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## HYPERLEDA

# II. The homogenized HI data^ 

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#### Abstract

After a compilation of HI data from 611 references and new observations made in Nançay, we produce a catalog of homogenized HI data for 16781 galaxies. The homogenization is made using the EPIDEMIC method from which all data are progressively converted into the adopted standard. The result is a catalog giving: 1) the logarithm of twice the maximum rotation velocity, $\log 2 V_{\mathrm{M}}^{\sin i}$, converted to the system of Mathewson et al. (1996). This quantity is given without correction for inclination; 2) the HI magnitude, $m_{21}$, (area of the $21-\mathrm{cm}$ line width expressed in magnitude) converted to the flux system of Theureau et al. (1998); 3) the HI velocity, $V_{\mathrm{HI}}$, expressed with the optical definition (i.e., using wavelengths instead frequencies). The typical uncertainties are: 0.04 for $\log 2 V_{\mathrm{M}}^{\sin i}, 0.25 \mathrm{mag}$ for $m_{21}$ and $9 \mathrm{~km} \mathrm{~s}^{-1}$ for $V_{\mathrm{HI}}$.


Key words. galaxies: general - catalogs

## 1. Introduction

Our first compilation of HI-data (Bottinelli et al. 1982) associated with optical ones was the origin of the LEDA database. The specificity of LEDA is to provide mean homogenized parameters, in the spirit of the series of Reference Catalogues initiated by G. and A. de Vaucouleurs (see the RC3 by de Vaucouleurs et al. 1991). The Tully-Fisher relation (Tully \& Fisher 1977) is the most straightforward application of such a large collection of HI and optical measurements.

The data are regularly maintained and the methods of homogenization are regularly revisited in order to take into account the evolution of measurements. In 1982 we collected HI measurements ( $21-\mathrm{cm}$ line width, HI flux or HI radial velocity) for 1210 galaxies (Bottinelli et al. 1982) and for 6439 galaxies in 1990 (Bottinelli et al. 1990). Today, we have measurements for 16781 galaxies. This increase could alone justify a new publication but another reason pushes us to revisit the method of homogenization: Many new measurements of logarithms of rotation velocity (hereafter, $\log 2 V_{\mathrm{M}}^{\sin i}$ ), obtained from rotation curves, are now available in the literature.

[^0]This gives us a new way to convert directly the observed $21-\mathrm{cm}$ line widths into the astrophysical parameter $\log 2 V_{\mathrm{M}}^{\sin i}$.

Further, the velocity resolution has been considerably improved (in 1982, 32 percent of our catalogue had measurements with a poor resolution of $63.5 \mathrm{~km} \mathrm{~s}^{-1}$ and only 50 percent of the sample were obtained with a resolution better than $25 \mathrm{~km} \mathrm{~s}^{-1}$. In the present catalogue, 92 percent of the data are obtained with a resolution better than $25 \mathrm{~km} \mathrm{~s}^{-1}$ and only a few galaxies (46) have no resolution better than $50 \mathrm{~km} \mathrm{~s}^{-1}$. This means that the correction for instrumental effects can be made effectively using a simple linear approximation.

In Sect. 2 we give a description of the present compilation of HI data in which new observations are included. The new observations are presented in the appendix. In Sect. 3 we briefly review the method of analysis (the EPIDEMIC method) applied to this study. Then, in Sects. 4-6, we study the homogenization of $21-\mathrm{cm}$ line widths, of HI radial velocities and of $21-\mathrm{cm}$ line fluxes, respectively. This allows us to produce a catalog of mean homogenized HI data, presented in the final section.

These data, together with the whole catalogue, are now available through the HYPERLEDA database (leda.univ-lyon1.fr), a project which aims at extending the capabilities of LEDA.

## 2. Main characteristics of the HI-compilation

We collected HI data from 611 papers. The full references with their code numbers are available in electronic form at the CDS.

Table 1. List of telescopes. Column 1: number of published 21-cm line widths; Col. 2: code number of the telescope in LEDA; Col. 3: Name of the telescope.

| $n$ | code | Radiotelescope |
| ---: | ---: | :--- |
| 19291 | 9 | Arecibo |
| 10089 | 1 | Nançay |
| 8702 | 7 | NRAO (91 m) |
| 3209 | 10 | Parkes |
| 2476 | 8 | Effelsberg |
| 1245 | 3 | NRAO (43 m) |
| 1062 | 5 | Jodrell B. (MarkI) |
| 816 | 2 | Westerbork |
| 244 | 14 | Jodrell B. (MarkII) |
| 114 | 11 | Owens Valley |
| 451 | 15 | VLA |
| 86 | 6 | Harvard |
| 54 | 12 | Dwingeloo |
| 38 | 4 | Cambridge |
| 37 | 13 | IAR (Argentina) |
| 3199 |  | Optical |
| 51113 |  | total |

Further, for 809 galaxies we added new measurements made with the radiotelescope of Nançay. These new data are presented in the appendix.

This compilation provides us with 51113 measurements of $21-\mathrm{cm}$ line widths or maximum velocity rotation, 31157 measurements of HI radial velocity and 25764 measurements of HI flux (area of the $21-\mathrm{cm}$ line) for 16781 galaxies. These data are characterized by some secondary parameters: telescope, velocity resolution, definition of the $21-\mathrm{cm}$ line width and bibliographic reference. Some additional parameters are also required: For instance the diameter and axis ratio of the observed galaxy (and sometimes the position angle) are needed to correct for the beam filling effect (see the section where the homogenization of fluxes is described). The optical radial velocity is also useful to provide a final check of velocities. These additional parameters are taken from LEDA.

Table 1 gives the list of telescopes with their code number and the number of $21-\mathrm{cm}$ line widths obtained by each.

Figure 1 gives the histogram of velocity resolutions. These resolutions will be shared in four classes designated by $r$, where $r$ represents the mean velocity resolution in $\mathrm{km} \mathrm{s}^{-1}$ (see Table 4):

$$
\begin{aligned}
& \text { - resol. } \leq 10 \mathrm{~km} \mathrm{~s}^{-1}\left(r=8 \mathrm{~km} \mathrm{~s}^{-1}\right) ; \\
& \text { - } 10 \mathrm{~km} \mathrm{~s}^{-1}<\text { resol. } \leq 18 \mathrm{~km} \mathrm{~s}^{-1}\left(r=16 \mathrm{~km} \mathrm{~s}^{-1}\right) ; \\
& \text { - } 18 \mathrm{~km} \mathrm{~s}^{-1}<\text { resol. } \leq 25 \mathrm{~km} \mathrm{~s}^{-1}\left(r=21 \mathrm{~km} \mathrm{~s}^{-1}\right) ; \\
& \text { - } 25 \mathrm{~km} \mathrm{~s}^{-1}<\text { resol. }\left(r=41 \mathrm{~km} \mathrm{~s}^{-1}\right)
\end{aligned}
$$

When the measurements are entered in our database, each $21-\mathrm{cm}$ line width receives a code defining how the measurement was done. These codes are given in Table 2 with their definitions and the number of occurences.

Finally, Table 3 presents the references giving more than 300 measurements. The reference numbers are arbitrary. They are internal numbers in LEDA. The full table, sorted by


Fig. 1. Histogram of velocity resolutions.
Table 2. Definition of the method of measurement of $21-\mathrm{cm}$ line width. The columns are arranged as follows: Col. 1: number of published $21-\mathrm{cm}$ line widths; Col. 2: code number of the method in LEDA; Col. 3: description of the method.

| $n$ | code | Definition of the level |
| ---: | ---: | :--- |
| 20684 | 5 | $50 \%$ of the maximum |
| 17546 | 1 | $20 \%$ of the maximum |
| 4702 | 8 | $50 \%$ of the mean |
| 3067 | 10 | $2 V_{\mathrm{M}}^{\sin i}$ from rotation curve |
| 1503 | 4 | $40 \%$ of the maximum |
| 1268 | 2 | $25 \%$ of the maximum |
| 1069 | 6 | $20 \%$ of the mean |
| 731 | 9 | mean of codes 1 and 8 |
| 395 | 7 | $25 \%$ of the mean |
| 139 | 0 | undetermined |
| 51113 |  | total |

the number of entries, is available in electronic form at the CDS. It is also available sorted in alphabetic order.

## 3. Description of the analysis: The EPIDEMIC method

At the time of the construction of our previous catalogs there was no standard system for $21-\mathrm{cm}$ measurements. This obliged us to refer the measurements to a mean system using the INTERCOMP method (see, Bottinelli et al. 1982). Today, thanks to some large homogeneous samples, it is possible to convert the measurements directly. This leads us to use a new method of analysis, the EPIDEMIC method (Paturel et al. 2003, Paper I). We start from a standard sample (a set of measurements giving a large and homogeneous sample). All other measurements are grouped into homogeneous classes (for instance, the class of measurements made at a given level and obtained with a given velocity resolution). Each class is crossidentified with the standard sample in order to establish the equation of conversion to the standard system. Then, the whole class is included in the standard sample which grows progressively (epidemic propagation). The order of inclusion follows the inverse of the quantity $t=\sigma / \sqrt{n}$, where $\sigma$ is the

Table 3. List of the richest references ( $n>300$ ). The columns are arranged as follows: Col. 1: number of measurements of 21-cm line width or $\log 2 V_{\mathrm{M}}^{\sin i} ;$ Col. 2: code number of the reference in LEDA; Col. 3: Full reference. The full table is available in electronic form at the CDS.

| $N$ | code | Reference |
| ---: | :--- | :--- |
| 4152 | 23289 | Theureau G., Bottinelli L., Coudreau-Durand N., et al. 1998, AAS, 130,333 |
| 2258 | 18000 | Fisher J. R., \& Tully R. B. 1981, ApJS, 47, 139 |
| 2048 | 23286 | Mathewson D. S., \& Ford V. L. 1996, ApJS, 107, 97 |
| 1392 | 23359 | (this paper) |
| 1361 | 23225 | Giovanelli R., Haynes M. P. 1993, AJ, 105(4), 1271 |
| 1267 | 16120 | Bottinelli L., Gouguenheim L., \& Paturel G. 1982, AA, 113, 61 |
| 1166 | 23296 | Giovanelli R., Avera E., \& Karachentsev I. D. 1997, AJ, 114(1), 122 |
| 1048 | 23200 | Schneider S. E., Thuan T. X., Mangum J. G., et al. 1992, ApJS, 81, 5 |
| 1036 | 24372 | Haynes M. P., Giovanelli R., Chamaraux P., Da Costa L. N., et al. 1999, AJ, 117, 2039 |
| 854 | 23151 | Giovanelli R., \& Haynes M. P. 1989, AJ, 97, 633 |
| 757 | 23122 | Haynes M. P., Giovanelli R., Starosta B. M., et al. 1988, AJ, 95, 607 |
| 726 | 23087 | Lewis B. M., Helou G., Salpeter E. E. 1985, ApJS, 59, 161 |
| 691 | 23070 | Giovanelli R., \& Haynes M. P. 1985, AJ, 90, 2445 |
| 636 | 23117 | Lewis B. M. 1987, ApJS, 515-542 |
| 628 | 23123 | Tifft W. G., \& Cocke W. J. 1988, ApJS, 67, 1-75 |
| 610 | 23152 | Giovanelli R., Haynes M. P., Myers S. T., \& Roth J. 1986, AJ, 92, 250 |
| 602 | 23192 | Haynes M. P., \& Giovanelli R. 1991, ApJS, 77, 331 |
| 546 | 16117 | Thuan T. X., \& Seitzer P. O. 1979, ApJ, 231, 327 |
| 535 | 23224 | Wegner G., Haynes M. P., \& Giovanelli R. 1993, AJ, 105(4), 1251 |
| 513 | 23092 | Bicay M. D., \& Giovanelli R. 1986, AJ, 91, 732 |
| 496 | 23091 | Bicay M. D., \& Giovanelli R. 1986, AJ, 91, 705 |
| 485 | 23226 | Lu N. Y., Hoffman G. L., Groff T., et al. 1993, ApJS, 88,383 |
| 482 | 23195 | Fouque P., Bottinelli L., Durand N., et al. 1990, AAS, 86, 473 |
| 481 | 23172 | Richter O. G., \& Huchtmeier W. K. 1991, AAS, 87, 425 |
| 474 | 23444 | Rosenberg J. L., \& Schneider S. E. 2000, ApJS, 130, 177 |
| 474 | 16093 | Shostak G. S. 1978, AA, 68, 321 |
| 472 | 23351 | Impey C. D., Burkholder V., \& Sprayberry D. 2001, AJ, 122, 2341 |
| 472 | 23067 | Bothun G. D., Aaronson M., Schommer R. A., et al. 1985, ApJS, 57, 423 |
| 448 | 23335 | Matthews L. D., \& van Driel W. 2000, AAS, 143, 421 |
| 439 | 23284 | Haynes M. P., Giovanelli R., Herter T., et al. 1997, AJ, 113(4), 1197 |
| 436 | 23333 | Huchtmeier W. K., Karachentsev I. D., Karachentseva V. E., et al. 2000, AAS, 141, 469 |
| 427 | 16057 | Rubin V. C., Ford Jr. W. K., Thonnard N., et al. 1976, AJ, 81, 687 |
| 423 | 16081 | Dickel J. R., \& Rood H. J. 1978, ApJ, 223, 391 |
| 367 | 12158 | Peterson S. D. 1979, ApJS, 40, 527 |
| 345 | 23113 | Hoffman G. L., Helou G., \& Salpeter E. E., et al. 1987, ApJS, 63, 247 |
| 344 | 23269 | Lu N. Y., \& Freudling W. 1995, ApJ, 449, 527 |
| 312 | 23121 | Staveley-Smith L., \& Davies R. D. 1988, MN, 231, 833-871 |
| $\ldots$ | $\ldots$ | .. |
| $\ldots$ | $\ldots$ | ... |
|  |  |  |

standard deviation of a preliminary comparison and $n$ the number of measurements. References having no intersection with the standard sample during the preliminary comparison are included following their population. All inclusions are made using weighted means. The standard sample receives the best standard deviation divided by $\sqrt{2}$. This assumes that it has the same scatter as the second best sample.

Obviously, the parameters used to make the classification must be relevant and we have to check that the equation of conversion is properly defined. Further, it is necessary to check that the conversion equation is reasonable (for instance, the conversion of flux should have a slope not very different from one in a $\log -\log$ scale). Let us see now the application of this method
to the homogenization of rotation velocities, heliocentric radial velocities and HI fluxes.

## 4. Homogenization of rotation velocities

### 4.1. First step

The $21-\mathrm{cm}$ line widths are used to derive the maximum rotation velocity $V_{\mathrm{M}}^{\sin i}$. It is important to note that in this paper, we consider always the quantity $\log 2 V_{\mathrm{M}}^{\sin i}$, that is twice the maximum rotation velocity in a logarithmic scale, in order to have a definition comparable with the $\log$ of a $21-\mathrm{cm}$ line width. Note that $\log 2 V_{\mathrm{M}}^{\sin i}$ is not corrected for the inclination.

The physical parameter measuring the maximum velocity rotation $V_{\mathrm{M}}$ is then related to $\log 2 V_{\mathrm{M}}^{\sin i}$ following the relation
$\log V_{\mathrm{M}}=\log 2 V_{\mathrm{M}}^{\sin i}-\log 2 \sin i$
where $i$ is the inclination of the galaxy (angle between the line of sight and the polar axis of the galaxy). Note that this relation implies a symmetric rotation curve. For non interacting galaxies this is reasonably justified.

The raw $21-\mathrm{cm}$ line widths are denoted $\log W$. They must be corrected for instrumental effects (velocity resolution $r$, standardization of the level of measurement, $l$ ) and then converted into $\log 2 V_{\mathrm{M}}^{\sin i}$ after correcting for internal velocity dispersion.

The adopted standard sample of $\log 2 V_{\mathrm{M}}^{\sin i}$ deduced from rotation curves is the sample by Mathewson et al. (1996) because it is the largest homogeneous one. Our new way of analysis allows us to convert directly $\log W(r, l)$ into $\log 2 V_{\mathrm{M}}^{\sin i}$ because we now have a large enough standard sample of $\log 2 V_{\mathrm{M}}^{\sin i}$. The correction for the resolution, the conversion to a standard level of measurement and the correction of internal dispersion are all done in a single operation provided that the data are distributed into homogeneous subsamples.

The choice of a $\log -\log$ conversion relation can be predicted in the case of good velocity resolution. However, we verified that it is linear. As an example, we show in Fig. 2 the relationship for the two richest subsamples.

It can be noted that some discrepant points are visible. They can result from an HI confusion by a close galaxy present in the lobe of the radiotelescope or from a bad detection. These discrepant objects are recorded for each regression in order to put a question mark on the final $\log 2 V_{\mathrm{M}}^{\sin i}$ value.

The total sample is divided into 26 subsamples defined by a given resolution $r$ and a given level $l$. After the preliminary test, the standard deviation of the initial standard sample of $\log 2 V_{\mathrm{M}}^{\sin i}$ is taken as $\sigma\left(\log 2 V_{\mathrm{M}}^{\sin i}\right)=0.027$. The results of the EPIDEMIC method is shown in Table 4. This table gives the slope $a$ and the intercept $b$ of the equation:
$\log 2 V_{\mathrm{M}}^{\sin i}=a . \log W(r, l)+b$.

It also gives the standard deviation $\sigma$, the number of points $n$ of the regression, and $N_{\mathrm{s}}$, the growing number of standard measurements.

This table gives the coefficients to convert any $21-\mathrm{cm}$ line width (measured with a resolution $r$ at a level $l$ ) into $\log 2 V_{\mathrm{M}}^{\sin i}$. It gives also the standard deviation $\sigma$ assigning a weight $1 / \sigma^{2}$ to the measurement.

It is visible from Table 4 that the level $l=5(50 \%$ of the maximum) gives a slope close to one, while the level $l=1$ ( $20 \%$ of the maximum) gives a slope significantly larger. One can see also that the low-resolution class ( $r=41 \mathrm{~km} \mathrm{~s}^{-1}$ ), for $l=1$ or $l=5$, gives systematically larger standard deviations. This reflects the fact that this class is less accurate and less homogeneous.


Fig. 2. An example of the regression used to convert log of $21-\mathrm{cm}$ line widths $\log W(r, l)$ into $\log 2 V_{\mathrm{M}}^{\sin i}$. Both quantities are observed ones, uncorrected for inclination. $\log W(r, l)$ refers to measurements made with a same class of velocity resolution and a same level of measurement. The regression is quite linear. The list of discrepant measurements (tiny dots) obtained with all references will be used to correct (or to erase) doubtful measurements.

Table 4. Correction of $\log W(r, l)$ to $\log 2 V_{\mathrm{M}}^{\sin i}$. This table gives the conversion coefficients and the mean error for the 26 subsamples (level and resolution). $l$ is the level code (see Table 2). $r$ is the mean velocity resolution in $\mathrm{km} \mathrm{s}^{-1}$. The full table is available in the electronic version of the journal.

| $l$ | $r$ | $a \pm \sigma_{a}$ | $b \pm \sigma_{b}$ | $\sigma$ | $n$ | $N_{\mathrm{s}}$ |
| ---: | ---: | ---: | ---: | ---: | ---: | ---: |
| 5 | 8 | $1.071 \pm 0.009$ | $-0.210 \pm 0.023$ | 0.044 | 928 | 2933 |
| 1 | 8 | $1.187 \pm 0.002$ | $-0.543 \pm 0.005$ | 0.036 | 5309 | 8952 |
| 1 | 21 | $1.193 \pm 0.005$ | $-0.563 \pm 0.012$ | 0.043 | 1438 | 9371 |
|  |  | $\ldots$ | $\ldots$ |  |  |  |
|  |  | $\ldots$ | $\ldots$ |  |  |  |
|  |  | $\ldots$ | $\ldots$ |  |  |  |

### 4.2. Test reference by reference

As a refinement ${ }^{1}$ we apply again the EPIDEMIC method using the bibliographic references to define the different homogeneous classes. The $\log 2 V_{\mathrm{M}}^{\sin i}$ measurements by Mathewson et al. (1996) are used as a standard. For this application, the $\log W(r, l)$ are first converted to $\log 2 V_{\mathrm{M}}^{\sin i}$ using the previous result (Table 4). This will allow us to improve the results for some references. If the first correction was perfect, the new

[^1]Table 5. Correction of $\log 2 V_{\mathrm{M}}^{\sin i}$ to the adopted standard (ref. 23286). This table gives the conversion coefficients and the mean error for the 65 references requiring an additional correction. The full table is available in the electronic version of the journal.

| Ref. | $a \pm \sigma_{a}$ | $b \pm \sigma_{b}$ | $\sigma$ | $n$ | $N_{\mathrm{s}}$ |
| ---: | ---: | ---: | ---: | ---: | ---: |
| 23289 | $1.000 \pm 0.000$ | $-0.007 \pm 0.002$ | 0.036 | 390 | 2439 |
| 18000 | $1.054 \pm 0.010$ | $-0.132 \pm 0.024$ | 0.024 | 158 | 4084 |
| 23121 | $1.026 \pm 0.010$ | $-0.058 \pm 0.024$ | 0.021 | 95 | 5064 |
| $\ldots$ | $\ldots$ | $\ldots$ |  |  |  |
| $\ldots$ | $\ldots$ | $\ldots$ |  |  |  |
| $\ldots$ | $\ldots$ | $\ldots$ |  |  |  |

slope should be $a=1$ and the zero-point should be $b=0$ for any references. We keep only those references for which the slope or the zero point are significantly different from that. A Student's $t$-test is first made on the slope $a$. If the slope is not significantly different from 1 at the 0.01 probability level, it is taken as 1 , exactly, and $b$ is recalculated. Another $t$-test is made on the new $b$ value. If it does not differ significantly from zero the reference does not need additional correction. Otherwise, the values of $a$ or $b$ are adopted, as given in Table 5. 65 references needed an additional correction. We established a list of 2087 galaxies that may have one discrepant measurement. This list will be used when we construct the final homogeneous catalog.

## 5. Homogenization of HI fluxes

### 5.1. Preliminary correction

The raw fluxes $F$ collected in the literature are converted into a logarithmic scale using $m_{21}$ magnitudes defined as (Vaucouleurs et al. 1991):
$m_{21}=-2.5 \log (0.2366 F)+15.84$
$F$ in Jy km s ${ }^{-1}$
Before applying the EPIDEMIC method, these magnitudes must be corrected for the beam filling effect. Sometimes radioastronomers publish corrected fluxes but, when possible, we prefer to collect raw fluxes in order to have a homogeneous conversion. The form of the correction is well established. We will use the correction established in our previous study (Bottinelli et al. 1990). Let us recall how the corrected HI magnitude, $m_{21}^{c}$, is calculated:
$m_{21}^{c}=m_{21}-1.25 \log [(1+x T)(1+x t)]$
$T=D_{25}^{2} / \theta^{2}$
$t=d_{25}^{2} / \theta^{2}$
$D_{25}$ and $d_{25}$ are respectively the major and minor axis diameter at the $25 B-$ mag $^{\operatorname{arcsec}^{-2}}$ and $\theta$ is the half power beam size of the considered radiotelescope. In this equation $D_{25}, d_{25}$ and $\theta$ are expressed with the same unit (i.e. arcmin). The parameter $x$ is $x=0.72 \pm 0.06$ (Bottinelli et al. 1990). We will keep this value.

Table 6. Correction of $m_{21}$ to the scale of Nançay. This table gives the the zero-point shift and the mean error for the different radiotelescope (Col. 1). By its adoption as a standard, the Nançay radiotelescope ( $t=$ $1)$ has $b=0$ and its standard deviation is $\sigma=0.24$. The full table is available in the electronic version of the journal.

| $t$ | $b \pm \sigma_{b}$ | $\sigma$ | $n$ | $N_{\mathrm{s}}$ |
| :---: | ---: | ---: | ---: | ---: |
| 7 | $0.052 \pm 0.016$ | 0.393 | 647 | 4182 |
| 9 | $-0.086 \pm 0.013$ | 0.415 | 1043 | 6413 |
| 3 | $0.051 \pm 0.012$ | 0.260 | 467 | 12813 |
|  | $\ldots$ | $\ldots$ |  |  |
|  | $\ldots$ | $\ldots$ |  |  |
|  | $\ldots$ | $\ldots$ |  |  |

For the Nançay radiotelescope the correction is more complex because the half power beam is not circular $\left(\approx 21^{\prime} \mathrm{NS} \times 4^{\prime} \mathrm{EW}\right)$. We adopted the the same correction (Rel. (4)) but:
$T=\left(D_{25}^{2} \sin ^{2} \beta+d_{25}^{2} \cos ^{2} \beta\right) / \theta_{\mathrm{EW}}^{2}$
$t=\left(D_{25}^{2} \cos ^{2} \beta+d_{25}^{2} \sin ^{2} \beta\right) / \theta_{\mathrm{NS}}^{2}$
where $\beta$ is the position angle of the considered galaxy (counted from North towards East) and where the ellipse of the half power bean size is defined by $\theta_{\mathrm{EW}}=4^{\prime}$ and $\theta_{\mathrm{NS}}=21^{\prime} . D_{25}$, $d_{25}$ and $\beta$ are calculated following Paturel et al. (2003, Paper I).

### 5.2. Homogenization by telescope

The EPIDEMIC method can now be applied to $m_{21}^{c}$. We expect a systematic effect depending on the radiotelescope itself. Thus, the classification of the EPIDEMIC method is built using the code for each radiotelescope. We use the Nançay radiotelescope $(t=1)$ as a standard because its large beam is less sensitive to the beam filling effect and because it constitutes one of the largest sample (see Table 1). The adopted equation of conversion is:
$m_{21}^{c}(t=1)=m_{21}^{c}(t)+b$.
This equation assumes that there is no significant scale error but only a zero-point shift. This is what we assumed in our previous HI catalogs, in agreement with the result (Paturel et al. 1991) that apparent differences in scale often reflect different limiting magnitudes. The result for each radiotelescope is given in Table 6.

## 5.3. homogneization reference by reference

We check reference by reference to find specific corrections of the zero-point of the $m_{21}^{c}$ magnitudes. They are first converted to the Nançay scale using the results from Table 6. For this second run, we adopted as a standard the reference 23289. Nevertheless, in order to reduce the number of references requiring a specific correction, it has been shifted by 0.2 mag (this correction has been obtained from a weighted mean of individual shifts, reference by reference). In other words,

Table 7. Correction of $m_{21}$ to the standard (ref. 23289). This table gives the zero-point shift and its mean error for the 46 references requiring an additional correction. The full table is available in the electronic version of the journal.

| Ref. | $b_{r} \pm \sigma_{b_{r}}$ | $\sigma$ | $n$ | $N_{\mathrm{s}}$ |
| ---: | ---: | ---: | ---: | ---: |
| 23195 | $-0.177 \pm 0.019$ | 0.147 | 61 | 2997 |
| 23321 | $-0.134 \pm 0.031$ | 0.183 | 41 | 3625 |
| 23359 | $-0.095 \pm 0.036$ | 0.266 | 56 | 3866 |
| $\ldots$ | $\ldots$ | $\ldots$ |  |  |
| $\ldots$ | $\ldots$ | $\ldots$ |  |  |
| $\ldots$ | $\ldots$ | $\ldots$ |  |  |

the actual zero-point of the magnitude scale is built from a mean system. The final $m_{21}^{F}$ magnitudes is then given by:
$m_{21}^{F}=m_{21}^{c}(t=1)+b_{r}-0.2$.
The value of $b_{r}$ is zero except for 46 references for which it is given in Table 7.

## 6. Homogenization of HI radial velocities

The EPIDEMIC method is now used for the heliocentric radial velocities. Note that all these velocities should be expressed with the optical convention $V=c\left(\lambda-\lambda_{0}\right) / \lambda_{0}$. We do not expect any sources of systematic effects. Thus we simply check that there is no reference effect. The homogeneous classes of the EPIDEMIC method are thus made using the reference code. The standard sample is the reference 23289 because it is the largest one. Two references (23000 and 23316) show a significant effect. It appears that they are expressed as $V=c\left(v_{\circ}-v\right) / v_{\mathrm{o}}$, where $v$ is the frequency. This effect disappears for these references when we applied the correction $V=V(v) /(1-V(v) / c)$. One reference (23286) showed a barely significant departure from the standard. It could have been corrected by applying a scaling factor $V_{\text {corr. }}=(1.0016 \pm 0.0008) \mathrm{V}$. Because reference 23286 was adopted as the standard for rotation velocity measured from rotation curves, one can ask if this effect could affect the accuracy of $\log 2 V_{\mathrm{M}}^{\sin i}$. One can calculate that this effect is at least 40 times smaller than the typical uncertainty and is completely negligible.

Finally, we have to estimate the standard deviation for each reference. As we did in our previous compilations, the best way consists of finding a correlation between the standard deviation and the resolution velocity. For this purpose, we apply the EPIDEMIC method using the true velocity resolution $R$ as a parameter of classification. For each regression we get the standard deviation and we plot it as a function of the resolution. The result is shown in Fig. 3. A direct regression leads to the relation: $\sigma=0.2 R+6$. This result is compatible with the result we obtained earlier (Bottinelli et al. 1990): $\sigma=0.15 R+7.4$. Nevertheless, we prefer to calculate the regression using only the resolution smaller than $24 \mathrm{~km} \mathrm{~s}^{-1}$ because the sample is dominated by small resolutions (see Fig. 1). The result is thus
$\sigma=(0.30 \pm 0.10) . R+(5.5 \pm 1.4)$.
This leads to increased $\sigma$ for large resolutions, i.e., to reduce the weight of old, poor resolution measurements. This justifies


Fig. 3. Empirical relation between the velocity resolution $r$ and the standard deviation $\sigma$ of the regressions $V(r)$ vs. $V($ ref $=23289)$. The adopted solution calculated with the resolution smaller than $24 \mathrm{~km} \mathrm{~s}^{-1}$ is represented by the full line.
also that we adopted this last solution. In Fig. 3 this solution is drawn.

We established a list of 710 galaxies that may have one discrepant velocity to be noted in the final homogeneous catalog.

As an additional test, we checked the homogenized HI velocities against the optical ones extracted from the LEDA database. The direct regression leads to the following result (obtained with 10019 galaxies):
$V($ optical $)=(1.0000 \pm 0.001) . V_{\mathrm{HI}}-(5.3 \pm 0.7)$.
Stricto Sensu, the zero-point is significantly different from 0. However, the standard deviation is $\sigma=39 \mathrm{~km} \mathrm{~s}^{-1}$ and the mean difference of both systems ( $\Delta V=5.3 \mathrm{~km} \mathrm{~s}^{-1}$ ) has no practical incidence. This comparison provided us with another list of 634 galaxies for which there is a discrepancy between optical and HI velocities.

## 7. The final catalogue of homogenized HI data

Using the results of the previous sections we produced a catalog of homogenized HI data for 16781 galaxies. The distribution of heliocentric radial velocities (Fig. 4) shows a bump around $5000 \mathrm{~km} \mathrm{~s}^{-1}$, although the completeness in volume is satisfied only up to $\approx 2000 \mathrm{~km} \mathrm{~s}^{-1}$. This corresponds to the nearest clusters surrounding the Local Super Cluster.

Let us recall briefly how the homogenized data are calculated. The $21-\mathrm{cm}$ line widths are first converted into $\log 2 V_{\mathrm{M}}^{\sin i}$ using empirical linear relations (Table 4) depending on the velocity resolution and on the level of measurement. For some references an additional correction is applied (Table 5). The mean error of each individual measurement comes from the standard deviation of the linear regression used for the last conversion (Tables 4 or 5 ). The HI magnitudes $m 21$ are first corrected for the beam filling effect (Rel. (4)) and then converted to the Nançay scale using linear relations (Table 6) depending on the radiotelescope used. For some references an additional correction is applied (Table 7). The mean error of each individual


Fig. 4. Distribution of heliocentric radial velocities.
measurement is also calculated as previously stated (Tables 6 or 7). The HI heliocentric radial velocities are not transformed (except one reference that is transformed into the optical definition and one that is corrected for a small scale effect). The mean error of each measurement is calculated as a funtion of the velocity resolution. For unknown velocity resolution we adopted a mean standard deviation of $15 \mathrm{~km} \mathrm{~s}^{-1}$.

Finaly, for each galaxy the weighted mean and its actual error are calculated. This actual error (Paturel et al. 1997) takes into account the accuracy of each individual source of data and the dispersion of the considered measurements. It gives a realistic description of the uncertainty attached to each mean value. Note that some data receive a flag (:) to show that they may include a discrepant measurement. As far as possible, discordant measurements are rejected (typically when there are several measurements in agreement and one discordant). In order to show the quality of the data the histograms of actual errors on $\log 2 V_{\mathrm{M}}^{\sin i}, m_{21}$ and $v_{\mathrm{HI}}$ are given in Fig. 5. The beginning of the final catalog is given in Table 8. This catalog is available in electronic form at the CDS.

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## Appendix A: New Nançay observations

In this section we present the HI follow-up of a set of galaxies selected in the infrared (IRAS) and near-infrared (DENIS) with the aim of starting to build of a TF catalogue in BIJHK bands. This survey is the preliminary step of the present cosmological key-project of the refurbished Nançay radiotelescope, KLUN+, which intends to collect new HI profiles of some 8000 galaxies in the period 2001-2005, on the basis of a DENIS and 2-MASS target selection.

Objects have been selected according to their IRAS flux, $S_{60}>0.6 \mathrm{Jy}$, and/or their $I$ magnitude, $I<14.5$. Radial velocities where known except for some galaxies at low galactic latitude (211 objects with $|b|<20^{\circ}$ ).


Fig. 5. Distribution of mean errors on $\log 2 V_{\mathrm{M}}^{\sin i}, m_{21}$ and $v_{\mathrm{HI}}$.

The present catalogue contains the HI profile and parameters (velocity measurements, $21-\mathrm{cm}$ line widths, HI fluxes, signal to noise ratio, and rms noise) for 817 spiral galaxies.

All these observations have been carried out between 1994 and 1998 with the old system of the meridian-transit Nançay radiotelescope (France). This instrument is a single dish antenna with a collecting area of $6912 \mathrm{~m}^{2}$ equivalent to that of a

Table 8. The final catalog of homogenized HI data. The catalog gives these data for 16781 galaxies. The full catalog is available in electronic form at the CDS. Col. 1: PGC number from LEDA; Col. 2: Alternate name from LEDA; Col. 3: Right Ascension and Declination for the equinox 2000, in hours, minutes, seconds and tenths, and degrees, arcminutes and arcseconds; Col. 4: Mean homogenized decimal logarithm of twice the maximum rotation velocity uncorrected for inclination ( $V_{\mathrm{M}}^{\sin i}$ in $\mathrm{km} \mathrm{s}^{-1}$ ); Col. 5: Mean homogenized 21-cm magnitude (see Eq. (3)); Col. 6: Mean homogenized HI heliocentric radial velocity (in $\mathrm{km} \mathrm{s}^{-1}$ ); Each parameter is given with a flag $f$ telling if it is uncertain $(f=1$ ).

| PGC | Name | RA(2000) DEC (2000) | $\log 2 \mathrm{~V}$ _M^sini | f | m21 | f | VHI | f |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| PGC0000002 | UGC12889 | J000001.6+471629 | $2.641+/-0.031$ | 0 | 15.30+/-0.39 | 0 | 5017.+/- | 9. 0 |
| PGC0000006 | CGCG456-13 | J000002.1+155253 | $2.312+/-0.024$ | 0 | 17.73+/-0.15 | 0 | 6001.+/- | 10. 0 |
| PGC0000004 | KUG2357+228 | J000003.4+230516 | $2.151+/-0.028$ | 0 | 16.69+/-0.25 | 0 | 4461.+/- | 8. 0 |
| PGC0000010 | MCG0-01-015 | J000007.7-000228 | $1.825+/-0.048$ | 0 | 16.55+/- 0.27 | 0 | 7099.+/- | 9. 0 |
| PGC0000012 | MCG-1-01-016 | J000008.6-062226 | $2.560+/-0.027$ | 0 | 15.97+/- 0.09 | 0 | 6547.+/- | 8. 0 |
| PGC0000016 | MCG-1-01-017 | J000011.3-050932 | $2.424+/-0.032$ | 0 | 17.09+/-0.25 | 0 | 5669.+/- | 8. 0 |
| PGC0000020 | MCG0-01-016 | J000012.7+010711 | $2.335+/-0.030$ | 0 | 16.57+/-0.32 | 0 | 7390.+/- | 8. 0 |
| PGC0000029 | CGCG517-10 | J000022.4+343658 | $2.101+/-0.057$ | 0 | 17.47+/-0.19 | 0 | 12702.+/- | 10. 0 |
| PGC0000035 | UGC12894 | J000022. 5+392944 | $1.500+/-0.070$ | 0 | 15.68+/-0.13 | 0 | 334.+/- | 5. 0 |
| . $\cdot$ |  |  |  |  |  |  |  |  |
|  |  | ... |  |  |  |  |  |  |
|  |  |  |  |  |  |  |  |  |

Table A.1. New observations. (The full table is available in electronic form at the CDS). Col. 1: PGC number from LEDA; Col. 2: Alternate name from LEDA; Col. 3: Right Ascension and Declination for the equinox 2000, in hours, minutes, seconds and tenths, and degrees, arcminutes and arcseconds; Col. 4: 21-cm line width at $20 \%$ of the maximum (in $\mathrm{km} \mathrm{s}^{-1}$ ); Col. 5: 21-cm line width at $50 \%$ of the maximum (in $\mathrm{km} \mathrm{s}^{-1}$ ); Col. 6: HI heliocentric radial velocity (in $\mathrm{km}_{\mathrm{s}} \mathrm{s}^{-1}$ ); Col. 7: HI flux (area of the 21-cm line) (in $\mathrm{mJy} \mathrm{km} \mathrm{s}^{-1}$ ); Col. 8: Signal to Noise ration; Col. 9: Quality of the line (see text) and notes (see Table A.3).


94 m -diameter parabolic dish. The half-power beam width at $21-\mathrm{cm}$ is $3.6 \operatorname{arcmin}(\mathrm{EW}) \times 22 \operatorname{arcmin}(\mathrm{NS})$ (at zero declination). Observations where limited to declination $\delta>-38.5^{\circ}$.

The minimal system temperature at $\delta=15^{\circ}$ was about 37 K in both horizontal and vertical polarizations. The spectrometer was a 1024-channel autocorrelator of 6.4 MHz bandwidth. The spacing of the channels corresponds to $2.6 \mathrm{~km} \mathrm{~s}^{-1}$ at 21 cm with a bank of 512 channels in each polarization. After boxcar smoothing the final resolution is typically $10 \mathrm{~km} \mathrm{~s}^{-1}$. In the velocity-search mode the 1024 channels are split in four banks of 256 channels leading to a range of $4800 \mathrm{~km} \mathrm{~s}^{-1}$ (generally from 400 to $5200 \mathrm{~km} \mathrm{~s}^{-1}$ or from 5200 to $10000 \mathrm{~km} \mathrm{~s}^{-1}$ ). The gain of the antenna has been calibrated according to Fouqué et al. (1990).

Table A.2. Periods of gain stability and correction factors $\left(=F_{0} / F_{\mathrm{HI}}\right)$.

| before June 30th 1991 | 0.943 |
| :---: | :---: |
| Jul. 1991-Oct. 1992 | 1.000 |
| Nov. 1992-Nov. 1994 | 1.143 |
| Dec. 1994-Nov. 1995 | 1.290 |
| Dec. 1995-Nov. 1996 | 1.449 |
| Dec. 1996-Jun. 1999 | 1.508 |

We used the Nançay processing package SIR (Système Interactif de Réduction). The processing chain consist of a selection of good observation cycles (one "observation" is a series of on/off observational sequences), the straightening of the


Fig. A.1. Sample of the $21-\mathrm{cm}$ line profiles. The full figure is availaible in the electronic version at http://www.edpsciences.org

Table A.3. Comments on some galaxies. The full table is available in electronic form at the CDS. The references concerning these comments are given at the end of the table.

| PGC | Name | \| Comments |
| :---: | :---: | :---: |
| PGC0000312 | 8ZW 3 | measured while pointing towards PGC351, 2.5 arcmin northwards |
| PGC0001431 | KUG 0019-016 | HI confusion with PGC1434 |
| PGC0001434 | UGC 212 | \| $\mathrm{HI} \sim$ confusion with PGC1431 |
| PGC0004878 | MCG +7-3-29 | HI confusion with PGC4901 at V=5913 and PGC4915 (Sc RC3) at V=5868 |
| PGC0002308 | IC 1562 | \| probably HI~confusion with PGC2305 (V=3897, barred Sb (DaCosta1998) or Sc (Loveday1996)) in the beam about 4 arcmin south |
| PGC0005335 | CGCG 385-135 | \| HI~confusion with PGC73964 (Sc Dressler1980) at V=5513 and PGC169815 at $\mathrm{V}=5586$ |
| PGC0005960 | ESO 297-11 | HI confusion with PGC5959 (S0-a Prugniel1998, Sc Loveday 1996) at V=5263 |
| . $\cdot$ | . $\cdot$ | ... |
|  | $\cdots$ | $\cdots$ |
|  |  | $\cdots$ |

base-line by a polynomial fit, and the application of a boxcar smoothing. The maximum of the line is determined by eye as the mean value of the maxima of each profile's horn after taking into account the medium noise (evaluated in the base-line). The widths, measured at the standard levels $20 \%$ and $50 \%$ of that maximum, correspond to the "distance" separating the two external points of the profile at these intensity levels.

Together with some other large HI surveys in Nançay, all the data are presented in the on-line Nançay HI extragalactic database (http://klun.obs-nancay.fr). The observed radial velocities are listed in Table A. 1 (Col. 6) and correspond to the median point of the $21-\mathrm{cm}$ line profile measured at $20 \%$ of maximum intensity and translated into the optical velocity scale. The average uncertainty on $V_{20}$ is about $8 \mathrm{~km} \mathrm{~s}^{-1}$ and does not exceed $20 \mathrm{~km} \mathrm{~s}^{-1}$.

The widths $W_{20}$ and $W_{50}$ are expressed in $\mathrm{km} \mathrm{s}^{-1}$ (Table A.1, Cols. 4 and 5) and correspond to direct measurements on the $21-\mathrm{cm}$ profiles.

HI-fluxes $F_{\mathrm{HI}}$ (Table A.1, Col. 7) are expressed in Jy km s${ }^{-1}$ and calibrated using as a reference a set of 9 calibrators regularly observed during the survey period (see Theureau et al. 1998). The evolution of the ratio $F_{\mathrm{HI}} / F_{0}$ with time allows us to supervise the Nançay system and to get for each observation the optimal flux measurement. In total, six periods of gain stability have been considered, with corresponding flux correction factors (see Table A.2). From this table and the number of ON/OFF integration cycles observed in each period for a given object, we have been able to correct the flux measurement for each galaxy. Each measurement has been corrected for the Nançay beam effect when position angle, diameter and axis ratio were known (beam-uncorrected flux are flagged with a star).

HI profiles have been classified in four classes according to their quality:
A. high signal to noise, symmetrical, double horn profile, best data for TF application;
B. robust HI parameters for TF application, more noisy than "A" galaxies;
C. well detected but noisy profile, generally HI confused, but OK for radial velocity measurement;
D. marginally detected, not useful.

General notes, mainly concerning possible or confirmed HI -confusion with some identified companion, are given in Table A.1.

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## Online Material

Table 4. Correction of $\log W(r, l)$ to $\log 2 V_{M}^{\sin i}$. This table gives the conversion coefficients and the mean error for the 26 subsamples (level and resolution). $l$ is the level code (see Table 2). $r$ is the mean velocity resolution in $\mathrm{km} \mathrm{s}^{-1}$.

| $l$ | $r$ | $a \pm \sigma_{a}$ | $b \pm \sigma_{b}$ | $\sigma$ | $n$ | $N_{\mathrm{s}}$ |
| ---: | ---: | ---: | ---: | ---: | ---: | ---: |
| 5 | 8 | $1.071 \pm 0.009$ | $-0.210 \pm 0.023$ | 0.044 | 928 | 2933 |
| 1 | 8 | $1.187 \pm 0.002$ | $-0.543 \pm 0.005$ | 0.036 | 5309 | 8952 |
| 1 | 21 | $1.193 \pm 0.005$ | $-0.563 \pm 0.012$ | 0.043 | 1438 | 9371 |
| 5 | 21 | $1.049 \pm 0.002$ | $-0.158 \pm 0.005$ | 0.033 | 3003 | 11381 |
| 5 | 16 | $1.048 \pm 0.003$ | $-0.156 \pm 0.006$ | 0.025 | 1591 | 11952 |
| 5 | 41 | $1.052 \pm 0.009$ | $-0.176 \pm 0.021$ | 0.051 | 751 | 14687 |
| 1 | 16 | $1.179 \pm 0.003$ | $-0.529 \pm 0.008$ | 0.038 | 2065 | 15028 |
| 1 | 41 | $1.156 \pm 0.010$ | $-0.521 \pm 0.025$ | 0.068 | 1047 | 15785 |
| 8 | 16 | $1.124 \pm 0.003$ | $-0.376 \pm 0.007$ | 0.028 | 2412 | 15815 |
| 6 | 8 | $1.146 \pm 0.005$ | $-0.450 \pm 0.010$ | 0.028 | 477 | 16212 |
| 4 | 41 | $1.120 \pm 0.011$ | $-0.384 \pm 0.027$ | 0.043 | 447 | 16408 |
| 2 | 8 | $1.145 \pm 0.003$ | $-0.430 \pm 0.007$ | 0.024 | 720 | 16408 |
| 4 | 21 | $1.093 \pm 0.005$ | $-0.280 \pm 0.012$ | 0.019 | 251 | 16412 |
| 4 | 8 | $1.088 \pm 0.006$ | $-0.266 \pm 0.014$ | 0.015 | 143 | 16412 |
| 8 | 8 | $1.084 \pm 0.006$ | $-0.266 \pm 0.015$ | 0.037 | 538 | 16412 |
| 7 | 21 | $1.278 \pm 0.020$ | $-0.827 \pm 0.051$ | 0.029 | 99 | 16443 |
| 2 | 21 | $1.279 \pm 0.059$ | $-0.786 \pm 0.149$ | 0.057 | 40 | 16447 |
| 4 | 16 | $1.087 \pm 0.006$ | $-0.260 \pm 0.013$ | 0.015 | 121 | 16447 |
| 7 | 16 | $1.226 \pm 0.016$ | $-0.677 \pm 0.041$ | 0.027 | 112 | 16447 |
| 8 | 21 | $1.126 \pm 0.007$ | $-0.390 \pm 0.017$ | 0.028 | 465 | 16457 |
| 9 | 8 | $1.107 \pm 0.012$ | $-0.339 \pm 0.031$ | 0.027 | 155 | 16467 |
| 9 | 16 | $1.148 \pm 0.011$ | $-0.431 \pm 0.027$ | 0.024 | 141 | 16554 |
| 6 | 16 | $1.149 \pm 0.016$ | $-0.466 \pm 0.041$ | 0.029 | 126 | 16621 |
| 9 | 21 | $1.181 \pm 0.015$ | $-0.524 \pm 0.038$ | 0.030 | 83 | 16643 |
| 2 | 41 | $1.043 \pm 0.027$ | $-0.205 \pm 0.066$ | 0.038 | 76 | 16686 |
| 7 | 8 | $1.082 \pm 0.039$ | $-0.302 \pm 0.098$ | 0.037 | 27 | 16721 |
|  |  |  | 0.050 |  |  |  |

Table 5. Correction of $\log 2 V_{M}^{\sin i}$ to the adopted standard (ref. 23286). This table gives the conversion coefficients and the mean error for the 65 references requiring an additional correction.

| Ref. | $a \pm \sigma_{a}$ | $b \pm \sigma_{b}$ | $\sigma$ | $n$ | $N_{\text {s }}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 23289 | $1.000 \pm 0.000$ | $-0.007 \pm 0.002$ | 0.036 | 390 | 2439 |
| 18000 | $1.054 \pm 0.010$ | $-0.132 \pm 0.024$ | 0.024 | 158 | 4084 |
| 23121 | $1.026 \pm 0.010$ | $-0.058 \pm 0.024$ | 0.021 | 95 | 5064 |
| 23152 | $1.000 \pm 0.000$ | $-0.011 \pm 0.004$ | 0.030 | 51 | 5121 |
| 24372 | $1.021 \pm 0.008$ | $-0.062 \pm 0.019$ | 0.028 | 523 | 5384 |
| 12320 | $0.958 \pm 0.018$ | $0.087 \pm 0.043$ | 0.047 | 119 | 6019 |
| 23171 | $1.000 \pm 0.000$ | $0.020 \pm 0.004$ | 0.023 | 28 | 6074 |
| 27036 | $1.000 \pm 0.000$ | $-0.010 \pm 0.004$ | 0.029 | 54 | 6081 |
| 16120 | $1.066 \pm 0.016$ | $-0.177 \pm 0.038$ | 0.048 | 257 | 6129 |
| 23195 | $1.000 \pm 0.000$ | $-0.016 \pm 0.003$ | 0.026 | 84 | 6277 |
| 23366 | $1.082 \pm 0.024$ | $-0.194 \pm 0.053$ | 0.041 | 57 | 6424 |
| 16121 | $0.902 \pm 0.016$ | $0.250 \pm 0.038$ | 0.030 | 76 | 7094 |
| 12310 | $1.000 \pm 0.000$ | $-0.026 \pm 0.008$ | 0.054 | 45 | 7268 |
| 23302 | $1.087 \pm 0.018$ | $-0.216 \pm 0.041$ | 0.030 | 56 | 7316 |
| 123303 | $1.000 \pm 0.000$ | $-0.029 \pm 0.005$ | 0.055 | 140 | 7337 |
| 12379 | $1.000 \pm 0.000$ | $-0.020 \pm 0.009$ | 0.055 | 36 | 7391 |
| 23296 | $0.968 \pm 0.011$ | $0.063 \pm 0.026$ | 0.017 | 68 | 7422 |
| 23284 | $1.031 \pm 0.011$ | $-0.089 \pm 0.026$ | 0.014 | 80 | 7931 |
| 23087 | $0.957 \pm 0.009$ | $0.097 \pm 0.022$ | 0.025 | 102 | 8666 |
| 23123 | $1.048 \pm 0.005$ | $-0.115 \pm 0.010$ | 0.033 | 467 | 8791 |
| 23092 | $1.022 \pm 0.009$ | $-0.056 \pm 0.022$ | 0.019 | 59 | 8877 |
| 23091 | $1.032 \pm 0.012$ | $-0.083 \pm 0.030$ | 0.016 | 63 | 9067 |
| 23444 | $0.954 \pm 0.022$ | $0.102 \pm 0.051$ | 0.063 | 100 | 9250 |
| 23067 | $1.000 \pm 0.000$ | $-0.024 \pm 0.003$ | 0.037 | 149 | 9713 |
| 16057 | $1.000 \pm 0.000$ | $0.015 \pm 0.006$ | 0.053 | 72 | 9961 |
| 16081 | $1.071 \pm 0.018$ | $-0.193 \pm 0.042$ | 0.044 | 96 | 10008 |
| 12158 | $1.053 \pm 0.018$ | $-0.126 \pm 0.044$ | 0.040 | 85 | 10013 |
| 23212 | $1.029 \pm 0.014$ | $-0.086 \pm 0.033$ | 0.022 | 85 | 10063 |
| 23324 | $1.000 \pm 0.000$ | $-0.013 \pm 0.004$ | 0.029 | 65 | 10329 |
| 23095 | $0.922 \pm 0.026$ | $0.177 \pm 0.062$ | 0.046 | 58 | 10776 |
| 23009 | $1.076 \pm 0.018$ | $-0.189 \pm 0.042$ | 0.044 | 136 | 11104 |
| 23251 | $1.075 \pm 0.011$ | $-0.190 \pm 0.026$ | 0.022 | 133 | 11132 |
| 23048 | $1.000 \pm 0.000$ | $-0.011 \pm 0.002$ | 0.023 | 99 | 11193 |
| 12327 | $0.805 \pm 0.053$ | $0.482 \pm 0.124$ | 0.058 | 35 | 11395 |
| 23232 | $1.000 \pm 0.000$ | $-0.027 \pm 0.003$ | 0.026 | 57 | 11454 |
| 23321 | $0.955 \pm 0.017$ | $0.105 \pm 0.040$ | 0.026 | 52 | 11506 |
| 23252 | $0.843 \pm 0.044$ | $0.378 \pm 0.108$ | 0.052 | 28 | 11549 |
| 16049 | $1.045 \pm 0.017$ | $-0.114 \pm 0.038$ | 0.021 | 41 | 11609 |
| 23000 | $1.084 \pm 0.014$ | $-0.208 \pm 0.033$ | 0.030 | 68 | 11672 |
| 23319 | $0.930 \pm 0.034$ | $0.156 \pm 0.080$ | 0.042 | 21 | 11725 |
| 23231 | $0.975 \pm 0.012$ | $0.055 \pm 0.031$ | 0.013 | 37 | 11782 |
| 23271 | $1.000 \pm 0.000$ | $-0.019 \pm 0.007$ | 0.036 | 25 | 11816 |
| 23234 | $1.000 \pm 0.000$ | $0.016 \pm 0.005$ | 0.036 | 49 | 11867 |
| 12346 | $0.867 \pm 0.035$ | $0.328 \pm 0.084$ | 0.061 | 50 | 11876 |
| 12393 | $1.000 \pm 0.000$ | $-0.009 \pm 0.003$ | 0.029 | 79 | 11921 |
| 16100 | $1.000 \pm 0.000$ | $0.029 \pm 0.013$ | 0.050 | 17 | 11988 |
| 23220 | $1.000 \pm 0.000$ | $-0.031 \pm 0.015$ | 0.050 | 11 | 12015 |
| 23295 | $1.000 \pm 0.000$ | $-0.041 \pm 0.009$ | 0.031 | 11 | 12081 |
| 23194 | $1.000 \pm 0.000$ | $-0.017 \pm 0.007$ | 0.039 | 29 | 12182 |
| 23043 | $1.000 \pm 0.000$ | $0.034 \pm 0.005$ | 0.026 | 32 | 12253 |
| 23147 | $0.895 \pm 0.031$ | $0.236 \pm 0.078$ | 0.031 | 43 | 12257 |
| 23360 | $1.302 \pm 0.062$ | $-0.639 \pm 0.118$ | 0.068 | 35 | 12287 |
| 23352 | $0.290 \pm 0.265$ | $1.738 \pm 0.674$ | 0.145 | 12 | 12311 |
| 23307 | $1.258 \pm 0.042$ | $-0.715 \pm 0.104$ | 0.061 | 56 | 12355 |
| 23101 | $0.917 \pm 0.029$ | $0.201 \pm 0.065$ | 0.049 | 36 | 12355 |
| 23316 | $0.895 \pm 0.052$ | $0.210 \pm 0.105$ | 0.043 | 22 | 12373 |
| 23332 | $1.000 \pm 0.000$ | $-0.078 \pm 0.027$ | 0.104 | 16 | 12397 |
| 23120 | $0.894 \pm 0.044$ | $0.262 \pm 0.098$ | 0.040 | 26 | 12405 |
| 23075 | $0.951 \pm 0.018$ | $0.119 \pm 0.048$ | 0.012 | 14 | 12426 |
| 23084 | $0.843 \pm 0.036$ | $0.385 \pm 0.089$ | 0.028 | 26 | 12434 |
| 23198 | $1.000 \pm 0.000$ | $-0.027 \pm 0.011$ | 0.042 | 15 | 12441 |
| 23237 | $1.067 \pm 0.030$ | $-0.138 \pm 0.066$ | 0.021 | 14 | 12478 |
| 12234 | $1.000 \pm 0.000$ | $-0.044 \pm 0.018$ | 0.071 | 15 | 12490 |
| 12275 | $1.000 \pm 0.000$ | $-0.112 \pm 0.021$ | 0.074 | 15 | 12490 |
| 23081 | $1.000 \pm 0.000$ | $-0.111 \pm 0.020$ | 0.079 | 18 | 12492 |

Table 6. Correction of $m_{21}$ to the scale of Nançay. This table gives the the zero-point shift and the mean error for the different radiotelescope (Col. 1). By its adoption as a standard, the Nançay radiotelescope ( $t=$ 1) has $b=0$ and its standard deviation is $\sigma=0.24$.

| $t$ | $b \pm \sigma_{b}$ | $\sigma$ | $n$ | $N_{\mathrm{s}}$ |
| ---: | ---: | ---: | ---: | ---: |
| 7 | $0.052 \pm 0.016$ | 0.393 | 647 | 4182 |
| 9 | $-0.086 \pm 0.013$ | 0.415 | 1043 | 6413 |
| 3 | $0.051 \pm 0.012$ | 0.260 | 467 | 12813 |
| 10 | $0.076 \pm 0.015$ | 0.324 | 448 | 13115 |
| 5 | $0.045 \pm 0.013$ | 0.261 | 443 | 13947 |
| 8 | $0.073 \pm 0.013$ | 0.317 | 594 | 13989 |
| 2 | $-0.026 \pm 0.016$ | 0.272 | 322 | 14417 |
| 14 | $-0.036 \pm 0.037$ | 0.305 | 68 | 14426 |
| 15 | $0.095 \pm 0.036$ | 0.405 | 136 | 14443 |
| 11 | $0.038 \pm 0.112$ | 0.515 | 23 | 14497 |
| 13 | $0.088 \pm 0.083$ | 0.371 | 32 | 14498 |
| 12 | $-0.036 \pm 0.075$ | 0.285 | 22 | 14499 |
| 4 | $-0.069 \pm 0.088$ | 0.403 | 20 | 14500 |
| 6 | $0.287 \pm 0.174$ | 0.655 | 13 | 14500 |

Table 7. Correction of $m_{21}$ to the standard (ref. 23289). This table gives the zero-point shift and its mean error for the 46 references requiring an additional correction.

| Ref. | $b_{r} \pm \sigma_{b_{r}}$ | $\sigma$ | $n$ | $N_{\mathrm{s}}$ |
| ---: | ---: | ---: | ---: | ---: |
| 23195 | $-0.177 \pm 0.019$ | 0.147 | 61 | 2997 |
| 23321 | $-0.134 \pm 0.031$ | 0.183 | 41 | 3625 |
| 23359 | $-0.095 \pm 0.036$ | 0.266 | 56 | 3866 |
| 23224 | $-0.060 \pm 0.027$ | 0.200 | 57 | 4505 |
| 23296 | $0.082 \pm 0.036$ | 0.247 | 47 | 4710 |
| 23226 | $0.060 \pm 0.025$ | 0.171 | 55 | 5770 |
| 23239 | $-0.252 \pm 0.041$ | 0.268 | 42 | 5960 |
| 23200 | $0.051 \pm 0.020$ | 0.263 | 181 | 6864 |
| 23271 | $-0.298 \pm 0.052$ | 0.206 | 17 | 7414 |
| 23238 | $-0.289 \pm 0.029$ | 0.267 | 81 | 7479 |
| 18000 | $0.039 \pm 0.007$ | 0.189 | 669 | 7525 |
| 23251 | $-0.065 \pm 0.015$ | 0.151 | 100 | 7972 |
| 27036 | $-0.098 \pm 0.047$ | 0.309 | 54 | 8097 |
| 16120 | $0.104 \pm 0.020$ | 0.306 | 305 | 9378 |
| 23444 | $-0.344 \pm 0.049$ | 0.472 | 94 | 9495 |
| 23097 | $-0.086 \pm 0.027$ | 0.303 | 123 | 9654 |
| 23009 | $-0.099 \pm 0.021$ | 0.255 | 150 | 9795 |
| 23337 | $-0.224 \pm 0.096$ | 0.468 | 24 | 9836 |
| 23067 | $-0.172 \pm 0.018$ | 0.217 | 144 | 10059 |
| 23048 | $0.272 \pm 0.031$ | 0.322 | 111 | 10147 |
| 23117 | $-0.090 \pm 0.037$ | 0.370 | 102 | 10268 |
| 12320 | $-0.083 \pm 0.024$ | 0.277 | 129 | 10543 |
| 23213 | $-0.236 \pm 0.081$ | 0.449 | 32 | 10610 |
| 12393 | $0.314 \pm 0.059$ | 0.276 | 86 | 10929 |
| 23318 | $0.075 \pm 0.037$ | 0.240 | 48 | 10933 |
| 23290 | $0.038 \pm 0.011$ | 0.109 | 95 | 11056 |
| 23338 | $-0.204 \pm 0.060$ | 0.248 | 20 | 11069 |
| 12310 | $0.286 \pm 0.042$ | 0.266 | 41 | 11145 |
| 16121 | $-0.134 \pm 0.027$ | 0.250 | 87 | 11196 |
| 23302 | $-0.248 \pm 0.044$ | 0.342 | 63 | 11361 |
| 23268 | $0.169 \pm 0.067$ | 0.266 | 32 | 11458 |
| 23231 | $0.078 \pm 0.035$ | 0.223 | 46 | 11506 |
| 23342 | $-0.095 \pm 0.029$ | 0.183 | 39 | 11630 |
| 23147 | $-0.385 \pm 0.061$ | 0.322 | 48 | 11664 |
| 23329 | $0.187 \pm 0.051$ | 0.315 | 40 | 11684 |
| 23279 | $0.129 \pm 0.055$ | 0.309 | 32 | 11714 |
| 23187 | $-0.094 \pm 0.041$ | 0.179 | 18 | 11840 |
| 23264 | $0.201 \pm 0.057$ | 0.250 | 18 | 11920 |
| 23265 | $-0.156 \pm 0.052$ | 0.272 | 30 | 11955 |
| 23234 | $-0.175 \pm 0.062$ | 0.368 | 35 | 12069 |
| 12232 | $0.274 \pm 0.081$ | 0.416 | 32 | 12131 |
| 23363 | $-0.219 \pm 0.050$ | 0.225 | 22 | 12173 |
| 16119 | $0.187 \pm 0.075$ | 0.372 | 27 | 12184 |
| 23365 | $-0.240 \pm 0.106$ | 0.511 | 23 | 12218 |
| 23316 | $-0.091 \pm 0.038$ | 0.182 | 26 | 12224 |
| 12157 | $0.197 \pm 0.063$ | 0.303 | 25 | 12277 |
|  |  |  |  |  |

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Fig. A.1. Sample of the $21-\mathrm{cm}$ line profiles.
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MCG -1-2-9












ESO 410-16


ESO 296- 2






Fig. A.1. continued.
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ESO 353- 9


ESO 353- 36




Fig. A.1. continued.
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Fig. A.1. continued.
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Fig. A.1. continued.
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Fig. A.1. continued.
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Fig. A.1. continued.
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Fig. A.1. continued.
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Fig. A.1. continued.
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Fig. A.1. continued.


















Fig. A.1. continued.
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Fig. A.1. continued.
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Fig. A.1. continued.


Fig. A.1. continued.
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Fig. A.1. continued.
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ESO 316-47






ESO 500- 5


Fig. A.1. continued.
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Fig. A.1. continued.
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Fig. A.1. continued.
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Fig. A.1. continued.
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Fig. A.1. continued.
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Fig. A.1. continued.
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Fig. A.1. continued.
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IRAS17306-0927








IRAS $19221+0019$


Fig. A.1. continued.
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Fig. A.1. continued.
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Fig. A.1. continued.


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Fig. A.1. continued.
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Fig. A.1. continued.


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    * Full Tables 4, 5, 6, 7 and Fig. A. 1 are available in electronic form at http://www.edpsciences.org.
    Full Tables 3, 8, A. 1 and A. 3 are available in electronic form at the CDS via anonymous ftp to
    cdsarc.u-strasbg.fr (130.79.128.5) or via
    http://cdsweb.u-strasbg.fr/cgi-bin/qcat?J/A+A/412/57

[^1]:    ${ }^{1}$ We search also for the effect of beam filling on $\log 2 V_{\mathrm{M}}^{\sin i}$. We found a very tiny effect $\delta\left(\log 2 V_{\mathrm{M}}^{\sin i}\right)=(0.0028 \pm 0.0013) \theta-(0.0024 \pm$ $0.0011),(\theta$ is the angular resolution of the radiotelescope, as seen in Sect. 5) that had no practical effect on the final result. Further, this effect disappeared for large galaxies $\left(\log D_{25}>1.5\right)$ where it is supposed to be more significant. We thus decided, for simplicity, to neglect it.

