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Kevin M. A. Ihaddadene, Joseph R. Dwyer, Ningyu Liu, Sébastien Celestin, Feng Shi. Modeling of a New Electron Acceleration Mechanism Ahead of Streamers. *Journal of Geophysical Research Space Physics*, 2019, 10.1029/2018JA026084 . insu-02263528

HAL Id: insu-02263528

<https://insu.hal.science/insu-02263528>

Submitted on 26 Aug 2019

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JGR Space Physics

RESEARCH ARTICLE

10.1029/2018JA026084

Key Points:

- Collisions between large negative streamers and small neutral plasmapatches produce thermal runaway electrons
- Single streamer-streamer head-on collision and multiple streamer-streamer head-on collisions are unlikely sources of runaway electrons
- X-ray emissions by laboratory spark discharges are unlikely to occur at the first burst of streamers

Supporting Information:

- Supporting Information S1
- Movie S1

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Citation:

Ihaddadene, K. M. A., Dwyer, J. R., Liu, N., Celestin, S., & Shi, F. (2019). Modeling of a new electron acceleration mechanism ahead of streamers. *Journal of Geophysical Research: Space Physics*, 124. <https://doi.org/10.1029/2018JA026084>

Received 12 SEP 2018

Accepted 24 JUL 2019

Accepted article online 4 AUG 2019

Modeling of a New Electron Acceleration Mechanism Ahead of Streamers

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Abstract Head-on collisions between negative and positive streamers have been proposed as a mechanism behind X-ray emissions by laboratory spark discharges. Recent simulations using plasma fluid and particle in cell models of a single head-on collision of two streamers of opposite polarities in ground pressure air predicted an insignificant number of thermal runaway electrons >1 keV and hence weak undetectable X-ray emissions. Because the current available models of a single streamer collision failed to explain the observations, we first use a Monte Carlo model coupled with multiple static dielectric ellipsoids immersed in a subbreakdown ambient electric field as a description of multiple streamer environment and we investigate the ability of multiple streamer-streamer head-on collisions to accelerate runaway electrons >1 keV up to energies ~200–300 keV instead of just one single head-on collision. The results of simulations show that the streamer head-on collision mechanism fails to accelerate electrons; instead, they decelerate in the positive streamer channel. In a second part, we use a streamer plasma fluid model to simulate a new streamer-electron acceleration mechanism based on a collision of a large negative streamer with a small neutral plasma patch in different Laplacian electric fields $|E_0| = (35, 40, 45)$ kV/cm, respectively. We observe the formation of a secondary short propagating negative streamer with a strong peak electric field >250 up to 378 kV/cm over a time duration of ~0.16 ns at the moment of the collision. The mechanism produces up to 10^6 runaway electrons with an upper energy limit of 24 keV.

1. Introduction

The mechanisms behind the production of runaway electrons, X- and gamma-rays by laboratory spark discharges, lightning leaders, and thunderstorms are still not yet understood. How such runaway electrons are produced is a key question in the field of atmospheric and space electricity. We know that negative lightning stepped leaders produce bursts of X-rays (Dwyer et al., 2003; Dwyer, Rassoul, Al-Dayeh et al., 2005; Moore et al., 2001), and many observation campaigns have been made to understand those emissions (e.g., Dwyer et al., 2011; Howard et al., 2008; Saleh et al., 2009; Schaal et al., 2014). The observations of X-rays were numerically analyzed by Babich et al. (2013). Laboratory spark discharges ignited by megavolt Marx generators produce X-rays as well (Dwyer, Rassoul, Saleh et al., 2005; Kochkin et al., 2012, 2015; March & Montanyà, 2011; Nguyen et al., 2008, 2010; Rahman et al., 2008). The presence of plasma streamer filaments in both lightning leader tips and laboratory spark discharges leads to the suggestion that streamers are behind those high-energy radiations, since streamers carry strong peak electric fields at their tips that may reach values higher than ~270 kV/cm at ground pressure under some circumstances. The value of 270 kV/cm is a deterministic electric field (e.g., Babich & Stankevich, 1973; Babich, 2003; Green & Peterson, 1968; Moss et al., 2006; Peterson & Green, 1968) necessary to accelerate the thermal low-energy electrons (tens of electron volts, e.g., Eichwald et al., 2006) present in the streamer head to high energies >120 eV. At this field the electric force exceeds the friction force caused by air molecules, resulting in so-called thermal runaway electron acceleration. The investigation of X-rays from laboratory sparks and explosive emission of electrons in vacuum have been a focus of many groups for many years with a significant body of literature published so far (e.g., Aleksandrov, 1966; Babich et al., 1990; Babich, 2005; Babich & Stankevich, 1973; Borukhov et al., 1973; Bugaev et al., 1975; Gurevich, 1961; Kremnev & Kurbatov, 1972; Litvinov et al., 1983; Mesyats et al., 1972; Stankevich & Kalinin, 1968; Stankevich, 1971; Tarasova & Khudyakova, 1970). Understanding the high-energy emissions from both lightning and laboratory sparks may be important for understanding terrestrial gamma ray flashes, which is the main focus of the recently launched ASIM

instrument (Neubert, 2009; Neubert et al., 2019) to the International Space Station, or missions under preparation such as the TARANIS satellite (Lefevre et al., 2008).

One of the mechanisms suggested to produce thermal runaway electrons is the head-on collision between negative and positive streamers, first proposed by Cooray et al. (2009). Head-on collisions may happen frequently in a laboratory spark discharge and during the connection between a main negative lightning leader and a bidirectional space leader. Cooray et al. (2009) used an electrostatic approach to study the problem of head-on collisions between a large number of positive and negative streamers. Processes associated with the dynamics of the streamers, ionization, attachment, conductivity, and relaxation were not considered in their study. Ihaddadene and Celestin (2015) reported the first numerical study of the production of X-rays by a single streamer head-on collision in a small gap of 8 mm, in an ambient field up to 60 kV/cm using a streamer fluid model and showed a significant increase of the electric field (~ 267 kV/cm) as the positive and negative streamer heads neutralize each other. However, such a high field exists over a very short time scale < 10 ps. They showed a production of maximum number of 2,000 thermal runaway electrons > 1 keV per single streamer collision. The 2,000 runaway electrons were estimated in two separate steps. They first modeled the increase of the electric field in a streamer head-on collision between negative and positive streamers, then they used the time-varying electric field ($E_{\text{max}} > 250$ kV/cm) in a Monte Carlo simulation as a homogeneously distributed electric field in space to maximize the production of thermal runaway electrons. Both the number and the energy of the runaway electrons were insufficient to explain the energy of the X-rays (50–500 keV) observed in the experiments (Dwyer, Rassoul, Saleh et al., 2005; Dwyer, Rassoul, Al-Dayeh et al., 2005). Kochkin, Köhn, et al., (2016) estimated a necessary number of 10^6 runaway electrons from head-on collisions with energies ~ 100 keV per one electron to explain the observations of 10^4 X-ray photons in a laboratory experiment. A criticism was given when the results of Ihaddadene and Celestin (2014) were first presented to the scientific community because the potential difference used ($60 \text{ kV/cm} \times 0.8 \text{ cm} = 48 \text{ kV}$) was not large enough to accelerate electrons to high energies compared to the 1 MV potential difference available in laboratory long gap experiments. In addition, the Ihaddadene and Celestin (2015) streamer head-on collision and the runaway electrons simulations were not self-consistent and the parallel plates boundary conditions were not a perfect representation of laboratory experimental conditions. Babich and Bochkov (2017) simulated a streamer head-on collision in a 5-cm gap with a 250-kV potential difference and they found similar conclusions as in Ihaddadene and Celestin (2015). Babich and Bochkov (2017) also used electron transport coefficient derived using the BOLSIG+ code (Hagelaar & Pitchford, 2005) instead of Morrow and Lowke (1997) along with open boundary conditions, and the effects did not change the conclusions of Ihaddadene and Celestin (2015). Moreover, da Silva et al. (2017, and references therein) detected X-ray photons ranging from 33 to 96 keV in a 4- to 8-cm short gap experiment with a lower applied voltage of 100 kV. The streamers they observed were 1-cm long and the X-rays were not associated with any streamer head-on collision (see also Kochkin et al., 2015, section 3.5.2). The X-ray bursts observed in long gap experiment are more intense than those observed in a short gap experiment presumably because the electrons experience a larger potential difference when they cross the gap (e.g., Kochkin et al., 2012, 2015; da Silva et al., 2017). Köhn et al. (2017) used full self-consistent particle in cell (PIC) simulations to investigate the acceleration of electrons during a single streamer head-on collision and they found no more than one electron with energy > 500 eV was produced. They drew the same conclusions as in Ihaddadene and Celestin (2015). Luque (2017) investigated the radio emissions from a single streamer head-on collision in long gap up to 30 cm using a full Maxwell's equations model in a protrusion-protrusion geometry and an applied electric potential of 750 kV. To estimate the energy of runaway electrons he used the one-dimensional deterministic equation for the acceleration of a single test electron in a homogeneous electric field. He assumed an already produced 1-keV electron in the high field zone and found a gain after the acceleration in few millimeters to 2 cm of 5 keV and 100 keV, respectively. The acceleration of an electron in a time-varying electric field (collapse) is different from that in a static electric field profile pulse because of the very short duration of the streamer head-on collision (picosecond).

Another mechanism that may produce thermal runaway electrons is the process of streamer branching (Arrayás et al., 2002; Kulikovskiy, 2000, 2001, 2002; Liu & Pasko, 2004; Moss et al., 2006; Pancheshnyi & Starikovskii, 2001; Rocco et al., 2002). Prior to the branching the streamer radius expands and reaches its maximum value of ~ 0.2 – 0.4 cm with an accumulated potential drop of ~ 100 kV (e.g., Celestin & Pasko, 2011). As a streamer approaches its maximum radius, the peak electric field in the streamer head is expected to increase in order to keep the preionization ahead of the streamer head on a similar level

(Liu & Pasko, 2004, sections 4.1 and 4.4). This may result in a peak streamer field >250 kV/cm. However, up to the present time no full physical 3-D simulation of streamer branching has been conducted to verify such expected streamer behavior, even though there have been great advances in 3-D streamer simulations in recent years (e.g., Li et al., 2009; Papageorgiou et al., 2011; Plewa et al., 2018; Shi et al., 2017; Teunissen & Ebert, 2017). This branching hypothesis has been assumed in Moss et al. (2006) to justify the high electric field used in their Monte Carlo simulations to study the production of thermal runaway electrons by a negative streamer. Moss et al. (2006) used a ~ 300 -kV/cm square pulse electric field in various ambient electric fields ~ 75 kV/cm to ~ 150 kV/cm. However, the description of a streamer by a square pulse in ambient electric fields higher than what may exist in a leader tip (~ 50 kV/cm; Raizer, 1991, p.363) may overestimate the energy and the flux of runaway electrons. Later on, Celestin and Pasko (2011) came up with a refined electric field shape of a streamer where a streamer simulation in an ambient electric field of 50 kV/cm was used. They estimated analytically the streamer parameters prior to the branching such as the radius and the potential drop, and then they injected in a Monte Carlo simulation the streamer electric field multiplied by a factor ~ 1.5 to reach a peak field of ~ 270 kV/cm required for the production of thermal runaway electrons. They accelerated the electrons in a maximum available potential difference in the streamer head of ~ 100 kV as they deduced from their analytical calculations. Celestin and Pasko (2011, section 4) came up with a runaway electron average energy value of ~ 60 keV.

Chanrion and Neubert (2010) proposed a single negative streamer mechanism that accelerates thermal electrons to explain terrestrial gamma ray flash observations. Their work is based on calculations by an advanced PIC-MCC code with an adaptive resampling scheme, the stochastic probability for an electron to runaway in a homogeneous electric field. They used that probability in a streamer PIC model to estimate in a stochastic manner the production of thermal runaway electrons. The probability was estimated in a homogeneous electric field that was not affected by the space charge. Based on their analysis, for every electric field value there is a certain probability for an electron to run away. The higher the number of available thermal electrons is, the higher the number of runaway electrons is generated $N_e \times P_e(E) = N_{\text{run}}$. For $E \geq 9E_k = 288$ kV/cm, where E_k is the breakdown electric field, the probability is equal to one. Because of the computational cost their streamer model does not include the photoionization, which is known to affect the propagation of both negative and positive streamers (e.g., Kulikovskiy, 2001, 2002; Liu & Pasko, 2004). Moreover, they used a background electric field up to ~ 100 kV/cm and a negative streamer peak electric field reached a value of ~ 320 kV/cm. When high ambient electric fields 100–1,000 kV/cm are used in experiments, a homogeneous ionized background develops and dominates the streamer regime (e.g., Mesyats & Bychkov, 1968). The first runaway electron with energy up to 171 keV appeared when the electric field peaked at ~ 157 kV/cm (Chanrion & Neubert, 2010, see Figure 3) instead of the deterministic value of ~ 250 kV/cm. From the point of view of the peak electric field the deterministic approach constrains more the ignition of the first population of thermal runaway electrons compared to the stochastic approach. For the case of a positive streamer, the acceleration is inefficient even though the peak electric field is higher than the negative one because the electrons flow into the weak field of the positive streamer body (Chanrion & Neubert, 2010, para. 34).

Ihaddadene (2016, pp.101–114) used a streamer plasma fluid model to reproduce the experimental results of Kremnev and Kurbatov (1972). Thermal runaway electrons and associated X-rays were produced by a single negative streamer in a 4-mm gap and a high ambient electric field 50 kV/cm at a standard pressure air (76–760 torr). Ihaddadene (2016) showed a production of runaway electrons between 1 and 6 keV and X-ray total energy between 0.00125 and 0.26 nJ in an ambient air density variation of 50% to 80%, respectively. They found a reasonable agreement between modeling and experiment. Recently, Köhn, Chanrion, Babich, et al. (2018) used a PIC code to simulate a double-headed streamer in a similar ambient field and size gap subjected to spatial sinusoidal air density variations along r and z axes between 15% and 80% to study the properties of streamers. Their simulation results showed the production of electron patches at the front of the main positive streamer possibly due to the photoionization. They stated that those patches acted as ignition points for negative streamers that collided afterward with the main positive streamer and then produced runaway electrons up to 3 keV in the 80% ambient air density variation case. The electron patches could be similar to the known stem structures or pilot systems usually observed in laboratory spark experiments described by, e.g., Kochkin et al. (2015) and Reess et al. (1995). However, as the model is a 2.5 Monte Carlo code, the patches have a cylindrical shape. Reducing the air density and keeping the same applied field is by definition an increase of the reduced electric field $\frac{E}{N}$ and equivalent to high ambient fields that were used by, e.g., Chanrion and Neubert (2010), Moss et al. (2006), and Ihaddadene (2016, pp.101–114).

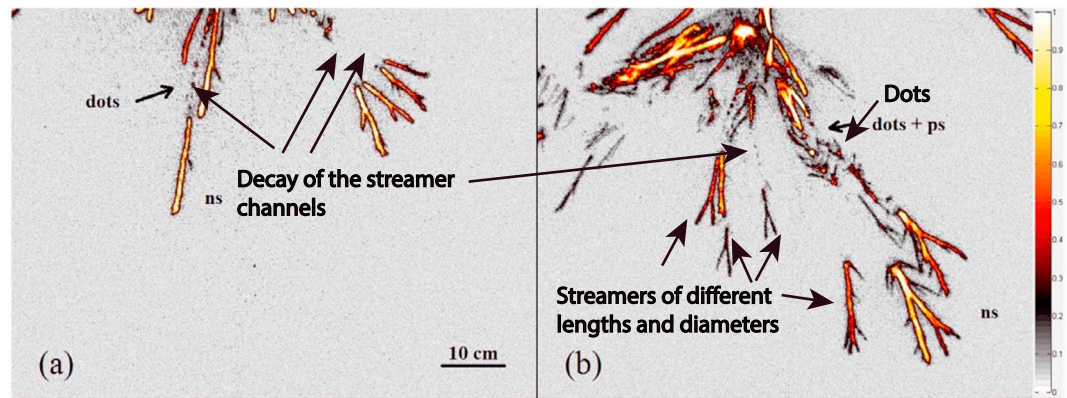


Figure 1. (a)–(b) Multiple visible dots, short and long streamers, streamer branching, and the complex morphology of discharge plasmas in a laboratory spark discharge. Adapted from (Kochkin et al., 2014).

Interestingly, if low ambient air density pockets occur in laboratory sparks, this indicates that the head-on collision is not required and a single streamer alone could produce thermal runaway electrons. It would be interesting to quantify at which value of low electric field and minimum density change the X-rays can be produced. Recently, Köhn, Chanrion and Neubert (2018) made a focused study on high-energy radiation from streamers and concluded that air density perturbations may be a candidate for the acceleration of runaway electrons.

As a single-streamer head-on collision poorly describes laboratory spark discharges because of the large number of streamers present (see Figure 1), in this work we investigate the multiple-streamer head-on collisions problem and the motion of a negative streamer into a neutral plasma patch. In the first part, we use a uniform distribution of multiple static dielectric ellipsoids as a streamer-like representation of the electric field in a Monte Carlo model to search for the configuration that efficiently accelerates the electrons. In the second part, we use a streamer plasma fluid model and simulate a new streamer-electron acceleration mechanism based on the collision between a large negative streamer and a small neutral plasma patch (see Figure 2). We will show that this mechanism is more efficient to accelerate electrons than streamer-streamer head-on collision.

2. Model Formulation

2.1. Monte Carlo Dielectric Ellipsoids Model

The Monte Carlo model we use in the present study is the 3-D runaway electron avalanche model that simulates the physics of runaway electron avalanche from ~ 1 keV to several hundred megaelectron volts. The model is fully described by Dwyer (2007, 2012). To study the dynamics of runaway electrons in a streamer head-on collision we implement in a Monte Carlo model an inhomogeneous electric field profile generated by one, two, or more static dielectric ellipsoids embedded in a homogeneous ambient electric field E_0 . The electrostatic analytic solution of the ellipsoid electric potential is well known (e.g., Landau & Lifshitz, 1960; Stratton, 1941, pp.207–217, pp.19–42), and we derive the electric field inside and outside the ellipsoid by taking the gradient of the following equations (1) and (2), respectively:

$$\phi = -\frac{E_0 z}{1 + \left(\frac{abc}{2}\right)\left(\frac{\epsilon_1 - \epsilon_2}{\epsilon_2}\right) \int_0^\infty \frac{ds}{(s+a^2)R_s}} \quad (1)$$

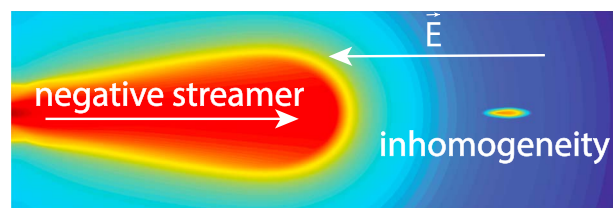


Figure 2. Illustration of a collision between a negative streamer (left) and a small plasma patch (right) in an external electric field \vec{E} .

$$\phi = \phi_0 \left(1 - \frac{\left(\frac{abc}{2}\right)\left(\frac{\epsilon_1 - \epsilon_2}{\epsilon_2}\right) \int_{\zeta}^{\infty} \frac{ds}{(s+a^2)R_s}}{1 + \left(\frac{abc}{2}\right)\left(\frac{\epsilon_1 - \epsilon_2}{\epsilon_2}\right) \int_0^{\infty} \frac{ds}{(s+a^2)R_s}} \right) \quad (2)$$

where $\phi_0 = -E_0 z$, $R_s = (s+a^2)^{\frac{1}{2}}(s+b^2)^{\frac{1}{2}}(s+c^2)^{\frac{1}{2}}$, $\zeta = \frac{1}{4}(r_1+r_2)^2 - a^2$, $r_1 = ((x-x_0)^2 + (y-y_0)^2 + (z-z_0-ae)^2)^{\frac{1}{2}}$, $r_2 = ((x-x_0)^2 + (y-y_0)^2 + (z-z_0+ae)^2)^{\frac{1}{2}}$, and $e = (1 - \frac{b^2}{a^2})^{\frac{1}{2}}$. The parameters a , b , and c are the semiprincipal axis of the ellipsoid along \vec{e}_z , \vec{e}_x , and \vec{e}_y directions; r_1 and r_2 are the distances of a given point to the foci of the ellipsoid and e its eccentricity. The coordinates (x, y, z) and (x_0, y_0, z_0) correspond to the positions of the electron and the ellipsoid, respectively. Based on different ellipsoid sizes, between 4-mm and 14-cm long, extremities of two ellipsoids separated by 1 to 4 mm along the z direction, and parallel ellipsoids by 10 to 30 mm along the x direction, in ambient electric fields $E_0 = 10$ –30 kV/cm, we test whether or not the runaway electrons continue to accelerate. We position the electrons in different locations between the ellipsoid heads, in the middle of the ellipsoid channel equivalently to an already collapsed collision, and in the background field in between the ellipsoid bodies. We calculate analytically the corresponding electric field at a given location of the electron (x, y, z) . In this study we take $c = b$. The parameter ϵ_1 is the permittivity of the ellipsoid that is embedded in a homogeneous medium of inductive capacity ϵ_2 , which is the air in this case. Equivalently to the streamer channel, ϵ_1 is chosen to adjust the value of the electric field inside the ellipsoid body to approximate the streamer channel.

2.2. Negative Streamer and Plasma Patch Model

To study the collision between an expanding negative streamer and a small neutral plasma patch, we use the streamer model based on the drift-diffusion equations for charged species coupled with Poisson's equation and which is fully described in Ihaddadene and Celestin (2015, 2017).

We choose to use a sphere-to-plane electrode configuration (e.g., Babaeva & Naidis, 1996a; 1996b) to initiate a single negative streamer connected to a point electrode with open boundary conditions. To obtain a maximum amplitude of the electric field of $| \sim 4E_k |$ at the surface of the sphere ($E_{\text{sph}} = \frac{\phi_0 + 3E_0 R}{R}$) (e.g., Liu et al., 2006, 2009) a sphere electrode of a radius $R_{\text{sph}} = 10^{-3}$ m is set to a potential $\phi_0 = -1.5$, 0, and 1.5 kV and placed in uniform electric fields $E_0 = -35$, -40 , and -45 kV/cm, respectively. The value of the breakdown electric field $E_k = 28.7$ kV/cm is calculated based on Morrow's coefficients (Morrow & Lowke, 1997). We have to mention that Morrow's coefficients underestimate the ionization frequency with respect to BOLSIG+ coefficients as demonstrated by Babich and Bochkov (2017, see Figure 1) and so one expects higher streamer electric fields than the results presented in this paper if one uses BOLSIG+ coefficients. In this study, we consider $E_0 = -35$ and $E_0 = -45$ kV/cm as lower and upper ambient electric fields. These electric field amplitudes are widely used in the streamer literature to describe the strength of the field in the streamer ignition zone in laboratory spark discharges and lightning leader tips (e.g., Babich et al., 2015; Babich & Bochkov, 2017; Bazelyan & Raizer, 1998; Rakov & Raizer, 2000; Chanrion & Neubert, 2010; Celestin & Pasko, 2011; Ihaddadene & Celestin, 2015; Köhn et al., 2017p. 76-84,p. 62-74). One case with a sphere of a radius $R_{\text{sph}} = 10^{-2}$ m, $\phi_0 = -6.5$ kV, and $E_0 = -28$ kV/cm over a short gap of 8 mm is also presented.

The negative propagating streamer is initiated by placing a Gaussian $n_e = n_{e_0} \exp(-(z-2\sigma_z)^2/\sigma_z^2) \exp(-r^2/\sigma_r^2)$ of neutral plasma cloud with characteristic sizes $\sigma_z = 2 \times 10^{-4}$ m, $\sigma_r = 2 \times 10^{-4}$ m, and $n_{e_0} = 10^{20} \text{ m}^{-3}$ near the vicinity of the sphere electrode. The size of the simulation domain varies from $1,001 \times 377$ to $2,625 \times 377$ regular grid points along z and r axis, respectively, with a spatial resolution defined by $\Delta z = \Delta r = 8 \times 10^{-6}$ m. The dimensions of the Gaussian neutral plasma patch are $\sigma_z = 2 \times 10^{-4}$ m, $\sigma_r = 15 \times 10^{-6}$ m, and $n_{e_0} = 10^{18} \text{ m}^{-3}$, and it is located at a distance d_{patch} from the sphere electrode.

We choose to use a weak peak electron density $n_{e_0} = 10^{18} \text{ m}^{-3}$ to prevent the initiation of streamers from the plasma patch and avoid conventional streamer-streamer collision. The patch is positioned at different locations $d_{\text{patch}} = (0.5, 1, 1.4)$ cm in -40 and in -45 kV/cm, at $(0.5, 1.4)$ cm in -35 kV/cm, and at 0.5 cm in -28 kV/cm.

3. Results

3.1. Streamer-Streamer Head-On Collision

The multistreamer discharge problem is complex to model if electromagnetic, hydrodynamic, and kinetic processes are included in one single model. In laboratory discharge experiments multiple streamers ($\sim 10^6$

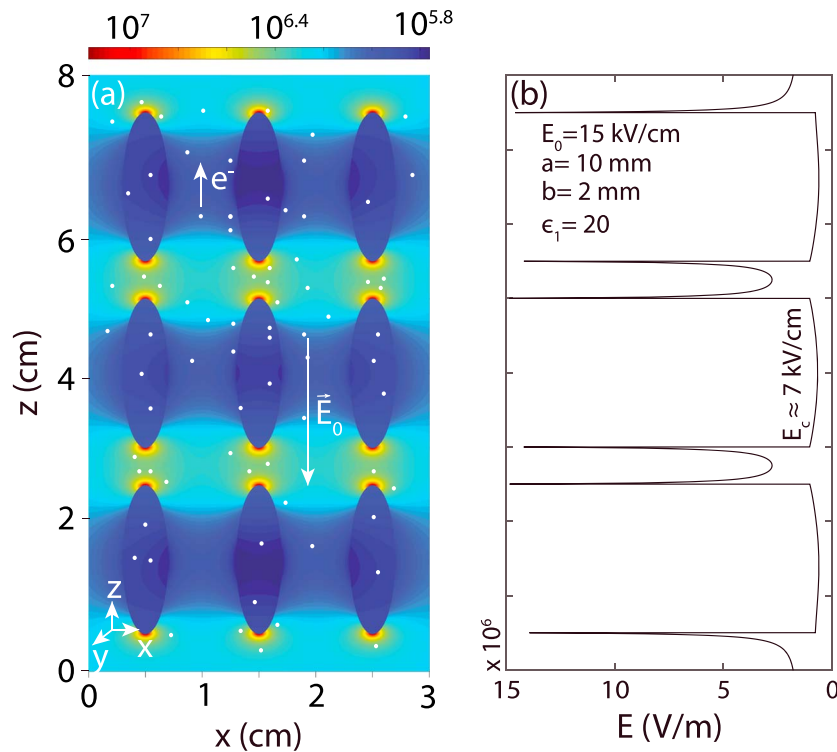


Figure 3. (a) Cross-sectional view of the electric field produced by several dielectric ellipsoids immersed in an ambient electric field $E_0 = 15$ kV/cm. The white dots illustrate electrons moving in between the ellipsoids. (b) Electric profile along the z axis (one streamer column).

Bazelyan & Raizer, 1998, section 2.4) with high localized electric fields are distributed inhomogeneously (see Figure 1) where the electric charge is a conserved quantity. From previous PIC simulations of a single-streamer head-on collision in air at ground pressure, results seem unlikely to explain the observations of X-ray emissions, with no more than one electron of 500 eV produced under an ambient electric field of 60 kV/cm (Köhn et al., 2017). In this work, we explore whether multiple streamers combined with multiple head-on collisions can possibly accelerate the electrons (≥ 1 keV) produced by a single-streamer head-on collision (Ihaddadene & Celestin, 2015; Köhn et al., 2017) to the higher values (50–500 keV) observed in the experiments. The hypothesis is that the electrons produced in the collision zone could cross the body of the positive streamer, reach another negative streamer head, then another positive streamer head, or another streamer head-on collision, and so on until finally the electrons would gain enough energy from multiple streamers to run away in a discharge gap. To test this hypothesis, we use a static distribution of multiple dielectric ellipsoids of different lengths and diameters equivalent in size to streamers we usually observe in laboratory experiments (see Figure 3) with peak and body electric fields reasonably close to streamer simulation results. Similar modeling approaches and different electric field configurations have been previously used to study runaway electron production (e.g., Köhn & Ebert, 2015; Kochkin, Köhn, et al., 2016; Kochkin, Lehtinen, et al., 2016). For an average applied electric field over a certain length the available potential difference must be conserved and the introduction of a dielectric ellipsoid approach is a good approximation for the multistreamer problem, as it modifies the uniform electric field to high and low field zones without introducing additional external charge or affecting the available potential difference. The ellipsoids are a simplified representation of a collection of streamers occurring in a laboratory discharge. The well-known analytic solution (equations (1) and (2)) makes it easy to implement in a Monte Carlo model.

Figure 3a shows the cross-sectional view of one of the configurations of multiple dielectric ellipsoids immersed in an ambient field $E_0 = 15$ kV/cm close to the average experimental value $\frac{1}{1}$ MV/m = 10 kV/cm. One can see the encounter of multiple ellipsoid heads forming multiple high field zones similar to streamer head-on collisions where runaway electrons can be produced. Based on the orientation of the ambient electric field, the upward and downward ellipsoid heads refer to negative and positive streamers, respectively.

Table 1

Runaway Electron Energy ϵ_{run} Required in Different Homogeneous Background Electric Fields, Deterministic (Equation (3)), and Stochastic Approach

E (kV/cm)	Deterministic ϵ (keV)	Stochastic ϵ (keV)
10	36	60
20	15	40
30	9	20
40	6	18
50	4	15
60	3	10
70	2.6	8
80	2	6
90	1.8	5
100	1.5	3

Note. REAM = 3-D runaway electron avalanche model.

From streamer simulations, we know that the positive streamer head field is higher than the negative one and inversely in their channels, but using ellipsoid approximation, we are unable to simulate this detailed property of streamers. The ellipsoid does not take into account the physical processes such as ionization, relaxation, production of electron density, or a continuous flowing charge through the channel. From Figure 3b one can see the nearly homogeneous field in the body of the ellipsoid and both heads possess almost similar maximum electric field. We use the superposition principle to estimate the total electric field produced by the ellipsoids. The sum of the electric field components do affect each other slightly and that is why we see a small increase of the field inside the body and the head of the ellipsoid in the middle of the figure.

The use of a simple linear superposition is justified by the fact that it leads to a reasonable representation of the electric field in streamer heads and channels as obtained in simulations. In this context the runaway electron acceleration can be easily maximized through a reorganization of the ellipsoids. We do not claim that the electric field obtained is a valid solution to the electrostatic problem involving real dielectric ellipsoids as the effect of mutual interactions should be taken into account.

To investigate the streamer head-on collision efficiency in the multiple streamer configuration, we assume there already exist 2,000 runaway electrons $\epsilon_i \geq 1$ up to 50 keV produced by a single collision (Ihaddadene & Celestin, 2015). The assumed initial energy ϵ_i is lower than the runaway threshold $\epsilon(E_0)$ for an electron to run away in an ambient field $E_0 < E_k$. After multiple combined tests we find that no electron was able to run away. The short ellipsoids do not provide enough potential drop in the high field region to accelerate the electrons and the long ones where the electrons gain energy in between the ellipsoid heads decelerate the electrons in the low body field of the long ellipsoid channel before they reach the other negative streamer or collision point.

A single long and large ellipsoid is more efficient than two colliding ellipsoids to accelerate a 1-keV electron produced at the tip (i.e., equivalent to a large streamer at the branching stage). Electrons run away easily ahead of a single ellipsoid in the ambient field in contrast to their deceleration in the channel field in a case of two colliding ellipsoids.

Table 1 shows the energy required for an electron to runaway in a different homogeneous electric fields established by a Monte Carlo and a one-dimension deterministic equation for a test electron ($m_e \frac{dv_e}{dt} = qE - F$. For details see section 3.3). As one can see the stochastic energy is greater than the deterministic one by up to ~50% due to the angular scattering which is not included in the one-dimension equation of motion. In reality, an electron does not travel in a straight line parallel to the electric field but goes through a longer path because of the scattering and requires more energy to runaway in that field. For instance, in a 10-kV/cm streamer channel field a streamer head-on collision must produce electrons >60 keV in the collision zone to allow them to cross the channel. This means a potential drop >60 kV in a high field zone must be available;

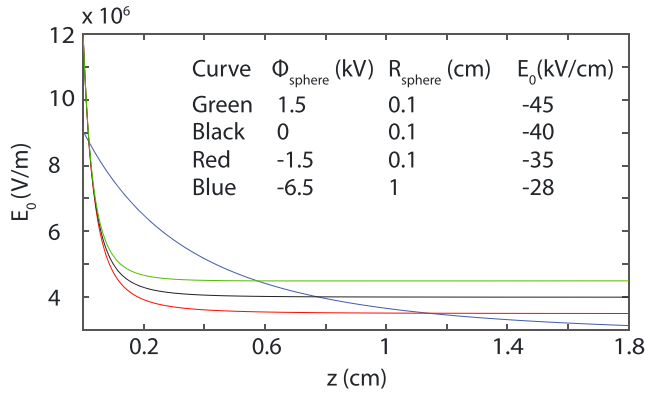


Figure 4. Ambient electric fields E_0 produced by a sphere-to-plane electrode.

otherwise, the electrons will decelerate in the channel. To achieve an efficient acceleration by a conventional streamer-streamer head-on collision, a positive streamer channel field must be strong enough to accelerate the runaway electrons that are produced at the collision point ($\epsilon_{\text{run}} > \epsilon_{\text{th}}^c(E_c)$).

3.2. Streamer-Patch Collision

From the results obtained in the previous section, the conventional streamer-streamer head-on collision appears unlikely to be an efficient electron acceleration mechanism even when considering reacceleration in a multiple-streamers configuration. In this section we develop a new picture of a streamer collision that involves a large negative streamer and a neutral plasma patch or inhomogeneity seen from experiments (see Figure 1). The figure is essentially a 2-D representation of the 3-D phenomenon. Thus, even though it appears that streamers might interact with either one another or with inhomogeneities, it could be that they are separated in depth and no interaction takes place (Nijdam et al., 2008).

This is an alternative solution to a conventional streamer-streamer colli-

sion channel problem as it does not involve a positive streamer but only a negative streamer moving forward with the electrons in the ambient field.

Figure 4 shows the Laplacian electric fields used in the current study. One clearly sees a highly enhanced electric field near the sphere electrode <0.2 cm and homogeneous >0.2 cm for the green, black, and red curves. The blue curve shows an inhomogeneous ambient field everywhere.

Figures 5a and 5b show the cross-sectional view of the electron density and the electric field at the moment of the collision between the negative streamer and the small neutral plasma patch in case of $E_0 = -40$ kV/cm. The figure shows a transition of the patch to a short negative streamer with an enhanced field at its tip due to the influence of the long negative streamer. Figures 5c and 5d show the electron density and the electric field along the z axis. One sees clearly a rightward propagating negative streamer with an unstable field and an increasing speed toward the patch. During the collision the initial electron density in the patch increases to a level that initiates a streamer because of the influence of the long streamer peak field. The charge flows from the large streamer head to the small patch, the potential drop transfers as well, and a short negative streamer starts to form and reaches a high electric field amplitude up to 325 kV/cm (see Movie S1 in the supporting information). As the short streamer develops, the radius increases causing the electric field to decrease, returning to the normal value of the previous long streamer. The maximum field value reached appears to be higher than the thermal runaway threshold 270 kV/cm as defined by the maximum friction

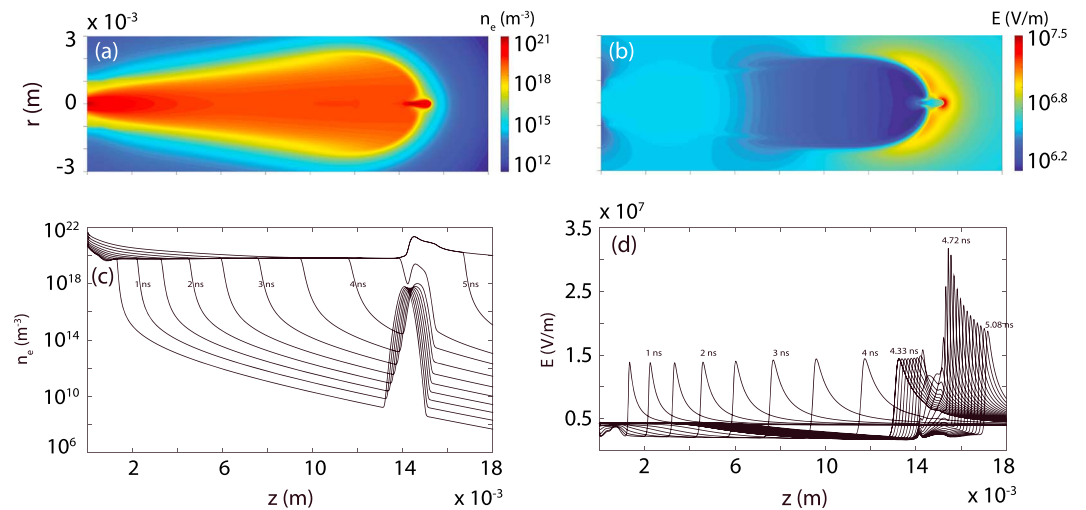


Figure 5. Cross-sectional view of the electron density (a) and electric field (b) at 4.72 ns in the case of $E_0 = -40$ kV/cm, respectively. Profile of the electron density (c) and the electric field along the z axis (d) in the case of $E_0 = -40$ kV/cm, respectively. In panel (d) results are shown with a time step of 0.5 and 0.03 ns after $t = 4.3$ ns.

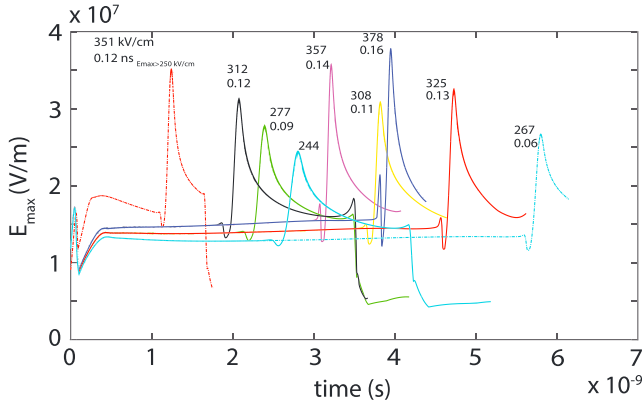


Figure 6. Evolution of the maximum electric field E_{\max} as function of time. Black, violet, and blue curves correspond to $(E_0, d_{\text{patch}}) = (-45 \text{ kV/cm}, 0.5 \text{ cm})$; $(-45, 0.5)$; $(-45, 1.4)$, respectively. Green, yellow, and red curves correspond to $(E_0, d_{\text{patch}}) = (-40, 0.5)$; $(-40, 1)$; $(-40, 1.4)$, respectively. Cyan and cyan dashed curves correspond to $(E_0, d_{\text{patch}}) = (-35, 0.5)$; $(-35, 1.4)$, respectively. Red dashed curve corresponds to $(E_0, d_{\text{patch}}) = (-28, 0.5)$.

force at $\sim 150 \text{ eV}$ (e.g., Babich & Stankevich, 1973; Babich, 2003; Green & Peterson, 1968; Moss et al., 2006; Peterson & Green, 1968).

Figure 6 shows the evolution of the maximum electric field in the simulation domain versus time for a variety of ambient electric fields and patch locations. The maximum electric field value obtained is between 244 and 378 kV/cm and it corresponds to the cases $(E_0 = -35 \text{ kV/cm}, d_{\text{patch}} = 0.5 \text{ cm})$ and $(-45 \text{ kV/cm}, 1.4 \text{ cm})$, respectively. We see clearly that the maximum electric field increases with increasing the distance of the plasma patch for a given ambient electric field. The explanation for this is that the maximum electric field in the patch is controlled by the amount of the available potential drop in the long streamer head. The longer the streamer channel the larger the quantity $\int_0^l (E_0 - E_c) dl$ to be transferred to the patch. A large amount of potential difference concentrated in a small volume creates a strong local electric field. This phenomena is equivalent to two conducting spheres of different radii $r_1 > r_2$ of potentials $U_1 > U_2$ connected with a wire and while at the same potential $U'_1 = U'_2$ the electric field ratio of the two spheres $\frac{E_2}{E_1} = \frac{r_1}{r_2}$ shows that the field is higher at the surface of the sphere of small radius. This is a good approach if one considers a streamer-sphere shape model (e.g., Gallimberti et al., 1974). However, in case of an accurate streamer description the geometry is more complex and a geometrical factor related to the streamer head shape must be taken into account in the above ratio. For instance, if a streamer is described as a long conducting filament (e.g., D'yakov & Kachorovskii, 1989) a logarithm of the length to the radius of the channel $\ln(\frac{l}{r})$ must be included in that ratio

as follows: $\frac{E_2}{E_1} = \frac{r_1}{r_2} \frac{\ln(\frac{l_1}{r_1})}{\ln(\frac{l_2}{r_2})}$, where (l_1, l_2) and (r_1, r_2) are the lengths and radii of the long and short streamers, respectively.

In the case of an inhomogeneous ambient field shown by the blue curve in Figure 4, the red dashed curve in Figure 6 shows an increasing streamer head field at $t < 0.5 \text{ ns}$ due to the large-sphere field enhancement and then a decreasing field as the streamer travel into a low ambient field region. The strong inhomogeneous ambient field produced by a large-sphere electrode seems efficient in reaching higher maximum electric fields of 351 kV/cm with a patch being located just at 5 mm than some of the homogeneous field cases.

Table 2 shows the streamer-patch collision parameters, the maximum electric field, the duration of the maximum field $> 250 \text{ kV/cm}$, the potential drop $\int_0^l (E_0 - E_c) dl$, the average streamer speed, and the ratio of the peak field E_{\max} to the streamer head field E_h . The reference value of E_h is taken as the value of streamer head electric field just before the encounter with the plasma patch. For a given ambient field the results clearly show an increase of the peak field and its duration versus the increase of the patch distance, the potential

Table 2 shows the streamer-patch collision parameters, the maximum electric field, the duration of the maximum field $> 250 \text{ kV/cm}$, the potential drop $\int_0^l (E_0 - E_c) dl$, the average streamer speed, and the ratio of the peak field E_{\max} to the streamer head field E_h . The reference value of E_h is taken as the value of streamer head electric field just before the encounter with the plasma patch. For a given ambient field the results clearly show an increase of the peak field and its duration versus the increase of the patch distance, the potential

Table 2

The Peak Electric Field E_{\max} , the Duration of the Peak Electric Field $E_{\max} > 250 \text{ (kV/cm)}$ t_{\max} , the Potential Drop $\int_0^l (E_0 - E_c) dl$, the Average Streamer Speed $V_{\text{str}} = \frac{l}{t}$, and the Ratio $\frac{E_{\max}}{E_h}$ for each Case in Figure 6

$E \text{ (kV/cm)},$ $d_{\text{patch}} \text{ (cm)}$	E_{\max} (kV/cm)	t_{\max} (ns)	$\int_0^l (E_0 - E_c) dl$ (kV)	$V_{\text{str}} \times 10^6$ (m/s)	$\frac{E_{\max}}{E_h}$
-35, 0.5	244	0	9	2	1.91
-35, 1.4	267	0.06	15	2.5	2.01
-40, 0.5	277	0.09	9.4	2.4	2.01
-40, 1	308	0.11	14.7	2.8	2.15
-40, 1.4	325	0.13	15.5	3	2.26
-45, 0.5	312	0.12	9.7	2.8	2.11
-45, 1	357	0.14	15.3	3.3	2.33
-45, 1.4	378	0.16	16.2	3.7	2.44
-28, 0.5	351	0.12	13	4.7	2.13

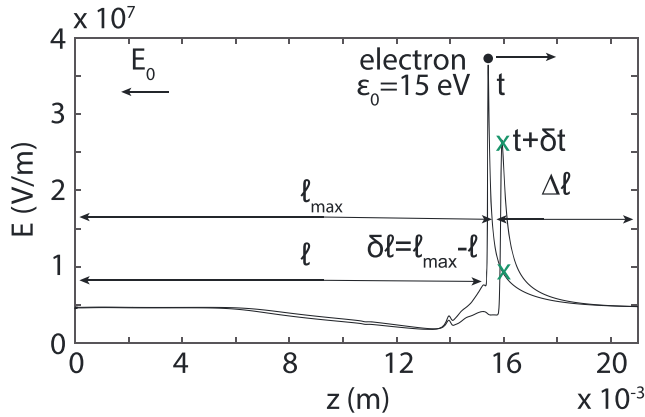


Figure 7. Illustration of the acceleration step of a test electron in a space-/time-varying streamer electric field. The results are taken from the case (−45 kV/cm, 1.5 cm) with a large time step 100 ps to show enough spatial scale. The quantities l_{\max} , l , δl , Δl , t , δt , and E_0 are the maximum streamer length when the maximum peak electric field is reached, the streamer reference length where the electric field suddenly increases, the streamer radius, the distance traveled by the runaway electron, the reference time for the streamer electric field data, the time step between two successive streamer electric field profiles, and the ambient electric field. The green cross shows two different values of the electric field at the same location when the dynamics of the streamer is taken into account.

drop, and the streamer speed. Note that the quantity $\int_0^l (E_0 - E_c) dl$ is estimated just before the collision stage and the integration is made from the surface of the sphere electrode up to $l = l_{\max} - \delta R$ where δR is by definition the streamer radius or the thickness of the streamer head charge layer.

3.3. Estimation of the Energy and the Number of Thermal Runaway Electrons in a Streamer-Patch Collision

To estimate the energy of thermal runaway electrons we solve numerically the deterministic equation (3) for a single test electron.

$$m_e \frac{dv_e(z_e, t)}{dt} = qE(z_e, t) - F(\epsilon(z_e, t)) \quad (3)$$

where the quantities m_e , $v_e(z_e, t)$, $E(z_e, t)$, $F(\epsilon)$, and z_e , are the electron mass, the electron velocity, the streamer electric field, the friction force, and the electron coordinate along the z axis of the streamer, respectively.

This deterministic approach has been used by several authors to study runaway electrons (e.g., Babich & Stankevich, 1973; Gurevich & Zybin, 2001; Gurevich et al., 1992, 2007; Luque, 2017; Zubarev et al., 2017). This approach maximizes the energy of the electron as it does not take into account the angular scattering, and the electron is assumed to travel parallel to the central axis of the streamer electric field. The effect of the angular scattering, the diffusion, and the cross section has been studied in detail in the framework of runaway electrons (e.g., Chanrion et al., 2016, and references therein). The estimation is still reasonable as it is made

over short distances of the order of 2–5 mm and in a nonrelativistic regime. The forward scattering for the electrons >50 eV prevail over the angular one (Pasko, 2006, pp. 261–265) especially in strong fields considered in this study (Babich et al., 1990, and references therein, p.523). Compared to the relativistic electrons traveling long distances the summation over all the small deviations is important and the angular scattering impacts the electron trajectory (Dwyer et al., 2012; Figure 7). As simulations are time consuming, we estimate just the lower and upper energy limits and we focus on the following cases (E_0 (kV/cm), d_{patch} (cm)) = (−45, 1.4), (−28, 0.5), and (−35, 1.4).

The friction force $F(\epsilon)$ that includes all the electron collision processes is tabulated as function of the electron energy based upon Moss et al. (2006). The maximum friction force is 260 keV/cm for an electron energy of ~ 120 eV. We use a numerical time step of $\delta t_{\text{num}} = 10^{-15}$ s to solve equation (3). A single test electron with an initial energy $\epsilon_0 = 15$ eV is placed at the streamer tip (Figure 7) and then it is accelerated in the streamer electric field $E(z_e, t)$ every step δt_{num} . At every step δt_{num} the electric field $E(z_e, t)$ and the friction force $F(\epsilon)$ are updated. To take into account the change of the electric field due to the moving streamer, we use streamer electric field data separated by a time step δt of the order of a picosecond (ps) to capture accurate streamer electric field changes. When the electron travels $n_{\text{iter}} \delta t_{\text{num}} > \delta t$ where n_{iter} is the number of numerical iterations, a streamer electric profile is update with a new one. The time steps are 0.2 ps for the cases (−35, 1.4) and (−28, 0.5), 1 ps for the cases (−45, 1.4) and (−28, 0.5). We choose 0.2 and 1 ps to run the fastest simulation (−28, 0.5) and to quantify the time step effect on the energy of the electron. Note that, if one uses the results of Figure 6 (the streamer head peak field) and assumes the electric field is distributed homogeneously following the approach of (Ihaddadene & Celestin, 2015, section 3.3) the electron synchronizes with the peak electric field and the energy is going to be overestimated. The runaway electrons produced at the streamer head exit the head and they follow a decreasing electric field $\sim E(\frac{1}{r})$.

Figure 8a shows the space and time-varying electric field traveled by a single test electron versus time. The electron acceleration is stopped when the electric field value falls off to 50 kV/cm. The black curve corresponds to the (−45, 1.4) case with a 1-ps time step between the data files, and blue and red curves correspond to the (−28, 0.5) case with 0.2- and 1-ps time steps, respectively. We first tested a single electron in a static streamer electric field and found that the electron failed to accelerate because of the lack of the available potential difference in the small radius of the short negative streamer as one can see from Figure 5d. However, once the streamer propagation is taken into account the electron is accelerated and gained energy. This

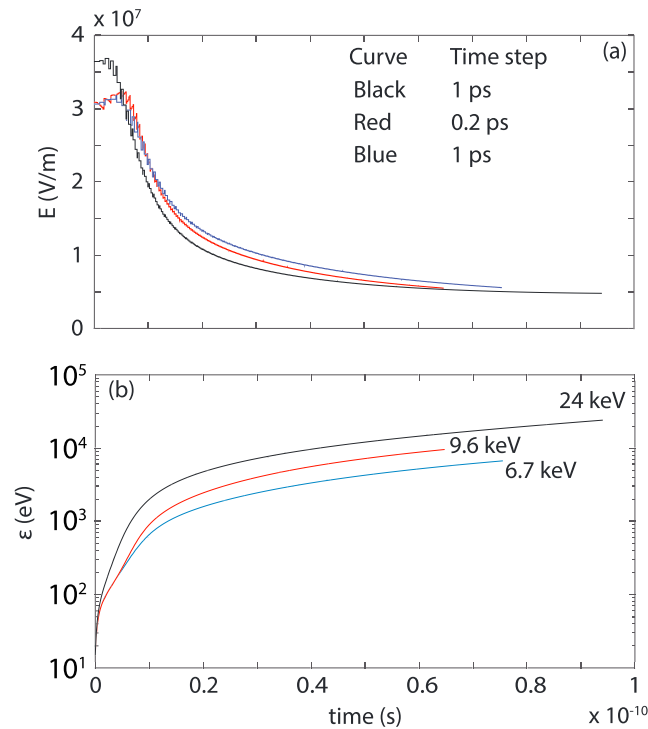


Figure 8. (a) Evolution of the space- and time-varying streamer electric traveled by a single test electron versus time. The black curve corresponds to the case $(E_0, d_{\text{patch}}) = (-45 \text{ kV/cm}, 1.4 \text{ cm})$, and the time step between two successive electric field data files is 1 ps. The blue and red curves correspond to $(E_0, d_{\text{patch}}) = (-28 \text{ kV/cm}, 0.5 \text{ cm})$, and time steps are 1 and 0.2 ps, respectively. (b) Evolution of the single test electron energy in the three cases versus time.

latter refers to a self acceleration mechanism introduced first by Babich (1982a, 1982b; Babich et al., 1990, p.534) and later on studied by Moss et al. (2006, section 4.2) who estimated a self-acceleration upper energy limit contribution no more than 9 keV by Monte Carlo simulations. The electron can gain more energy than the available potential difference if that electron stays surfing in the streamer head for a sufficient time. A streamer speed of $\sim 10^7 \text{ m/s}$ corresponds to an electron energy of $\sim 280 \text{ eV}$. Some of the electrons >120 up to 280 eV amplify their energy gain while traveling within the high field region of the streamer head.

At the end of the electron acceleration we estimate 124 eV, 6.7 keV, 9.6 keV, and 24 keV for the cases $(-35, 1.4)$, $(-28, 0.5)$, $(-28, 0.5)$, and $(-45, 1.4)$, respectively (Figure 8b). The values 124 eV, 6.7 keV, 9.6 keV, and 24 keV were computed with $\delta t = 0.2$ -, 1-, 0.2-, and 1-ps data time steps, respectively. The 24-keV electron reaches 200 keV in an ambient field of 45–30 kV/cm over 4–7 cm, respectively. We stop the electron acceleration when the electron reaches a location ahead of the streamer where the electric field is $\sim 50 \text{ kV/cm}$. We do not show the results for the case $(-35, 1.4)$ because the electron reaches 124 eV and then loses its energy. If a very small time step is used the above energy increases by 43% up to 177 eV as one can see from the

Table 3

Estimation of the Quantities $U_1 = \int_0^{t_f} E v_e dt$ and $U_2 = \int_{l_{\text{max}}}^{l_{\text{max}}+\Delta l} E dl$ where $\Delta l = \int_0^{t_f} v_e dt$, in Case (1): $50 \text{ kV/cm} < E < E_{\text{max}}$ and Case (2): $250 \text{ kV/cm} < E < E_{\text{max}}$

	(E (kV/cm), d (cm))			
	(-45, 1.4) Case (1)	(-45, 1.4) Case (2)	(-28, 0.5) Case (1)	(-28, 0.5) Case (2)
U_1 (keV)	39	2.8	21.9	2.13
U_2 (keV)	37	2.4	19.4	1.35

Table 4

The Number of Runaway Electrons N_{run} , the Flux of Runaway Electrons Φ , and the Number of Streamer-Patch Collisions N_{coll} Required to Explain the Experimental Results for each Case in Figure 6

E (kV/cm), d_{patch} (cm)	−35, 0.5	−35, 1.4	−40, 0.5	−40, 1	−40, 1.4	−45, 0.5	−45, 1	−45, 1.4	−28, 0.5
$N_{\text{run}} \times 10^3$	0	0.105	0.324	13	68	14	550	1870	280
$\Phi = \frac{N_{\text{run}}}{t_{\text{max}}} (\text{s}^{-1}) \times 10^{12}$	0	1.75	3.6	181.2	523	116.7	3,437.5	1,168.8	233.3
$N_{\text{coll}} = \frac{10^6}{N_{\text{run}}}$	—	10^4	10^3	77	15	71	2	1	4

results of (−28, 0.5) case. The electron energy increased from 6.7 to 9.6 keV when changing the time step from 1 to 0.2 ps.

We estimate the quantity $U_1 = \int_0^{t_f} E v_e dt$, where E and v_e are the space- and time-varying streamer electric field and the electron velocity, respectively. The quantity t_f stands for the end time of the electron acceleration and can be read from Figure 8. We compare U_1 with the quantity $U_2 = \int_{l_{\text{max}}}^{l_{\text{max}} + \Delta l} E dl$ where E is a single space-varying streamer electric field profile. The quantity $\Delta l = \int_0^{t_f} v_e dt$ is the distance traveled by an electron when it is accelerated. In the case (−45, 1.4), we find $U_1 = 39$ kV versus $U_2 = 7$ kV, and 2.8 kV versus 2.4 kV if we estimate both quantities in the following electric field interval: $250 \text{ kV/cm} \leq E \leq E_h \text{ kV/cm}$. In the case (−28, 0.5) with a 1-ps time step we find 21.9 kV versus 19.4 kV, and 2.13 kV versus 1.35 kV. One sees differences between the quantities U_1 and U_2 of 2, 2.5, and 0.4 kV, 0.78 kV when $250 \leq E \leq E_h \text{ kV/cm}$ for the cases (45, 1.4) and (28, 0.5), respectively. The values of U_1 and U_2 are also summarized in Table 3. From the estimations above, one can find an average energy loss over the ~ 5.5 -mm distance traveled by the electron of $\epsilon_{\text{loss}} = \frac{(39 \text{ keV} - 24 \text{ keV})}{5.5 \text{ mm}} \simeq 2.7 \text{ keV/mm}$ and so ~ 190 eV per streamer radius $\delta l \sim 7 \times 10^{-5} \text{ m}$. The average energy gain is $\epsilon_{\text{gain}} = \frac{24}{5.5} \simeq 4 \text{ keV/mm}$ and it corresponds to 280 eV per streamer radius. It appears that the energy difference of 400 eV up to 800 eV (in the case $E_h \geq 250 \text{ kV/cm}$) supplied by the streamer dynamics is a helpful boost for the low-energy electrons ~ 120 eV to escape the thermal zone to the side of the runaway one. Figure 7 clearly shows how tiny the streamer radius δl is and that the electric field at the same location marked by a green cross can drastically jump from 100 up to 250 kV/cm while the streamer dynamics is taken into account.

The number of thermal runaway electrons is estimated based on the equation (4) (e.g., Babich et al., 2015; Bakhov et al., 2000).

$$N_{\text{run}} = \int_0^{t_{\text{max}}} \int_V v_{\text{run}} n_e dV dt \quad (4)$$

where $v_{\text{run}}(E) = 3.5 \times 10^{-24} \exp(-(2.166 \times 10^{-7} \times E)^2 + 3.77 \times 10^{-6} \times E)$ is the runaway electron frequency calculated by a Monte Carlo simulation (Bakhov et al., 2000).

The integration is made over a field region $E \geq 250 \text{ kV/cm}$ and over an interval of time t_{max} (see Table 2). The time step between data files used to make this estimation is 0.01 ns. Table 4 lists the estimated number of runaway electrons, their average flux, and the number of streamer-patch collisions required to explain the experimental results of Kochkin, Köhn, et al. (2016) for every case considered in the simulation. From the results we estimate an upper limit of the runaway electrons self-current and electric field created at the streamer front. We focus on the case (−45, 1.4) and we find a self-current $I = \frac{e N_{\text{run}}}{t_{\text{max}}} \simeq 0.002 \text{ A}$ and an average self-electric field $E = \frac{e N_{\text{run}} r_{\text{str}}}{(\frac{4}{3} \pi r_{\text{str}}^2) 3 \epsilon_0} = \frac{e N_{\text{run}}}{4 \epsilon_0 S_{\text{str}}} \simeq 6 \text{ kV/cm}$ assuming a spherical geometry. A runaway electron at the front of the streamer creates low-energy electrons along its trajectory. This additional ionization due to runaway electrons combined self consistently with the ionization due to low-energy thermal electrons may increase the speed of the ionization front. The increase of streamer speed by the runaway electrons was already suggested and discussed by Aleksandrov et al. (1995, p.72, and references therein).

4. Discussion

Since the suggestion of the streamer head-on collision as the process behind the production of thermal runaway electrons and X-ray emissions by laboratory spark discharges and by lightning stepped leaders (e.g., Cooray et al., 2009; Kochkin et al., 2012), it became a hot topic within the community of atmospheric and space electricity. It may also be a source of VHF-UHF emissions from lightning (e.g., Montanyà et al., 2015;

Luque, 2017; Shi et al., 2018, 2019). Recently, Babich and Bochkov (2017) made an interesting analysis of Ihaddadene and Celestin (2015) results from the point of view of a possible underestimation effect of the calculated electric field by the applied conducting boundary conditions as suggested by Kochkin, Lehtinen, et al. (2016) and by the used transport coefficients data from Morrow and Lowke (1997). Babich and Bochkov (2017) came up with no effect on the conclusion that conventional streamer head-on collision is unlikely a source of runaway electrons. Kochkin, Lehtinen, et al. (2016), Lehtinen, Kochkin, and Ostgaard (2016), and Lehtinen and Ostgaard (2016) used point-to-point and point-to-plane electrodes discharge 2-D and 1-D models and came to a similar conclusion. Several modeling attempts of a single head-on collision have been conducted (e.g., Babich & Bochkov, 2017; Ihaddadene & Celestin, 2015; Köhn et al., 2017; Lehtinen, Kochkin & Ostgaard 2016; Lehtinen & Ostgaard 2016; Luque, 2017) with a variety of modeling complexities and it appears that the general idea does not seem to be explaining the experimental observations.

The maximum electric field in a head-on collision can be approximated as $E_{\max} \simeq \frac{\left(\int_0^{l^-} (E_0 - E_c^-) dl + \int_0^{l^+} (E_0 - E_c^+) dl\right)}{(\delta R)}$, with $\delta R = \delta R^+ + \delta R^-$, l , and δR are the size of the high field region, the streamer length, and the streamer radius, respectively. The plus and minus signs stand for positive and negative streamers. As the streamers approach each other they squeeze the available potential drop into a short region δR and the electric field is increased. However, the field increase has a physical limit because of the local conductivity and the size of the high field zone. When the Maxwell time τ_m becomes smaller than the ionization time τ_i , then the field starts to collapse (Celestin & Pasko, 2011, see Figure 7). If one assumes $\tau_i = \frac{1}{\nu_i}$ and $\tau_m \sim \frac{\epsilon_0}{q\mu_e n_{e0}} \exp(-\nu_i t)$, one sees that the Maxwell time decreases faster than the ionization time at a certain level of the electron density n_{e0} in between the streamer heads. The photoionization rate also affects the ionization time as long as the distance between the two streamer heads is still larger than the photoionization characteristic length. The field increase is streamer speed dependent, and the faster the collision the faster the field increases and vice versa. At extremely high field values such as 400 kV/cm, to produce a 100-keV electron, the average characteristic distance required is 4 cm, which is higher than a realistic maximum high field zone of 0.4 cm formed by two large streamers at their branching stage (e.g., Celestin & Pasko, 2011). Even with large colliding streamers, 100-keV electrons are unlikely to be produced. Moreover, the ~1-keV electron energy found by Monte Carlo and PIC simulations of streamer head-on collisions (e.g., Ihaddadene & Celestin, 2015; Köhn et al., 2017) requires a channel electric field $E_c = 100$ kV/cm to run away. The runaway electron moves against the approaching high field front and their residence time in the high field zone is always shorter than the maximum field duration t_{\max} .

The multistreamer head-on collisions problem is beyond the single-streamer head-on collision complexity. It is a combination of high number of streamers propagating in an electromagnetic, kinetic, and hydrodynamic environment that we are not able to simulate with the current available models if simplifications are not made. That is the reason why in the present work we use a simplified multiple ellipsoids approach to test how likely the head-on collision is a source of runaway electrons. Indeed, the dynamics of the electric field in a meter-scale discharge is very complex, and our model does not capture the details of streamer collective effects. However, from the present results we find two issues that prevent the electrons from running away when short and large dielectric ellipsoids representative of streamers are used. The short ellipsoids possess a limited potential drop and electrons do not gain enough energy to run away, and the long ellipsoids possess a large potential drop but then electrons fall into the long channel's low electric field and they lose the energy they gained at the collision point. Based on Table 1, we would require a single head-on collision to produce electrons $\epsilon_{\text{run}} > 60$ keV at the collision point to allow them to runaway in $E_c = 10$ kV/cm positive streamer channel, if one assumes an average discharge ambient field of $\simeq 20$ kV/cm. From Table 1, one sees that the low electric field requires a high-energy ϵ_{run} for an electron to run away. A self-consistent particle model taking into account the ionization from runaway electrons (Köhn et al., 2017, see Figure 4) found one electron with $\epsilon \sim 500$ eV that is unlikely to runaway in ~ 20 kV/cm positive streamer channel's field, which requires $\epsilon_{\text{run}} > 40$ keV. Unlike a short gap simulation of Köhn et al. (2017), Luque (2017) conducted a long plasma fluid streamer simulation over ~ 30 -cm gap to achieve a large potential drop, and their one-dimensional deterministic acceleration produced an electron with $\epsilon > 5$ keV that is also not going to runaway in $E_c \sim 12$ kV/cm, which requires $\epsilon_{\text{run}} > 60$ keV. We emphasize that there is not only the parameter of the available potential drop to take into account at the head-on collision location for the production of runaway electrons, but others, such as the size of the high field zone, the duration of the peak field, and the energy of the runaway electrons, when they come out from the collision zone to flow through the channel,

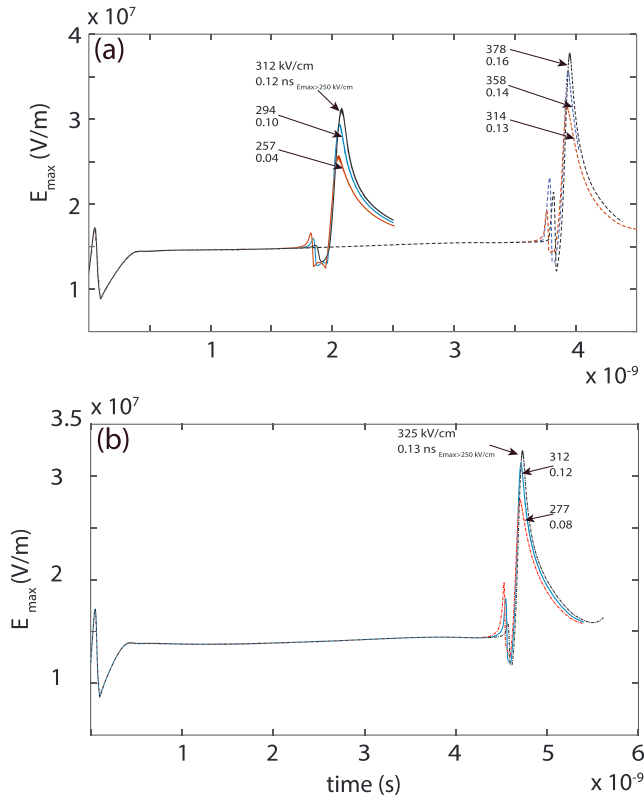


Figure 9. (a) Evolution of the maximum electric field E_{\max} as function of time. Black, dashed black, blue, dashed blue, and red, dashed red, curves correspond to patch sizes σ_r (μm) = 15, 40, 80, and E_0, d_{patch} = -45 kV/cm, 0.5 cm; -45, 1.4, respectively. (b) Evolution of the maximum electric field E_{\max} as function of time. Dashed black, blue, and dashed red, curves correspond to patch sizes σ_r (μm) = 15, 40, 80, and E_0, d_{patch} = -40 kV/cm, 1.4 cm, respectively.

are also important. Streamer head-on collisions are believed to happen in the middle of the discharge gap during the encounter of long negative and positive streamers (e.g., Kochkin et al., 2012). However, there is no strong experimental evidence for streamer collisions to occur and to produce X-rays (e.g., da Silva et al., 2017). The detector response time is much longer than the picosecond interaction times of a single-streamer head-on collision, so the detector cannot pinpoint the source. At the present time only a statistical correlation between the X-ray emissions and the encounter of long streamers has been reported as a possible evidence (e.g., Kochkin et al., 2012).

Figure 1 shows a complex morphology of the streamer zone near the point electrode when X-ray emissions are produced. The figure clearly shows different streamer channel sizes, and small aligned isolated dots probably formed after the decay of the streamer channel. Some of the dots act as inhomogeneities initiating upward propagating positive streamers. From the results presented in this work, we believe that a collision between a large negative streamer and a plasma patch similar to those in Figure 1 is a better candidate for explaining the production of runaway electrons and the associated X-ray emissions. The plasma patches used in this study have a size $\sigma_z = 200$ and $\sigma_r = 15$ μm . We may ask the question why the patch is so thin radially and also whether a numerical resolution of 8 μm could resolve accurately such a thin patch. We have run a numerical test with $\sigma_r \sim 5 \times \Delta_r$ ($\Delta_r = 3$ μm), and the results do not show a change that would affect the principal conclusions of this work. The maximum field is increased by 1% at 3- μm resolution with respect to the 8- μm resolution case. In addition to that, we have also conducted additional simulations with different patch sizes $\sigma_r = 40$ and 80 μm (same $\sigma_z = 200$ μm) at locations $z = 0.5$ and 1.4 cm in -45 and 40 kV/cm ambient electric fields (sphere-to-plane electrode configuration), and the results are shown in Figure 9. The results clearly show strong electric fields similar to the values obtained using a patch size $\sigma_r = 15$ μm that are sufficient to generate thermal runaway electrons. One sees that the maximum electric field is decreased by ~ 15 –20% when increasing the patch radial size from 15 to

80 μm . The origin and the upper and lower size limits of the patch are both uncertain. We believe that the patch might be produced either by a decay of former streamer channels created when streamers propagate into a low ambient field volume (e.g., ~ 95 μm radius for a thinnest streamer (Briels et al., 2006)) or by a development of a group of polarized electron avalanches/pilots that act as initiation points for positive and negative streamers (e.g., Reess et al., 1995, section 3). As far as we know, the characteristic size that corresponds to an avalanche to streamer transition for atmospheric air is about 30–50 $\mu\text{m}/2$ to 15–25 μm (e.g., Naidis, 1996; Qin, 2013, p.28–31).

We choose a sharper ellipsoidal shape for the plasma patch because it produces a stronger peak electric field than a spherical shape. The analysis by Østgaard et al. (2016) showed that runaway electrons are likely to be produced at tens of centimeters near the electrode zone where the electric field is believed to be $\geq E_k$. From the Figure 2 of Kochkin et al. (2012) one can estimate a 50 kV/cm to span over 16 cm at ~ 0.8 MV applied voltage when X-rays are produced. Those studies support our selection for the ambient electric fields used in the present paper. We conducted simulations in a low ambient electric field ~ 20 –25 kV/cm, in a ~ 2 -cm gap, and we obtained with a ~ 12 -kV potential drop a peak electric field up to ~ 180 kV/cm, which is below the 250 kV/cm required for the production of thermal runaway electrons. The low ambient field cases ($< E_k$) require longer streamers to achieve larger potential drops and higher runaway threshold energy ($\epsilon > 40$ keV, see Table 1). For instance, a potential drop $U = 40$ –100 kV requires a streamer length $l \approx \frac{U}{(E_0 - E_c)} \approx 5$ –10 cm that is beyond our simulation capabilities. Furthermore, according to Celestine and Pasko (2011, section 2) when the field in the streamer channel approaches the stability field (10–15 kV/cm) in the case of a negative streamer, the streamer head moves with a constant potential drop and does not expand. In addition, the streamer speed is lower in ambient fields $< E_k$, which results in a slower and softer collision with a weaker

peak electric field. In low ambient electric fields a long streamer develops a large radius to host a large potential drop and keep a stable head electric field as it propagates (see properties of stable streamers; e.g., Qin & Pasko, 2014). As the electric field at the front of the streamer head follows $\sim \frac{1}{r_s}$ profile, the plasma patch is expected to polarize and initiate streamers in a quasi-homogeneous streamer electric field before the full collision occurs, which will prevent the short streamer from reaching high peak electric fields. The streamer residence time in the plasma patch $\frac{l_{\text{patch}}}{V_{\text{str}}}$ must be short enough to develop a fast ionization growth to increase the peak electric field in the patch. The thermal runaway electrons are easy to produce in this fashion in ambient electric fields $\geq E_k$ where streamers expand and carry large potential differences. The low ambient field case deserves to be explored in future studies.

Based on the results of Kochkin et al. (2014) X-rays are not produced at the first burst of the streamer but at the fourth burst. This may indicate that a specific set of conditions needs to be filled. The first burst of streamers propagates short distances with low speed and dies off (e.g., Morrow & Lowke, 1997; Qin & Pasko, 2014) as the ambient field far away from the electrode approaches ~ 0 kV/cm. The first burst of streamers creates a preionized medium for the next coming bursts. Both voltage and current increase continuously while the previous streamer channels decay over the attachment time scale. If one assumes a streamer channel electron density 10^{20} m^{-3} , a streamer channel field of 25 kV/cm, and an ambient field 50 kV/cm, it takes 50 ns for a streamer to travel ~ 15 cm (3×10^6 m/s) and for the density to go down to 10^{18} m^{-3} and to form plasma inhomogeneities. Within the decay process the field rises again with the voltage, and if a collision occurs, it is going to be the adequate moment for the X-ray emissions to be produced. We do not expect the X-ray emissions to be produced within the first burst of streamers because the plasma patches do not exist yet. The frequency of streamer-patch collisions determines the number of both the runaway electrons and the associated X-ray photons. From Table 4, one collision in a case (-45 kV/cm, 1.4 cm) is sufficient to produce the 10^6 runaway electrons with ~ 100 keV per one electron to explain the 10^4 X-ray photons based on 100 electrons per a single X-ray photon (Kochkin, Köhn, et al., 2016). The lower ambient field demands a higher number of collisions. Tens of collisions is a reasonable number to consider. However, it is hard to imagine a higher number because of the many parameters that control the streamer-patch mechanism to occur such as the streamer length, the speed, the homogeneous or inhomogeneous ambient field, the streamers density per burst, the characteristic decay time of the streamer channel, and the shape of the plasma patch, make it unlikely for all the collisions to produce runaway electrons and the associated X-rays.

The estimation of the energy and the number of thermal runaway electrons used in this work is an acceptable approximation to quantify the two quantities and computationally not expensive. We realize that the estimation is sensitive to the initial location of the test electron and the local high electric fields. We position the electron in the back ($E > 250$ kV/cm) of the peak streamer field (378 kV/cm) in the case (-45 kV/cm, 1.4 cm), and we find that the electron can gain without the streamer dynamics up to 20 keV. We do not expect a drastic change in the energy of the electron using a Monte Carlo approach as we already explained in section 3.3. The energies considered in the current study are at 4% maximum of the relativistic electron mass energy (~ 500 keV). From Table 1, the Monte Carlo simulation requires $\epsilon_{\text{run}} = 15$ keV in $E_0 = 50$ kV/cm and if one applies up to 30% less energy gain to the estimated value $\epsilon = 24$ keV due to Monte Carlo considerations one finds ~ 17 keV $> \epsilon_{\text{run}} = 15$ keV. The Monte Carlo model would give a better distribution of electrons for a broad range of energies as it has the ability to deal with a high number of electrons instead of a one single test electron method. However, the separate treatment of runaway electrons and streamer simulations makes the estimation not self consistent as it is in a particle code (e.g., Chanrion & Neubert, 2010; Köhn et al., 2017). In case of high fields > 250 – 300 kV/cm the drift-diffusion local field approximation is no longer valid and the relation $v_e(E) = \mu_e(E)E$ is violated as electrons are no longer in equilibrium in the local electric field; hence, the electron transport coefficient and the electron distribution are not well resolved. In addition, the amount of ionization due to the runaway electrons affects the streamer local electric field (Köhn et al., 2017, see Figure 2), the collision speed, the duration of the peak field, and the rate of runaway electrons with the associated X-ray photons. Li et al. (2007, 2008) set a limit between the self consistent particle and fluid models at 200 kV/cm. At that field value the particle code describes better the streamer ionization front with a 60% higher ionization rate than the fluid models due to electrons > 15 and < 120 eV that are not runaway and create more ionization. The nonlocal approximation of Naidis (1997) may be an alternative solution in between the local approximation and particle approach to describe the electron distribution in this case. We also point that PIC techniques such as resampling or superparticle treatment of

high number of electrons beyond the computation capabilities produce numerical enhanced ionization and branching effects (Chanrion & Neubert, 2010, paragraph 36).

How the runaway electrons affect the streamer properties at high-energy regimes is still not fully understood. The runaway electrons may create electron avalanches at the front of the negative streamer head (e.g., Chanrion & Neubert, 2008; Köhn, Chanrion, Babich, et al., 2018) and affect its velocity. Advanced simulations are required to investigate these effects and understand the mechanisms of the runaway electrons and X-rays in laboratory discharges and lightning leaders. The ellipsoid model used in this study may not be an accurate description of the complex geometry of the electric field produced by a large number of streamers. The estimation of the energy and the number of runaway electrons produced by the streamer-patch mechanism is one dimensional and not self-consistent. A PIC simulation of multiple-streamer (>2) head-on collisions to test the possibility of a continuous accelerations of thermal runaway electrons by other neighboring streamers and/or a reproduction of one of the streamer-patch configurations would be an interesting future work to do and compare with the results of this paper.

5. Conclusions

The main conclusions of this work can be summarized as follows:

1. We have combined dielectric ellipsoids with Monte Carlo simulations to study the conventional multiple-streamer head-on collisions and we have found that the runaway electrons lose their total energy in the positive streamer channel's low electric field. The positive streamer channel is an additional factor of energy loss for the electrons to overcome if they are produced by a streamer head-on collision.
2. We have simulated a collision between a large negative streamer and a small neutral plasma patch, and we have shown it to produce thermal runaway electrons in ambient electric fields $> E_k$.
3. We have found a number of runaway electrons in the streamer-patch simulations between 10^2 and 10^6 in ambient electric fields $|E_0| = 35$ and 45 kV/cm, with an average energy per electron up to 24 keV.
4. We have found that the streamer self-acceleration contributes to a runaway acceleration process.

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Acknowledgments

This material is based upon work supported by the Air Force Office of Scientific Research under Award FA9550-16-1-0396. Sebastien Celestin's research is supported by the French space agency CNES in the framework of the satellite mission TARANIS. This work complies with the AGU data policy. All data used in this paper are available at <https://zenodo.org/DOI:10.5281/zenodo.2362416/zenodo.2565279/zenodo.3239526>. The authors thank L. P. Babich and the anonymous reviewers for their comments that improved this article.

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