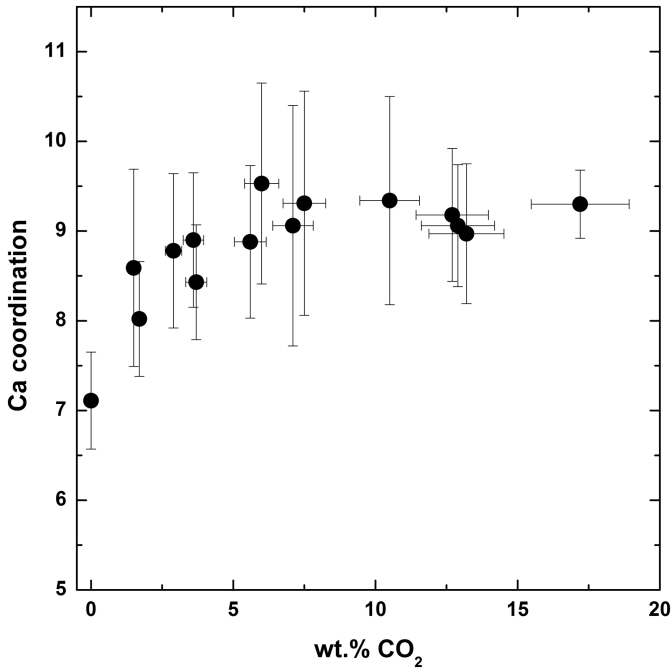


Figure 7