

# Influence of Extreme Ultraviolet Irradiance Variations on the Precipitating Ion Flux From MAVEN Observations

Antoine Martinez, François Leblanc, Jean-Yves Chaufray, Ronan Modolo, Olivier Witasse, Y. Dong, T. Hara, J. Halekas, R. Lillis, J. Mcfadden, et al.

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1	Influence of Extreme Ultraviolet irradiance variations on the precipitating ion
2	flux from MAVEN observations.
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## 29 Abstract:

30 We study the influence of the solar EUV flux intensity on the precipitating ion fluxes as seen by 31 MAVEN/SWIA, an energy and angular ion spectrometer. We defined three periods with significantly 32 different EUV flux intensity (1.6 and 3.2 times the lowest EUV intensity) and compare the 33 precipitating ion flux measured by MAVEN/SWIA during each period. At low energy [30-650] eV, we 34 find that the median (average) precipitating ion flux during the medium and low EUV periods are 35 respectively 1.7 (2.1) and 3 (3.5) times more intense than the flux during the high EUV period. At high 36 energy [650-25000] eV, a similar trend in the intensity of the precipitating ion flux is observed but 37 with an increase by 50% (46%) and 70% (79%) respectively. A larger EUV flux does therefore not 38 seem to favour heavy ion precipitation into Mars' atmosphere, contrary to modelling prediction and 39 overall expectations.

#### 41 I. Introduction

42 Mars' upper atmosphere is constantly bombarded by energetic ions from either the solar wind or from Mars' exosphere (planetary pickup ions) (Lillis et al. 2015). Planetary ions created by the 43 44 ionization of exospheric neutral particles can be accelerated by the motional solar wind electric field. 45 Some of these ions escape (Brain et al. 2016), but a significant portion can impact Mars' atmosphere, 46 leading to collisional cascade in Mars' thermosphere and to the ejection of atmospheric particles into 47 the exosphere (Luhmann and Kozyra, 1991). This process, named "atmospheric sputtering" (Johnson, 48 1994), could have induced a significant atmospheric loss along Mars' history (Luhmann et al., 1992; 49 Leblanc, and Johnson, 2001 and 2002).

50 Atmospheric escape induced by sputtering at present epoch is expected to be small compared to 51 other mechanisms (Leblanc et al. 2017) and so difficult to measure. However, since the main driver 52 of sputtering is ion precipitation (Johnson et al, 2000; Wang et al, 2014;2015), it is crucial to constrain 53 the dependence of the precipitating ion flux on present solar wind conditions. While many studies 54 have been carried out on the influence of the Extreme-Ultraviolet/Ultraviolet (EUV/UV) flux on the 55 Martian environment as well as on ion escape (Modolo et al., 2005; Ma and Nagy, 2007; Terada et al. 56 2009; Edberg et al., 2009; Lundin et al, 2008,2013; Ramstad et al., 2015; Dong et al., 2017), the 57 relation between ion precipitation and EUV flux has only been studied in theoretical works (Chaufray 58 et al., 2007; Leblanc and Johnson, 2002; Wang et al., 2014,2015) and briefly discussed in Nilsson et 59 al., 2012. We therefore describe here the first empirical study on the influence of the EUV flux on ion 60 precipitation by analyzing MAVEN measurements (Jakosky et al. 2015). In section II, we describe the 61 set of data used in this work. In section III, we carefully analyzed the various potential solar wind 62 drivers to extract three sets of measurements obtained during different EUV flux intensity but for similar solar wind conditions and geographical coverage. In section IV, we summarize and discuss the 63 64 main results of this study and conclude in section V.

#### 66 II. Data set used for this analysis

We use measurements performed by the Solar Wind Ion Analyzer (SWIA) (Halekas et al., 2015), the 67 68 Magnetometer (MAG) (Connerney et al., 2015a; 2015b) and the Solar Extreme Ultraviolet Monitor 69 (EUVM) (Eparvier et al., 2015). The MAVEN/SWIA is an energy and angular ion spectrometer covering 70 an energy range between 25eV/q and 25keV/q with 48 energy steps logarithmically spaced, a field of 71 view (FOV) of 360°x90° on 64 angular bins and 4s time resolution. We based our study on 72 MAVEN/SWIA despite the lack of mass resolution because MAVEN/STATIC (which has mass 73 resolution) has much more restricted coverage of the precipitating flux. The MAVEN/EUVM measures 74 the solar irradiance in three bands from the soft X-ray to the EUV range (in three spectral bands 0.1-75 7nm, 17-22nm and 121-122nm) with a temporal resolution of 1s. The solar wind density, speed, and 76 IMF vector are measured by MAVEN/SWIA and MAVEN/MAG and averaged on an orbit-by-orbit basis 77 in order to characterize the solar wind conditions (Halekas et al., 2017). We then reconstruct the 78 Mars Solar Electric (MSE) coordinate system in which the X-axis is toward the Sun, the Z-axis points 79 along the solar wind motional electric field and the Y-axis completing the right hand system (Fedorov 80 et al., 2006).

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82 In order to reconstruct the precipitating ion flux, we follow the method described in Leblanc et al. 83 (2015) and Martinez et al (2019). We select all MAVEN measurements performed between 200 and 84 350km. Within such altitude range, any ion which velocity direction is at less than 75° with respect to 85 the local nadir direction has a very large probability to impact Mars' atmosphere. Therefore, to 86 reconstruct the precipitating flux, we sum all measurements of SWIA anodes which FOV is at less 87 than 75° from the local zenith direction. Moreover, in order to exclude reconstructed precipitating flux with poor coverage, we only consider measurements during which the total FOV of SWIA anodes 88 89 is more than 65% of the 75° solid angle cone centered on the zenith direction.

In this study, to characterize the EUV flux during the declining solar cycle from medium to minimum
activity, we use the MAVEN/EUVM channel A which measures solar irradiance between 17 and

92 22nm. We sort the MAVEN data from December 2014 to July 2018 into three periods: high EUV 93 (From 2 December 2014 to 10 March 2015) with an average EUV flux of 0.42mW.m<sup>-2</sup>, medium EUV 94 (From 4 June 2015 to 22 October 2015, from 23 December 2015 to 1 June 2016, and from 4 95 December 2016 to 20 July 2017) with an average EUV flux of 0.21mW.m<sup>-2</sup> and low EUV (from 28 96 November 2017 to 31 July 2018) with an average EUV flux of 0.13mW.m<sup>-2</sup>. These three intervals 97 represent several different Mars seasons. Table 1 summarizes the main characteristics of each 98 period.

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	High Medium		Low	
[Date begin- Date End] (dd/mm/yyyy)	[02/12/2014- 10/03/2015]	[04/06/2015-22/10/2015] [23/12/2015-01/06/2016] [04/12/2016-30/07/2017]	[28/11/2017 31/07/2018]	
Heliocentric distance (AU)	[1.38-1.43]	[1.54-1.66] [1.49-1.66] [1.39-1.63]	[1.40-1.65]	
Solar Longitude (LS)°	[244.3°-305.2°] (northern hemisphere winter)	[352.5°-57.9°] [85.1°-161.5°] [273.5°-35.7°] (one Martian year)	[93.4°-221.1°] (northern hemisphere summer and autumn)	
Mean EUV flux ( <i>mW</i> . <i>m</i> <sup>-2</sup> ) from channel A of MAVEN/EUVM	Mean EUV flux $(mW. m^{-2})$ from channel A of0.420.21		0.13	
Average Precipitating ion flux $(eV. cm^{-2}. sr^{-1}. s^{-1})$ [30,650] eV	$1.18 (\pm 0.11) \times 10^7$	$2.49 (\pm 0.12) \times 10^7$	$4.05 (\pm 0.33) \times 10^7$	
Median Precipitating ion flux $(eV. cm^{-2}. sr^{-1}. s^{-1})$ [30,650] eV	$0.62 \times 10^{7}$	$1.07  imes 10^7$	$1.83 \times 10^{7}$	
Average Precipitating ion flux ( <i>eV. cm</i> <sup>-2</sup> . <i>sr</i> <sup>-1</sup> <i>s</i> <sup>-1</sup> ) [650,25000] eV	$15.9 (\pm 0.7) \times 10^7$	23.9 (±0.7) × 10 <sup>7</sup>	26.9 (±1.1) × 10 <sup>7</sup>	
Median Precipitating ion flux ( <i>eV. cm</i> <sup>-2</sup> . <i>sr</i> <sup>-1</sup> <i>s</i> <sup>-1</sup> ) [650,25000] eV	$11.0 \times 10^{7}$	16.1 × 10 <sup>7</sup>	$19.7 \times 10^{7}$	
Number of measurements	425	1205	374	

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**Table 1:** Summary of each period of our study. In brackets, the 1-sigma standard deviation

(uncertainty) of the precipitating ion flux values.

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### 103 III. Upstream solar wind conditions and planetary coverage

104 To investigate the Solar EUV dependence of the precipitating ion flux, we first need to reconstruct the solar wind parameters. Lillis et al. (2015), Edberg et al. (2009) and Ramstad et al., (2017) 105 106 suggested that the dynamic pressure compresses Mars' magnetosphere and favors the acceleration 107 and precipitation of the planetary pickup ions. Martinez et al. (2019) and Hara et al. (2017a) showed 108 that the precipitating ion flux can be organized with respect to the orientation of the solar wind 109 convective electric field. Also, the presence of crustal magnetic fields (Acuña et al., 1999; Connerney 110 et al., 2005) can also influence locally the precipitating ion flux, as shown by Hara et al. (2017b). For 111 each orbit, during which the precipitating flux was measured, we therefore used the part of MAVEN 112 orbit in the solar wind to infer the solar wind. We use such solar wind measurements prior to our 113 precipitating flux measurement if the time difference between the two measurements is less than 4 hours and 30 minutes (the MAVEN orbital period). We assume that the solar wind conditions do not 114 115 change significantly between MAVEN measurements of the solar wind and MAVEN measurements of 116 the precipitating flux (in average 2 hours and 20 minutes).

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118 In order to verify the validity of this assumption, we compare the solar wind density and speed 119 measurements from Mars Express(MEX)/ASPERA-3 (Barabash et al., 2006) with those measured by 120 MAVEN. Starting from  $n_{MAVEN/SW}$  and  $V_{MAVEN/SW}$  the solar wind density and speed as measured by 121 MAVEN/SWIA before a measurement of the precipitating flux below 350 km, we determine 122  $n_{MEX/SW}$  and  $V_{MEX/SW}$ , as measured by MEX/ASPERA-3 (Barabash et al., 2006) when MAVEN's altitude was within the time interval during which SWIA measured the precipitating flux. We only 123 consider ASPERA-3 measurements if the measured density was larger than 0.01cm<sup>-3</sup> and the quality 124 125 flag larger than 0.6.



Figure 1: Comparison between the measured solar wind parameters by MAVEN/SWIA (blue) and MEX/ASPERA-3 (orange) for the high EUV period (from December 2014 to March 2015): **(a)** the solar wind density, **(b)** Solar wind speed.

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130 With such criteria, we did not identify any measurement for the low EUV period and ASPERA-3 131 measurements for 4.2% and 8.0% of MAVEN precipitating flux measurements for the medium and 132 high EUV periods (see also the supplementary materials, Table S.1). The set of measurement 133 comparisons displayed in Figure 1 for the high EUV period, show that the solar wind speeds as 134 measured by SWIA and ASPERA-3 are usually in very good agreement (MAVEN:  $403 \pm 12$ km/s and 135 MEX:  $396 \pm 12$  km/s, for 85 measurements). Contrary to the solar wind speed, the measurements of 136 the density display a relatively poor agreement (Figure 1b) but still acceptable in an average and within the one-sigma dispersion (MAVEN:  $3.31\pm0.35$  cm<sup>-3</sup> and MEX:  $3.07\pm0.31$  cm<sup>-3</sup>, for 85 137 138 measurements).

The approach to select the set of measurements used for our study follows the method developed inDong et al. (2017). As in Dong et al. (2017), we restrict the set of precipitating measurements to

141 those obtained during solar wind conditions with dynamic pressure  $P_{dyn} = m_p n_{SW} V_{SW}^2$ , (where m<sub>p</sub> is the mass of the proton,  $n_{SW}$  and  $V_{SW}$  the solar wind density and velocity) between 0.3 and 2.6nPa 142 143 (moderate solar wind dynamic pressure) and with IMF strength between 2.2 and 6.7nT. We also restrict our sample of precipitation measurements to those obtained in an SZA interval between 70° 144 145 and 130°. Furthermore, in order to avoid the potential influence of the crustal fields (Leblanc et al. 146 2017; Hara et al. 2017b), we also only consider measurements performed when the average 147 magnetic field between 200 and 350km is less than 60nT. For each set of solar wind parameters, we 148 calculate the solar wind dynamic pressure, the norm of the solar wind motional electric field defined as  $E = \left| \vec{V}_{SW} \times \vec{B}_{IMF} \right|$  where  $B_{IMF}$  is the interplanetary magnetic field (IMF), the Alfvén Mach number 149  $M_A = V_{SW}/V_A$  where V<sub>A</sub> is the Alfvén speed, the solar wind flux defined as  $F_{SW} = n_{SW}V_{SW}$  and the 150 pickup  $O^+$  gyroradius in the solar wind  $R_g = \frac{m_O V_{SW} \sin \theta_{cone}}{q B_{IMF}}$  where  $m_O$  is the mass of an  $O^+$  ion, 151  $\theta_{cone}$  is the angle between the IMF direction and the MSO X-axis and q is the electron charge value. 152 153 We also define the MSE angle as the anticlockwise angle between the vector formed by the latitude 154 and the longitude in MSE of the position where the ion precipitation is being measured and the East 155 MSE direction (MSE longitude equal to +180° and latitude equal to 0°) (Martinez et al., 2019).





Figure 2: Distribution of the set of measured solar wind parameters for the high EUV period (black), the medium EUV period (dark blue) and the low EUV period (red): **(a)** Values of the EUV photon flux as measured by MAVEN/EUVM instrument, **(b)** Solar wind dynamic pressure, **(c)** Solar wind convective electric field, **(d)** Interplanetary Magnetic Field as measured by MAVEN/MAG, **(e)** pickup 0<sup>+</sup> ion gyroradius calculated in the solar wind and **(f)** solar wind speed determined by MAVEN/SWIA. The green zone in 2.e and 2.f correspond to the portion of the distribution common to the three periods.

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166 In order to demonstrate that the selected sets of measurement for the three EUV periods are similar 167 in terms of solar wind conditions and planetary geographic coverage (MSO and MSE frames, see also 168 the supplementary materials, Figures S.1 and S.2), we calculate for each parameter, the area (green 169 part in Figure 2.e and 2.f) common to the three distributions. The ratio between this common area 170 and each of the three areas provides a measure of the overlap between the three sets of 171 measurement for a given solar wind parameter. In Table 2, for 10 solar parameters, we provide the 172 minimum percentage of the distribution areas in common to the three periods as well as the first and 173 third quartile, the mean and 1-sigma standard deviation of each parameter.

Figure 2 displays the distribution of the values of each of these parameters for the three EUV periods and illustrates our method to determine the similarity between these sets of measurement. As an example, in the case of the gyroradius (Figure 2.e), at least 83% of each sets of measurement were performed for similar values of this parameter. This percentage decreases to 62% for the solar wind speed (Figure 2.f). On the other hand, the three EUV periods are clearly distinct (Figure 2.a) and correspond to a null common area (Table 2).

		High	Medium	Low	Common coverage (%)
EUV Flux	$\mu \pm \sigma$	$0.42 \pm 0.08$	$0.21 \pm 0.04$	$0.13 \pm 0.02$	0.0

$mW.m^{-2}$	Q25	0.37	0.19	0.11	
	Q75	0.46	0.22	0.14	
Dynamic	$\mu \pm \sigma$	$0.87 \pm 0.44$	$0.85 \pm 0.46$	$0.90 \pm 0.43$	
pressure	Q25	0.56	0.52	0.56	86.0
nPa	Q75	1.08	1.05	1.14	
Electric field	$\mu \pm \sigma$	$1.18 \pm 0.50$	1.15 ± 0.51	$1.00 \pm 0.43$	
$mV m^{-1}$	Q25	0.83	0.80	0.74	85.4
ntv.nt	Q75	1.48	1.42	1.27	
INAC	$\mu \pm \sigma$	3.84 <u>+</u> 1.16	3.43 ± 1.04	3.25 ± 0.98	
	Q25	2.89	2.59	2.52	74.5
101	Q75	4.74	4.06	3.67	
Currente di une D	$\mu \pm \sigma$	$4.36 \pm 1.72$	5.22 ± 1.94	$5.08 \pm 2.02$	
Gyroradius K <sub>g</sub>	Q25	2.87	3.80	3.61	83.2
<b>N</b> <sub>Mars</sub>	Q75	5.51	6.54	6.41	
Alfrée Mach	$\mu \pm \sigma$	8.89 <u>±</u> 2.71	9.68 <u>+</u> 2.71	10.6 ± 3.0	
Alfven Wach	Q25	6.87	7.88	8.40	74.2
number	Q75	10.1	11.3	12.3	
Crood	$\mu \pm \sigma$	378 <u>+</u> 78	422 <u>+</u> 74	$406 \pm 70$	
speed	Q25	317	369	353	61.6
кт. 5	Q75	413	464	443	
Density	$\mu \pm \sigma$	$4.00 \pm 2.60$	3.02 ± 1.86	3.63 ± 2.25	
cm <sup>-3</sup>	Q25	2.21	1.67	1.89	76.5
Cm	Q75	4.89	3.70	4.99	
	$\mu \pm \sigma$	185 ± 112	$202 \pm 111$	172 ± 115	
Angle MSE (°)	Q25	143	149	38.8	67.1
	Q75	306	326	208	
	$\mu \pm \sigma$	$1.42 \pm 0.77$	$1.21 \pm 0.67$	$1.38 \pm 0.73$	
$10^8  {\rm s}^{-1}  {\rm cm}^{-2}$	Q25	0.88	0.72	0.81	79.0
10 3 CIII	Q75	1.70	1.50	1.80	

181 **Table 2:** Mean  $\mu$ , standard deviation  $\sigma$ , first quartile Q25 and third quartile Q75 of each solar 182 parameter distribution for the three sets of precipitating flux measurement. The last column provides 183 the percentage of the area in common between the 3 distributions.

184 If we now only consider the medium and low periods, for all parameters listed in Table 2, the 185 distribution area in common to these two sets of measurement represents more than 75% of each 186 distribution. According to these percentages, we conclude that the three sets of precipitating ion 187 measurements were obtained under similar solar wind conditions.

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To further demonstrate this conclusion, considering the solar wind speed parameter with the less similar distributions between the three periods, we reduce each sample so that the solar wind speed measured during each period was restricted to values between 370 and 500km/s. In that case, the percentage increases to 81%. We then calculate the values of the precipitating ion flux with these reduced samples (see also the supplementary material, Table S.2) and found that all measured precipitating flux values are equivalent for the two different samples except during the medium period when the flux values differ slightly more than one sigma. In another words, the reduction of the sampling improves the similarity between the three sets of measurement and does not change significantly the measured precipitating ion flux

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## 199 IV. <u>Results and Analysis</u>

200 Our results show that the average ion precipitation decreases when the EUV flux increases (see Table 201 1 and Figure 3). The energy dependence of the differential flux is similar from one period to the 202 other, but we can note that the increase in precipitating ion flux is more significant for energies below 1 to 3 keV than above. This conclusion is not consistent with studies modelling the 203 204 precipitating ion flux for different solar EUV fluxes (Chaufray et al., 2007) which showed that the flux 205 of  $O^{\dagger}$  which re-impacts the atmosphere should be smaller at solar minimum than at solar maximum. 206 However, they also concluded that the differences between the two periods were smaller than 207 predicted by Luhmann et al. (1992) using simple gas dynamic description of Mars' interaction with 208 the solar wind. These authors explained this difference by the deceleration of the solar wind due to 209 the mass-loading (Dubinin and Lundin, 1995) which happens further away from Mars at high solar 210 activity and reduces therefore the probability of accelerated pickup ions to re-impact Mars' 211 atmosphere.

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213 Mass-loading depends to first order on the ionization rate of the Martian exosphere. The exospheric 214 density at a given altitude also depends on the EUV/UV flux which heats the Martian thermosphere 215 whereas the ionization rate directly depends on the solar radiation flux in the wavelength range 216 below 91nm. Forbes et al. (2008) found a very good correlation (0.96) between the long-term EUV 217 flux (proxy F10.7) and the Martian exospheric density at 390km in altitude. Seasonal variations of the 218 exosphere might also significantly impact the density at a given altitude as studied in Bhattacharyya

et al, (2015) and Chaufray et al, (2015), as well as the position of the bow shock (Halekas et al, 2017). A peak of exospheric hydrogen density during the Martian dust season was reported by Chaffin et al. (2014), Bhattacharyya et al, (2015) and Halekas (2017) between Ls=200° and 300°. Therefore, an increase in mass-loading of the solar wind might be related to this seasonal variability but would only explain the decrease of the precipitating flux during the high EUV period. During the low and medium periods, there are no significant change in season so that the difference in precipitating flux cannot be associated to seasonal variability.

As originally suggested by Johnson & Luhmann (1998), an increase of the EUV/UV flux should increase the number of pickup exospheric ions but might also load more efficiently the solar wind and reduce the percentage of these ions able to re-impact the atmosphere. What our study seems to suggest is that this effect might be much more efficient than shown in Chaufray et al. (2007), perhaps due to the limited spatial resolution of their simulation at this time (~300km).

Another potential effect is the MSO latitude of the crustal magnetic fields which is solar dependent. 231 232 Brain et al. (2005) showed that the Magnetic Pile-Up Boundary altitude changes significantly with 233 season because of the latitude variation of the main crustal field structure (centered at 180° GEO 234 East longitude and -50° GEO latitude) along one Martian year. Such effect might be actually due to an 235 increase exospheric hydrogen density rather due to the crustal fields latitudes according to Chaufray 236 et al. (2015). Actually, Hall et al. (2016) shows that an increase in the EUV flux is associated with a 237 position of the bow shock further from Mars, suggesting also an influence of the exosphere on the 238 induced magnetosphere. In any case, we cannot exclude that part of the variation in precipitating 239 flux described here might be due to the crustal field, as actually shown by Hara et al. (2017b). We 240 calculate the average position of the main crustal field structure for the three sets of measurements and found  $-47 \pm 17^{\circ}$ ,  $-45 \pm 16^{\circ}$ , and  $-46 \pm 17^{\circ}$  latitude MSO and  $27 \pm 106^{\circ}$ ,  $-20 \pm 107^{\circ}$  and 241  $-3 \pm 84^{\circ}$  longitude MSO for the low, medium and high EUV periods. Therefore, the relative variation 242 243 of the precipitating ion flux between the three periods selected in our paper might be also related to 244 the variation of the crustal field latitudinal positions. In order to test this possibility, we further

reduce our set of measurements in order to only use measurements of the precipitating ion flux performed when the main crustal field structure was at a SZA larger than 120°. The reconstructed differential precipitating ion fluxes display very similar relative variation showing, here again, that the main driver for this variation should be the EUV flux.

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Dong et al. (2017) found that the total escape rate of oxygen ions with  $E_i \ge 6$  eV increases by 1.5 when the EUV flux increases by a similar factor. Our analysis suggested that an increased EUV flux induces a decrease of the ion precipitation. In the case of the high EUV period (similar to that in Dong et al (2017)), for the ions with  $E_i \ge 25$  eV, we have an integrated precipitating flux equal to  $1.3 \times 10^{22} \ particles. s^{-1}$  which should be compared to the estimate by Dong et al. (2017) of an escape rate of  $1.8 \times 10^{24} \ particles. s^{-1}$  at  $0.2R_{Mars}$  in altitude. Therefore, the precipitating flux represents only 1% of the escaping flux in the energy range considered in this paper.



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Figure 3: Precipitating ion differential energy spectra as measured by MAVEN/SWIA during high EUV
(black), medium EUV (dark blue) and low EUV periods (red): (a) Average value (with 1-sigma standard
deviation) and (b) Median value (with first and third quarter).

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### 262 V. Summary and Conclusion

We study the precipitating ion flux from MAVEN instruments over more than three years of data divided into three intervals of solar EUV/UV intensity: between 2 December 2014 and 10 March 2015 (Ls from 244.3° to 305.2°), between 4 June 2015 and 1 June 2017 (two Ls periods: from 273.5° to 57.9° and from 85.1° to 161.5°) and between 28 November 2017 and 31 July 2018 (Ls from 93.4° to 221.1°). By constraining the solar wind conditions, we show that these three samples were obtained under similar solar wind conditions and cover similar regions in a MSE frame.

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We reconstruct the measured precipitating ion flux by MAVEN/SWIA. The increase of the EUV flux intensity from 0.13 to 0.42 mW.m<sup>-2</sup> is clearly associated with a decrease of the precipitating ion flux from 18.3 to  $6.2 \times 10^{6}$  eV. cm<sup>-2</sup>. sr<sup>-1</sup>. s<sup>-1</sup> (by almost a factor 3) for the energy range between 30 and 650 eV and a decrease by 18% to 44% at higher energies ([650-25000] eV).

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275 This somewhat surprising result, in contradiction with previous theoretical studies (Chaufray et al., 276 2007), might illustrate the complex relation between the variations of the exosphere as induced by 277 the solar EUV/UV flux and the solar wind interaction with Mars. Dong et al. (2017) reported an 278 almost linear dependence of the ion escape on the EUV/UV flux, whereas we have found an almost 279 linear anti-correlation between the EUV flux and the precipitating ion flux. In others words, even if 280 the EUV flux increase the ionization rate in the exosphere, the probability of a pickup ion to escape or 281 to precipitate into the atmosphere should also be changed, illustrating a much more complex 282 dependence of the fate of Mars' planetary ions on the solar conditions than anticipated.

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293

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Figure1.



Figure2.



Figure3.

