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# Xenon and iodine behaviour in magmas

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## Abstract

Iodine (I) and xenon (Xe) are two key elements that trace Earth's differentiation (e.g. atmosphere formation) and dynamics (e.g. volcanism and recycling at subduction zones). Iodine and Xe abundances are linked through the decay of the extinct  $^{129}\text{I}$  that produced  $^{129}\text{Xe}$ , which is today depleted in the Earth's atmosphere compared to the composition of the solar system (i.e. chondrites). Iodine and Xe cycles and storage in the deep Earth are almost unknown, which is in large part due to the fact that their behaviour in magmas and fluids, key agents of mass transfer through planetary envelopes, are poorly known. Here, the solubility of Xe and I in melts is measured under high pressure ( $P$ ) and temperature ( $T$ ) conditions using large volume presses, and Xe and I behaviour in melts and fluids is monitored in situ under high  $P$ - $T$  conditions using resistive heating diamond anvil cells combined with synchrotron x-ray fluorescence (XRF) and Raman spectroscopy. Xenon, I and H ( $\text{H}_2\text{O}$ ) contents were measured in quenched glasses by particle x-ray Emission (PIXE) and Elastic Recoil Detection Analysis (ERDA). Solubility, speciation and degassing processes are investigated for two different compositions: haplogranitic melt (HPG analogue for crustal melts) and basaltic melts (MORB and IAB). Experimentally measured solubilities for both elements are much higher than their natural abundances in terrestrial magmas. Xenon solubility at 3.5 GPa reaches 4.00 wt.% in HPG and 0.40 wt.% in basalts. Iodine solubility is 0.46 wt.% at 0.4 GPa on average in HPG, and reaches 1.42 wt.% in basalts at 2 GPa. The *in situ* Raman spectroscopic study shows that I forms I-I bonds in hydrous high P fluids/melts unlike Xe that was previously shown to oxidize in high P melts. The XRF monitoring of I and Xe partitioning

between aqueous fluids and silicate melts during decompression (i.e. water degassing) shows that Xe degassing is strongly  $P$ - $T$  dependent and can be retained in the melt at deep crust conditions,, while I is totally washed out from the silicate melt by the aqueous phase. Xenon and I degassing processes are based on different mechanisms, which implies that the atmospheric isotopic signature of Xe cannot be inherited from a process involving volcanic water degassing. Instead,  $^{129}\text{Xe}$  depletion may originate from a separation of both elements at depth, by deep fluids, a proposition that agrees with a deep storage of Xe in minerals.

**Keywords** iodine, xenon, silicate melt, solubility, speciation, volcanic degassing

## 1. Introduction

Iodine (I) and xenon (Xe) are highly volatile elements, differing only by one proton, but **due to their difference in electron configuration**, I is a very reactive halogen element while Xe is a noble gas. They share a very low abundance compared to chondrites in the Bulk Silicate Earth (BSE) and low temperatures of condensation (see Marty, 2012). Both I and Xe are potentially “missing” in the BSE compared to solar abundances. Iodine and Xe elements are neighbours in the periodic table, and are linked by the radioactive decay of the today extinct  $^{129}\text{I}$  to  $^{129}\text{Xe}$  (half life 16.1 million years).

Yet unsolved issues are: (1) why is the Earth’s atmosphere depleted in elemental Xe and  $^{129}\text{Xe}$ ? (2) how are I and Xe recycled back to the mantle at subduction zones? (3) whether or not I and Xe are stored in the deep Earth, and if so, by which processes and to what extent?

Iodine is the heaviest stable halogen element, having the largest ionic radius, and for this reason it is much less frequently incorporated in minerals than chlorine (Cl) and bromine (Br) (McDonough and Sun, 1995), whereas they shares many other characteristics. Iodine is mostly found at the ppb level in magmatic rocks and volcanic glasses (Muramatsu and Wedepohl, 1998; Schilling et al., 1980; Déruelle et al., 1992; Kendrick et al., 2012). It is more concentrated in sediments than in any other geologic materials (Muramatsu and Wedepohl, 1998). Being hydrophilic and biophilic, it is enriched in marine organic matter and associated fluids (Fehn 2012).  $^{129}\text{I}$  is used to trace crustal

79 and volcanic fluids in subduction contexts (Fehn, 2012), where serpentinites are efficient traps for  
80 subducted I (John et al., 2011; Kendrick et al., 2012). Experimental studies have demonstrated that  
81 heavy halogens Br and I may significantly degas from hydrous silicic magmas (Bureau et al., 2000;  
82 2010; 2016). The uncertainty on I outgassing fluxes is due to the lack of constraints on: (1) the real  
83 I content of magmas; (2) the storage capacity of magmas with respect to I, and more specifically I  
84 solubility in melts relevant for subduction-related volcanism; (3) the real abundance of I in the bulk  
85 silicate Earth, (4) how much of subducted I is deeply recycled. Mass balance calculations based on  
86 the I content of MORBs (Mid Ocean Ridge Basalts) have been used to predict the total amount of  
87 I in the bulk silicate Earth (McDonough and Sun, 1995), resulting in a strong depletion of the BSE  
88 compared to solar compositions (i.e. CI chondrites). Two scenarios are proposed: (1) I loss by a  
89 strong degassing event of the magma ocean during the Hadean eon (Bureau et al., 2016); (2) I  
90 storage in the core (Armytage et al., 2013) or in a hidden deep mantle reservoir(s). Alternatively,  
91 there could be no I depletion as indicated by new measurements of halogen contents in chondrites  
92 (Clay et al. 2018).

93 As most noble gases, Xe is a tracer of Earth's dynamics. Xenon isotopes are important markers of  
94 planetary processes, and the I-Pu-Xe system is used to date the formation of the atmosphere (Ozima  
95 and Podosek, 1999, Pepin 2006, Avice and Marty 2014). About 90% of Earth's primordial Xe is  
96 missing from the atmosphere of the Earth (Anders and Owen 1977, Ozima and Podosek, 1999),  
97 this is the "missing Xe" issue. Xenon is depleted by one order of magnitude relative to other volatile  
98 elements (water, carbon, neon, argon and krypton) when normalized to the chondritic composition,  
99 and atmospheric Xe is enriched in heavy isotopes by 3–4%  $\text{amu}^{-1}$  (unified atomic mass unit)  
100 relative to chondritic and solar compositions, this is the "Xe paradox" (Krummenacher et al. 1962).  
101 The Earth is also depleted in radiogenic Xe isotopes, i.e. I-derived  $^{129}\text{Xe}$  and Pu- and U-derived

fissiogenic isotopes (Tolstikhin et al., 2014). Recent measurements on deep fluids and fluid inclusions trapped in Archean crustal minerals show that Xe isotopic fractionation was still occurring throughout the Archean era (Pujol et al. 2011, Avice et al. 2017). Modelling Xe degassing and re-gassing over geological times reveals that Xe recycling has been significant after 2.5-2.8 Ga (Holland and Ballentine 2006, Avice et al., 2017; Parai and Mukhopadhyay, 2018).

Xe escape from the atmosphere could explain this evolution (Pepin, 2006; Hebrard and Marty 2014, Zahnle et al., 2019). That requires either some degree of storage of Xe at depth (Pepin 2006) or the photoionization of Xe (Hebrard and Marty 2014, Zahnle 2019) providing conditions: the presence of an atmospheric organic haze and strongly nonlinear photoionization process to compensate for the weak UV flux in the Archean (Hebrard and Marty 2014), or an unlikely large amount of H<sub>2</sub>O reduction by FeO during the Archean to provide sufficient H<sup>+</sup> escape flux or yet unexplained Xe escape bursts (Zahnle et al. 2019).

Alternatively, the missing Xe could be stored at depth, in the deep crust or mantle either in high *P* melts (Leroy et al., 2016) and/or minerals as shown for quartz (Sanloup et al., 2005, Crepisson et al. 2019) and olivine (i.e. Sanloup et al., 2011; Crepisson et al., 2018), or in the core (Lee and Steinle-Neumann, 2006, Zhu et al. 2014).

In the case of Mars' atmosphere, in which Xe is also depleted both in bulk element content and in light isotopes, it was proposed that I separated from Xe through a process involving water (Musselwhite and Drake, 1991) as on a chemical point of view I is hydrophilic whereas Xe is hydrophobic. Liquid water, supposed to be present at that time, would have trapped I (including <sup>129</sup>I later transformed in <sup>129</sup>Xe) while Xe was lost to space. This is the model proposed to explain that for Mars (<sup>129</sup>Xe/ <sup>132</sup>Xe)<sub>Mars'satmosphere</sub> > (<sup>129</sup>Xe/ <sup>132</sup>Xe)<sub>Marsmantle</sub> whereas this is the opposite for

the Earth:  $(^{129}\text{Xe}/^{132}\text{Xe})_{\text{Earth's atmosphere}} < (^{129}\text{Xe}/^{132}\text{Xe})_{\text{MORB}}$ . Iodine and Xe should exhibit different behaviours during an early major water degassing event that may have occurred during the Hadean magma ocean era.  $^{129}\text{I}$  became extinct about 100 million years after the formation of the solar system. If I was degassed from the magma ocean and lost from the atmosphere together with water (Bureau et al., 2016), but not Xe, such a process would have affected the atmosphere only, and may have contributed to the isotopic atmospheric depletion in  $^{129}\text{Xe}$  observed today. If Xe was lost together with I, magma degassing is not the process to explain its atmospheric isotopic depletion.

We test the effect of an early water degassing process of the magma ocean on a potential separation of I from Xe. For this purpose, we developed a multi-diagnostic experimental approach to measure the solubility of both Xe and I in melts using large volume presses, and monitor Xe and I partitioning between hydrous melts and aqueous fluids in situ at high pressure ( $P$ ) and temperatures ( $T$ ) during decompression (i.e. degassing of an ascending magma) combining diamond anvil cells experiments with synchrotron x-ray fluorescence XRF and, with Raman micro-spectroscopy devoted to measure I speciation in high  $P$  melts/fluids. We propose that Xe and I could be stored at depth, possibly in deep magma or melt layers and their crystallisation products.

## 2. Methods

### 1. Synthesis of the starting materials



Iodine and Xe bearing glasses are synthesized to obtain starting materials for diamond anvil cells experiments, and determine solubilities in high P silicate melts.

We use a natural Island Arc Basalt (IAB, basaltic sample STV301, from St Vincent volcano, Pichavant et al., 2002), relevant for subduction zones context; a natural Mid Ocean Ridge Basalt (MORB, BN06-11 from Juan de Fuca Ridge, Prouteau and Scaillet, 2013), and a synthetic haplogranite (HPG, Bureau and Keppler, 1999), a composition relevant for continental crust melts (e.g. Schmidt et al., 1997), also used as a proxy for subduction zones magmas (e.g. Bureau et al., 2003, 2010; 2016).

The IAB and MORB powders are obtained by crushing the natural lava samples. The HPG glass is prepared from a mixture of pure oxides and carbonate powders ( $\text{SiO}_2$ ,  $\text{Al}_2\text{O}_3$ ,  $\text{Na}_2\text{CO}_3$ ,  $\text{K}_2\text{CO}_3$ ). The HPG powder is dehydrated and degassed from  $\text{CO}_2$  by a slow ramp heating from room temperature to  $1200^\circ\text{C}$  during 8 hours, it is melted in an atmospheric furnace for one hour at  $1600^\circ\text{C}$ , quenched in water, crushed and melted again and crushed to a fine homogeneous powder.

### *Iodine bearing glasses*

To limit iron loss during experiments some platinum (Pt) capsules are doped in iron before the high *P-T* experiments. The IAB powder is loaded in open Pt capsules and melted in an atmospheric furnace at the same T and for the same duration as for the high P experiments. Capsules are then cleaned in hydrofluoric acid before being loaded and sealed. HPG and basalt powders together with NaI salt (3wt.%) are loaded in Pt capsules without or with distilled water. Pt capsules are welded and inserted either (1) in a 3.2 cm long (1/2 inch) cell-assembly composed of a graphite heater surrounded by pyrex and talc sleeves for piston cylinder experiment; or (2) in an internally heated

pressure vessel pressurized with an Ar-H<sub>2</sub> gas mixture (ISTO, Orléans). High  $P$ - $T$  conditions of 2-3.5 GPa and 1400-1600°C are generated using the piston cylinder apparatus for 60 minutes; low  $P$ - $T$  conditions of 0.2-0.4 GPa and 1300°C are generated in the internally heated pressure vessel for 240 minutes.  $T$  is monitored with a Type C thermocouple. The samples are quenched as glasses either by turning-off the power supply for the piston cylinder or using a fast quench device for the internally-heated pressure vessel.

### *Xenon bearing glasses*

For glasses not doped in I, hydration is done at 2 GPa and 1400°C using a piston–cylinder apparatus with a half-inch talc pyrex assembly, a graphite heater, and dried MgO powder packed around the Pt capsule. Two hydration levels are targeted, 5 wt% and 7 wt%. The glass is sealed with the proper amount of distilled water in a Pt capsule. Xe doping is done subsequently by gas-loading in another Pt capsule partially filled with the recovered hydrous HPG glasses, and brought to 3.5 GPa and 1600°C for one hour, followed by rapid  $T$  quenching and decompression at room  $T$ . For I-Xe bearing glasses the I-glasses are Xe-doped following the same procedure. Experimental conditions are summarized in **Table 1**.

### *2.2 Diamond anvil cells experiments*

The degassing process (*i.e.* the separation of a fluid phase from a silicate melt) is monitored *in situ*, through decompression induced by  $T$  decrease in Bassett-modified externally heated diamond anvil cells (DAC) equipped with  $P$ -driving membranes (Munsch et al., 2015). The partitioning of Xe and I between the two phases is measured *in situ* by using synchrotron XRF analysis during decompression following the same strategy as in previous studies (Bureau et al., 2010; 2016). The

DACs are equipped with large culet size (0.8 mm to 1 mm) and 2 mm thick diamonds, and externally heated by furnaces attached to both anvils.  $T$  is controlled by two K-type thermocouples touching each diamond anvil, and previously calibrated. The sample chamber is a 300  $\mu\text{m}$  to 500  $\mu\text{m}$  hole drilled in a Re gasket (initial thickness: 250  $\mu\text{m}$ ) compressed between the two diamonds. Observations are done along the  $P$  axis through the diamond anvils, and the loaded aqueous fluid is used as  $P$ -medium. Density of the aqueous fluid is calculated by measuring the homogenization  $T$  of a water vapor bubble in liquid water. Knowing the density,  $P$  can be calculated for a given  $T$  from the equation of state of the fluid phase (see Munsch et al., 2015). We obtain a  $P$ -range of 1.21 GPa - >1.68 GPa with a precision of 0.2 GPa for a  $T$ -range of 684°C - >851°C.

### 2.3 *In situ* SXRF

Synchrotron XRF quantitative analysis is performed at the beamline P02.2, Petra III (DESY, Hamburg). We use a 42.7 keV monochromatic x-ray beam focused down to 4  $\mu\text{m}$  x 6  $\mu\text{m}$ . For each experiment, the DAC is positioned perpendicularly to the beam, XRF spectra are detected in transmission geometry with an angle of 8° with the incident beam, with a Silicon Drift x-ray detector. Iodine and Xe can be detected *in situ* in a DAC as their high  $K\alpha$  and  $K\beta$  Rays (*i.e.* I:  $K\alpha_{\text{I}}$  28.61 keV;  $K\beta_{\text{I}}$  32.99 keV, Xe:  $K\alpha_{\text{Xe}}$  29.78 keV;  $K\beta_{\text{Xe}}$  33.62 keV) are not absorbed by the 2 mm thick diamond anvils. We use the I-Xe doped HPG glasses as magma analogues (HPGXe-C11, HPGXeI-04). For the HPG composition, the glass transition occurs at lower  $T$  compared to basaltic or peridotitic melts, which makes possible a study in resistive heating DAC whose  $T$  are limited to 1000°C in routine mode (max 1100°C). For each experiment, the sample chamber is half-filled with pieces of glass, and further loaded with pure water. When possible,  $T$  is increased up to total

miscibility, and then decreased by steps causing decompression; at each  $T$  step, co-existing aqueous fluid and silicate melt are analyzed (Bureau et al., 2010; 2016). At the very last stages of decompression, water bubbles nucleate in the melt and form fluid inclusions until the glass transition  $T$  is reached, compromising the analysis of the quenched glasses *in situ* in the DAC; the quenched aqueous fluid is also characterized. XRF analyses are calibrated *in situ* in the DAC at room conditions using a NIST 610 glass (e.g. Rocholl et al., 1997). For XRF analysis it is not necessary to use I- and Xe-enriched glasses as standards, because the whole NIST 610 spectrum is calibrated for multi elemental analysis using glasses containing elements with  $K\alpha$  ray energies close to that of I and Xe , such as Sb  $K\alpha$  = 25.2 keV,  $405 \pm 10$  ppm in NIST610 (Rocholl et al., 1997).

#### 2.4 *In situ* Raman spectroscopy

Raman spectra are collected *in situ* in a DAC on a Jobin Yvon® HR-460 spectrometer equipped with an Andor® CCD camera at IMPMC using the 1500 grooves/mm grating mode. Raman signal is excited using the 514.5 nm wavelength of an  $Ar^+$  laser, delivering 300 mW focused into a 2  $\mu$ m spot by a long-working distance Mitutoyo® x20 objective and collected in back-scattering geometry. Backscattered Rayleigh light is filtered with a set of holographic Notch filters. Spectral resolution is about  $2.5\text{ cm}^{-1}$ . *In situ* measurement collection times are 60 s or 120 s per spectra in order to get reasonable signal-to-noise ratios. The DAC is loaded with pieces of HPG glass and a solution of  $H_2O$ -NaI (20 g/l), heated up to total miscibility between the aqueous fluid and silicate melt, and the first Raman measurements are performed in the supercritical fluid at 860°C.  $T$  is then

decreased, and Raman spectra are acquired after 10 min of equilibration in co-existing aqueous fluid and silicate melt.

## *2.5 Analyses of quenched samples*

Quenched glasses from large volume press and DAC experiments are recovered, mirror polished (except for DAC samples) for *ex situ* scanning electron microscopy, electron microprobe and ion beam analyses. We use a Zeiss Ultra 55 field emission scanning electron microscope (IMPMC) with a working distance between 3 mm and 7.5 mm, and a voltage of 15 kV to observe crystal-free glasses texture. Major elements are analyzed using a Cameca SX FIVE electron microprobe at CAMPARIS facility (Sorbonne Université). Acceleration voltage is set to 15 kV, current to 4 nA and beam is focused to 15  $\mu\text{m}$  diameter at the surface of the sample. Counting times are 10 s per element on both peak and background. We use two internal standards (EtC, Bureau et al., 2000 ; KE12, Bureau and Metrich, 2003). About 10 to 30 analyses are performed for each sample. Major elements compositions for recovered high P-T samples and starting glasses are listed in **Table 2**.

Particle induced x-ray emission (PIXE), Rutherford backscattering spectrometry (RBS) and Elastic Recoil Detection Analysis (ERDA) are performed at the nuclear microprobe, CEA Saclay NIMBE / LEEL. Samples are analyzed by PIXE for I and Xe. The  $3 \times 3 \mu\text{m}^2$  proton beam of 2 MeV, 600 pA is scanned on large areas ( $250 \times 250 \mu\text{m}^2$  to  $50 \times 50 \mu\text{m}^2$ ) during one to a few hours, corresponding to deposited charges ranging from 0.6  $\mu\text{C}$  to 2.4  $\mu\text{C}$ . Emitted x-rays are collected using a Silicon-drift x-ray detector. Simultaneous RBS analysis allows to check the major element compositions of the glasses, and to map fluid inclusions trapped in the glass. Elemental maps are processed in order to select the areas of interest (e.g. bubble-free glass areas). Concentrations are obtained as

described in Bureau et al. (2010, 2016). Data analysis is calibrated with the NIST 610 Glass, including the calibration of the PIXE solid angle done on Sb L $\alpha$  line at 3.44 keV (Sb content = 405  $\pm$  10 ppm). The quantification process is cross-checked using the well-known reference glasses EtC and KE12. To quantify I content of HPG glasses, we use the L $\alpha$  (3.92 KeV) and L $\beta$  rays (4.22 KeV), both rays giving similar concentrations (**Table 2**). To quantify the I content of basaltic glasses, we use only the L $\beta$  lines due to interferences between the I L $\alpha$  line at 3.94 keV and the calcium Ca-K $\beta$  line at 3.69 keV (**Figure 1**). To quantify Xe in all glasses, we use the L $\alpha$  line at 4.1 keV.

Water contents are obtained by ERDA, a technique based on the detection of protons after their ejection from the samples through elastic collisions at a grazing angle of 15° induced by a  $^4\text{He}$  4x16  $\mu\text{m}^2$  incident beam of 3 MeV. RBS measurements are performed using an annular detector positioned at 170° with respect to the incident beam direction in order to monitor the charge delivered to the sample, necessary to allow hydrogen quantification. Details about the procedure can be found in Bureau et al (2009).

### **3. Results**

#### *3.1 I-Xe doped Glasses*

All glassy samples are chemically homogeneous with respect to major elements, water, and I thus attesting that chemical equilibrium was reached. For the samples of HPG composition synthesized at 0.4 GPa, we observe large I-rich bubbles, that are excluded from the maps for PIXE data processing. The presence of I-rich water bubbles may be due to the fact that at this  $P$  condition, water solubility in HPG melts is lower than the initial amount of water loaded in the capsule ( $> 5$  wt.% at 0.2 GPa, 800°C, Holtz et al., 1992), water and I degas simultaneously during decompression (Bureau et al., 2016). Bubbles are not observed in samples of IAB composition as we use initial water contents ( $< 2$  wt.%) lower than the solubility of water in basalts at high  $P$  (up to 8 wt%  $\text{H}_2\text{O}$  at 0.4 GPa, Sisson and Grove 1993). Xe-doped glasses also exhibit large bubbles due to Xe oversaturation from the gas loading system.

Iodine contents range from  $0.46 \pm 0.07$  wt.% I to  $1.33 \pm 0.15$  wt.% I (in the presence of Xe) in both hydrous HPG, and from  $0.48 \pm 0.07$  wt.% to  $2.97 \pm 0.58$  wt.% I for IAB. At 3.5 GPa (see below), there is no excess I, added at the level of 2.54 wt.% I (3 wt.% NaI) in the starting material. This suggests that the solubility at 3.5 GPa may exceed 2.97 wt.%. Solubility expresses the ability of a solute to dissolve in a solvent; it is given by the maximum concentration of the solute in a solvent at equilibrium. Therefore we measured I solubility only for experiments performed up to 2 GPa.

These concentrations are much higher than the I content measured in basaltic volcanic glasses, usually a few ppb to a few hundreds of ppb (i.e. Schilling et al., 1980; Déruelle et al., 1992; Kendrick et al., 2012). We observe a positive trend between I content and  $P$  for the basaltic composition (**Figure 2A**), similar to that observed for the solubility of chlorine (Webster, 1997) and bromine (Bureau and Métrich, 2003) in hydrous melts in the absence of aqueous brines and/or

hydrosaline fluids (see Webster and Mandeville 2007). For  $P$  up to 2 GPa, I solubility values in silicate melts are comparable to those of chlorine and bromine (**Figure 2C**). Because the amount of I dissolved in the basaltic melt at 3.5 GPa (i.e.  $2.97 \pm 0.58$  wt.% and  $2.67 \pm 0.58$  wt.%) is equal to the initial amount loaded in the starting sample (about 2.54 wt.%) the melts may not be saturated with respect to I at these high  $P$ , which means that I solubility is probably higher than 2.97 wt.% for  $P > 3.5$  GPa. The Fe content of the samples is variable due to Fe loss to the Pt capsules (1.05 wt.% to 8.51 wt.% FeO). Even experiments performed at high  $P$  (2 GPa and 3.5 GPa) and  $T$  (1400-1600°C) with Fe-doped Pt capsules do not prevent partial Fe loss from the silicate melt (see samples BA-PC88 and BA-PC89, respectively at 1.75 wt.% and 5.77 wt.% FeO while sample BA-C7 synthesized at 0.2 GPa-1300°C exhibits the right FeO content at 8.51 wt.%). Because the initial water content is not equivalent for all samples we cannot discuss potential relationship between I and water contents.

The Xe solubility value at 3.5 GPa is much higher for HPG than for basalts: 4 wt.% for HPG falling down to 0.4 wt.% for IAB and to 0.27 wt.% for MORB (**Figure 2B**), in agreement with data for tholeiitic melts (Schmidt and Keppler, 2002).

Two experiments are performed for basaltic compositions doped in both Xe and I at 3.5 GPa (**Table 2**), i.e. IAB (BA-PC73 1.81 wt.% I, 0.19 wt.% Xe) and MORB (BM-PC73 2.00 wt.% I, 0.33 wt.% Xe). These concentrations are slightly lower than for experiments performed in the presence of a single element. **Figure 2A** and **2B** summarize Xe and I concentrations at saturation conditions with respect to  $P$ . Both Xe and I solubility increase with  $P$  in melts, I solubility is possibly affected by the composition of silicate melts, which remains to be tested by structural studies, while Xe solubility is 10-fold higher in HPG than in basalts. However, such high concentrations are not expected in natural rocks, in the view of their very low terrestrial abundances.



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319 3.2. *In situ* high  $P$ - $T$  partition coefficients of iodine and xenon between melts and fluid

320

321 Xe and I concentrations measured *in situ* at high  $P$  and  $T$  in melts and fluids (SXRF), *ex situ* in  
322 glassy phases (PIXE) along with calculated partition coefficients are reported in **Table 3**. *Ex situ*  
323 data illustrate the crucial need of *in situ* data, due to strong Xe and I degassing during  
324 decompression in the aqueous phases that produces Xe-I-rich bubbles (**Figure 3**) and modifies  
325 dramatically the I and Xe concentration in quenched glasses compared to the silicate melts. *If*  
326 *remaining in the glass I content is too low to be detected*. Therefore, in the present case the analysis  
327 of both phases is necessary (silicate melt and aqueous fluid), as illustrated in the SXRF spectra  
328 presented in **Figure 4**. For experiment #02, the quantification of Xe in the supercritical fluid allows  
329 mass balance calculation, at high  $P$ : 45% melt and 55% aqueous fluid.

330 These partitioning experiments are performed between two phases, they differ from solubility  
331 experiments performed for a single phase. The presence of a water phase in the system strongly  
332 modifies the behaviour of both Xe and I in the silicate melt, the concentrations of these elements  
333 measured in aqueous fluids and in water-saturated silicate melts, in equilibrium with each other,  
334 cannot be compared with the solubility values measured in glasses quenched from high  $P$ - $T$  liquid  
335 silicates.

336

337 For a  $P$  range of 1.64 GPa to 1.22 GPa and a  $T$  range from 860 °C to 650 °C,  $D^{\text{Xe}}_{\text{f/m}}$  ranges from  
338 0.15 to 2748 (including Xe-I doped experiments), while I partitioning could only be measured at

1.5 GPa and 794°C with  $D_{f/m}^I = 0.97$  (**Figure 5**). Partition coefficients increase with decreasing  $P$ , implying strong degassing processes for both Xe and I. The comparison with previous results for I, i.e.  $D_{f/m}^I$  from 1.92 to 41, at  $P$  from 1.8 GPa to 0.1 GPa (Bureau et al., 2016) shows that Xe is more efficiently lost from the silicate melt during decompression than I. A relationship is observed between the partitioning of Xe and its melting curve as shown in **Figure 6**, where the melting curve is presented together with the  $P$ - $T$  conditions for each experiment. For experiment #02, the first partition coefficient ( $D_{f/m}^{Xe} 0.15$  at 851°C-1.68 GPa) was measured 10°C below total miscibility conditions (i.e. when the silicate melt and aqueous fluid separate from a single supercritical fluid). For that experiment only the  $D_{f/m}^{Xe}$  is below unity at 0.15, which implies more Xe dissolved in the melt than exsolved. It is observed in **Figure 6** that the  $P$ - $T$  conditions of this experiment fall close to the melting curve of Xe, while those of all other experiments are below this curve, and correspond to situations where Xe is largely dominant as a fluid phase ( $D_{f/m}^{Xe} \gg 1$ ). It shows that the exsolution of Xe from the silicate melt during both decompression and decrease of  $T$  is not promoted by the water phase as is the case for I (highly hydrophilic, Bureau et al., 2016), but is only determined by the  $P$ - $T$  conditions, in other words by the melting of Xe that is  $P$ - $T$  dependent. This observation is consistent with previous reports of Xe retention in silicates occurring systematically above the melting curve of Xe (Sanloup et al. 2005, Sanloup et al. 2011, Crépisson et al. 2019), a reaction likely driven under  $P$  by the important volume reduction between the reactants (i.e. silicates and molten Xe with a consequently large molar volume) and the product (i.e. Xe-doped silicates).

### 3.3. Iodine speciation in hydrous silicate melts

361

362 *In situ* Raman measurements are performed in the supercritical fluid enriched with respect to I at  
363 860°C, and in the melt and aqueous fluid in equilibrium close to miscibility at 850°C and 1.5 GPa  
364 (**Figure 7**). We observe a Raman band at about 134 cm<sup>-1</sup> in the aqueous fluid at both room and  
365 high *P-T* conditions, that progressively also appears in the HPG silicate melt albeit with a smaller  
366 intensity close to miscibility conditions, and in the single fluid resulting from total miscibility. We  
367 attribute this band to I species. The 110-174 cm<sup>-1</sup> range is indeed known to correspond to I-I bond  
368 stretching vibration of polyiodides moieties in aqueous solutions, as the Raman band at 180 cm<sup>-1</sup>  
369 for solid I<sub>2</sub> is expected to move towards lower frequencies when di-I interacts with donors to form  
370 complexes (Deplano et al., 1999). We do not observe oxidized I species (IO<sup>3-</sup>) at 779 cm<sup>-1</sup> (Shen  
371 et al., 1937). Spectroscopic data agree well with *in situ* partitioning measurements of Figure 6,  
372 showing a transfer of I to the aqueous fluid ( $D_{f/m}^I > 1$ ) when both *P* and *T* decrease.

373

## 374 **4. Discussion**

375

### 376 *4.1 Iodine and Xenon in silicate melts*

377

378 We measured I solubility in silicate melts from various settings (deep continental crust, subduction  
379 zone). Iodine contents for quenched hydrous glasses are up to 1.33±0.15 wt.% at 3.5 GPa in HPG,  
380 and up to 1.42±0.23 wt.% at 2 GPa for IAB. In all cases we show that I can be significantly  
381 incorporated in the silicate networks, meaning that I may be locally enriched in silicate melts at  
382 depth. This may have consequences on the real I content of the bulk silicate Earth that is believed

to be very low, based on measurements performed in basalts recovered at ambient conditions (McDonough and Sun, 1995).

We measured Xe solubility in silicate melts at 3.5 GPa for both HPG (up to 4 wt.% Xe) and basalts (tholeiites 0.4wt.%, Arc basalts 0.27 wt.% Xe), in agreement with data on tholeiitic melts from Schmidt and Keppler (2002). All results show that the amount of Xe dissolved in high  $P$  melts at depth may be significant, provided that  $P$  and  $T$  conditions remain above Xe melting curve. It is also observed that Xe and Ar have similar solubility values in wt% in silicate melts (this study, Leroy et al., 2018, Schmidt and Keppler 2002). The high solubility of Xe in HPG melts was shown to be due to Xe insertion in the melt ring structure (6-membered rings) through the formation of Xe-O bonds (Leroy et al., 2018), it is suggested that the same process may be at stake for basaltic melts that have much less 6-membered rings than polymerized silica-rich melts such as HPG, and a consequently much lower Xe solubility. The incorporation of Xe through  $P$ -induced oxidation process in the melt differs from the incorporation of I through I-I bonds of polyiodides, and/or potentially associated to alkali cations.

The *in situ* monitoring of Xe and I behaviours in fluids at conditions analogous to magma degassing show that degassing of I and Xe from silicate melts to the aqueous fluid involve different processes. I is decoupled from Xe because it systematically degasses with water (forming  $I_2$  moieties in the aqueous fluid), while Xe only degasses if its retention in the silicate melt is not thermodynamically favored. Similarly, mantle re-gassing processes at subduction zones could be different for both elements (see 4.3). We note a slope change in the pressure evolution in  $D_{f/m}^I$  circa 0.5 GPa, flattening at  $D_{f/m}^I = 4$ ; it might relate to a change of retention mechanism of iodine in the silicate melt as reported for bromine under pressure (Cochain et al., 2015).

#### 4.2 Iodine and Xenon in subduction zones

Although I is present at the ppb level in igneous rocks, the potential content of I in silicic melts and natural basaltic melts at high  $P$  can reach a few weight percent, similarly to bromine. This may have important consequences in subduction zones.

About 70% of the I present in the crust is believed to originate from sea water sediments of organic origin, indicating potentially a significant recycling of I to the mantle. I mobility in subduction zones is believed to be promoted by the dehydration of serpentinites, because serpentines have the ability to significantly incorporate I (John et al., 2011).

When subducted with marine sediments, part of I is mobilized in fluids and released in brines from the forearc areas, together with I from old organic-rich marine formations located in the upper plate of the forearc regions (i.e. Japanese forearc, Tomaru et al.; 2007). Volcanic gases and forearc fluids from subduction zone volcanoes are enriched in I with respect to other settings, as for example the Central American volcanic arc (Witt et al., 2008) or in Japan (Tomaru et al.; 2007). The abundance of I found in the volcanic arc material may therefore be controlled by fluid release from deserpentinizing slab mantle, a conclusion also supported by the analysis of I in back-arc basin basalts (Kendrick et al., 2012). The high solubility of I in natural basaltic melts reported here indicates that most of I released from serpentinites can be incorporated in magmas formed in the mantle wedge. The first melt produced in the slab at typical  $T$  for sub-arc depths are silica-rich with high alkali contents such as HPG (e.g. trondhjemites, Prouteau & Scaillet, 2013), possibly enriched

in I. Iodine volcanic injection to the atmosphere is potentially having significant consequences on the stratospheric ozone chemistry (Bureau et al., 2000).

While the major part of I present in subducted slabs is reinjected to the atmosphere, another part may be transported at depth in slabs. This is in agreement with the study of intraplate mantle xenoliths from various settings, some of them exhibiting I/Cl ratios more than one order of magnitude higher than the mantle value. It is proposed that I may be transported into the mantle at least beyond sub-arc depth (Kobayashi et al., 2019). This constrains the amount of subducted I into the convecting mantle to a few % of the subduction input. A deep storage of I is in agreement with the presence of I-rich aqueous fluid inclusions in lithospheric diamonds (Burgess et al., 2002), with the high halogen concentrations of the de-serpentinized harzburgite residues from serpentine breakdown, that are returned to the deep mantle (John et al., 2011) or in the seawater-derived I preserved in exhumed mantle wedge peridotite (Sumino et al., 2010). The processes suggested from our preliminary results may be migration and storage in deep melts.

We show the appearance of polyiodides moieties in hydrous silicate melt and supercritical fluid at 850°C and >1.5 GPa pressure. In a recent study of bromine in high *P* hydrous melts of HPG composition, Cochain and co-workers (2015) show that significant amounts of Br (up to 2.2 wt.%) can be retained in silicate melts at 3.4 GPa in the presence of water (up to 4.4 wt.%). They observe a change of Br coordination shell from a  $\text{Br}[\text{Na}_x(\text{H}_2\text{O})_y]$  complex at low *P* to an alkali structure at high *P* (at least from 2 GPa on although it remains to be investigated at lower *P*), whereby Br is surrounded by two oxygens at 1.9 Å and a mean of 6.6 Na atoms at 3.7 Å. A similar process may occur for I at high *P*, allowing a storage at depth in hydrous high *P* melts.

Xenon is also believed to be significantly recycled in the mantle (e.g. Holland and Ballentine 2006; Parai and Mukhopadhyay, 2018). Different processes are involved for I and Xe “ingassing” of the

mantle. During slab dehydration, a significant part of I would be efficiently washed out from the slab by aqueous fluids, and degassed back to the surface, while Xe is also recycled through aqueous fluids (**Figure 6**) but then it would remain in the mantle where it is expected to be stored in deep silicate melts/minerals. The recent examination of noble gases (including Xe) and halogens in serpentinites and secondary peridotites illustrates the key role played by serpentinites and other equivalent metamorphosed minerals where all these elements are progressively less abundant during dehydration processes and then recycled in the mantle (Kendrick et al., 2018). For Xe and possibly I, en route to the deep Earth, successive incorporations in mineral lattices and silicate melts may occur such as at the boundary between the transition zone and the lower mantle during dehydration melting process (Schmandt et al., 2014).

#### *4.3. The depletion in light xenon of the Modern Earth's atmosphere and the missing xenon*

The monitoring in real time of I and Xe partitioning during conditions analogous to magma ascent reveals two distinct behaviours. I is efficiently washed out of the silicate melt by water ( $D_{f/m}^I \gg 1$  see Bureau et al., 2016). Xenon is lost from the silicate melts much more significantly ( $D_{f/m}^{Xe}$  up to 2748) but only below its melting curve. Water can therefore not be involved in the depletion in light Xe of the Modern Earth's atmosphere, except in the eventuality of deep hydrous fluids trapping I, implying a deep storage of  $^{129}\text{Xe}$ , daughter isotope of  $^{129}\text{I}$ , that would cause an increased ( $^{129}\text{Xe}/^{132}\text{Xe}$ ) ratio in mantle partial melts (MORB). If the Early Earth was wet, both elements should have been totally degassed from the upper magma ocean, and potentially from the whole

magma ocean through effective mixing by convection. Fractional crystallization processes are proposed to account for the missing Xe, Xe can be stored in olivine up to 0.4 wt.% (Crepisson et al., 2018) through Si substitution in the olivine lattice and consequent oxidation of Xe, such as observed in high  $P$  melts (Leroy et al., 2018). Fractional crystallization of the magma ocean may have separated I from Xe by keeping Xe in minerals at depth while I (including  $^{129}\text{I}$ ) is concentrated in the melt and later lost to space.

In the current day Earth both I and Xe, can be stored in supercritical fluids if the  $T$  is high enough. Xe can be efficiently stored in deep crustal melts (i.e. very high Xe solubility in haplogranitic melts) and in the mantle, and/or in melt layers, such as at the lithosphere asthenosphere boundary (Crepisson et al., 2014), at 350 km depth (Tauzin et al., 2010), at the top of the transition zone at 410 km depth or below the TZ at 660 km depth (Schmandt et al. 2014), at the core mantle boundary CMB (Labrosse et al., 2007). This would contribute to the “missing Xe” and would imply a significant storage of I at depth, questioning in turn our knowledge of the abundance of I in the bulk silicate Earth.

## Conclusions

We have studied Xe and I behaviour at magmatic conditions, from room conditions up to 3.5 GPa and 1600°C, and under hydrous conditions for two magmatic systems, haplogranite (magma analogue for silicic melts) and basalts (MORB and IAB). We show that both elements reach high solubilities at high  $P$ , albeit with a strong composition dependence for Xe only. This strongly



contrasts with the low abundances of these elements in mantle products. Results illustrate that despite the proximity of Xe and I in the periodic table, their respective behaviour in silicate melts involve different mechanisms. I is incorporated as polyiodides (I-I bonds) while Xe is oxidized at depth (Xe-O bonds, Leroy et al., 2018). The *in situ* monitoring of shallow water degassing process reveals that both elements totally degas during decompression, but while water enhances the degassing of I, it has not effect on the degassing of Xe. Instead Xe degassing is strongly T-dependent with a partition coefficient one order of magnitude higher than for I at low  $T$ , but lower than unity below the Xe melting curve. Water can therefore not be involved in the depletion in light Xe of the Modern Earth's atmosphere, except for deep aqueous fluids [in the eventuality of a deep storage, at conditions below the Xe melting curve. Storage of Xe in deep melt layers melts layers, such as at the lithosphere-asthenosphere boundary, or deeper, may account for the missing Xe, possibly associated to fractional crystallization processes as](#) Xe incorporation in olivine can be significant (up to 0.4 wt.%, Crepisson et al., 2018).

Iodine is not expected to enter the lattice of nominally anhydrous phases, it is concentrated in silicate melts, causing the separation of  $^{129}\text{I}$  (later  $^{129}\text{Xe}$ ) from Xe. Deep silicate melts may favour the storage of I (and Xe) in the bulk silicate Earth, either upon early Earth's differentiation or upon modern recycling at subduction zones. The high capability of magmas to incorporate I at depth and to degas I during ascent in subduction zones is also a potential significant source of I for the atmosphere, an output that still lacks quantification today.

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**Table 1: Experimental conditions**

Sample	P (GPa) ( $\pm 0.05$ GPa)	T ( $^{\circ}$ C) ( $\pm 50^{\circ}$ C)	Duration (min)	Starting materials	Description
HPGI-03	0.4	1300	240	HPG + NaI + H <sub>2</sub> O	internally heated pressure vessel
HPGI-05	0.4	1300	240	HPG + NaI + H <sub>2</sub> O	internally heated pressure vessel
HPGXe-02	3.5	1400	60	HPG + Xe + H <sub>2</sub> O	Piston cylinder
HPGXe-C11	3.5	1600	60	HPG + Xe + H <sub>2</sub> O	Piston cylinder
HPGXe-C17	3.5	1400	240	HPG + Xe + H <sub>2</sub> O	Piston cylinder
HPGXe-C6	3.5	1400	240	HPG + Xe + H <sub>2</sub> O	Piston cylinder
HPGXeI-04	3.5	1600	240	HPG + Xe + NaI + H <sub>2</sub> O	Piston cylinder
BA-C7	0.2	1300	240	STV301 + NaI + H <sub>2</sub> O	internally heated pressure vessel
BA-PC82	3.5	1600	60	STV301 + NaI + H <sub>2</sub> O	Piston cylinder
BA-PC83	3.5	1600	60	STV301 + NaI + H <sub>2</sub> O	Piston cylinder
BA-PC87	2	1400	60	STV301 + NaI + H <sub>2</sub> O	Piston cylinder
BA-PC88	2	1400	60	STV301 + NaI	Piston cylinder - Fe-doped Pt capsule
BA-PC89	2	1400	60	STV301 + NaI	Piston cylinder - Fe-doped Pt capsule
BA-PC71	3.5	1600	60	STV301 + Xe + H <sub>2</sub> O	Piston cylinder
BM-PC71	3.5	1600	60	MORB + Xe + H <sub>2</sub> O	Piston cylinder
BA-PC73	3.5	1600	60	STV301 + NaI + Xe + H <sub>2</sub> O	Piston cylinder
BM-PC73	3.5	1600	60	MORB + NaI + Xe + H <sub>2</sub> O	Piston cylinder

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Sample names: HPG stands for Haplogranite composition, BA for Basalt of Arc (IAB) composition, BM for Basalt of MORB composition



**Table 2: Major elements (EMPA), iodine and xenon (PIXE), and water (ERDA) compositions.**

Starting materials and samples wt. %	SiO <sub>2</sub>	TiO <sub>2</sub>	Al <sub>2</sub> O <sub>3</sub>	FeO	MnO	MgO	CaO	Na <sub>2</sub> O	K <sub>2</sub> O	I Lα	I Lβ	Xe Lα	H <sub>2</sub> O	Total
HPG (1)	73.65		11.77					4.66	4.26				5.02	99.36
IAB (2)	47.01	1.07	15.28	8.79	0.16	12.5	10.95	2.23	0.47					98.47
N-MORB (3)	49.97	2.27	13.68	12.69	0.22	6.83	10.56	2.69	0.27					99.54
HPGI-03	75.17 (0.54)		12.10 (0.22)					4.11 (0.11)	4.04 (0.14)	0.47 (0.09)	0.47 (0.07)		4.04 (0.61)	99.92
HPGI-05	74.43 (0.55)		11.91 (0.31)					4.10 (0.18)	3.93 (0.19)	0.46 (0.06)	0.45 (0.07)		5.48 (0.82)	100.3
HPGXe-02	71.33 (0.66)		11.99 (0.22)					4.61 (0.07)	4.13 (0.17)			4.00 (0.80)	5.06 (0.75)	100.8
HPGXe-C11	72.03 (0.67)		11.86 (0.27)					4.44 (0.12)	4.06 (0.12)			3.54 (0.50)	4.66 (0.70)	100.3
HPGXe-C17	69.91 (0.63)		11.67 (0.13)					4.23 (0.11)	3.95 (0.13)			1.89 (0.38)	7.75 (1.1)	99.36
HPGXe-C6	69.57 (0.44)		11.41 (0.20)					4.26 (0.13)	3.89 (0.24)			3.72 (0.53)	7.74 (1.02)	100.6
HPGXeI-04	71.48 (0.50)		11.75 (0.20)					4.35 (0.15)	3.81 (0.13)	1.33 (0.15)		3.64 (0.52)	4.37 (0.65)	100.7
BA - C7	45.72 (0.43)	1.12 (0.04)	14.91 (0.18)	8.51 (0.30)	0.17 (0.05)	11.85 (0.12)	10.65 (0.16)	2.60 (0.08)	0.32 (0.04)		0.48 (0.07)		3.97 (0.60)	100.3
BA-PC82	46.91 (0.51)	1.37 (0.07)	15.16 (0.27)	0.95 (0.09)	0.16 (0.05)	11.58 (0.15)	10.58 (0.18)	2.98 (0.08)	0.50 (0.03)		2.97 (0.58)		3.82 (0.57)	96.98
BA-PC83	47.86 (0.45)	1.47 (0.09)	15.63 (0.28)	1.20 (0.07)	0.15 (0.07)	11.79 (0.15)	10.82 (0.22)	2.87 (0.08)	0.49 (0.04)		2.67 (0.58)		1.64 (0.25)	96.59
BA-PC87	49.83 (0.55)	1.36 (0.10)	16.34 (0.23)	1.05 (0.06)	0.12 (0.05)	12.40 (0.12)	11.42 (0.28)	2.82 (0.08)	0.40 (0.03)		1.42 (0.23)		2.33 (0.35)	99.49
BA-PC88*,**	50.83 (0.49)	1.29 (0.11)	16.54 (0.17)	1.75 (0.08)	0.15 (0.07)	12.89 (0.12)	11.43 (0.17)	2.66 (0.08)	0.45 (0.03)		0.66 (0.13)		0.66 (0.10)	99.31
BA-PC89*	48.51 (0.45)	1.25 (0.07)	16.04 (0.20)	5.77 (0.16)	0.17 (0.05)	12.23 (0.12)	11.09 (0.25)	2.88 (0.10)	0.36 (0.03)		1.02 (0.17)		0.13 (0.02)	99.45
BA-PC71	51.41 (0.67)	1.20 (0.06)	16.91 (0.27)	0.29 (0.09)	0.15 (0.04)	13.34 (0.12)	12.07 (0.17)	2.48 (0.06)	0.57 (0.05)			0.27 (0.10)	1.26 (0.20)	99.96
BM-PC71	55.74 (0.67)	2.65 (0.07)	14.78 (0.28)	1.02 (0.26)	0.22 (0.04)	7.47 (0.13)	11.98 (0.23)	2.93 (0.10)	0.33 (0.05)			0.40 (0.21)	1.89 (0.28)	99.41
BA-PC73	50.04 (0.48)	1.41 (0.04)	16.35 (0.28)	0.15 (0.08)	0.09 (0.07)	12.32 (0.13)	11.65 (0.19)	2.79 (0.09)	0.48 (0.04)		1.81 (0.19)	0.19 (0.16)	1.48 (0.42)	98.75
BM-PC73	54.77 (0.35)	2.81 (0.06)	14.43 (0.19)	0.34 (0.10)	0.17 (0.06)	7.36 (0.08)	12.07 (0.12)	3.35 (0.12)	0.27 (0.05)		2.00 (0.21)	0.33 (0.19)	1.46 (0.41)	99.36

(1) Bureau and Keppler, 1999; (2) STV301, Pichant et al., 2002; (3) BN06-11, Prouteau and Scaillet 2013. \* Fe-doped Pt capsules, \*\*water leak during the run. Sample names: HPG stands for Haplogranite composition, BA for Basalt of Arc (IAB) composition, BM for Basalt of MORB composition

**Table 3:** Concentrations and partition coefficients calculated from data collected in situ at pressures and temperatures

Exp.	Starting glass	T °C	P GPa	Xe <sub>fluid</sub> wt. %	Xe <sub>melt or glass</sub> wt. %	D <sup>Xe</sup> <sub>f/m</sub>	I <sub>fluid</sub> wt. %	I <sub>melt or glass</sub> wt. %	D <sup>I</sup> <sub>f/m</sub>
#01	HPGXe-C11	761	1.44	3.78	0.0148	255			
		692	1.31	3.88	0.0111	348			
		684	1.21	3.93	0.0014	2748			
		RT	RP	4.33					
		RT**	RP**		0.0283*			< DL	
		RT**	RP**		0.0345*			< DL	
#02	HPGXe-C11	>851 <sup>S</sup>	>1.68	2.915					
		851 <sup>CS</sup>	1.68	3.74	0.5645	0.15			
		791	1.50	5.16	0.1778	6.63			
		741	1.40	4.74	0.0252	188			
		692	1.31	4.011	0.0191	210			
		RT	RP	4.97					
		RT**	RP**		< DL				
		RT**	RP**		< DL				
#03	HPGXeI-04	794	1.51	4.6	0.053	86.4	1.72	1.77	0.97
		789	1.50		0.039			1.256	
		RT**	RP**		0.9682*			0.2547*	
		RT**	RP**		0.9953*			0.3153*	

Exp. Experiment. <sup>S</sup>supercritical fluid, in agreement with complete miscibility data for HPG-water after Bureau and Keppler (1999); <sup>CS</sup> close to supercritic \*quench bubbles included, \*\*from PIXE; Analytical uncertainties are 10% for SXRF, 5% for PIXE analysis. DL detection limit. Uncertainty on pressure is 0.2 GPa. Temperatures are corrected from the calibration performed at room pressure from chemicals of known melting points.

## Figures Captions

**Figure 1:** PIXE spectrum for A-HPGI-03 sample, B-BA-PC82 sample. x-rays are labelled. The dark line represents the experimental spectrum, dotted and dashed lines correspond to the theoretical fits for potassium K, calcium Ca and iodine I x-rays (see legends).

**Figure 2:** Xe wt.% (A) and I wt.% (B) concentrations measured in haplogranite (HPG) and basaltic glasses versus Pressure (GPa). These concentrations reflect the solubility data, except for I at 3.5 GPa, because the total amount of I in the starting compositions corresponds to the concentrations measured after the experiment. The concentration for Xe are in good agreement with the solubility measured in tholeiites (Schmidt and Keppler, 2002). Dark squares: HPG enriched in Xe, open squares: HPG enriched in both Xe and I, open circles: basalts enriched in Xe, dark circle: Basalts enriched in both Xe and I, diamonds: data from Schmidt and Keppler (2002). (C) Halogen element (I, Br, Cl) solubility in wt.% versus pressure in GPa for HPG melt and I solubility in IAB melts for comparison at low pressures (P max 0.4 GPa). Data for I (HPG: dark squares, IAB: open square) are from this study, data for Br (HPG: dark triangle) are from Bureau and Métrich (2003), data for Cl (HPG black circles) are from Webster (1997). Uncertainties are within the size of the symbols. All halogens exhibit similar ranges of solubility values in both melt compositions.

**Figure 3:** SEM images of the three quenched samples recovered from HP-HT DAC experiments. A. Sample chamber of exp #01, the arrow indicates the area zoomed in fig. 5.B ; B. Zoom on the large glassy globules showing the high density of Xe-I-rich bubbles that contaminate the glass analysis; C High density of Xe-I-rich bubbles in the glass that strongly contaminate the analysis. D. exp #02, small glassy globule, without bubbles.

**Figure 4:** In situ SRXF spectra of Xe content in DAC, in supercritical fluid (green) at  $T > 851^\circ\text{C}$ , aqueous fluid (red) and silicate melt (black) at 1.68 GPa and  $851^\circ\text{C}$  for experiment #02. The  $K\alpha$  and  $K\beta$  rays of Xe are observed at 29.78 keV and 33.62 keV respectively, measurements are performed in the supercritical fluid, then in the aqueous fluid and silicate melt at equilibrium at  $860^\circ\text{C}$  and 1 GPa. Xe is immediately lost from the silicate melt to the aqueous fluid formed upon cooling the supercritical fluid.

**Figure 5:** Pressure versus partition coefficient,  $D^{f/m}$ , for HPG starting glasses enriched in I and Xe. I data: this study in the presence of Xe (grey square), from Bureau et al., 2016 (open squares ); Xe

data: this study (dark circles grey circle in the presence of I). The cross represents the uncertainties. The inset represents all the data, including high  $D_{f/m}$ . The diagram shows a stronger affinity of Xe for the fluid phase compared to I with respect to pressure (i.e. magma ascent), except at high P where the opposite is observed, and the consequent decoupling of both elements. Grey symbols represent partition coefficient of I (square) and Xe (circle) when both elements are present in the system. They show that the presence of Xe has no influence on I partitioning and reciprocally. The slope change in the pressure evolution in  $D_I^{f/m}$  at 0.5 GPa, might relate to a change of retention mechanism of iodine in the silicate melt (see text).

**Figure 6:** Pressure versus temperature conditions of the three series of partitioning experiments (Table 3): #01 (open squares), #02 (open triangles), #03 (open circles Xe and I present), performed with starting glasses HPGXe-C11, HPGXeI-04 (see Table 1). Labels are the partition coefficients measured for Xe ( $D_{Xe}^{f/m}$ ). The grey curve is the melting curve of Xe after Belonoshko et al. (2006). The dashed curve is the P-T regime of crustal xenoliths after Rodriguez-Vargas et al. (2005). The P-T regime of the oceanic slab is out of the scale of the diagram. Results show a good agreement between the  $P$  and  $T$  conditions of one experiment performed 10°C below total miscibility (02, 851°C, 1.68 GPa) where most of the Xe is still in the silicate melt ( $D_{Xe}^{f/m}=0.15$ ). All the other experiments are located below the Xe melting curve, when most of the Xe, if not all, is transferred to the fluid phase. It suggests that Xe is recycled in the mantle in aqueous fluids where it can be incorporated in silicate melts.

**Figure 7:** In situ Raman spectroscopy in the DAC in the supercritical fluid (860°C, 1.6 GPa), and in the aqueous fluid and silicate melt (850°C, 1.5 GPa). This experiment was performed with a I-free HPG glass and I-enriched aqueous fluid (H<sub>2</sub>O-NaI 20 g/l) as starting materials following the protocol defined in Bureau et al. (2016). We observe a band at around 134 cm<sup>-1</sup> corresponding to I-I bond stretching vibrations (Deplano et al., 1999).

Figure 1

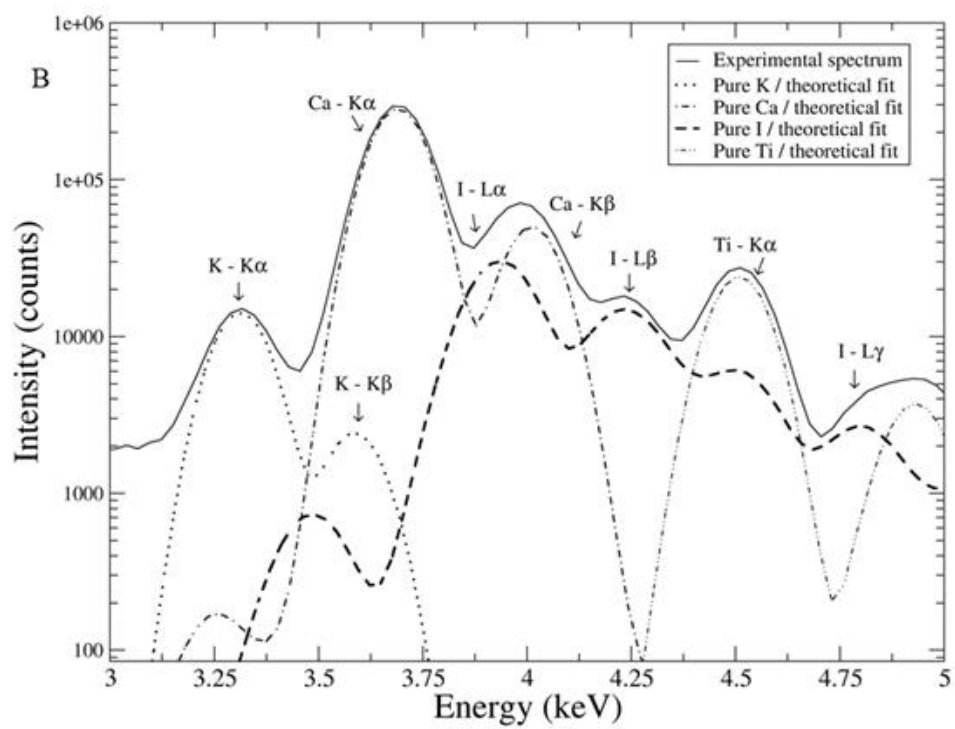
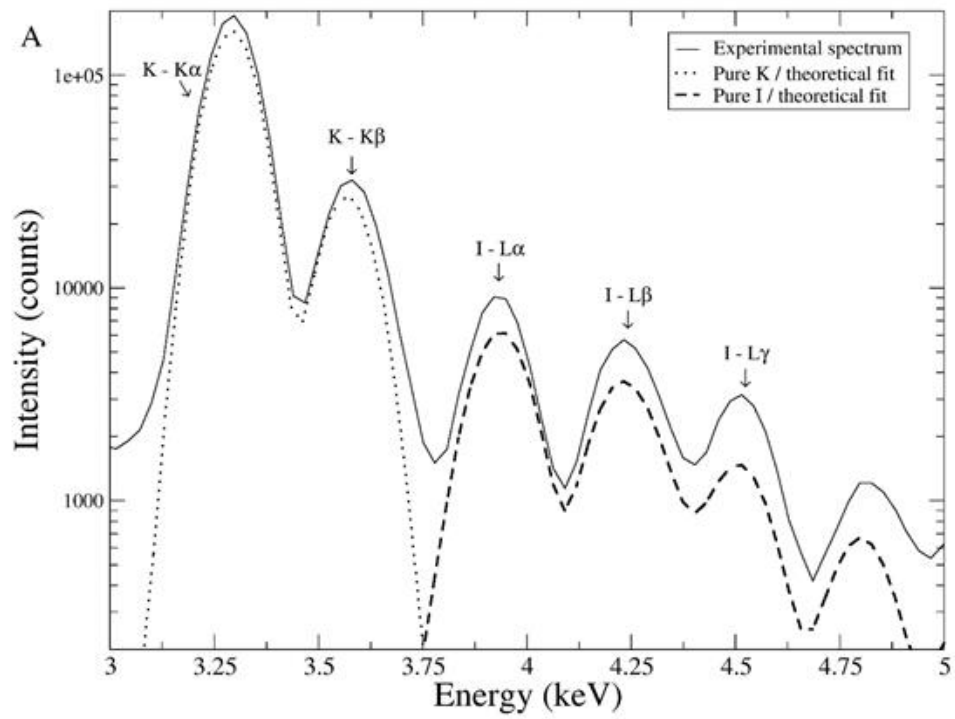


Figure 2

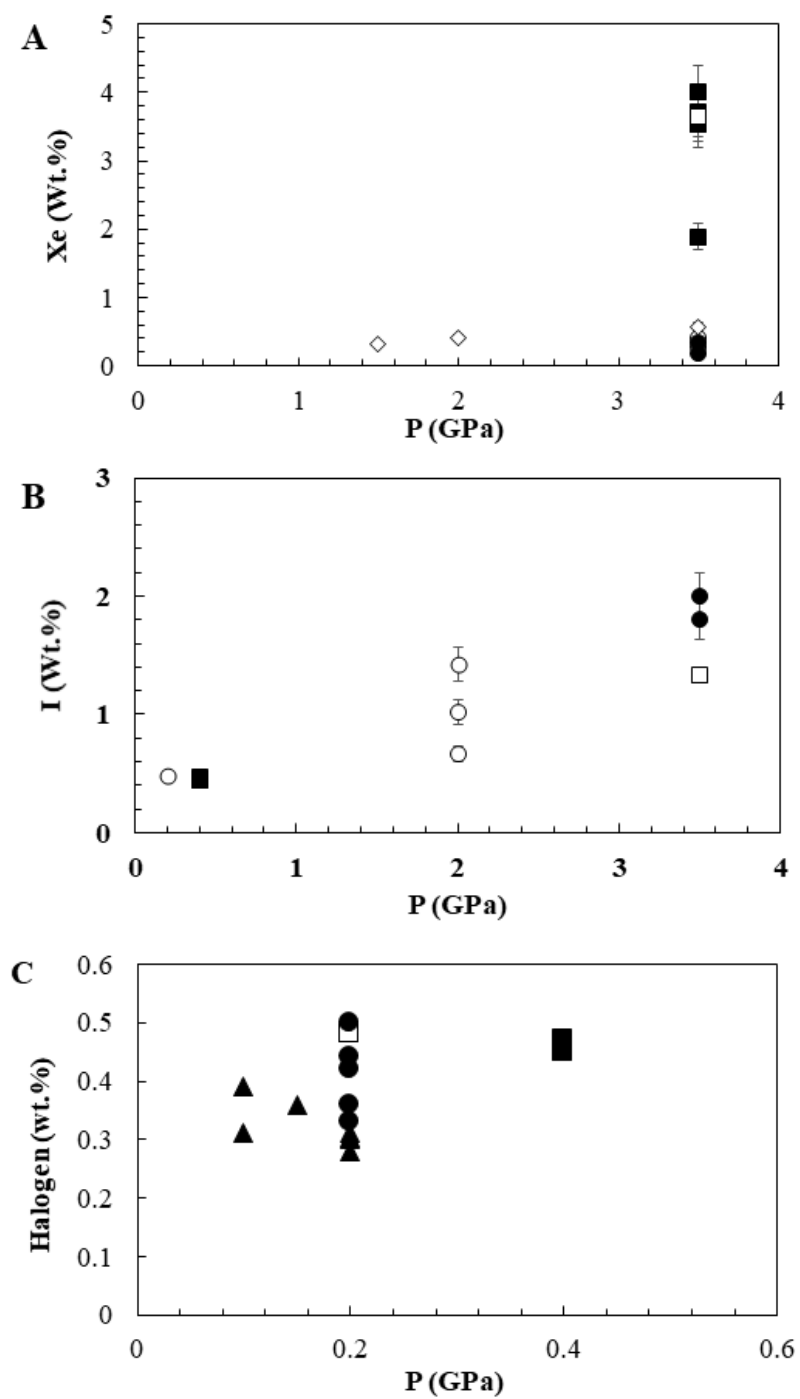


Figure 3

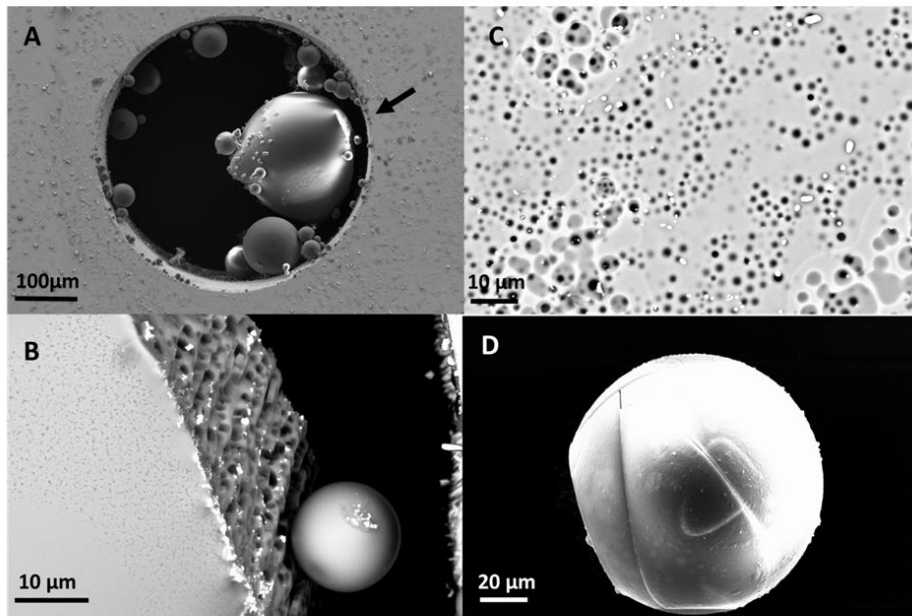


Figure 4

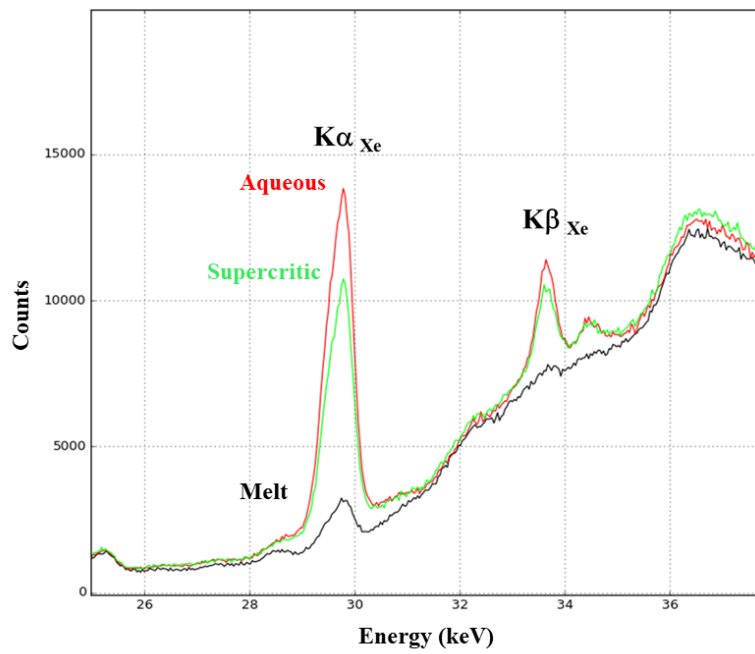




Figure 5

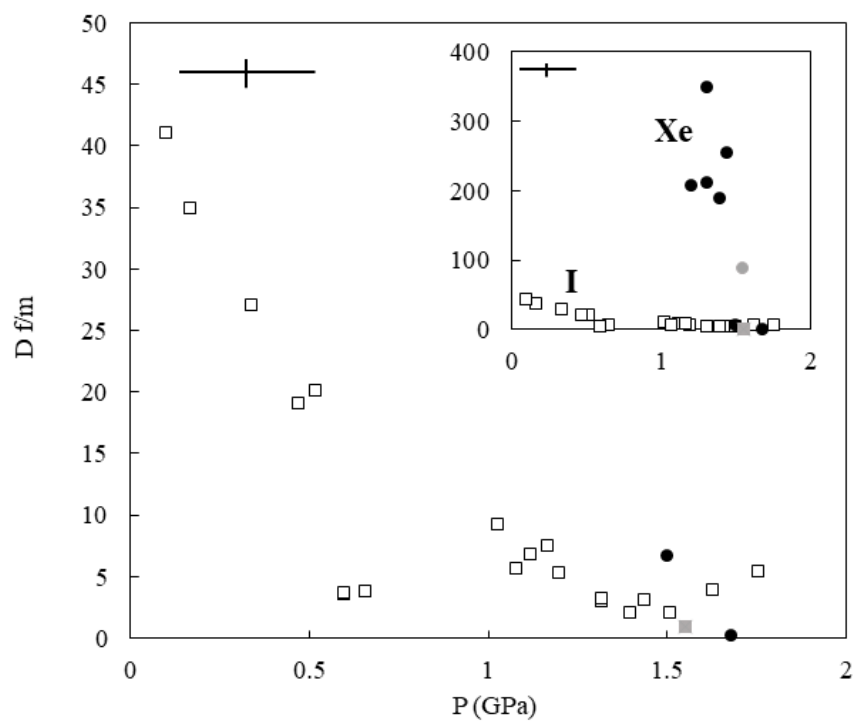


Figure 6

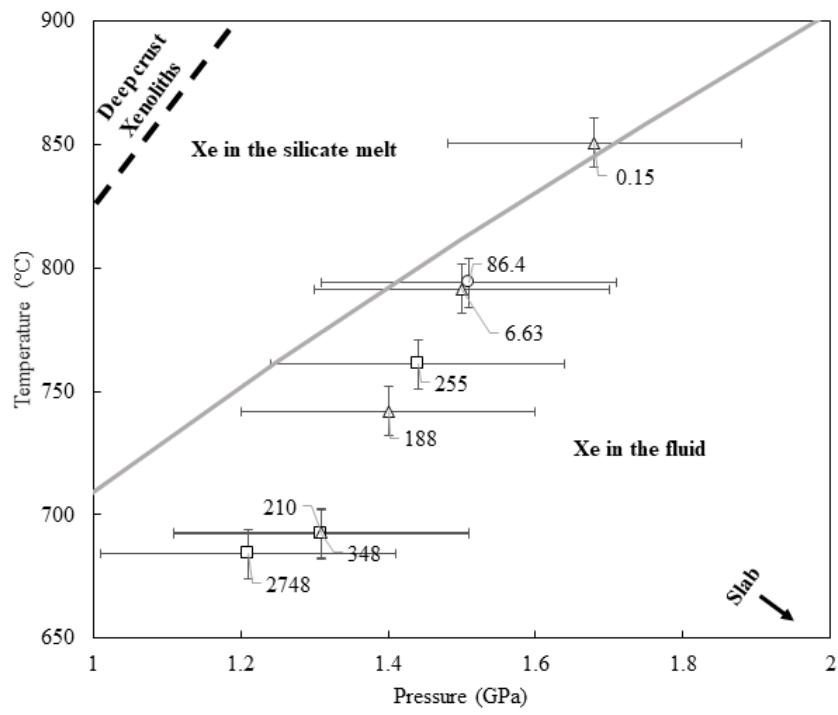


Figure 7

