

# Detection of aerosols in Antarctica from long-range transport of the 2009 australian wildfires

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1	DETECTION OF AEROSOLS IN ANTARCTICA FROM LONG-RANGE
2	TRANSPORT OF THE 2009 AUSTRALIAN WILDFIRES
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4	M. Fromm <sup>4</sup> , and P. Keckhut <sup>1</sup>
5	
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12 13 14 15 16 17 18 19 20 21 22	<ul> <li>Keypoints:</li> <li>1. First lidar identification of smoke aerosols originating from the 2009 wildfires above the French Antarctic Dumont d'Urville station</li> <li>2. A dedicated simple microphysical and transport model uses CALIOP data as constraints to successfully model the aerosol overpass above Antarctica</li> <li>3. Positive ozone anomaly is reported using lidar measurements</li> </ul>
23 24	A series of bushfires ignited in the Australian state of Victoria on February 7th 2009 and ended
25	up being the most devastating fire hazard in Australia before the recent 2019/2020 fires. Active remote
26	sensing monitoring instruments are deployed on the French Antarctic station Dumont d'Urville. For
27	the first time, the station recorded presence of aerosols having originated from this biomass burning
28	event at stratospheric altitudes using the atmospheric laser sounding technique (lidar). We combine
29	model calculations to space-borne and ground-based measurements to track the long-range transport
30	of a small filament of the aerosol plume down to the Antarctic station to highlight the possible global
31 32	impact of such events.
52	

### 35 Abstract

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37 We analyze the long-range transport to high latitudes of a smoke particle filament originating 38 from the extra-tropics plume after the Australian wildfires colloquially known as 'Black Saturday' on 39 February 7<sup>th</sup> 2009 and report the first Antarctic stratospheric lidar characterization of such aerosols. 40 Using a high-resolution transport/microphysical model, we show that the monitoring cloud/aerosol 41 lidar instrument operating at the French Antarctic station Dumont d'Urville (DDU-66°S-140°E) 42 recorded a signature of those aerosols. The 532 nm scattering ratio of this filament is comparable to 43 typical moderate stratospheric volcanic plume, with values between 1.4 and 1.6 on the 1<sup>st</sup> and 3<sup>rd</sup> days 44 of March above DDU station at around the 14 and 16 km altitude respectively.

A dedicated model is described and its ability to track down fine optical signatures is validated against Antarctic lidar elastic aerosol and DIAL ozone measurements. Using one month of tropical CALIOP data to support a relatively simple microphysical scheme, we report modeled aerosol presence above DDU station after advection of the aerosol size distribution. In situ measurements also report associated positive ozone anomaly.

50 This case study provides evidence that biomass burning events injecting significant amounts 51 of material up to stratospheric altitudes can be transported towards high latitudes. We highlight a 52 potential imprint of smoke particles on the Antarctic atmosphere over larger time scales. Any 53 underestimation of the global impact of such deep particle transport will lead to uncertainties in 54 modeling the associated chemical or radiative effects, especially in polar regions where specific 55 microphysical and chemical processes take place.

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#### 58 **1. Introduction**

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60 Aerosols are known to cause significant effects on the radiative balance of the Earth, both 61 directly, through scattering and absorption of short-wave and long-wave radiation [Charlson et al., 62 1992], and indirectly, by acting as condensation nuclei [Penner et al., 2001]. As a matter of fact, long-63 term observations of stratospheric aerosols are needed for understanding the global atmospheric 64 temperature and ozone layer variability [Thomasson and Peter 2006]. The role of particles, being either 65 sulfated background/volcanic aerosols or carbon-based smokes as cloud condensation nuclei, is 66 especially important in Antarctica. Recent studies also tend to demonstrate that polar aerosols can be 67 characterized by large variations in their optical, physical and chemical composition parameters, 68 causing a wide variety of particulate radiative properties throughout the year [Wang et al., 2005; 69 Tomasi et al., 2007]. Besides, Polar Stratospheric Clouds (PSCs) grow from stratospheric sulfate 70 aerosols under given temperature thresholds and have a crucial impact on stratospheric ozone 71 depletion [Adriani et al., 2004; Spinhirne et al., 2005; WMO, 2018], and on radiative balance, which in 72 turn modulate the local circulation [Liou, 1986; Ramanathan et al., 1989; Lachlan-Cope, 2010]. Polar 73 clouds are also involved in the equilibrium of the largest freshwater reservoir in the world [Arthern et 74 al., 2006].

76 Because of the low aerosol number concentration, their direct effect was first thought to be 77 negligible in Antarctic regions [Bodhaine, 1995]. But, above snow- or ice-covered areas, multiple 78 scattering between these particles, cloud surface, and ground surface may still lead to some degree of 79 warming [Randles et al., 2004; Liu et al., 2013]. Aside from the background sulfate aerosols, smoke 80 particles also act as nucleation sites, favoring the growth of sulphuric acid droplets. Smoke particles 81 also have a greater absorption/heating potential than sulfate particles [Burton et al. 2012, Weller et 82 al., 2013]. Therefore, the global effect of a given stratospheric aerosol load has yet to be precisely 83 estimated as it is as highly dependent on latitude as it is on chemical composition.

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85 Worldwide information on stratospheric aerosols is now available from remote sensing 86 instruments, especially with the numerous space missions [SPARC, 2006]. Nevertheless, the Antarctic 87 particle budget remains hard to monitor due to an evident lack of instrumental capabilities beside the 88 sparse ground stations which are mainly located at the edge of the continent. Pollution originating 89 from trans-continental transport is increasing at high latitudes [Stohl and Sodemann, 2010]. Bullard et 90 al. [2016] highlighted the significant impact of dust event and processes at high latitudes. Antarctica is 91 more isolated from human activities than the Arctic and this geographical situation results in a lower 92 density of aerosol particles relative to other regions [Ito, 1986].

93

94 Exchanges between troposphere and stratosphere have important impacts on atmospheric 95 chemistry by mixing air masses containing potent greenhouse gases like ozone and water vapor [Gauss 96 et al., 2003; Forster et al., 2007] and by air masses from lower latitudes entering the Antarctic 97 circulation. Particle sedimentation and aerosol deposition then cause ice cores to act as climate proxies 98 keeping track of past aerosol concentration levels [Wolff and Peel, 1985; McConnell et al., 2007]. Ice 99 core samplings may also contain information on the frequency and intensity of past volcanic eruptions 100 [Sigl et al., 2014]. Around 1.5% of stratospheric air reaches the surface after 10 days which is around 101 the average age of air mass at the ground level in spring [Stohl and Sodemann, 2010]. However, the 102 extent, scales and impact (either volcanic or biomass burning) occurring at lower latitudes on the 103 Antarctic atmosphere through deep transport is still unclear. The particular processes taking place in 104 Antarctica make the assessment of this impact a critical point in the interplaying interactions between 105 ozone and climate [WMO, 2018].

106

107 This case study reports evidence of another source of stratospheric Antarctic aerosols related 108 to biomass burning. Biomass burning may spread across large areas and significantly increase 109 temperature and generate some explosive troposphere to stratosphere transport favored by intense 110 forest fire. Convective processes associated with fire activity are referred to as pyroconvection and 111 under specific heat and moisture conditions, fire-generated aerosol release can lead to thunderstorm-112 like convective cloud known as PyroCumulonimbus (PyroCb). This feature has been reported and 113 investigated, with many events spanning from 1950 up to now [Livesey et al., 2004; Waibel et al., 1999; 114 Fromm et al., 2000, 2005 ; Siebert et al., 2000 ; Fromm and Servranckx, 2003 ; Jost et al., 2004 ; Tupper et al., 2005 ; Fromm et al.; 2006, 2008 ; Peterson et al., 2015; Kablick et al., 2018 ]. One of the 115 116 interesting features of the PyroCb events is that the microphysics and dynamics of associated plumes 117 are most often very close to characteristics observed in the case of volcanic eruption [Peterson et al., 118 2018]. Modeling studies have shown that tropospheric plumes above thunderstorms are very likely to 119 enter the lower stratosphere, from which they may reach latitudes far removed from their emission 120 location [Wang, 2003]. In the case of an underestimated frequency of such events, this process could 121 be considered a new alternative to efficiently redistribute gases throughout the Upper Troposphere / 122 Lower Stratosphere range (UTLS). While it is an efficient source process to inject particles up into the 123 stratosphere, further studies are still required to better identify the convective storms responsible for 124 this injection, with both observations and modeling required [Fromm et al., 2005]. In terms of optical 125 measurement analysis, this also questions any misleading attribution of such carbonated aerosol 126 signatures to volcanic activity. Discrimination of aerosols as of either volcanic or biomass burning origin 127 requires microphysical information that can be provided by highly sensitive measurement capabilities 128 such as those of lidar instruments.

129

130 The transport of material over long distances, and especially smoke, is not often directly 131 observed in the Antarctic UTLS. A few individual events have been documented, such as the transport 132 of dust from Patagonia [Gassó et al., 2010] or smoke from biomass burning in South America 133 [Evangelista et al., 2007; Pereira et al., 2006] to the tip of the Antarctic Peninsula. Fiebig et al. [2009] 134 revealed transport of biomass burning aerosol from South America to an inland site on Queen Maud 135 Land. Statistical studies of transport down to Antarctica have also been performed using trajectory 136 calculations for single measurement stations or ice core drilling sites [Kottmeier and Fay, 1998]. Weller 137 et al. [2013] investigated seasonal variations of black-carbon particles and albedo assessment with a 138 maximum variability in spring time from decadal time series using absorption photometers. However, 139 microphysical modeling and transport of particles from the lower latitudes down to Antarctica remains 140 difficult due to scale issues within the models: large scale is needed for longer runs and small scale 141 resolution needed to resolve the numerous fine filaments. Models have often been used in conjunction 142 with ground-based measurements to attribute pollution in Antarctica to its sources [Stohl and 143 Sodemann, 2010]. As for ground-based detection, the lidar instrument still remains one of the best 144 suited to monitor and detect small changes in particle load in the UTLS, with a vertical resolution often 145 under 100m and very high sampling frequency.

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147 This paper reports observations along with microphysical and transport simulation of a specific 148 smoke filament originating from the particle plume which was largely confined within the tropics one 149 month after the Australian bushfires. The overall complex structure of the different plume layers has 150 been used as an additional constraint in our model both upon initialization and during the simulation 151 we present here. We describe a methodology based on a combination of local and global lidar 152 measurements along with a simple coupling of optical and transport modeling. We aim at taking 153 advantage of local measurements to efficiently report on similar fine events (or those involving long 154 range transport of volcanic aerosols). We also highlight the importance in having ground reference 155 monitoring measurement stations to complement and validate global measurements, model results 156 or even to some extent microphysical schemes through case studies like this one. Section 2 provides 157 additional information on the event. Observations and numerical tools are described in section 3. 158 Section 4 details the transport method and simulation setup. Section 5 discusses the model results 159 against observations before drawing conclusions in section 6.

163

## 2. The 2009 Australian Bushfires

#### 2.1. Event Description

164 In the Southern Hemisphere, biomass burning activity is minimal in February. The major fire 165 sources of Indonesia, Africa and South America become significant later in the year [Edwards et al., 166 2006]. Particle transport from these areas has a strong seasonality with maxima in winter for Southern 167 Africa and spring from South America and Australia [Stohl and Sodemann, 2010]. A series of bushfires 168 ignited in the Australian state of Victoria on February 7<sup>th</sup> 2009 and ended up being the most 169 devastating fire hazard in Australia up to the recent 2019/2020 event and are now referred to as the 170 'Black Saturday' bushfires. Over 4500 km<sup>2</sup> of land burned to ashes, and more than 170 fatalities were 171 recorded with major fire spots identified over 2 weeks from the Black Saturday [BOM, 2009]. The 172 intensity of the Victorian bushfires was such that pyroconvection uplifted gas species along with 173 particles up to stratospheric altitudes. Such typical transport of gas and particles has been reported by 174 Wang [2003], Fromm et al. [2005] and Field et al. [2016]. On 'Black Saturday', the combination of typical 175 meteorological synoptic features with lack of precipitation and localized convective processes rose the 176 fire danger level to extreme and several PyroCb were identified [Reeder et al., 2015; Tosca et al., 2015; 177 Field et al., 2016]. Recent investigations of the responses between the fires and the atmosphere lead 178 to a better understanding of the pyroconvection processes and especially the risk of new fire ignitions 179 using pyrogenic lightning. Dowdy et al. [2017] gathered evidence of pyroconvection in the Black 180 Saturday events using a number of distinct electrified PyroCb clusters, therefore retro-validating the 181 studies on the deep imprint of these plumes on the Southern Hemisphere atmosphere during several 182 months [Siddaway and Peletina, 2011; Pumphrey et al., 2011; Glatthor et al., 2013].

183

184 The main smoke plume was first observed at stratospheric altitudes the day after the outbreak 185 and encircled the globe approximately 6 weeks after being first detected [Pumphrey et al., 2011]. 186 Similar to typical volcanic behaviour, smoke aerosols remained in the extratropical channel. From mid-187 February to June, it was mostly located between 18-22 km altitudes (see Section 5). Space-borne 188 observations acquired onboard the CALIPSO satellite reported both the main part of the plume within 189 tropical latitudes and various filaments reaching higher latitudes but backscatter values after 2 or 3 190 months often reached the lower detection limit of the lidar instrument. Siddaway and Peletina [2011] 191 used OSIRIS observations to report a predominant westward direction in the smoke plume transport 192 and characterized advections from around 19 to 22km up before the end of April, and background 193 radiance levels on OSIRIS observations by mid-June 2009.

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#### **3. Observations**

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#### 3.1. Ground-based lidar

Lidar observations have been carried out at the French station Dumont d'Urville [DDU] (66°S -140°E), since April 1989 within the framework of a French (LATMOS/IPSL) Italian (ISAC/CNR) cooperation. The station is a primary site of the Network for Detection of Atmospheric Composition Changes (NDACC). Except for year 2000 and 2005, aerosols and PSCs measurements have been routinely carried out at DDU every year since 1989, from roughly mid-March to late September 202 [David et al., 1998, 2005] providing the only continuous aerosol lidar time series over 30 years in 203 Antarctica. The DDU Rayleigh/Mie lidar was originally designed to operate in two exclusive modes, one 204 for ozone and one for aerosol/PSC and temperature measurements. The ozone mode includes light 205 emission and reception at wavelengths of 308 and 355 nm and additional reception at 332 and 387 nm. 206 The aerosol mode includes emission and reception at wavelengths of 532 and 1064 nm. Lidar 207 measurements are performed at DDU on a yearly winter campaign basis; the main focus being on polar 208 stratospheric clouds monitoring during the winter and early spring months. For calibration purposes 209 and additional aerosol records, a lighter measurement calendar also plans lidar acquisitions from early 210 March, when the sky gets dark enough for the lidar to safely operate for extended measuring sessions. 211 Within the operational frame of the NDACC, some alerts regarding specific events like volcanic 212 eruption and specific air mass intrusions may trigger additional acquisitions.

213

214 A complete description of both the instrumental design and inversion procedure is featured in 215 [David et al., 2012]. The Nd:YAG laser source emits 500mJ per pulse at 1064nm at 10 Hz frequency, 216 doubled in frequency to 532 nm. A collocated Excimer laser emits 55 mJ per pulse at 308 nm at 80 Hz 217 frequency. Backscattered signal is collected by an 80 cm telescope. Aerosol vertical density profile is 218 defined as the backscattering ratio or Scattering Ratio (hereafter called SR) expressed as the ratio of 219 extra scattering of aerosols to the molecular scattering of air at the same altitude. After a 220 preprocessing phase removing potential saturation effects and background noise level (low at this time 221 of the year in Antarctica), lidar signals are inverted using the Klett-Fernald formalism [Klett 1981,1985; 222 Fernald, 1984], to derive individual SR profiles. The zero-reference altitude is set between 28 and 223 32 km depending on signal-to-noise ratio. The extinction-to-backscatter ratio is set to 50 sr which is a 224 commonly assumed value for small aerosol loading and even periods of moderate volcanic eruptions 225 [Ridley et al., 2014]. Cumulative uncertainties of the backscatter measurements induced by random 226 detection processes, possible presence of aerosol at the reference altitude and the error in lidar ratio 227 value do not exceed 10% as reported by Chazette et al., [1995]. Another major source of uncertainty 228 is the molecular number density used in Rayleigh calculations derived from atmospheric pressure and 229 temperature profiles. The lidar inversion is particularly sensitive to the molecular density and the clear-230 air reference altitude. Collocated radiosondes are used to estimate reference atmospheric density and 231 interpolation and scaling to National Centers for Numerical Prediction (NCEP) daily meteorological 232 data is done at balloon burst altitude.

233

Ozone profiles are retrieved using the DIfferential Absorption Lidar (DIAL) technique. The ozone number density can be derived from two signals, one being strongly absorbed by ozone (308 nm), the other (355 nm) being only weakly absorbed. The Nd:YAG frequency is tripled to 355 nm to provide the reference signal whereas the Excimer laser directly emits at the 308 nm. Pre-processing of the lidar signal in ozone mode is similar to the aerosol mode. Ozone concentration is retrieved without external assumptions by comparing the number of laser shots between the absorbed and reference wavelengths [David et al., 2012].

241

242To complement the simulations presented in this paper, we also include a lidar profile acquired243at the Australian Davis station (-68.6°S, 78°E). The 532nm SR is obtained with the Klett-Fernald244formalism after normalization at 30-35km altitude. The molecular signal is derived from density profile

inferred from measurements by the AIRS instrument onboard the Aqua satellite. As for DDU, the lidar ratio of the background aerosol is assumed to be 50 sr. For observations used here, a mechanical chopper reduced signal levels below approximately 11 km altitude. Further details of the system and analysis procedure are provided by Alexander et al. [2011]. Observations were made with the Davis lidar system throughout the autumn during quiescent tropospheric weather, however only data collected on February 26<sup>th</sup> are relevant for comparisons with the simulations discussed below.

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#### **3.2. Satellite aerosol sounders**

253 The CALIOP (Cloud-Aerosol Lidar with Orthogonal Polarization) instrument onboard the 254 CALIPSO satellite platform is a nadir-viewing active sounder [Winker et al., 2010]. Operational since 255 June 2006, CALIOP provides range-resolved measurements of elastic backscatter at 532 nm and 1064 256 nm with a vertical resolution of around 200 m at stratospheric altitudes. CALIOP lidar makes use of a 257 Nd:YAG laser operating at 20.2 Hz with a 110 mJ/pulse power and a 0.78 m<sup>2</sup> telescope. In this study 258 we use data based on the night-time 532 nm level 2 product, post-processed using a treatment 259 described by Vernier et al. [2009]. We assume backscatter measurement uncertainties around 20% as 260 stated in Winker et al. [2010]. We also use the level 2 Cloud-Aerosol discrimination score to 261 characterize the smoke particles. The CALIOP 532 nm SR is obtained after applying an extinction 262 correction on the attenuated backscatter with a generic non-background regime stratospheric aerosol 263 lidar ratio of 50 sr. We focus on UTLS altitudes, where attenuation is lower than in the lower 264 troposphere, especially from a space-borne acquisition geometry. The molecular atmospheric profile 265 is derived from temperature and pressure provided by Goddard Earth Observing System [Yasunari et 266 al., 2011].

267

268 The Atmospheric Chemistry Experiment (ACE) was launched in August 2003 onboard the 269 Canadian scientific satellite SCISAT-I. The high-resolution spectrometer measures during sunrise and 270 sunset infrared absorption to provide vertical profiles of atmospheric compounds (gases and particles). 271 Aerosols and clouds are monitored using the solar radiation extinction measured by two filtered 272 imagers at 527 (VIS) and 1020 nm (NIR) [Dodion et al., 2007]. The satellite operates in a circular orbit 273 at an altitude of 650 km. Aerosol extinction profiles are obtained by subtracting the gas extinction 274 contribution (air, O<sub>3</sub>, NO<sub>2</sub>) derived from the Fourier Transform spectrometer from the total 275 atmospheric extinction profiles. On average, the NIR ACE imager gives results that are 20% higher than 276 SAGE II instrument [Gilbert et al., 2007].

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### 4. Numerical simulations

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### 4.1 The aerosol microphysical and transport model

The Modèle Isentropique de transport Mésoéchelle de l'Ozone Stratosphérique par Advection (MIMOSA) model [Hauchecorne et al., 2001] is originally a Potential Vorticity (PV) advection model running on isentropic surfaces (surfaces of constant potential temperature). The advection scheme is semi-Lagrangian with a time step of 1 hour. Regridding onto the original orthonormal grid is performed every 6 hours. The model resolution is 0.5°x0.5°. The advection is driven by European Centre for Medium Range Weather Forecast (ECMWF) meteorological analyses at a resolution of 0.5°x0.5°. The accuracy and mass conservation of the model have been evaluated by Hauchecorne et al. [2001] and validated against airborne lidar ozone measurements using the strong correlation between PV and
 ozone, a quasi-conserved chemical tracer on a week timescale within most of the lower stratosphere
 [Heese et al., 2001].

291

292 A sectional aerosol microphysical module has been coupled to MIMOSA to simulate the 293 formation, transport and microphysical evolution of stratospheric aerosol particles [Jumelet et al., 294 2009]. The size distribution is discretized into a number of size bins using a geometrically increasing 295 volume scale. The number of size bins is 50 with the bin particle radius ranging from 0.01 to 5  $\mu$ m. The 296 model prognostic variables are the number density of each size bin and chemical composition. The 297 module also includes an optical routine that calculates theoretical properties such as volume 298 backscatter and/or extinction coefficients at given wavelengths (in the visible and near-infrared 299 domains e.g. the 532 nm aerosol wavelength) from the model-simulated particle size distribution and 300 the composition-dependent refractive index.

301

302 We consider here two different aerosol types within the model, one being the background 303 sulfate stratospheric aerosols and the other one being the injected smoke particles. Since background 304 sulfate aerosols are not relevant to our long-range transport case study, only the results for smoke 305 particles are presented and discussed here. In order to simulate smoke particles, the original aerosol 306 scheme is extended. Smoke particles can be complex objects, i.e. a homogeneous mixing of organic 307 and inorganic aerosol components with black carbon coating in some more or less complex core/shell 308 structure often modelled using fractal aggregates [Smith and Grainger, 2014]. However, our approach 309 is simple because the objective is to calculate the aerosol optical parameters (properties integrated 310 over the entire distribution) instead of the detailed behaviour of the size distribution and the 532nm 311 SR will be driven by observations. Dählkotter et al. [2014] have shown that the optical properties of 312 aged forest fire smokes (having experienced intercontinental transport) can be estimated with Mie 313 Theory without any major uncertainty. The smoke particles are treated here as a mixture of 314 sulfate/black carbon with overall real refractive index at 1.85 [Smith and Grainger, 2014, Mischenko et 315 al., 2014]. The smoke SR is calculated using the Mie scattering code with aspherical aspect ratio at 1.50 316 [Scarnato et al., 2015]. There are still computational uncertainties in the optical calculations due to the 317 limited number of bins covering the size range that drive the overall uncertainty estimated around 20% 318 on the backscatter coefficient.

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### 4.2 Initialisation & CALIOP SR insertion into the model

321

322 The initialization of aerosol size distributions is based on measurements. There was no 323 significant volcanic influence on aerosol loading at the time of the Australian Bushfires in the southern 324 hemisphere. For this reason, the sulfuric acid aerosol size distribution was initialized over the entire 325 model domain with a typical background (non-volcanic) lognormal distribution:  $N_0=10$  cm<sup>-3</sup>, 326  $r_m$ =0.07  $\mu$ m and  $\sigma$ =1.8 [Pinnick et al., 1976; Larsen et al., 2000]. To our best knowledge, no information 327 on the Australian smoke particle size distribution are available. The smoke size distribution is therefore 328 initialized with a lognormal distribution based on data collected in the northern hemisphere on a 329 comparable PyroCb particle plume after a similar event of smoke long-range transport (N=50 cm<sup>-3</sup>,  $r_m$ =0.12 µm and  $\sigma$ =1.3) [Dahlkötter et al., 2014]. This size distribution can also be typical of the fine mode of volcanic aerosol plumes.

332

333 The initialization of the stratospheric aerosol spatial distribution is based on the insertion of 334 CALIOP SR measurements in the model. The insertion is straightforward as compared to other 335 approaches where CALIOP backscatter products have already been successfully assimilated in several 336 studies using full assimilation schemes [Sekiyama et al., 2010; Campbell et al., 2010]. From a data 337 assimilation point of view, our insertion scheme is equivalent to a sequential assimilation assuming 338 model errors to be much greater than observational errors and using our optical model as 339 observational operator. Our scheme remains local because the insertion of one single observation only 340 impacts one model grid cell, meaning the length-scale of the spatial correlation is taken to be less that 341 a size of a grid cell, i.e. less than 0.5°.

342

343 The insertion procedure is the following: one variable of the aerosol size distribution 344 parameter (the so-called control variable) at each grid cell is modified iteratively in order for the model-345 calculated SR to match the CALIOP-observed SR. The choice of the control variable depends on the 346 value of the observation, controlling the aerosol type. When CALIOP SR is smaller than the background 347 value (taken as 1.025, from simple Mie Calculations on H<sub>2</sub>SO<sub>4</sub>/H<sub>2</sub>O droplets), one can assume that the 348 measurements correspond to stratospheric sulfuric acid aerosols and the control variable is the total 349 aerosol concentration N. The median radius and standard deviation remain unchanged ( $r_m$ =0.12  $\mu m$ 350 and  $\sigma$ =1.3). Note that, in that case, the optical module uses the aerosol optical index of sulfuric 351 acid/water solution in equilibrium with the gas phase [Jumelet et al., 2009]. When CALIOP SR is greater 352 than 1.025, the measurement is taken as indicative of the presence of smoke particles and the control 353 variable is the mode radius. N and standard deviation remain unchanged. In that case, the optical 354 module uses the smoke refractive index. This choice of control variable for smoke type is based on the 355 fact that there is little information on smoke particle size and, within the range of smoke size 356 distribution parameters, model-calculated SR is assumed to be much more sensitive to the mode 357 radius than to the total number concentration, the latter variable displaying a linear behavior regarding 358 the aerosol backscatter coefficient.

359

The first 2 weeks of the simulation in January, before the fire outbreak, define the spin-up phase. In this phase, continuous high-frequency CALIOP SR insertion on the whole grid slowly shifts modelled stratospheric aerosol SR towards CALIOP measurements. After the spin-up, the average aerosol levels do not change much beyond sedimentation, and the model reaches some degree of equilibrium regarding the forcing of satellite observations.

- 365 366
- 4.3 Model set-up and simulations
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A key aim of the present study is to test the model ability to forecast the position of plumes of either volcanic or biomass origin. For this reason, simulating these long-range transport events use aerosols as tracers. After the initialization/spin-up phase, CALIOP SR data are inserted in the model during the run, but only outside a large free-running model domain. The CALIOP insertion carries on in order to maintain realistic boundary conditions on this free running domain. The sensitivity of the model results has been tested by varying the latitude of the free running domain boundary beyond which the aerosol particles are freely advected (no CALIOP SR insertion). Stable results were obtained when the CALIOP insertion is confined to the region equatorward of 50°S latitude. Of course, the CALIOP insertion domain covers the Australian source region. We keep the latitude of the limit between the free running insertion zones relatively low to highlight the efficiency of the transport model.

- The model is integrated from January 1<sup>st</sup> to May 1<sup>st</sup>, with an advection step of 1 hour, and a regridded output is dumped every 6 hours. The model is run at several isentropic levels every 10 K over the 400-450 K potential temperature range on a southern hemispheric grid. It is forced with ECMWF temperature and pressure fields at a 0.5°x0.5° resolution. The spin-up phase covers the first 2 weeks of January. After that, CALIOP SR data are only inserted outside the large free-running model domain, equatorwards of 50°S. The signature of the Australian fire plumes at stratospheric altitudes is clearly detected after 2 days after the outburst on February 7<sup>th</sup> above the 410 K level.
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- **5. Results**
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#### 5.1 CALIOP and lidar observations

392 Two weeks of 532 nm CALIOP SR averaged data are presented on Figure 1. The figure shows 393 the evolution of smoke plume distribution between February 1<sup>st</sup> and 16<sup>th</sup>. Recall that the fires started 394 on February 7<sup>th</sup>. To avoid any false detection, only SR above 1.1 (clearly indicative in this study of smoke 395 particles) are considered. The SR averaged over the 15-25 km altitude range show the plume in the 396 South Pacific (Figure 1a). The SR vertical distribution as a function of longitude (data averaged over the 397 25°S-50°S latitudinal band) also confirm the presence of smoke particles in the 15-22 km altitude range 398 (cf. Figure 1b). Figure 1c shows the zonal mean SR distribution (data averaged along latitude circles) 399 with isentropic levels superimposed; smoke particles remain mostly confined between the 380 K and 400 450 K isentropic levels during the first 2 weeks of February, although an analysis at higher time 401 resolution revealed that few smoke particles can be detected up to the 500 K level a few days after the 402 fire outbreak (not shown).

403

404 The plume is identified above DDU station during a short time span in early March, with 405 unambiguous signatures on the 1<sup>st</sup> and 3<sup>rd</sup>. At that time, no direct comparison with CALIOP is available. 406 CALIOP data and results of model simulations (described in the next section) indicate more overpasses 407 of the smoke plume above DDU, but no lidar measurements were available due to bad local weather 408 conditions. Both vertical ground-based SR profiles obtained on the 1<sup>st</sup> (Figure 2a) and 3<sup>rd</sup> of March 409 (Figure 2b) reveal clear signals of smoke particles at 13-14 and 16-17 km respectively. SR peaks at 410 K 410 and 450 K on the March 1<sup>st</sup> (Figure 2a) and March 3<sup>rd</sup> respectively. Note that the lidar data are averaged 411 over ~160 min on March 1<sup>st</sup> and ~210 min on March 3<sup>rd</sup>. SR peaks above 1.4 in both cases which is 412 higher than SR signals of minor volcanic eruptions [Vernier et al., 2011]. An additional weaker layer 413 can also be seen just above 15km (Figure 2a). It is worth pointing out that measurements at this time 414 of the year (at the end of continuous Antarctic daylight) were performed for calibration purposes in 415 March, related to nighttime sky conditions.

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#### 5.2 Model simulations versus DDU lidar detection

Figure 3a shows the plume filament on February 26<sup>th</sup>, a few days prior detection above DDU on March 3<sup>rd</sup>, near the location of the Davis station. Figure 3b shows the associated lidar profile on the same day, with a broad but weak peak of about 1.05 around 20km. This small SR value may be indicative of background aerosols, especially as 2009 is a volcanic quiescent time period. Besides, lidar measurements have been integrated on several hours. Even if the plume position is at the Davis

longitude, identification is not clear looking at the lidar profile, especially considering that for small SR
values, the molecular scattering derived from the pressure and temperature profiles is a sensitive
parameter. Still, aerosol loading is present and there is a strong agreement between Figure 3a and 3b,
in that the simulated plume does not reach Davis with scattering ratios greater than 1.1.

428

429 Figure 4 shows modeled SR fields between February 10<sup>th</sup> (Figure 4a) and March 10<sup>th</sup> (Figure 4i) 430 on the 410 K level corresponding to the altitude of the ground-based SR peak on March 1<sup>st</sup>. The highest 431 CALIOP SR values are detected during the first half of February, with values peaking slightly above 1.4 432 and gradually decreasing below the 1.1 threshold during the first half of April (not shown). A weak 433 plume spanning both sides of South America is visible in the modeled SR field in February. This is not 434 related to the Australian fires and will not be discussed here. A week after the fire outbreak, a small 435 patch of high SR values appears on the 10<sup>th</sup> of February east of Australia (Figure 4a). This patch, 436 indicative of the smoke plume, carries on being transported eastward. At the same time, it is very 437 strongly stretched by wind shear, forming a large filament extending over thousands of kilometers. 438 Filamentary structures quickly form at the subtropical barriers. Between the 20<sup>th</sup> and 25<sup>th</sup> of February, 439 the plume filament split into two main parts with the main part remaining in the tropics and a smaller 440 filament being transported towards higher latitudes.

441

Around the time of the fires and during the ascension between Feb 8<sup>th</sup> and 16<sup>th</sup>, part of the CALIOP data may have been misflagged so the overall aerosol load near the tropics may not be accurately depicted. This issue is still out of focus as we aim at identifying the small filament transported at high latitudes and clearly separated from the main bulk (visible from Figure 4d). Besides, the consistency of SR values (excluding high cloud scattering) has been checked prior to integration within the model.

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449 Overall, several smoke plume filaments were observed over a wide altitude range (13-22km) 450 with CALIOP till the end of May. Even though the bulk of the smoke plume did not cross the subtropical 451 barrier, smaller filaments occurred throughout the extratropics due to efficient quasi-horizontal 452 transport (along isentropic surfaces) in the lower stratosphere. After one month of long-range 453 transport, the plume is detected above DDU on the 1<sup>st</sup> of March during the time of the measurements, 454 between 13:33 and 16:12 UT (see Figure 2). The model indicates the passing of a thin filament just 455 above the station on the same day (see Figure 4f); model outputs are only available every 6 hours. In 456 order to investigate the passing of the filaments above DDU, Figure 5 displays a close-up of the 457 simulation at 410 K and 450 K, respectively on March 1<sup>st</sup> and 3<sup>rd</sup> when the DDU lidar clearly detected

- the plume. There is a reasonably good agreement in terms of timing. The ground-based measurements in Figure 1a and 1b identified 2 layers peaking at 410 K and 450 K. The lower layer is detected on March 1<sup>st</sup>, the higher layer is detected on March 3<sup>rd</sup>. This is consistent with the model simulation which shows a DDU overpass at 410 K level only on March 1<sup>st</sup> (Figure 5c) but not on March 3<sup>rd</sup> (Figure 5d); it is very close to the time of the afternoon lidar plume measurements whereas the model output is set at 12:00 UT. The model simulation also appears to be in agreement with the DDU lidar at 450 K level with a DDU overpass only on March 3<sup>rd</sup> (Figure 5b) but not on March 1<sup>rst</sup> (Figure 5a).
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466 As expected, the modeled SR values are in excellent agreement with CALIOP SR. However, the 467 peak modeled SR is slightly higher than the corresponding maximum values of the DDU ground-based 468 lidar. The instantaneous model outputs (every 6 hrs) show the maximum SR values within the smoke 469 filament whereas the ground-based lidar data correspond to around a 3 hr time average of the passing 470 filament as seen from a local atmospheric column above the station. Model fields show that there are 471 sharp SR gradients within the plume. It is unlikely for the ground-based instrument to probe only within 472 the most dense parts of the filament. Furthermore, space-borne and ground-based lidars have 473 different acquisition fields of view and different data processing (calibration, inversion, spatio-474 temporal averaging), and small discrepancies cannot be ruled out. In view of these caveats, the 475 agreement between the ground-based lidar measurements and the CALIOP-driven model simulation 476 can be viewed as good. It indicates that the setup is able to accurately model long-range transport of 477 the information provided by CALIOP measurements from 50°S (the edge of the free-running model 478 domain) down to Antarctica. This also suggests that the transport scheme and the ERA-Interim winds 479 driving the advection are accurate enough if the model is used as a forecasting tool. After the DDU 480 overpass, during the first 10 days of March, the plume is stretched and breaks into smaller parts that 481 are scattered throughout the extratropics and high latitudes.

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#### 5.3 Model simulations versus ACE measurements

485 ACE imager measurements of extinction ratio on March 4<sup>th</sup> were used as a comparison to the 486 modeled SR field. Figure 6. displays the model SR with crosshairs at the location of the ACE 487 measurements. Overall, SR values peak at slightly above 1.6 as forced by CALIOP measurements. The 488 main bulk of the plume remains confined within tropical latitudes, even though the largest SR values 489 are found in a filament crossing out of the tropics at the square location. Figure 7 displays ACE imager 490 measurements of extinction ratio on the same March 4<sup>th</sup>, when the plume has been clearly identified, 491 peaking respectively at 19 and 21 km. Figure 7b displays associated VIS and NIR imager extinction 492 coefficients. In early March, the total AOD fell to around 0.002 and 0.006 for the NIR and VIS 493 respectively whereas it was estimated around 0.03 in February. Once again, the SR values on Figure 6 494 follow the shape of very fine dynamic structures that can be difficult to model without dedicated 495 microphysics and resolution but a large amount of particles seems to be able to cross the tropics. 496 Burton et al. [2012] provides an assessment of the 532 nm extinction to scattering coefficient (lidar 497 ratio) for smoke particles of around 60-80 sr and associated IR/Visible color ratio of around 1.5-3. Using 498 reference stratospheric values of temperature and pressure around 200-210K and 80-100 hPa 499 respectively, we calculate a Rayleigh contribution to the scattering ratio of around 1.7-1.9E-4 km<sup>-1</sup>.sr<sup>-</sup> 500 <sup>1</sup> and using the 60-80 sr lidar ratio, we derive a 1.3-1.5 532 nm scattering ratio from the ACE measurements of Figure 7b which is in good agreement with the modelled 1.4-1.6 values of Figure 6.
 This agreement between ACE measurements of Figure 6 and 7 suggests that present plume modeling
 is accurate for identifying particles at high latitudes.

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#### 5.4 Local impact of the smoke plume on stratospheric ozone

506 The DDU lidar was also used for measurements of ozone vertical profiles up to 2012. Such 507 measurements are available on the 3<sup>rd</sup> of March, but not on the 1<sup>st</sup> of March. Figure 8a presents the 508 lidar ozone profile measured on the same day, a few hours after the measurement of the aerosol 509 profile. The dotted black line indicates the 3-year (2006-2008) average lidar ozone profile for this same 510 period of the year. This background averaged profile agrees with a typical zonal mean climatology 511 [Bencherif et al., 2011] for the relevant latitudinal range [75°S-85°S] which is indicated by the dotted 512 red line. The interesting feature is the presence of a 1 km-thick enhanced ozone layer at about the 513 altitude of the smoke layer.

514

515 However, similar ozone enhancements have been reported in smoke layers in previous 516 studies. Fromm et al. [2005] studied a PyroCb injection of smoke to the stratosphere in August 1998 517 above Northwestern Canada using balloon profiles. They reported the first highly enhanced ozone 518 concentration at the smoke altitude. These reports remain rare because simultaneous smoke and 519 ozone measurements in the stratosphere are not common. It would be worthwhile to consider more 520 cases to establish unequivocally a link between smoke and ozone enhancements in the stratosphere. 521 It has been shown in tropospheric biomass burning plumes that ozone production can be expected 522 and is observed [Pickering et al., 1996; Andreae et al., 2001] but at stratospheric altitudes, no clear 523 process is identified to our knowledge to explain this positive correlation between ozone and smoke 524 particles. Yu et al. [2019] used photochemistry modeling to assess the impact of smokes on  $O_3$  on the 525 Canadian 2017 fire event, only attributing negative anomalies observed by satellite to transport of air 526 masses from the troposphere, with no in situ chemistry.

527

528 Figure 9 shows the ozone mixing ratio on the 395 K and 430 K isentropes from ERA-Interim. A 529 filament of ozone-rich air from higher latitude appears to be lying close to DDU. It is possible that the 530 ozone feature, which is mainly below and inside the smoke layer, could simply show the demarcation 531 between air that has originated from further south (having higher ozone, and being below the smoke 532 layer) from air that is more representative of the zonal average. Further studies combining 533 measurements and chemistry-transport model on the air mass composition inside and outside the 534 plume would be needed to state on this positive correlation because different mechanisms could lead 535 to opposite ozone variations (like decrease from photolysis on water vapor and increase from 536 CO concentrations). Besides, investigations using the observations like ACE or OSIRIS at high latitudes 537 is difficult as the extinction coefficients associated to these aerosols are relatively small (compare with 538 Figure 7).

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#### 6. Summary and concluding remarks

543 An event of stratospheric biomass burning aerosols was detected in the lower stratosphere on 544 the coast of Antarctica using lidar measurements at the end of the Southern Hemisphere summer. Our 545 investigations show that these aerosols originate from the 2009 Australian 'Black Saturday' smoke 546 plume. From this detection, we develop a transport/microphysic modeling tool taking advantage of 547 space-borne measurements to be a useful complement to the groundbased measurements in a region 548 where instrumental capabilities are low as compared to lower latitudes. Due to this low instrumental 549 coverage and fine optical signature of such events, we therefore highlight the underestimated 550 occurrences of such polar detections in the many reported PyroCb events, and therefore link the 551 PyroCb impacts to scientific questions more specific to high latitude regions.

552

Our global simulations of the microphysical/transport model indicated that following the 553 554 February 7<sup>th</sup> Black Saturday event, particles reached the stratosphere within 2 days after the fire 555 outbreak and crossed the tropical dynamical barriers down to the high latitudes in less than a week. 556 Lidar measurements are routinely performed from the early days of March at the French Dumont 557 d'Urville station, around one month after the fires in Australia. We identified a clear optical signature. 558 An average scattering ratio slightly above 1.4 is identified at the 532nm wavelength, which is 559 comparable to the optical response of moderate volcanic plume after several months of stratospheric 560 transport.

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562 A first step towards a more in-depth parametrization of the different particle types present in 563 the UTLS has been developed and applied to this case study. The modeling approach allows for high 564 resolution hemispheric transport of a smoke plume proxy to which optical properties are fit at the 565 microphysical level to the optical properties measured by the CALIOP space-borne lidar. The dedicated 566 MIMOSA- $\mu\phi$  model was run over a 5-month time period, as the overall global signature of the fire 567 plume is reported up to June.

568

569 Without implementation of a full data assimilation scheme, this simulation is built on a 2-step 570 procedure. A first spin-up step where the model grid is filled with particles, being essentially 571 background sulfate aerosols with an initial carbonated smoke cluster, using the CALIOP 532nm SR as 572 the variable controlling the evolution of the bin-resolved size distribution from its reference 573 background parameters. This first step runs through the first month of the run time and leads to a 574 realistic pattern of the global particle stratospheric load on the southern hemisphere. In the second 575 step, the model runs on both a forced and free regime, considering the 50°S parallel as the latitude 576 limit beyond which the advection is freely performed. After more than one month of long-range 577 transport, our proxy of the smoke plume is found to be at a time and place globally consistent with the 578 lidar measurements on both March the 1<sup>st</sup> and 3<sup>rd</sup>. This simulation highlights the opportunity for air 579 masses coming from lower latitudes to deeply imprint the Antarctic stratosphere, as only some polar 580 cloud types and similar transport of any strong volcanic plume would feature similar SR values. The 581 quantitative assessment of the carbonated aerosol mass and smoke composition will need further 582 microphysical investigations and many improvements on the model are necessary. 583

584 This work reports to our knowledge the first lidar identification of stratospheric smoke aerosols, 585 the observed signature observed at the French Dumont d'Urville station originates from the 2009 586 wildfire event. We highlight the benefits of an efficient yet simple model-instrumental coupling. 587 Accurate plume characterization can be made from a small sample of measurements and relatively 588 simple modeling tools. Thus, we see three different outlines in this study: first, the scientific and 589 confirmed pathway to a global impact (yet to be estimated) of long-range troposphere to stratosphere 590 transport through pyroconvection and deep stratospheric transport for smoke. Second, the need for 591 better characterization of the microphysical properties of such smoke, as their trapping above the 592 Antarctic area is likely to have an impact on the ground level once deposited on ice shelves, especially 593 over longer time scales. Third, the needed effort towards a better coupling of models and observations 594 at the global scale. This can be achieved through assimilation of diagnostic variables such as the 595 scattering ratio, especially when considering the overall success of space-borne lidar instruments and 596 the need for maintaining routine lidar operations in polar stations.

597

598 In this study few compromises are made with the model grid size, resolution, and level of 599 microphysical detail. This highlights the importance of dedicated data assimilation tools. Satellite 600 measurements combined with global and regional models can be an efficient complementary data 601 source to document transport of fine volcanic or biomass burning plumes, especially considering the 602 2019/2020 Australian wildfire event, which strength and amplitude is already acknowledged to be 603 much greater than the 2009 event.

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608

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616 The MIMOSA transport model outputs are available online at the french ESPRI data archive :

617 https://cds-espri.ipsl.upmc.fr/etherTypo/index.php?id=1663&L=1).

- 618 Measurements from the ESA SCISAT-1/ACE instrument are available online at:
- 619 https://databace.scisat.ca/level2/
- 620 621 622 623 Measurements from the CALIPSO/CALIOP space-borne lidar are available online at:
- https://www-calipso.larc.nasa.gov/tools/data avail/
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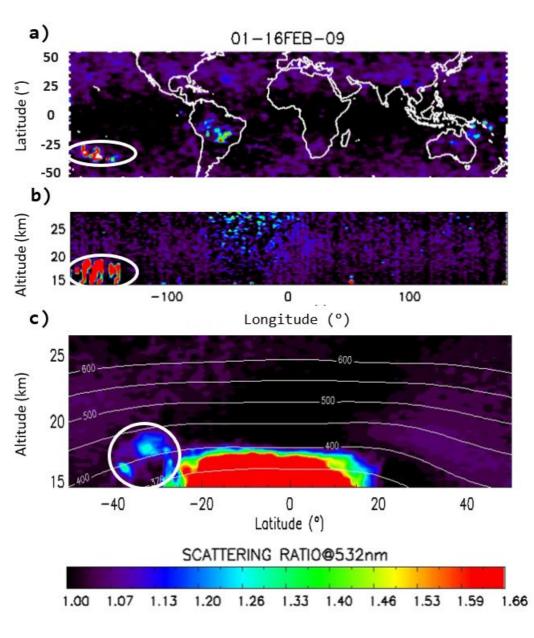
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**Fig. 1.** 532nm CALIOP Scattering Ratio in the first 16 days of February plotted as longitude vs latitude, 15-20 km average (a), longitude vs altitude, averaged in the -25° to -50° latitudinal band (b), and latitude vs altitude (c). The smoke plume cluster is encircled in white and is clearly separated from the global tropospheric cloud signature. Contours of constant potential temperature (K) are shown.

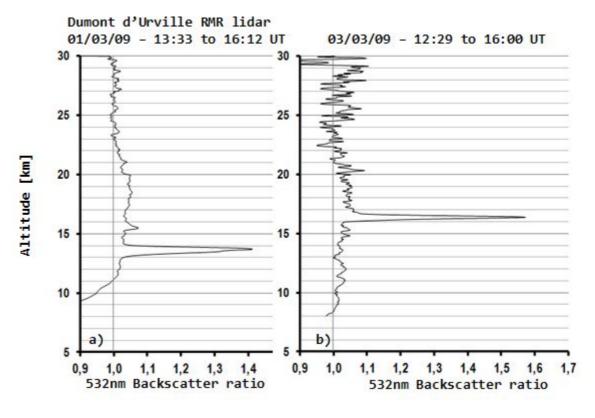
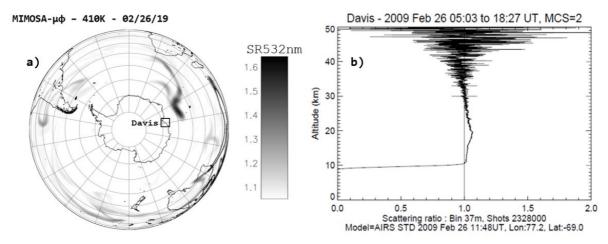




Fig. 2. Dumont d'Urville 532nm Scattering Ratio on March 1<sup>st</sup> (a), and 3<sup>rd</sup> (b) with smoke layers
 respectively visible at 13.5 km and 16.4 km.

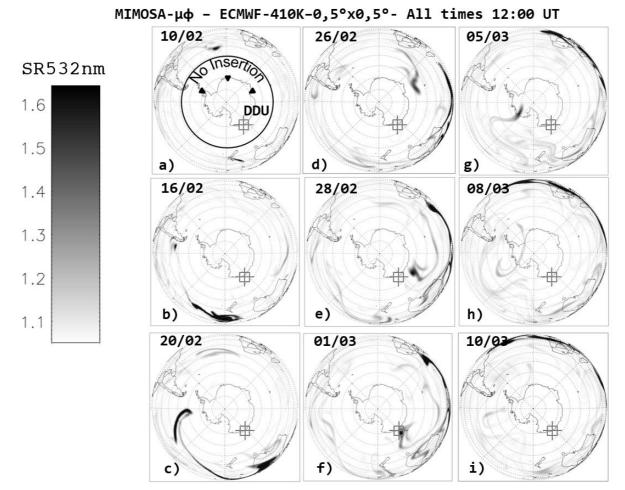
846



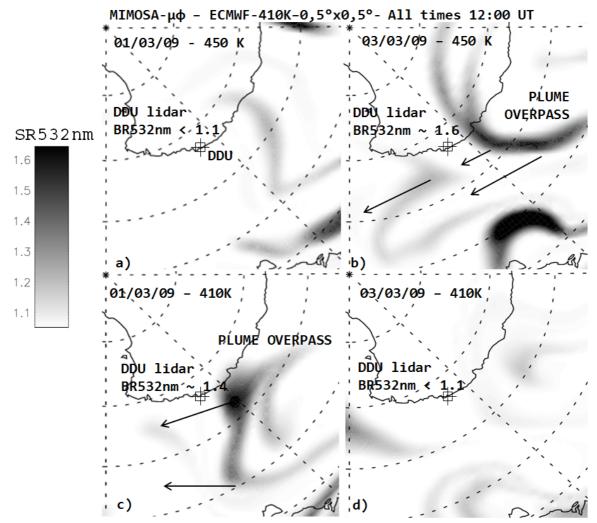
847



849 acquired at Davis Station.



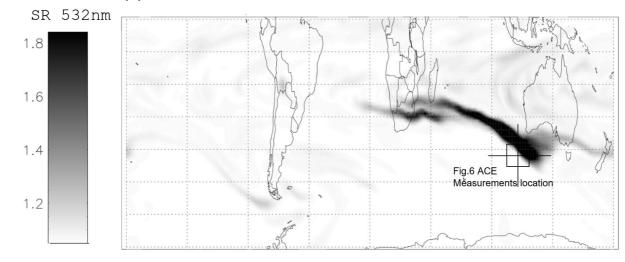
850 851 **Fig. 4.** 532nm MIMOSA-μφ SR calculations from February 10<sup>th</sup> to March 10<sup>th</sup>. Spatial resolution is 0.5°x0.5° and the potential temperature level is 410 K. CALIOP SR is integrated only up to 852 853 the 16<sup>th</sup> of February and in the (-50°;50°) latitudinal range. The location of Dumont d'Urville station is indicated by the square. 854



**Fig. 5.** 532nm MIMOSA- $\mu\phi$  Scattering Ratio calculations and detailed outputs above the vicinity of DDU on March 1<sup>st</sup> at 450 K (a), 410 K (c) potential temperature levels and March 3<sup>rd</sup> at 450 K (b), 410 K (d) potential temperature levels. Plume overpass is detected in Figure 4b and 4c consistently with lidar detections in Figure 1.

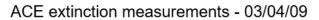
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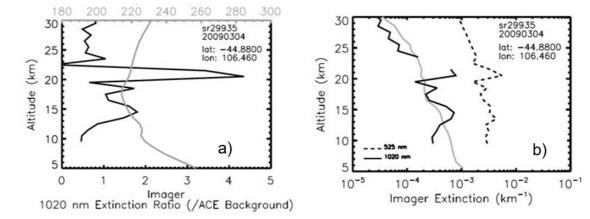
MIMOSA-µφ – 532nm SR - 550K - 03/04/09



863Fig. 6. 532nm MIMOSA-μφ SR calculations on March 4th (b). Crosshairs locate the ACE864measurements displayed in Figure 7.

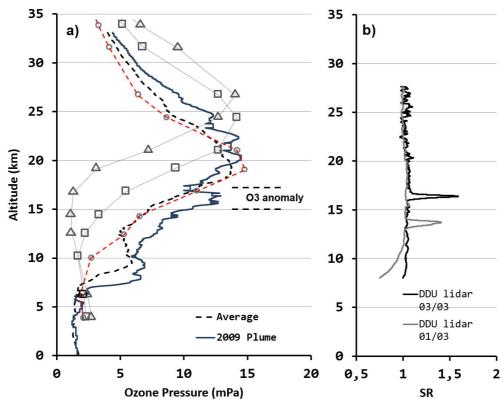
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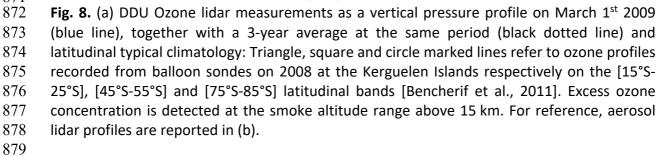


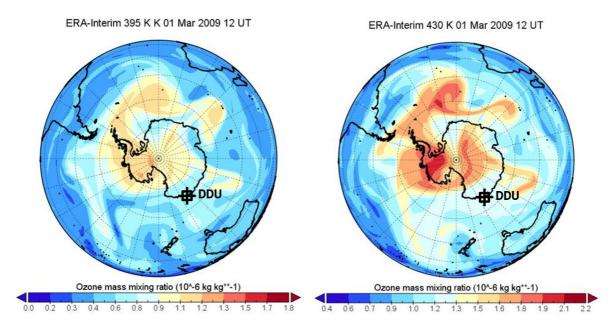
866
 867 Fig. 7. Visible and Near Infra Red ACE imager measurements on March 3<sup>rd</sup>. (a) Extinction ratio

868 is reported as compared to the ACE reference background. Imager extinction is in (b).











**Fig. 9.** Ozone mass mixing ratio from ERA-Interim at 12:00 UT on 01 March 2009 on the isentropic surfaces of (left) 395 K and (right) 430 K.