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To cite this version:

Jon Blundy, Helena Melekhova, Luca Ziberna, Madeleine C.S. Humphreys, Valerio Cerantola, et al.. Effect of redox on Fe–Mg–Mn exchange between olivine and melt and an oxybarometer for basalts. Contributions to Mineralogy and Petrology, 2020, 175 (11), $10.1007/s00410-020-01736-7$. insu-03092886

HAL Id: insu-03092886 <https://insu.hal.science/insu-03092886v1>

Submitted on 3 Jan 2021

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ORIGINAL PAPER

Efect of redox on Fe–Mg–Mn exchange between olivine and melt and an oxybarometer for basalts

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Received: 28 April 2020 / Accepted: 11 September 2020 / Published online: 15 October 2020 © The Author(s) 2020

Abstract

The Fe–Mg exchange coefficient between olivine (ol) and melt (m), defined as Kd_{Fe}_{T−Mg}=(Fe_{ol}/Fe_m)·(Mg_m/Mg_{ol}), with all Fe^T expressed as $Fe²⁺$, is one of the most widely used parameters in petrology. We explore the effect of redox conditions on Kd_{Fe}^{*T*}−Mg</sub> using experimental, olivine-saturated basaltic glasses with variable H₂O (≤7 wt%) over a wide range of *f*O₂ (ironwüstite buffer to air), pressure (\leq 1.7 GPa), temperature (1025–1425 °C) and melt composition. The ratio of Fe³⁺ to total Fe (Fe3+/∑Fe), as determined by Fe K-edge µXANES and/or Synchrotron Mössbauer Source (SMS) spectroscopy, lies in the range 0–0.84. Measured Fe³⁺/ Σ Fe is consistent (\pm 0.05) with published algorithms and appears insensitive to dissolved H₂O. Combining our new data with published experimental data having measured glass Fe³⁺/ Σ Fe, we show that for Fo_{65–98} olivine in equilibrium with basaltic and basaltic andesite melts, Kd_{Fe}^{*T*}−Mg</sub> decreases linearly with Fe³⁺/∑Fe with a slope and intercept of 0.3135 ± 0.0011 . After accounting for non-ideal mixing of forsterite and fayalite in olivine, using a symmetrical regular solution model, the slope and intercept become 0.3642 ± 0.0011 . This is the value at Fo₅₀ olivine; at higher and lower Fo the value will be reduced by an amount related to olivine non-ideality. Our approach provides a straightforward means to determine Fe³⁺/∑Fe in olivine-bearing experimental melts, from which *f*O₂ can be calculated. In contrast to Kd_{Fe} T_{Mg} , the Mn–Mg exchange coefficient, Kd_{Mn–Mg}, is relatively constant over a wide range of P–T–*f*O₂ conditions. We present an expression for Kd_{Mn-Mg} that incorporates the effects of temperature and olivine composition using the lattice strain model. By applying our experimentally-calibrated expressions for Kd_{Fe} T _{−Mg} and Kd_{Mn−Mg} to olivine-hosted melt inclusions analysed by electron microprobe it is possible to correct simultaneously for post-entrapment crystallisation (or dissolution) and calculate melt Fe³⁺/ Σ Fe to a precision of ≤ 0.04 .

Keywords Olivine · Experiments · Melt inclusions · µXANES · Mössbauer · Oxygen fugacity

Communicated by Dante Canil.

Electronic supplementary material The online version of this article [\(https://doi.org/10.1007/s00410-020-01736-7\)](https://doi.org/10.1007/s00410-020-01736-7) contains supplementary material, which is available to authorized users.

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Introduction

The exchange of iron and magnesium between olivine and coexisting melt bears directly on the generation and chemical evolution of basaltic magmas. Consequently, the

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Fe–Mg exchange coefficient, Kd_{Fe-Mg} , is one of the mostly widely used parameters in petrology.

 $\text{Kd}_{\text{Fe}-\text{Mg}}$ is related to the equilibrium constant for the exchange reaction:

$$
Fe2SiO4(melt) + Mg2SiO4(olivine) = Fe2SiO4(olivine) + Mg2SiO4(melt)
$$
\n(1)

and is defned as

$$
Kd_{Fe^{2+}-Mg} = \frac{(Fe^{2+}/Mg)_{ol}}{(Fe^{2+}/Mg)_{melt}},
$$
\n(2)

where $Fe²⁺$ and Mg are expressed as atomic concentrations. At equilibrium $Kd_{Fe^{2+}-Mg}$ is expected to vary with the free energy of exchange reaction (1) for the pure end-members, as well as any non-ideal interactions of $Fe²⁺$ and Mg dissolved in melts and in olivine. As the enthalpies, entropies and volumes of fusion of the olivine end-members forsterite and fayalite are somewhat diferent from each other (e.g., Lange and Carmichael 1990) one would expect, a priori, some temperature and pressure dependence of $\text{Kd}_{Fe^{2+}-\text{Mo}}$. In addition, the non-ideality of silicate melts and, to a lesser extent, of olivine solid solutions would be expected to confer significant compositional dependence on $Kd_{Fe^{2+}-Mg}$. However, despite these expectations, the seminal work of Roeder and Emslie (1970) showed that $\text{Kd}_{Fe^{2+}-Me}$ is remarkably constant over a wide range of pressure, temperature and composition (P–T–X), such that a value of 0.30 ± 0.03 can be used with some confdence to describe melting and crystallisation in a wide variety of olivine-bearing systems. We refer to this as the 'canonical' value of $\text{Kd}_{\text{Fe}^{2+}-\text{Mg}}$.

There have been many subsequent studies of olivinemelt equilibrium (e.g., Ulmer 1989; Beattie et al. 1991; Beattie 1993; Herzberg and O'Hara 2002; Toplis 2005; Mysen 2006; Matzen et al. 2011; Putirka 2016) and associated attempts to refne the canonical value or establish its sensitivity to P–T–X, but it has remained one of the most durable underpinning tenets of basalt petrology, used in a wide variety of ways, from tests of the primitive (i.e., mantle-derived) character of basaltic magmas, to corrections for post-entrapment crystallisation of melt inclusions, to fractionation of basaltic magmas in the crust and mantle.

A particular challenge with using Eq. (1) is the need to know the $Fe³⁺$ content of the silicate melt, which is not readily measurable by conventional electron microprobe techniques (see review by Hughes et al. 2018). Olivine contains negligible $Fe³⁺$ (less then a few thousand ppm and not more than a few percent of the total Fe, Fe^T ; Ejima et al. 2018) so that where Fe^{3+} in the melt is low, i.e., in relatively reduced systems, $Kd_{Fe^{2+}-Mg}$ can be used with confidence assuming that $Fe^{2+} = Fe^{T}$. In such cases

a variant of Eq. (2) , with all Fe expressed as Fe^T, i.e., $Fe^T = \Sigma Fe = Fe²⁺ + Fe³⁺$, may be used instead:

$$
Kd_{Fe^T - Mg} = \frac{\left(Fe^T / Mg\right)_{ol}}{\left(Fe^T / Mg\right)_{melt}}.
$$
\n(3)

Most natural magmatic systems contain some $Fe³⁺$, thus Kd_{Fe} $^{T}-Mg$, as expressed in (3), will be sensitive to redox state. Fe³⁺/ΣFe ratios vary widely in basaltic magmas, from almost zero in the case of lunar basalts to > 0.5 in the case of some oxidised, hydrous subduction-related igneous rocks (Stolper and Bucholz 2019). Moreover, even at constant P, T and fO_2 , Fe³⁺/ΣFe is known to vary with melt composition (Kress and Carmichael 1991; Putirka 2016; Borisov et al. 2018). For example, elevated Fe³⁺/ΣFe occurs in alkaline basaltic magmas due to the stabilising effect of $Na⁺$ and K^+ on Fe³⁺ (Mysen and Virgo 1989; Kress and Carmichael 1991). Thus, in some tectonic environments, Kd_{Fe}^{*T*}−Mg</sub> can be lower than $Kd_{Fe^{2+}-Mg}$ by as much as a factor of two or more, with far-reaching implications for olivine-melt equilibrium in basaltic systems.

If the availability of Fe^{2+} in the melt is the dominant control on Fe–Mg exchange, then Kd_{Fe}^{*T*}–M_g should vary systematically with the redox state of the system; the relationship between Kd_{Fe}^{*T*}−Mg</sub> and Kd_{Fe²⁺−Mg} will reflect the proportion of total iron in the melt that is trivalent (Fe^{3+}) at the pressure, temperature and melt composition of interest:

$$
\mathrm{Kd}_{\mathrm{Fe}^T - \mathrm{Mg}} = \mathrm{Kd}_{\mathrm{Fe}^{2+} - \mathrm{Mg}} \times \left(1 - \frac{\mathrm{Fe}^{3+}}{\Sigma \mathrm{Fe}}\right). \tag{4}
$$

Recognising the problem of $Fe³⁺$ in the melt Roeder and Emslie (1970) attempted to account for its effect by determining, via wet chemistry, the Fe³⁺/ΣFe ratio of their bulk experimental charges and making a correction for any contained olivine crystals using a simple mass balance. The behaviour predicted in Eq. (4) is apparent in the original Roeder and Emslie (1970) dataset (Fig. 1), although the uncertainty in the Fe³⁺/ΣFe ratio of the glass (as opposed to the bulk) precludes any meaningful conclusions. For example, it is unclear if the scatter in Fig. 1 arises due to temperature or compositional effects, the use of different capsule materials (and attendant Fe loss from the glass), or the presence of crystals in the aliquot of the experimental charge used for measuring $\text{Fe}^{3+}/\text{2Fe}$. The scatter is not removed even when more sophisticated mass balance techniques are used, taking into account other iron-bearing mineral phases in the glass, e.g., clinopyroxene, magnetite (Matzen, 2012). However, in principle, if we know the redox state of a magmatic system, usually defned in terms of an oxygen fugacity (*f*O₂), and the relationship between *f*O₂ and Fe³⁺/ΣFe, then we should be able to use Eq. (4) as an oxybarometer. The potential of this approach was recognised by Putirka (2016)

Fig. 1 Variation of olivine-melt Kd_{Fe^T-Mg} with $\text{Fe}^{3+}/\text{2Fe}$ in the bulk experimental run product from the seminal experimental study of Roeder and Emslie (1970). Results for three diferent capsule materials are shown. The data hint at the operation of a relationship akin to Eq. (4), but the measurement of $Fe^{3+}/\Sigma Fe$ is insufficiently precise for a more thorough treatment because of the need to correct the bulk $Fe³⁺/\Sigma Fe$ analysis for included, Fe-bearing crystalline phases. The line and equation correspond to fts to Eq. (4), for comparison to Fig. 6a

who developed an expression for Kd_{Fe}^{*T*}−Mg that explicitly includes an $fO₂$ term. However, the form of his Eq. (9b) is not optimised for oxybarometry, partly because of the way the ftting was performed, and partly because of the dearth of experimental olivine-melt data with measured $Fe^{3+}/\Sigma Fe$.

Application of Eq. (4) as an oxybarometer requires, of course, that the olivine-melt pair of interest is in equilibrium. When Fe³⁺/ΣFe in the melt is unknown, the problem of testing for equilibrium and calculating *f*O₂ becomes circular. To resolve this circularity requires understanding the behaviour of a component that is not redox-sensitive, such as exchange of Mn and Mg between olivine and melt, Kd_{Mn-Mg} . If the diference between the behaviour of Fe and Mn can be quantifed, it has the potential to be used to test for olivine-melt equilibrium and so enable oxybarometry. The aim of this study is to (a) compare methods for the reliable, high-spatial resolution measurement of the $Fe³⁺$ content of hydrous and anhydrous glasses of broadly basaltic composition synthesised over a range of fO_2 ; and (b) determine the partitioning of Fe, Mg and Mn between olivine and basaltic melt across a wide range of redox conditions. The ultimate objective is to establish whether Eq. (4) provides a useful means to correct Kd_{Fe} T _{−Mg} for Fe³⁺-rich systems, and whether it is possible to recover *f*O₂ from determinations of Kd_{Fe} r _{−Mg} in natural or experimental basaltic systems by exploiting the redox insensitive exchange of Mn and Mg between olivine and melt.

Experimental methods

This study utilises 72 experimental samples, mainly from our previously published studies with existing, modifed or new determinations of the Fe³⁺/ΣFe ratio in the glass to explore the effect of redox on Kd_{Fe}^{*T*}−Mg</sub>. Our experimental starting materials comprised eight basalts, two basaltic andesites and one andesite, all based on natural rock compositions from subduction-related magmatic systems: Lesser Antilles arc (St. Vincent, Grenada, Martinique, St. Kitts, Montserrat), Central American arc (Masaya), Aeolian arc (Stromboli), and the post-collisional, calc-alkaline Adamello Batholith, Italy. For St. Vincent and Martinique we used three and two diferent starting materials, respectively; these are referred to as St. Vincent series 1, 2 and 3, and Martinique series 1 and 2. MgO contents of the starting materials range from 2.3 to 17.1 wt%; total alkali contents range from 1.6 to 4.7 wt%. Mg#, expressed in terms of Fe^T (i.e., molar Mg/[Mg+Fe*^T*]) ranges from 0.34 to 0.76. Compositions, on an anhydrous basis, of all eleven starting materials are given in Table 1. None of the studied glasses contain sulphur, to prevent possible modifcations of Fe3+∕ΣFe ratios due to the homogeneous reaction S^{2-} + 8Fe³⁺ = S^{6+} + 8Fe²⁺ during quenching of the glass (Nash et al. 2019).

With the exception of the three 'XANES' experiments (see Supplementary Method in supplementary material 3) all of the experimental runs are taken from the published studies listed in Tables 1 and 2, where the experimental techniques are described in detail, and summarised as follows. One-atmosphere experiments (Adamello series) were run in vertical quench furnaces at the Geophysical Laboratory, Carnegie Institution of Washington, with CO–CO₂ gases to regulate $fO₂$ using methods described in Ulmer (1989) and Kägi et al. (2005). One-atmosphere experiments (Grenada series) were run in vertical quench furnaces at University of Bristol and Australian National University with $CO-CO₂$ gases to regulate $fO₂$ using methods in Stamper et al. (2014). Internally-heated pressure vessel (IHPV) runs (St. Vincent 2 and 3, St. Kitts, Martinique 1 and 2, and Montserrat series) were performed at Université d'Orléans using techniques reported by Pichavant et al. (2002) and Melekhova et al. (2017) with $fO₂$ controlled by $H₂$ added to the argon pressurising gas, and $fO₂$ monitored using an NiPd sensor. IHPV experiments at Leibnitz University of Hannover (Stromboli and Masaya series) were run using the methods described in Lesne et al. (2011) , without H₂ control. Piston cylinder (St. Vincent 1 and Grenada series) runs were performed in half-inch pressure cells at University of Bristol using

1. Ulmer (1989); Kägi et al. (2005)

2. Stamper et al. (2014)
3. Melekhova et al. (2015) 2. Stamper et al. (2014)

4. Pichavant and Macdonald (2007) 4. Pichavant and Macdonald (2007) 3. Melekhova et al. (2015)

5. Lesne et al. (2011) 5. Lesne et al. (2011)

6. Pichavant (unpublished) 6. Pichavant (unpublished)

7. Pichavant et al. (2002) 7. Pichavant et al. (2002)

8. Melekhova et al. (2017) 8. Melekhova et al. (2017)

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techniques described by Melekhova et al. (2015) and Stamper et al. (2014). Those experiments used a double capsule technique, with sample materials placed in both inner and outer capsules to minimise H loss or gain by diffusion. In some St. Vincent series 1 runs a sensor capsule of pure Pd, loaded with a mixture of Ni, NiO and H₂O, was placed inside the outer capsule to monitor $fO₂$ by means of analyses of NiPd alloys. All 72 experiments have the Fe³⁺/ΣFe ratio of their glass determined by one or multiple techniques (see Table 3); two samples have replicate measurements. $fO₂$ was known most precisely in the 18 one-atmosphere runs. In high-pressure runs, performed at water-undersaturated conditions in IHPV or piston cylinder with NiPd sensors, $fO₂$ can be calculated only by taking account of the reduced H_2O activity

 $(aH₂O)$ in the melt at run conditions. This was done using the method of Burnham (1979). The fO_2 of the run is then that of the NiPd sensor plus $2log(aH₂O)$. This approach, which involves a greater uncertainty than the one-atmosphere experiments, provides $fO₂$ estimates for a further 15 runs. The overall range in measured $fO₂$ these 33 experiments is 9.3 log units relative to the NNO bufer (O'Neill and Pownceby 1993) at experimental P and T (i.e., from NNO $- 2.8$ to NNO + 6.5). In the remaining 40 experimental runs the experimental $fO₂$ was not precisely constrained.

In addition, we considered another \sim 100 experiments with measured $\text{Fe}^{3+}/\text{\Sigma}\text{Fe}$ ratios for the glass (see below) and a database of over 1000 published olivine-bearing experiments conducted at known *f*O₂ above NNO − 3, to test various aspects of our parameterizations.

Analytical methods

Electron probe microanalysis (EPMA)

Major and minor elements in glasses and minerals in the Adamello (RC158c), St. Vincent 1 (RSV49, XANES) and Montserrat series of experiments were analysed using a Cameca SX100 electron microprobe at University of Bristol. Analytical conditions were: olivine—20 kV primary beam, 10 nA beam current and 1 µm beam diameter; glass—20 kV, 4 nA and 10 µm beam. Calibration was performed on a range of mineral, oxide and glass standards. Secondary standards used were: St. John's Island olivine, Kakanui hornblende, diopside, Columbia River Basalt glass (USGS) and an inhouse synthetic amphibolite glass (#3570). All other glasses and minerals are those reported by the original authors (Table 1). At least 5 analyses were made of each olivine and glass. The subset of the 72 experiments containing analysed olivine crystals is 52.

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Volatiles measured by SIMS or Karl-Fischer titration (Martinique series) except where estimated from glass fraction and starting material (or volatiles by difference) shown in italics fVolatiles measured by SIMS or Karl–Fischer titration (Martinique series) except where estimated from glass fraction and starting material (or volatiles by diference) shown in italics

eCalculated at experimental P, T,

fOlivine-melt exchange coefficient

'Olivine-melt exchange coefficient

*f*O2 and glass composition: Kress and Carmichael (1991) (KC91), Borisov et al. (2018) (BBH18), O'Neill et al. (2018) (OBM18)

Calculated at experimental P, T, fO₂ and glass composition: Kress and Carmichael (1991) (KC91), Borisov et al. (2018) (BBH18), O'Neill et al. (2018) (OBM18)

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Micro X‑ray absorption near‑edge spectroscopy (µXANES)

Fe3+∕ΣFe ratios were measured in 66 experimental glasses by µXANES at the Diamond Light Source synchrotron facility, UK, using techniques described in Stamper et al. (2014), with some modifications, as summarized here. μ XANES spectra were collected at the Fe K-edge in fuorescence mode on Beamline I18 using the Si(111) monochromator. The beam size at the sample was an ellipse with principal axes approximately 2.5×4.5 microns. The incident beam flux was reduced by placing one or more 50 μ m-thick aluminium foil sheets in front of the sample in addition to the fxed 15 µm-thick flter used for all analyses (Table 3). The total fux density (defned as total photons delivered per second per square micrometer; Cottrell et al. 2018) is 6.8×10^{10} (with 15 μ m Al filter) or 2.5 $\times \times 10^{10}$ photons s⁻¹ μ m⁻² (with 15+50 µm Al flters). One sample (PU58) was analysed twice with both 65 µm and 115 µm thickness of Al flters (flux density of < 1.1×10^{10} photons s⁻¹ µm⁻²); the two measured Fe³⁺/ΣFe ratios are 0.761 and 0.777; the mean is reported in Table 3.

Fluorescence counts were normalized to the incident beam fux at every energy step and collected using a 9-element solid-state Ge detector. The energy was calibrated by defning the frst peak of the frst derivative of Fe foil to be 7112 eV. Each spectrum was collected using four sets of energy acquisition conditions, giving good resolution over the pre-edge region and sufficient post-edge detail to allow high-quality normalization and background ftting. Each point was analyzed for 1000 or 2000 ms, giving total acquisition times of approximately 10–15 min. Calibration was performed on ten basalt glass standards loaned by the Smithsonian Institution (Cottrell et al. 2009), using the revised and updated Fe³⁺/ΣFe values of Zhang et al. (2018). Data processing is described in more detail in Stamper et al. (2014), but their original values have been updated in Table 3. Multiple (2 or 3) analyses were made of each glass; mean Fe³⁺/ΣFe and 2 standard deviations (s.d.) are reported in Table 3.

Secondary ion mass spectrometry (SIMS)

Dissolved volatiles $(H₂O, CO₂)$ in seven high-pressure experimental glasses from Montserrat and St. Vincent 1, 2 and 3 series were measured by SIMS using a Cameca ims4f instrument at the NERC Edinburgh ion-microprobe facility (EIMF), UK. Samples were gold-coated for analysis. The primary beam was 10 keV O− ions with net impact energy of 14.5 keV (4.5 kV secondary voltage). Beam current was 5 nA, focussed to $a \sim 15 \mu m$ spot at the sample surface. Prior to analysis the sample surface was sputtered with a 25×25 µm raster for 3 min to eliminate surface contamination. Positive

secondary ions of ${}^{1}H$ and ${}^{12}C$ were collected with the 25 μ m imaged field and 150 µm field aperture in two separate analytical routines with offset voltages of 75 (¹H) and 50 V (12) . To avoid 24 Mg²⁺ interference 12 C⁺ secondary ions were collected at high mass resolution (*M*/∆*M* = 1200). ${}^{1}H^{+}$ secondary ions were collected at lower mass resolution ($M/\Delta M$ = 500). A total of 15 analytical cycles was used, with count times of 5 s per cycle for ${}^{1}H$ and 12 s for ${}^{12}C$. Only the final 5 cycles (1 H) or 7 cycles (12 C) were processed to remove any lingering efects of surface contamination. ³⁰Si was used as an internal standard, based on prior EPMA analyses. Calibration was performed on a suite of H_2O - and $CO₂$ -bearing synthetic and natural glass standards. Backgrounds, monitored on natural quartz grains co-mounted with the experimental glasses, were 2.8 ± 0.8 counts per second (cps) ¹²C and 1600 ± 700 cps ¹H. Minimum detection limits, calculated as 3 s.d. on the blanks, were 140 ppm H_2O and 26 ppm $CO₂$. At least three analyses were made of each glass except for St. Vincent series 2 run 7–3, where only one sufficiently large glass pool could be found.

 H_2O and CO_2 contents of a further 32 Grenada, St. Kitts, Stromboli, Masaya and St. Vincent 1 series glasses analysed by SIMS using similar techniques can be found in the original sources. Pichavant et al. (2002) report H_2O , but not CO_2 , in fve St. Vincent 2 and Martinique series glasses. Melekhova et al. (2015) report $H₂O$ contents for 14 St. Vincent series 1 glasses. Uncertainties on H_2O and CO_2 (Table 2) are based on 1 s.d. of multiple analyses for new SIMS data, or are taken from original data sources.

In‑house Mössbauer spectroscopy

Mössbauer spectroscopy measurements of two experimental glasses were made at Bayerisches Geoinstitut, Bayreuth, using a constant acceleration Mössbauer spectrometer in transmission mode with a nominal 370 MBq $57C_O$ point source in a 12 μm Rh matrix. Active dimensions of the point source were $500 \times 500 \mu m^2$. The velocity scale was calibrated relative to α -Fe and line widths of 0.36 mm s⁻¹ were obtained for outer lines of α -Fe at room temperature. Experimental glasses were at room temperature during data collection. The glass samples were 200 μm thick and Mössbauer spectra were collected over a 500 μm diameter region in the middle of each sample. Data collection took 7–12 days.

Mössbauer spectra were ftted using MossA (Prescher et al. 2012) with a linear baseline to account for shadowing. We adopted the xVBFmodel (see Alberto et al. 1996; Lagarec and Rancourt 1997) for the fit, with the full transmission integral to account for thickness efects of the source and absorber (Rancourt 1989), and conventional constraints (doublet components with equal widths and areas). The $Fe²⁺$ doublet ft used the xVBFmodel with correlation. One extra $Fe²⁺$ doublet was added to improve the fit of the $Fe²⁺$ absorption envelope; we used pseudo-Voigt line shape to minimize the number of extra parameters. The $Fe³⁺$ doublet fit used a pseudo-Voigt line shape to approximate the xVBFmodel with no correlation (see Partzsch et al. 2004; McCammon 2004). The Fe³⁺/ΣFe ratio was determined from relative areas. Spectra were ft with diferent models to assess dependence of $Fe^{3+}/\Sigma Fe$ on fitting model and error bars were estimated accordingly $(\pm 0.03$ in the ratio).

Synchrotron Mössbauer Source (SMS) spectroscopy

Energy-domain SMS measurements on 37 experimental glasses were conducted at the Nuclear Resonance Beamline ID18 (Rüfer and Chumakov 1996) at the European Synchrotron Radiation Facility (ESRF), Grenoble, France, operating in multibunch $(7/8 + 1)$ mode. SMS is based on a nuclear resonant monochromator employing pure nuclear refections of an iron borate $({}^{57}FeBO_3)$ single crystal (Potapkin et al. 2012). This source provides ⁵⁷Fe resonant radiation at 14.4 keV within a bandwidth of 6 neV which is tuneable in energy over a range of ± 0.6 μeV (Potapkin et al. 2012). Sample glasses were prepared as $150-350 \mu$ m-thick, doublypolished wafers. After polishing, the transparent wafers were checked under a microscope to locate areas free of bubbles or crystals.

The X-ray beam was focused at the sample surface to an ellipse with principal axes 17×18 µm at the full-width half-maximum (FWHM). Before and after each sample measurement, SMS linewidth was determined using a $K_2Mg^{57}Fe(CN)_6$ reference single-line absorber. The velocity scale (\pm 5 mm s⁻¹) was calibrated relative to a 25 µm-thick natural α-Fe foil. The small cross section, high brilliance and fully resonant and polarized nature of the beam allowed for rapid spectrum collection (approximately 2 h). Slightly longer run times (up to 6 h) were required for Fe-poor samples. Note that the glasses measured in this study contain only natural abundances of 57 Fe-atoms, i.e., ~2% of total Fe, thus the total radiation dosage, defned as photon delivered per μ m² of the sample is 150, compared to 10¹² for μ XANES at Diamond.

Typical samples for SMS analysis are shown in Fig. 2; SMS spectra (and fits) for these samples are shown in Fig. 3. All SMS spectra consist of two broad symmetric doublets, typical of basaltic glasses. The spectra were ftted with a full transmission integral and Lorentzian line shape using the software package MossA (Prescher et al. 2012). The single line spectra were ftted with a normalized Lorentzian-squared source line shape. A linear function was applied to model the background. To obtain the maximum amount of photons we used the confocal Be-lenses installed at the ESRF beamline ID18. Be-lenses always bear traces of iron that result in the presence of two components

Fig. 2 Transmitted light photomicrographs of typical run products prepared for SMS analysis: **a** run RSV49_4; and **b** run HAB23. Doubly polished glass chips show clear glass pools suitable for SMS analysis. Crystalline phases are olivine in **a** and magnetite in **b**. The locations of the SMS analyses reported in Table 3 are shown as red circles

(Supplementary Fig. 1). The signal for iron in Be-lenses is easily corrected for and defned in all of the glass SMS spectra (Fig. 3).

The approach we adopt for the ftting uses only two distinct components represented by two doublets, one for $Fe³⁺$ and one for $Fe²⁺$. Attempts using a model involving an additional $Fe²⁺$ component did not improve the quality of the fitting nor change the final $\text{Fe}^{3+}/\text{\Sigma}\text{Fe}$ results. Hyperfine parameters (centre shift, CS, and quadrupole splitting, QS) were frst determined from the spectra where the two components were easily identifed. For the other samples, the hyperfne parameters were allowed to vary within those ranges. $Fe³⁺/$ ΣFe values were obtained from the relative areas of the two components. The errors (2 s.d.) on the Fe3+*/*ΣFe ratios were obtained by normal error propagation.

Hyperfine parameters, $Fe^{3+}/\Sigma Fe$ values and related errors for all SMS analyses are reported in Supplementary Table 1. Centre shift varies in the range 1.00–1.12 mm s⁻¹ for Fe²⁺ and $0.27-0.41$ for Fe³⁺. Quadrupole splitting varies in the range 1.89–2.60 for Fe²⁺ and 0.74–1.29 mm s⁻¹ for Fe³⁺. With one exception (run HAB23; Supplementary Table 1), absolute errors on $Fe^{3+}/\Sigma Fe$ range from 0.02 to 0.05. Highest errors are typically associated with samples having either

Fig. 3 Mössbauer spectra of experimental samples **a** HAB23 and **b** RSV49_4 (spot a) obtained using synchrotron Mössbauer source (SMS) spectroscopy at the beamline ID18, ESRF. Green and blue areas represent the fitted Fe^{2+} and Fe^{3+} components, respectively, while the grey area represents the contribution of Fe in Be-lenses (see text and Supplementary Material 1). Red curve represents the sum of all components (i.e., the modelled ft)

low or high Fe³⁺/ΣFe, for which the Fe³⁺ or Fe²⁺ component is less easily resolved.

Colorimetry

Four glasses of the Stromboli and Masaya series of Lesne et al. (2011) were analysed previously for Fe³⁺/ΣFe using colorimetric wet-chemistry, following the method of Schuessler et al. (2008).

Results

Experimental run conditions and products are presented in Table 2, $Fe^{3+}/\Sigma Fe$ ratios and exchange coefficient in Table 3, and analyses of glasses and olivines in Tables 4 and 5, respectively. On an anhydrous basis glasses are basalts

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Blank denotes not analysed; bd is below detection

All Fe expressed as FeO

All Fe expressed as FeO

^aTotal includes H₂O and CO₂ content of glasses as reported in Table 2

 $^{\rm a}$ Total includes H2O and CO₂ content of glasses as reported in Table 2

 $(n=61)$ and basaltic andesites $(n=10)$ and a single dacite (HAB-26). H_2O and CO_2 range from zero (nominally) up to 8.3 and 1.2 wt%, respectively. Glass MgO contents range from 2 to 17 wt%; olivines are Fo_{86} to Fo_{98} with 0.09–0.59 wt% MnO (one at $F_{0,65}$ and 0.56 wt% MnO) and 0.02–0.94 wt% CaO. Total alkalis $(Na₂O + K₂O)$ in the glasses range from 0.1 to 4.9 wt%.

Ferric–ferrous ratios

Fe³⁺/ΣFe ratios range from 0.06 to 0.84 (μ XANES) and 0.04 to 0.80 (SMS). Typical uncertainties (2 s.d.) in Fe3+∕ΣFe, propagated through the various sources of analytical error, are in the range 0.001–0.32 (mean=0.016) for μ XANES; 0.02–0.07 (mean = 0.029) for SMS; and ~ 0.03 for in-house Mössbauer and colorimetry. To test the homogeneity of individual glasses we used SMS to analyze three separate chips of water-poor glass RSV49_4 from the centre (RSV49_4a, b) and from the edge (RSV49_4c) of the experimental charge. All three measurements lie within 2 s.d. of each other (Table 3). Similarly, two SMS analyses of hydrous glass BM46 from the centre (BM46a) and periphery (BM46b) of the same glass chip give values of Fe³⁺/ΣFe that agree within 2 s.d. (BM46c analysis is described below).

Two experimental glasses (XANES9, 6-4) analysed by SMS and in-house Mössbauer displayed consistency to within 2 s.d. (Table 3). SMS and colorimetry agree within 2 s.d. for three of the four glasses analysed by both techniques; the fourth (MAS1_B4) has a signifcantly higher Fe³⁺/ΣFe by SMS (0.289) than by colorimetry (0.18). The cause of this discrepancy is unclear. However, it is noteworthy that the Fe³⁺/ΣFe from colorimetry is signifcantly lower than other three values from glasses synthesised under similar conditions in the study of Lesne et al. (2011) , i.e., 0.32–0.37, suggesting a potential problem with colorimetric analysis of MAS1_B4. For the 27 glasses analysed by both µXANES and SMS, the former method gives significantly higher $\text{Fe}^{3+}/\text{2Fe}$ for all but ten (Table 3). This discrepancy is well outside the analytical uncertainty and is suggestive of oxidation by the X-ray beam during µXANES analysis (Cottrell et al. 2018). The problem of oxidation is found to be most acute in hydrous glasses (> 0.5 wt% H₂O) containing a significant proportion of the oxidizable species, $Fe²⁺$. This is evident from a plot of the difference between $\text{Fe}^{3+}/\text{\textstyle{\Sigma}}\text{Fe}$ by SMS and Fe³⁺/ΣFe by μXANES against the Fe³⁺/ΣFe (by SMS) in glass (Fig. 4). For anhydrous glasses $(< 0.5 \text{ wt\% H}_2O$) the two methods agree across a broad range of $\text{Fe}^{3+}/\text{2Fe}$; for hydrous glasses the scale of the mismatch increases roughly linearly with decreasing $Fe^{3+}/\Sigma Fe$ (i.e., increasing Fe3+∕ΣFe). This is consistent with the fndings of Cottrell et al. (2018), albeit here based on a much wider range of glass compositions and Fe³⁺/ΣFe. Cottrell et al. (2018) conclude that oxidation is most acute when photon fux densities are as high as those used in our analyses; they recommend using fux densities some two orders of magnitude lower to minimise the oxidative efects of the beam.

To explore further the infuence of photon dosage on in situ oxidation during µXANES analysis, we ran a series of continuous time-scans at Diamond at energies corresponding to the peaks of the Fe^{2+} and Fe^{3+} regions of the preedge, without attenuating Al foils in front of the beam. This enables a semi-quantitative, real-time assessment of maximum beam damage. We analysed synthetic hydrous glasses MAS1_B4 and MAS1_B5 from Lesne et al. (2011) with 3.0 and 2.6 wt% H_2O , respectively, alongside anhydrous glass standard LW_10 from Cottrell et al. (2009). Both hydrous glasses have ~ 10.7 wt% FeO^T; LW has 10.2 wt% FeO^T (Cottrell et al. 2009). All samples are moderately oxidised with $Fe^{3+}/\Sigma Fe$ of 0.18 (MAS1_B4) and 0.32 (MAS1_B5), as determined by colorimetry (Lesne et al. 2011), 0.289 $(MAS1_B4)$ by SMS (Table 3), and 0.235 (LW_10) by inhouse Mössbauer (Cottrell et al. 2009). The sample shutter was kept closed until the beginning of counting, and counts were collected every 5 s for 500 s, which is approximately the same length of time taken to reach the pre-edge region in our quantitative μ XANES analyses. For LW_10, Fe³⁺ count rates (normalised to I_0) showed only a minimal increase during the analysis, whereas MAS1_B4 and MAS1_B5 showed relatively small increases in apparent $Fe³⁺/\Sigma Fe$ (calculated from peak height alone, not from a full area-weighted centroid ft) of ~ 0.02 and 0.05, respectively (Supplementary Fig. 2).

This experiment confirms that beam damage during µXANES is occurring in our samples, and that the extent of damage increases with higher H_2O concentrations and lower initial redox state, as described in Cottrell et al. (2018) . However, regardless of H₂O concentration, the measured deviations are very weak in oxidised glasses with $\text{Fe}^{3+}/\text{2Fe} \geq 0.5$ and strongest in reduced glasses with Fe³⁺/ΣFe ≤ 0.5 (Fig. 4). In unpublished μ XANES data from more evolved (rhyolitic) compositions we have also observed greater beam damage in reduced, H_2O -poor glasses than in oxidised, H_2O -rich equivalents. We suggest, therefore, that the dominant factor controlling beam-damage susceptibility is the initial redox state of the glass. However, as anhydrous glasses experience little or no damage, the presence of H_2O is clearly important (Cottrell et al. 2018). During electron beam damage of hydrous glasses, H_2O appears to facilitate damage by enhancing the difusion of alkalis and OH towards or away from the centre of electron deposition (Humphreys et al. 2006). We speculate, following Cooper et al. (1996), that a similar efect may occur during X-ray irradiation, with Fe oxidation generated as a result of alkali migration towards the site of maximum X-ray fux. Such mechanisms could be tested in future by detailed mapping

 ${}^{\text{a}}$ Molar percentage of forsterite (Fo), fayalite (Fa) and tephroite (Te)

Fig. 4 Comparison of µXANES and SMS methods used to measure Fe3+∕ΣFe in our experimental glasses. The diference between XANES and SMS measurements of $Fe³⁺/\Sigma Fe$ is plotted versus that measured by SMS. μ XANES tends to overestimate Fe³⁺/ΣFe due to beam damage at the high photon fuxes used. As the amount of oxidisable $Fe²⁺$ in the glass increases, so the tendency to oxidise during µXANES increases for all but the most dry $(< 0.5$ wt% H₂O) glasses. Only for very oxidised (Fe³⁺/ΣFe>0.5) or dry glasses do μ XANES and SMS agree within experimental uncertainty. Data are taken from Table 3; error bars are 2 sd

of cation distribution in the regions surrounding irradiated sites (Cottrell et al. 2018).

To directly simulate possible beam damage caused by µXANES analysis at the high photon fux densities used at Diamond we analysed by SMS a glass chip from experiment BM46 previously exposed to comparable beam dosage to µXANES. A spot on the sample (BM46c in Table 3) was first irradiated for 10 min with a beam flux of 10^{14} photon s⁻¹, corresponding to a total photon dosage of 8×10^{11} . A dark spot appeared at the glass surface after irradiation. This spot was then analysed by SMS, using the same low photon fux as previously. The analysis revealed signifcantly higher Fe³⁺/ΣFe (0.196±0.017) compared to un-irradiated spots (BM46a,b) on the same chip $(0.130 \pm 0.023$ and 0.104 ± 0.020 , confirming that high photon-flux irradiation of hydrous glass, such as that during µXANES at Diamond, causes appreciable oxidation. The extent of the oxidation is, however, somewhat less than we observed for the typical mismatch between SMS and µXANES (Fig. 4), probably due to the smaller spot size of µXANES (i.e., greater fux density) compared to the SMS simulation. The mechanism of oxidation is not clear, but may involve the formation of tiny magnetite nanolites (Di Genova et al. 2017; Hughes et al. 2018), as suggested by Raman spectra of some µXANES spots and the darkening of the glass after irradiation. This mechanism may operate in tandem with the alkali migration proposed above.

In light of the likely oxidation of some hydrous glasses during µXANES analysis, we favour SMS data over µXANES data for all hydrous glasses. For glasses with < 0.5 wt% H₂O, but lacking SMS analyses, we have adopted the μ XANES Fe³⁺/ΣFe values. Our 'preferred' Fe³⁺/ΣFe values for each experiment are given in Table 3; these are the Fe³⁺/ΣFe values used in all subsequent calculations and plots. Note that some hydrous glasses without SMS analysis have no preferred $Fe^{3+}/\Sigma Fe$ value and are not considered further. The total number of diferent glasses with preferred Fe³⁺/ΣFe values is 47, of which 28 were synthesised at known $fO₂$ and 20 coexist with olivine.

We have supplemented our new experimental data with a further 108 experimental glasses of broadly basaltic (or haplobasaltic) and basaltic andesite (≤ 60 wt% SiO₂) composition with measured $\text{Fe}^{3+}/\text{\Sigma}\text{Fe}$ performed at controlled (or measured) fQ_2 . This dataset was selected to include experiments that contain olivine and/or experiments where the glass contains known amounts of dissolved H_2O . This provides, respectively, an additional set of data with which to explore olivine-melt partitioning and an additional set of Fe³⁺/ΣFe data for hydrous systems. The studies used are: Mysen (2006), Matzen et al. (2011), Partzsch et al. (2004), Botcharnikov et al. (2005), Vetere et al. (2014), Wilke et al. (2002), Moore et al. (1995), and Schuessler et al. (2008). The number and type of measurements from these studies are provided in Supplementary Table 2.

Calculated Fe³+∕*6***Fe in hydrous basaltic melts**

In Fig. 5, we compare the measured $\text{Fe}^{3+}/\text{2Fe}$ in our 28 glasses from experiments with known $fO₂$, together with the additional 108 glasses described above, to the calculated Fe³⁺/ΣFe at the experimental P–T–*f*O₂ using three different algorithms: the widely used algorithm of Kress and Carmichael (1991), which includes a pressure term; and the recent algorithms of Borisov et al. (2018) and O'Neill et al. (2018) that do not.¹ None of these algorithms includes an explicit H₂O term, even though O'Neill et al. (2018) did use the hydrous data of Moore et al. (1995) in their calibration. The agreement for all three algorithms is encouraging (average absolute deviation, *aad*=0.06–0.08 in Fe3+∕ΣFe), with Borisov et al. (2018) performing marginally better than the other two. The observed *aad* is in excellent agreement with that claimed by Borisov et al. (2018) for their much larger, anhydrous calibrant dataset (\pm 0.05 at Fe³⁺/ΣFe=0.25). The good performance of both algorithms is despite the fact that the combined dataset in Fig. $5(n=136)$ includes a large number $(n=79)$ of hydrous glasses, whereas only that of O'Neill et al. (2018) included any hydrous glass in their calibrations. For our dataset, plots (not shown) of the deviation of Fe³⁺/ΣFe from the values calculated using any of the three algorithms against wt% $H₂O$ (or $aH₂O$, calculated at

 1 For a useful comparison of a variety of pre-2016 algorithms for the relationship between Fe^{3−}/ΣF*e* and *f*O₂ see Putirka (2016).

Fig. 5 Tests of the ability of three widely used algorithms to recover Fe³⁺/ΣFe from experiments at know temperature and fO_2 . **A** Kress and Carmichael (1991); **b** Borisov et al. (2018); **c** O'Neill et al. (2018). Only Kress and Carmichael (1991) include a pressure term.

Plotted are new data from Table 3 (red symbols), alongside published data for hydrous basalts (grey symbols). A total of 151 glasses are plotted; the average absolute deviation (*aad*) from the analysed values is given in each panel

experimental conditions using the method of Burnham 1979) does not reveal any systematic behaviour. This fnding is in agreement with previous studies that conclude that dissolved H₂O has negligible influence on Fe³⁺/ΣFe in silicate melts (e.g., Sisson and Grove 1993; Moore et al. 1995; Botcharnikov et al. 2005). If there is an influence on Fe³⁺/ΣFe of H_2O it is subtle and non-systematic, and requires further, targeted experimental exploration.

It is noteworthy that the two persistent outliers in all panels of Fig. 5 are both highly oxidised glasses from our dataset with > 8wt % H₂O (HAB-20, HAB-23). These are the wettest glasses yet analysed for $Fe^{3+}/\Sigma Fe$, so it is unclear whether their failure to lie on the 1:1 line refects an efect of very high H_2O on $Fe^{3+}-Fe^{2+}$ equilibria or a consequence of analytical challenges in such unstable glasses even using SMS. We conclude that, except possibly for very H_2O -rich basaltic glasses, there is no appreciable interaction between Fe species and H_2O , such that the algorithms of Borisov et al. (2018), O'Neill et al. (2018) and Kress and Carmichael (1991) can be used with some confdence. For reference, in a typical basalt the average absolute deviation of ± 0.05 in calculated Fe³⁺/ΣFe equates to approximately \pm 0.6 log units of $fO₂$ at NNO + 1 and 1250 °C. This gives an indication of the accuracy available to recover $fO₂$ from a hydrous mafic melt with known $\text{Fe}^{3+}/\text{\Sigma}\text{Fe}$ and equilibration temperature. In all subsequent calculations we will adopt, for convenience,

Fig. 6 Fe–Mg exchange between olivine and melt in new experiments and published studies given in Supplementary Table 2. **a**
Kd_{Fe}*T*_{−Mg} versus measured Fe³⁺/ΣFe showing the diluting effect of Fe³⁺. For clarity we distinguish the experiments of Stamper et al. (2014) that are included in Table 2. The line has a slope and inter-

the Borisov et al. (2018) algorithm, although, according to Fig. 5, the results would be broadly similar for the other two algorithms, as well as that of Putirka (2016).

Olivine‑melt partitioning of Fe2+, Mn and Mg

Olivine-melt KdFe*T*−Mg shows a considerable range in our experimental dataset, as would be predicted from the presence of Fe³⁺ in the glass. As expected, Kd_{Fe} r _{−Mg} decreases with increasing Fe³⁺/ΣFe as the availability of Fe²⁺ diminishes and, as a consequence, olivine becomes more Fo-rich (Fig. 6a). A weighted ft to Eq. (4) yields

cept of 0.3135 ± 0.0011 . Kd_{Fe²⁺−Mg} versus **b** temperature; **c** pressure; and **d** olivine forsterite content (mol%). The dashed lines are shown for guidance only. Note the two persistent outlier experiments (runs BM34 and D-7) in all panels; labelled in **c**

$$
Kd_{Fe^T - Mg} = 0.3135 \pm 0.0011 \times \left(1 - \frac{Fe^{3+}}{\Sigma Fe}\right).
$$
 (5)

This is consistent with a simple dilution of Fe^{2+} by Fe^{3+} in the melt, in the manner hinted at from the data of Roeder and Emslie (1970; Fig. 1). The linearity of the relationship across a very wide range of $\text{Fe}^{3+}/\text{\textstyle{\Sigma}}\text{Fe}$ argues strongly against significant non-ideal interactions between $Fe²⁺$ and $Fe³⁺$ in the melt, of the type proposed by Jayasuriya et al. (2004). The intercept at Fe³⁺/ΣFe equates to Kd_{Fe²⁺-Mg} and is in excellent agreement with the canonical value of 0.30 ± 0.03 of Roeder and Emslie (1970). This observation Mysen (2006)

0.35

0.40

Fig. 7 Comparison of $Kd_{Fe^{2+}-Mg}$ calculated using the algorithm of Toplis (2005) with the experimental data in Table 3. The 1:1 line is shown. The two persistent outliers in our dataset (runs BM34 and D-7) are highlighted. Error bars are 1 sd for the experimental measurements and ± 0.02 for the Toplis (2005) method

confirms Roeder and Emslie's (1970) finding that $fO₂$ has little effect on $Kd_{Fe^{2+}-Mg}$. Consequently, the measured value of KdFe*T*−Mg can be used to recover Fe3+∕ΣFe, provided that the potential role of olivine non-ideality can be disregarded or suitably accounted for. We return to this issue in a later section.

We explore the influence of other parameters on $\text{Kd}_{\text{Fe}^{2+}-\text{Me}}$ in Fig. 6b–d. There is no discernible influence of temperature (Fig. 6b) suggesting that over the studied temperature range the change in free energy of reaction (1), ΔG_1 , is relatively small. According to the standard state free energy of fusion data compiled by Toplis (2005) an increase in temperature from 1100 to 1500 °C would equate to an increase in ΔG_1 of 2.9 kJ mol⁻¹, corresponding to a 14% increase in $Kd_{Fe^{2+}-Mg}$, broadly consistent with that observed. There is also no effect of dissolved H_2O (as previously discussed) or pressure (Fig. 6c). Using the volume change for reaction (1), ΔV_1 , as presented by Toplis (2005), we would expect very little change in $Kd_{Fe^{2+}-Mg}$ over the pressure range studied. For example, at 1200 °C, $\overline{Kd}_{Fe^{2+}-Me}$ should increase by just 4% from one atmosphere to 1400 MPa, again broadly consistent with that observed.

Our observations raise the possibility that much of the pressure dependence claimed by Ulmer (1989) results simply from the efect of pressure on ferric–ferrous ratio in basaltic melts. For a given fO_2 , at higher pressures Fe²⁺ in the melt is stabilised relative to $Fe³⁺$ up to at least 3 GPa (Kress and Carmichael, 1991). Thus, as pressure increases the availability of Fe^{2+} will increase, leading to an increase in KdFe*T*−Mg. This is precisely the efect seen by Ulmer (1989); from 1 atmosphere to 3 GPa his experiments show an increase in Kd_{Fe}^{*T*}−Mg</sub> from 0.30 to 0.37. Assuming that $fO₂$ in these experiments remains constant, relative to a solid state buffer such as NNO, the $Fe³⁺/\Sigma Fe$ ratio, calculated using Kress and Carmichael (1991) for Ulmer's (1989) starting material and a reference temperature of 1200 °C, ranges from 0.191 at atmospheric pressure to 0.105 at 3 GPa. From the one atmosphere value of 0.30, this equates to an increase in Kd_{Fe} r_{-Mg} to 0.33, approximately 50% of the observed increase. However, Ulmer's (1989) higher pressure experiments used graphite-lined Pt capsules, and consequently are signifcantly more reduced than NNO, approximating the iron-wüstite buffer. Thus, the higher pressure melts will contain negligible $Fe³⁺$, leading to a further increase in Kd_{Fe}^{*T*}−Mg</sub> to 0.37, relative to the 1 atmosphere value of 0.30, as observed. We suggest that the apparent pressure increase in Kd_{Fe}^T_{−Mg} observed by Ulmer (1989) is a combination of two effects: the decrease in $\text{Fe}^{3+}/\text{\Sigma}\text{Fe}$ that occurs with increasing pressure at fixed temperature along an $fO₂$ buffer; and the more reduced nature of his higher pressure experiments compared to those at one atmosphere.

Finally, there is a small negative dependence of $\text{Kd}_{\text{Fe}^{2+}-\text{Mo}}$ on olivine composition (Fig. 6d), consistent with modest non-ideality on the forsterite–fayalite join, as previously noted by Toplis (2005). For a symmetrical regular solution, characterised by a binary interaction parameter $W_{\text{Fe}-\text{Mg}}$, Kd_{Fe²⁺–Mg} will decrease from Fo₅₀ towards both the fayalite (Fo₀) and forsterite (Fo₁₀₀) end-members. With a value of $W_{Fe-Mg}=2.6 \text{ kJ mol}^{-1}$, as determined experimentally O'Neill et al. (2003), Kd_{Fe²⁺−Mg} (at 1200 °C) would be expected to decrease by ~16% from Fo_{60} to Fo_{100} . This is a small effect, but not inconsistent with the variation seen in Fig. 6d. In a subsequent section we incorporate olivine nonideality into our expression for $Kd_{Fe^{2+}-Mg}$.

The fact that Eq. (5) holds across a wide range in pressure, temperature and melt composition (especially H_2O) contrasts with the strong melt compositional dependence of $\text{Kd}_{\text{Fe}^{2+}-\text{Mg}}$ that was proposed by Toplis (2005) and Putirka (2016) based on large published experimental datasets. In Fig. 7, we compare our measured values of $\text{Kd}_{\text{Fe}^{2+}-\text{Mg}}$ to those calculated using the method of Toplis (2005) at our experimental conditions. There is broad agreement (within about ± 0.05 in $\text{Kd}_{\text{Fe}^{2+}-\text{Me}}$) despite the fact that Toplis (2005) does not include an explicit compositional term for $Fe³⁺$ in his parameterisation. Instead he used the parameterisation of Kilinc et al. (1983) to calculate $\text{Fe}^{2+}/\text{Fe}^{3+}$ at the run T–*f*O₂ conditions for his experimental dataset. Kilinc et al's (1983) parameterization was updated by Kress and Carmichael (1991) using additional experimental data, with signifcantly diferent temperature and composition parameters as a result. It is unclear to what extent

Fig. 8 Mn–Mg exchange between olivine and melt in new experiments and published studies given in the text—for clarity we distinguish the experiments of Stamper et al. (2014) and Melekhova et al.

(2015). Kd_{Mn−Mg} versus: **a** measured Fe³⁺/ΣFe; **b** fO_2 (expressed as log units relative to NNO buffer); c temperature; and **d** olivine forsterite content (mol%). The dashed lines are shown for guidance only

some of the compositional dependence $\mathrm{Kd}_{\mathrm{Fe}^{2+}-\mathrm{Mg}}$ identified by Toplis (2005) is aliased to compositional dependence of $Fe^{2+}/$ $Fe³⁺$ above and beyond that captured by the parameterisation of Kilinc et al. (1983). In this respect, it is notable that Toplis' (2005) compositional terms include a strong infuence of molar $Na₂O$ and $K₂O$ in the melt, the two components known to exert the greatest influence on $\text{Fe}^{2+}/\text{Fe}^{3+}$ in melts at fixed P, T and $fO₂$ (Kress and Carmichael 1991; Borisov et al. 2018). It is very difficult to disentangle these two effects in the absence of measured Fe³⁺/ΣFe in experimental glasses.

In contrast to Fe–Mg exchange between olivine and melt, we would not expect Mn–Mg exchange to be redox-sensitive under most terrestrial fO_2 conditions. In Fig. 8, we plot Kd_{Mn-Mg} for our experimental dataset against a variety of intensive and compositional parameters. As expected, a strong variation with Fe³⁺/ΣFe (Fig. 8a) is not observed, because Mn is divalent across almost all of the $fO₂$ range considered (Stokes et al. 2019). Nonetheless, at the highest $Fe^{3+}/\Sigma Fe$, i.e., the most oxidised conditions, any Mn that is trivalent should lead to a small reduction in $\text{Kd}_{\text{Mn}-\text{Mg}}$. In Fig. 8b, we show Kd_{Mn−Mg} plotted against $\triangle NNO$ (the deviation in log units relative to the NNO buffer at P and T). Any reduction in $\text{Kd}_{\text{Mn}-\text{Mg}}$ at the highest ∆NNO is evidently very small, consistent with the observation that $Mn^{3+}/\Sigma Mn$ in lime–alumina–silica melts is low, even in air (Tamura et al. 1987). Moreover, $Mn^{3+}/\Sigma Mn$ decreases signifcantly with decreasing molar Na/Si ratios (Schreiber et al. 1994), such that for basalts (Na/Si \approx 0.1) equilibrated in air at~1150 °C Mn³⁺/ΣMn will be less than

Fig. 9 Models for Fe–Mg and Mn–Mg exchange coefficients. a Kd_{Fe}^{*T*}−Mg</sub> (after correction for olivine non-ideality, as described in text) versus Fe³⁺/ΣFe. Data labelled 'this study' include experiments presented in Table 2 plus those additionally plotted in Fig. 6a for which $\text{Fe}^{3+}/\text{\textstyle{\Sigma}}\text{Fe}$ is independently measured (see Supplementary Table 2). 'Literature data' are taken from the compilation of Matzen et al. (2011), whence original sources can be found, plus additional data from Laubier et al. (2014), Gaetani and Grove (1997), Mallman and O'Neill (2009, 2013), Canil (1997), and Canil and Fedortchouk

0.02. Thus the effect of redox on $\text{Kd}_{\text{Mn}-\text{Mg}}$ is expected to be minimal under all conditions investigated experimentally and likely to be found in nature (e.g., Stokes et al. 2019).

Kd_{Mn−Mg} shows minor dependence on temperature (Fig. 8c) and olivine composition (Fig. 8d), as expected for trace element substitution. To a large extent these efects are interlinked in the experimental dataset and should be considered in combination. Mn^{2+} exchange for Fe²⁺ or Mg^{2+} in olivine is likely to be highly non-ideal, due to the large sixfold ionic radius diferences (Shannon, 1976) between Mn^{2+} (0.83 Å), Fe²⁺ (0.78 Å) and Mg²⁺ (0.72 Å). The ease of inserting an Mn^{2+} ion into the olivine lattice increases with temperature at fxed olivine composition, and with decreasing forsterite content at fxed temperature. This behaviour is consistent with the lattice strain model of trace element partitioning (Blundy and Wood, 1994), as developed further below.

Discussion

The experimental data presented above show the strong dependence of Fe–Mg exchange on $fO₂$, but the relative insensitivity of Mn–Mg exchange to the same parameter. This raises the possibility of using Fe–Mg partitioning systematics to determine $fO₂$ from equilibrium pairs of olivine and silicate melt, for example in high-pressure experiments

(2001). Glass Fe³⁺/ΣFe was not measured in these experiments; the plotted value is calculated using Borisov et al. (2018) at the experimental $T-fO₂$ and, therefore, subject to greater uncertainty than the measured values. The line (red) ftted to the data from this study data has a slope and intercept of $Kd_{Fe^{2+}-Mg}^{0} = 0.3642 \pm 0.0011$. The line (black) through the entire literature dataset has a very similar slope of 0.349±0.002. **b** *RT*lnKd_{Mn - Mg} (in kJ mol⁻¹) plotted against molar fraction Fo in olivine. Data sources as in **a**. The black line is a ft to the lattice strain Eq. (10) with the fit parameters given

where $fO₂$ is hard to control/monitor, and using the redoxinsensitive Mn–Mg exchange as a means to test for olivinemelt equilibrium. Ultimately our goal is to introduce a simple method to correct olivine-hosted melt inclusion compositions for post-entrapment crystallization (using Mn–Mg) and then recover the magmatic $fO₂$ (using Fe–Mg). We are helped in this enterprise by the fact that existing models for ferric–ferrous equilibria in silicate melts as a function of redox (e.g., Kress and Carmichael 1991; Borisov et al. 2018; O'Neill et al. 2018), have proven surprisingly accurate, even for hydrous basaltic compositions (Fig. 5). The frst step is to develop models for Fe–Mg and Mn–Mg exchange between olivine and melt, and then to test these expressions against large datasets of experimental olivine-melt pairs.

A model for Fe–Mg partitioning in basaltic melts

Ignoring, for the moment, non-ideality in the melt, $Fe^{2+}-Mg$ exchange can be described by the following relationship (cf. Toplis 2005):

$$
\mathrm{Kd}_{\mathrm{Fe}^{2+}-\mathrm{Mg}} = \exp\Bigg\{\frac{-\Delta H_1 + T\Delta S_1 - P\Delta V_1 + W_{\mathrm{FeMg}}^{\mathrm{olivine}}\left(1 - 2X_{\mathrm{Fo}}\right)}{RT}\Bigg\},\tag{6}
$$

where the terms ΔH_1 , ΔS_1 and ΔV_1 denote, respectively, the enthalpy, entropy and volume changes of reaction (1), and $W_{\text{FeMg}}^{\text{olivine}}$ is a symmetrical regular solution interaction parameter for the forsterite–fayalite binary solid solution. *W*^{olivine} has been constrained experimentally to be in the region of a few kJ mol⁻¹; we will adopt the value of 2.6 kJ mol⁻¹ of O'Neill et al. (2003). As noted above, our experimental data provide some constraints on ΔH_1 , ΔS_1 and ΔV_1 over the P–T–X studied. Given the limited temperature dependence (Fig. 6b) ΔH_1 is expected to lie close to zero. We did explore the use of ΔH_1 = +6.766 kJ mol⁻¹ derived from the free energy of fusion of forsterite and fayalite (Toplis 2005),² but it led to no improvement in our modelled fts, suggesting that the effective ΔH_1 is smaller than this value, probably due to fortuitous cancelation by non-ideal interactions of $Fe²⁺$ and Mg in the melt. Beattie (1993), for example, obtains a value for ΔH_1 of + 3.2 ± 0.8 kJ mol⁻¹ using experimental olivinemelt pairs. The limited pressure dependence (Fig. 6c) suggests that ΔV_1 is also small; Toplis (2005) proposes a value of − 0.35 kJ GPa⁻¹ mol⁻¹, while Beattie (1993) obtains a value of -0.09 ± 0.11 kJ GPa⁻¹ mol⁻¹ from experimental olivine-melt data. The remaining term, ∆*S*1, embraces the entropy of mixing of Fe^{2+} and Mg in both olivine and melt. The quantity $exp(\Delta S_1/R)$ in Eq. (6) equates to the 'ideal' value of $\text{Kd}_{\text{Fe}^{2+}-\text{Mg}}$ (termed $\text{Kd}_{\text{Fe}^{2+}\text{Mg}}^{0}$) prior to any correction for olivine non-ideality. The value of $exp(\Delta S_1/R)$ using the derived thermodynamic data of Beattie (1993) is 0.121; the equivalent value from Toplis (2005) is 0.414. We can use our experimental data to derive a value for $\text{Kd}_{\text{Fe}^{2+}-\text{Mg}}^{0}$ by modifying Eq. (6) to

$$
Kd_{Fe^{2+}-Mg} = Kd_{Fe^{2+}-Mg}^{0} \times \exp\left(\frac{\Delta W_{Fe^{2+}Mg}^{\text{olivine}}(1-2X_{Fo})}{RT}\right), \quad (7)
$$

where $\text{Kd}_{\text{Fe}^{2+}-\text{Mg}}^{0}$ can then be obtained from a weighted regression of Kd_{Fe}^{*T*}−Mg</sub>/ex*p* $\left(\frac{\Delta W_{\text{Fe}^{2+}Mg}^{\text{olivine}}(1-2X_{\text{Fe}})}{RT} \right)$ against Fe3+∕ΣFe. For the 52 experiments with known Fe3+∕ΣFe the regressed value of $\text{Kd}_{\text{Fe}^{2+}-\text{Mg}}^{0}$ is 0.3642 ± 0.0011 at $Fe³⁺/\Sigma Fe = 0$ (Fig. 9a). To obtain this value we eliminated two persistent outliers in the dataset: experiments D-7 of Stamper et al (2014) and BM34 of Melekhova et al (2015). Note that for a symmetrical, regular solution $\text{Kd}_{\text{Fe}^{2+}-\text{Mg}}^{0} = 0.3642$ gives the Fe²⁺–Mg exchange coefficient at $F_{\sigma_{50}}$; it will be lower for olivines with higher or lower Fo than 0.5. This accounts for the diference from the value of $\text{Kd}_{\text{Fe}^{2+}-\text{Mg}} = 0.3135$ (Fig. 6a) presented above, which takes no account of olivine non-ideality. For example, for a Fo_{70}

olivine at 1200 °C, the bracketed term in Eq. (7) equals 0.92; for Fo_{90} at 1450 °C it equals 0.86. It is worth noting that the strong dependence of the olivine-melt Mg partition coefficient (D_{Mg}) on olivine Fo content will lead to an apparent negative correlation between $\text{Kd}_{\text{Fe}-\text{Mg}}$ and D_{Mg} , such as that observed by Herzberg and O'Hara (2002). Evidently, any such correlation should be carefully evaluated in the light of olivine non-ideality, as discussed above in the context of Fig. 6c.

An assumption in using Eq. (7) to describe $Fe^{2+}-Mg$ exchange between olivine and melt is that melt composition does not play a role, i.e., a single value of $Kd_{Fe^{2+}-Mg}^{0}$ captures the full variation in the data provided that $Fe^{3+}/\Sigma Fe$ is known (measured) and can be treated as an independent variable. Toplis (2005) and Putirka (2016) have argued, on the basis of an experimental dataset for which $\text{Fe}^{3+}/\text{\Sigma}\text{Fe}$ was not measured, that additional compositional terms are important. Unfortunately, the experimental, olivine-bearing dataset with measured glass Fe³⁺/ΣFe is too small ($n=55$) to explore fully such additional melt compositional effects. Thus, as noted by Matzen et al. (2011) in their Appendix, some of the compositional dependence in the $\text{Kd}_{\text{Fe}^{2+}-\text{Mg}}$ expressions of Toplis (2005) may arise through any compositional dependence of melt $\text{Fe}^{3+}/\text{2Fe}$ not fully captured by the algorithm of Kilinc et al. (1983). To illustrate this point, we show also in Fig. 9a a database of over 1000 experimental olivine-melt Kd_{Fe}^{*T*}−Mg, for which fO_2 is known but Fe³⁺ /ΣFe is not. The data are fitted to an adapted version of Eq. (5) that allows for olivine non-ideality, using the expression in Eq. (7), to give

$$
Kd_{Fe^T - Mg} = 0.3642 \times \left(1 - \frac{Fe^{3+}}{\Sigma Fe}\right) \times \exp\left(\frac{312.7(1 - 2X_{Fo})}{T}\right),\tag{8}
$$

where T is in Kelvin and X_{F_0} is the molar fraction for sterite in olivine. In fitting this equation, we calculated $Fe^{3+}/\Sigma Fe$ at the experimental $T-\gamma O_2$ conditions using the expression of Borisov et al (2018). The bulk of the data lie close to the regression line; 57% of all of the calculated Kd_{Fe}^{*T*}−Mg</sub> values lie within 0.02 of the experimental values, and the average absolute deviation (*aad*) of all 1135 experiments is 0.030. If we consider only the natural compositions, i.e., those with non-zero alkalis, *aad* reduces to 0.022. The *aad* is also dependent on the total alkali contents ($Na₂O + K₂O$). For example, melts with Na₂O + K₂O ≥ 6 wt% have *aad* = 0.040, while those with $Na₂O + K₂O \ge 9$ wt% have *aad* = 0.063. This behaviour cannot be attributed to olivine non-ideality; it must instead be the result of either non-ideality in the melt [the interpretation preferred by Toplis (2005), and Putirka (2016)] or a failure of the Borisov et al (2018) algorithm to capture fully the efect of alkalis on Fe3+∕ΣF*e*. In all likelihood the deviation is a combination of these two factors, but without determinations of $\text{Fe}^{3+}/\text{\textstyle{\Sigma}}\text{Fe}$ in these

 2 Note the change in sign because Toplis (2005) provides data for the reverse of reaction (1).

alkali-rich experiments it is not possible to disentangle the two effects. We conclude that Eq. (8) is best suited to basalts with between 0 and 8 wt% $Na_2O + K_2O$ (*aad* = 0.019); at higher alkali contents its precision is reduced slightly. As noted by Toplis (2005), the analytical precision of Kd_{Fe}^{*T*}−Mg. determinations in most experiments is of the order 0.02.

The observed scatter in the data at low $Fe^{3+}/\Sigma Fe$ (Fig. 9a) is curious. It cannot arise through olivine nonideality or temperature. It may arise instead from the failure of the Borisov et al (2018) algorithm to capture precisely Fe³⁺/ΣFe at low fO_2 , as suggested by Fig. 5b. Such a situation likely has its origins in the difficul of determining $Fe^{3+}/\Sigma Fe$ precisely by spectroscopic or wet chemical means when the $Fe³⁺$ content is very low (O'Neill et al. 2018). This problem warrants further investigation, with additional data on Fe³⁺/ΣFe at very reducing conditions. For the time being we consider that Eq. (8) provides an adequate description of Kd_{Fe}^{*T*}−Mg</sub> for basaltic magmas under most terrestrial redox conditions.

The advantage that Eq. (8) offers over that of Toplis (2005) is that it contains no empirical ft parameters to describe the effect of melt non-ideality on Kd_{Fe}^{*T*}−Mg</sub>. Any

advocate caution in applying Eq. (8) at pressures higher than this. Likewise, as noted above, we do not see an apparent increase of $\text{Kd}_{\text{Fe}^T-\text{Mg}}$ with melt SiO₂ content, in contrast to Toplis (2005) and Putirka (2016). This again likely refects the limited compositional range of experimental glasses from our olivine-bearing experiments $(45.9-56.9 \text{ wt\% SiO}_2)$ on an anhydrous basis; Table 4) and those in the other data listed in Supplementary Table 2. For this reason we would not recommend using Eq. (8) for olivine-bearing andesitic liquids with more than ~ 60 wt% SiO₂.

A model for Mn–Mg partitioning in basaltic melts

Unlike Fe^{2+} , Mn is a minor (or trace) element in most natural olivines, allowing for a diferent thermodynamic treatment. Based on the observations in Fig. 8, we will deploy the lattice strain model for trace element partitioning (Blundy and Wood 1994), as recast for the case of a 'proxy' element (Eq. 8 of Blundy and Wood 2003). In this formulation, which is well suited to describing exchange coefficients (*Kd*s) between trace and major elements, the proxy is Mg and the trace is Mn. This yields

$$
Kd_{Mn-Mg} = \exp\left[\frac{-4\pi N_A E_M^{2+}}{RT} \left\{ \frac{r_{0(M)}^{2+}}{2} \left(r_{Mg}^2 - r_{Mn}^2 \right) - \frac{1}{3} \left(r_{Mg}^3 - r_{Mn}^3 \right) \right\} \right],
$$
\n(9)

melt compositional dependencies arise exclusively from those in the algorithm used for $Fe^{3+}/\Sigma Fe$, in this case Borisov et al (2018); the term for olivine compositional dependence is independently derived. Equation (8) offers a straightforward means to determine the composition of olivine crystallising from a melt at known temperature and $fO₂$. Conversely, from a measured value of Kd_{Fe}^{*T*}−Mg</sub> at known temperature and olivine composition (X_{F_0}) , Fe³⁺/ΣFe (and $fO₂$) can be recovered by rearrangement of Eq. (8) and use of an algorithm such as that of Borisov et al (2018) or Putirka (2016).

Our data do not suggest a pressure dependence of Kd_{Fe}^{*T*}−Mg</sub> in contrast to several previous studies (e.g., Ulmer 1989; Herzberg and O'Hara 1998; Putirka 2005, 2016; Toplis 2005). We speculate above that changing Fe³⁺/ΣFe in the melt with increasing pressure may account for some apparent pressure dependence. Unfortunately there are no high pressure experimental data with measured $\text{Fe}^{3+}/\text{2Fe}$ at known *f*O₂ with which to test this idea, and our experimental dataset does not extend above 1.7 GPa. For this reason we would

where N_A is Avogadro's number, R is the gas constant, T is temperature in Kelvin, $E_{\rm M}^{2+}$ is the effective Young's Modulus of the VI-fold M-site in olivine, and $r_{0(M)}^{2+}$ is the optimum radius of that site. r_{Mg} and r_{Mn} are the ionic radii of Mg²⁺ and Mn^{2+} in VI-fold coordination, 0.72 and 0.83 Å, respectively (Shannon 1976). The ideal M-site cation radius changes along the forsterite–fayalite solid solution, because Fe²⁺ (0.78 Å) is larger than Mg²⁺ (0.72 Å). Thus $r_{0(M)}^{2+}$ is not constant. In principle we can use the experimental data, from this study and from the literature, to derive values for $E_{\rm M}^{2+}$ and for $r_{0(M)}^{\tilde{2}+}$ as a function of olivine composition. However, as noted by Wood and Blundy (1997), there is considerable trade-off between these two parameters from least squares fitting. For that reason, we fix $E_M^{2+} = 150$ GPa, following Beattie (1994) and Purton et al. (2000), and fit only for $r_{0(M)}^{2+}$. We do this by assuming that $r_{0(M)}^{2+}$ increases linearly from forsterite $(r_{0(F_0)}^{2+})$ to fayalite $(r_{0(F_a)}^{2+})$. Thus, Eq. (9) can be rewritten and linearised in such a way that a plot of RT lnKd_{Mn−Mg} versus X_{F_0} can be used to obtain $r_{0(\text{F}_0)}^{2+}$ and $r_{0(Fa)}^{2+}$.

$$
RT\ln \mathrm{Kd}_{\mathrm{Mn-Mg}} = -150 \times 4\pi N_{\mathrm{A}} \left\{ \left(\frac{X_{\mathrm{Fo}} r_{0(\mathrm{Fo})}^{2+} + (1 - X_{\mathrm{Fo}}) r_{0(\mathrm{Fa})}^{2+}}{2} \right) \left(r_{\mathrm{Mg}}^{2} - r_{\mathrm{Mn}}^{2} \right) - \frac{1}{3} \left(r_{\mathrm{Mg}}^{3} - r_{\mathrm{Mn}}^{3} \right) \right\}.
$$
 (10)

Fig. 10 Fully propagated uncertainty estimates for ∆NNO calculations using Kd_{Fe}^{*T*}−Mg</sub>. Calculated values are strongly dependent on ∆NNO, and greatly magnified at low fO_2 . Shown are curves for the typical, propagated analytical uncertainty (red), assuming 5% relative error in Kd_{Fe}^{*T*}−Mg</sub>, and the total uncertainty (blue) that includes 0.38 log units in fO_2 as given by Borisov et al. (2018) for their Fe³⁺/ΣFe algorithm. The shaded region illustrates an uncertainty of ± 0.6 log units in fQ_2

Using the literature experimental dataset described above, as well as the new data presented here, we fnd that this equation describes well the Mn–Mg exchange coefficient for the entire experimental dataset (Fig. 9b), returning values of $r_{0(Fo)}^{2+}$ = 0.565 ± 0.002 Å and $r_{0(Fa)}^{2+}$ = 0.713 ± 0.007 Å. A more complex expression, that takes into account the possible temperature dependence of $r_{0(F_0)}^{2+}$ and $r_{0(F_a)}^{2+}$ is not warranted. There is no evidence for strong melt compositional dependence of *RT*lnKd_{Mn−Mg}; most of the scatter in Fig. 9b arises through analytical uncertainty on the MnO contents of experimental olivines, which are more difficult to measure precisely than major elements (for instance MnO is usually reported to two signifcant fgures compared to four for MgO). For example, a typical uncertainty on Kd_{Mn−Mg} of 15% relative, translates to an uncertainty of ± 2 kJ mol⁻¹ in RT lnKd_{Mn−Mg}, which encompasses over 80% of the observed scatter in Fig. 9b. Note that this approach takes into account both the temperature and olivine composition dependence of $\text{Kd}_{\text{Mn-Mg}}$, in contrast to the expression of Matzen et al (2017) that considers only the efect of melt MgO content. For olivine-saturated systems, the latter is, of course, a function of both temperature and olivine composition.

Regarding the infuence of melt composition on Mn–Mg partitioning, there is no evidence for such an efect in our basaltic experimental dataset. In contrast, Kohn and Schofeld (1994) in their study of Mn partitioning between olivine (Fo_{100}) and melt in the system forsterite–albite–anorthite show that *RT*lnKd_{Mn−Mg} varies with the degree of melt polymerisation, increasing from~ 0.12 to 0.23 over a wide range in NBO/T (0.09–0.37). However, the overall variation in RT lnKd_{Mn−Mg} for their experiments is -20.6 to

Fig. 11 Calculated ∆NNO for olivine-melt pairs from Table 2 and in the dataset of published experiments introduced in Fig. 8a. Only experiments run at $fO₂$ above NNO -4.5 are plotted. The 1:1 line and error envelopes from Fig. 10 are shown as black lines

 $-$ 22.9 kJ mol⁻¹, compared to the value of $-$ 20.4 kJ mol⁻¹ calculated for Fo_{100} olivine using Eq. (10). The mean deviation of Kohn and Schofeld's (1994) data from this value is only \pm 1.3 kJ mol⁻¹, despite the wide range in NBO/T. This gives a sense of the relatively small infuence of melt composition on RT lnKd_{Mn−Mg} compared to olivine composition, which changes by 4.3 kJ mol⁻¹ from Fo₁₀₀ to Fo₇₀ (Fig. 9b).

Applications

Equation (8) obviously has wide utility in modelling the crystallisation and melting of olivine-bearing systems. By way of example, we focus on two novel oxybarometry applications using the derived expressions for Kd_{Fe}^{*T*}−Mg</sub> and Kd_{Mn−Mg}. The first involves determination of *f*O₂ for experimental, olivine-bearing melts using Kd_{Fe}^{*T*}−Mg</sub>; the second involves determination of $\text{Fe}^{3+}/\text{2Fe}$ in olivinehosted melt inclusions that have frst been corrected for post-entrapment modifcation (crystallisation or dissolution) by the host olivine using Kd_{Mn-Mg} .

Oxygen fugacity estimates from olivine‑bearing experiments

Equation (8) can be rearranged to obtain $\text{Fe}^{3+}/\text{\textstyle{\Sigma}}\text{Fe}$ of an experimental glass in equilibrium with olivine of a known **Table 6** Calculated fO_2 for Lesser Antilles experiments using olivine Fe–Mg exchange oxybarometry

a ∆NNO reported in the original source using the method indicated: µXANES; olivine-spinal oxybarometry; Fe in AuPd alloy

b ∆NNO calculated using method described in text in combination with algorithm of Kress and Carmichael (1991) or Borisov et al. (2018)

 $\text{c}_{\text{wt}}\%H_2O$ by SIMS, except where calculated by mass balance (italics). Figure in parentheses is 1 sd in terms of least significant figures

composition at a known (experimental) temperature. This value of Fe³⁺/ΣFe can then be used to derive fO_2 using the algorithm of Borisov et al (2018), which we adopt here in light of its slightly greater ability to recover experimental fO_2 for our sample set (Fig. 5b). Determination of fO_2 for experiments is particularly useful at high pressure, when bufering redox and/or preventing the egress or ingress of hydrogen are perennial problems.

Fig. 12 Calculated ∆NNO for 40 high-pressure experiments on hydrous Lesser Antilles basalts from Melekhova et al. (2015) noted as M15 and Stamper et al. (2014) noted as S14. The horizontal axis shows the values determined from olivine-melt Fe–Mg exchange, while the vertical axis shows values obtained by three independent methods, described in the text and reported by the original authors. µXANES analyses suffer from oxidation due to the high photon flux used (see also Fig. 4)

Before testing our approach against the experimental olivine-melt pairs produced at known redox conditions, it is instructive to gauge the likely precision of the method. The uncertainty on the calculation is not linear with $fO₂$. The method relies, in effect, on the difference between Kd_{Fe}^{*T*}−Mg and $Kd_{Fe^{2+}-Mg}$, after first accounting for olivine non-ideality. As Fe³⁺/ΣFe approaches zero, so Kd_{Fe}^{*T*}−M_g approaches $Kd_{Fe^{2+}-Mg}$. This leads to significant uncertainty on Fe3+∕ΣFe, which propagates to considerable uncertainty on calculated $fO₂$. The same problem is not encountered under more oxidising conditions, when $\text{Fe}^{3+}/\text{\Sigma}\text{Fe}$ approaches 1 and Kd_{Fe}^{*T*}−Mg</sub> tends to zero. To capture this non-linearity in the uncertainty we have propagated the errors from the various fts performed on the data through the analytical uncertainty, assuming a 5% relative uncertainty on the measured Kd_{Fe} $r_{\text{–Mg}}$, as well as that arising from the Borisov et al (2018) calibration (0.38 log units in $fO₂$). The resulting behaviour is shown in Fig. 10. As expected, uncertainty is greatly magnifed under reducing conditions, exceeding 0.6 log units around NNO − 1, and clearly becoming meaningless below NNO − 4. The uncertainty contracts considerably at more oxidising conditions, attaining the limit provided by the Borisov et al (2018) calibration above NNO + 2. Over almost all of the $fO₂$ range, this calibration uncertainty dominates over that due to analytical precision, emphasising the

need for continuous refinement of the Fe³⁺/ΣFe algorithm, especially at very low fO_2 .

In Fig. 11, we plot the calculated $fO₂$ (relative to NNO) for all experiments in the dataset conducted above NNO − 4; we eliminate only the subset of data where Kd_{Fe}^{*T*}−Mg>Kd_{Fe}²⁺−Mg equating to a spurious negative value of $Fe^{3+}/\Sigma Fe$. The observed behaviour is in line with the expectations from Fig. 10, namely that precision deteriorates dramatically below NNO − 1. Nonetheless, in 72% of all experiments at NNO -2 and above we recover $fO₂$ to within ± 1.2 log units; 44% are recovered within \pm 0.6 log units. In addition to the factors discussed above, uncertainty also accrues from the effects of alkalis on Kd_{Fe}^{*T*}−Mg, either directly through melt non-ideality, or indirectly through the algorithm used to recover fO_2 from Fe³⁺/ΣFe. Once again we emphasise that our method is best applied only to relatively oxidised melts (above~NNO – 1) with less than about 8 wt% $Na₂O+K₂O$.

We test our approach on a set of real high-pressure experiments where the fO_2 was either not constrained directly, e.g., by solid-state bufers, or may have been compromised by hydrogen difusion through the capsule walls. We have selected 40 experiments from Stamper et al. (2014) and Melekhova et al. (2015) on magnesium-rich basalts from the Lesser Antilles with varying amounts of added H_2O (Table 6). In these experiments various alternative strategies were employed to recover fO_2 , including μ XANES, olivine-spinel oxybarometry (O'Neill and Wall 1987), and the solubility of iron in AuPd alloy capsules (Barr and Grove 2010). In Fig. 12, we compare these diferent methods. As expected, we fnd that the µXANES analyses of the glasses from these hydrous experiments yield higher $fO₂$ estimates than all other methods, due to the oxidative efects of the high photon fux used in the analyses. We can also see the wide variation in $fO₂$ that results from varying amounts of Fe redox due to hydrogen difusion, a phenomenon hard to avoid in most piston-cylinder experiments where, unlike with internally-heated gas-pressure vessels, fH_2 cannot be easily maintained without external (solid-state) buffers. The methodology developed here reproduces very well the estimates from olivine-spinel barometry in those runs where spinel was large enough to analyse. The Fe–in–AuPd approach also yields results that lie close to our new method, albeit with much greater scatter, primarily due to the difficulty of analysing Fe at very low concentrations in Au–Pd close to the melt–capsule interface where secondary fuorescence can be a problem. We conclude that our olivine-melt oxybarometer has the potential to recover $fO₂$ from olivine-bearing experiments that are either lacking spinel or where the runproduct spinel is too small to analyse. However, care must be taken to screen for Fe-loss to the capsule material during the experiment. This problem is especially acute under reducing conditions and in Au alloy or Pt capsules (Barr and Grove

Fig. 13 Illustration of the effect of olivine addition on calculated $fO₂$ (expressed as ∆NNO) for a melt inclusion from St. Vincent (MI14 of Bouvier et al. 2008). The optimal correction to match the calculated Kd_{Mn−Mg} is – 0.5 wt% olivine (i.e., olivine dissolution). An analytical uncertainty of $\pm 10\%$ relative on this value is shown as the grey bar, which translates to an uncertainty in fO_2 of \pm 0.7 log units. With increasing olivine addition the calculated $fO₂$ drops significantly, emphasising the need for precise Mn measurements of olivine and glass for this method

 2010). Uptake of metallic iron (Fe⁰) by the capsule occurs via disproportionation of melt Fe^{2+} to $Fe^{0} + Fe^{3+}$ leading to an apparent increase in glass $\text{Fe}^{3+}/\text{\Sigma}\text{Fe}$ and hence calculated $fO₂$. To ensure reliable application of our oxybarometer the efects of Fe-loss must be fully accounted for.

Fe³+∕*6***Fe estimates from olivine‑hosted melt inclusions**

Olivine-hosted melt inclusions are widely used to determine volatile contents of pre-eruptive magmas, with implications for volatile recycling at subduction zones or magma storage conditions, for example. Most olivine-hosted melt inclusions have experienced some degree of post-entrapment modifcation, for example by crystallisation of olivine on the inclusion wall, dissolution of the olivine host, or difusive exchange between melt and host (Danyushevsky et al. 2002). Standard practice in these situations is to correct the melt inclusion composition for these efects by adding (or subtracting) some amount of host olivine (or FeO and MgO) to the melt inclusion to bring the pair back into equilibrium. This inevitably leads to some modification to Kd_{Fe}^{*T*}−Mg, the parameter required for oxybarometry. To obviate this circularity, we

Fig. 14 Calculated Fe³⁺/ΣFe for basaltic melt inclusions using the technique described in the text for six volcanoes for which $Fe^{3+}/\Sigma Fe$ was determined by μ XANES. The average and sd for each volcano are plotted; the number of individual analyses is given in parentheses. Data sources are: Erebus (Moussallam et al. 2014); Agrigan, Mariana (Kelley and Cottrell 2012; Brounce et al. 2014); Sarigan, Mariana (Brounce et al. 2014); Kilauea (Moussallam et al. 2016); Andaman Islands, and Paricutin, Mexico (Kelley and Cottrell 2009)

propose to use the systematics of redox-insensitive Mn–Mg exchange to perform the post-entrapment corrections, and then use the Kd_{Fe} r _{−Mg} of the corrected melt and its host olivine to recover $Fe^{3+}/\Sigma Fe$. To make the post-entrapment correction requires some knowledge of the equilibrium temperature, which can be obtained, for example, through olivine-melt thermometry (e.g., Beattie 1993; Putirka et al. 2007).

Before testing our approach on natural melt inclusions, it is instructive once again to get a sense of the expected uncertainty of the method. In addition to the other factors considered above, the correction for post-entrapment modifcation adds further uncertainty; the more olivine that is added back into the inclusion, the higher the resulting Kd_{Fe}^{*T*}−Mg</sub> and the lower the calculated Fe³⁺/ΣFe (and hence *f*O₂). To illustrate this we have calculated the effect on calculated $fO₂$ due to varying amounts of olivine addition (and subtraction) from an Mg-rich basaltic melt inclusion (#14) from St. Vincent (Bouvier et al. 2008). We use a temperature of 1200 °C and an initial (measured) Kd_{Fe} r_{-Mg} =0.205 for the calculations (Fig. 13). The calculated, equilibrium $Kd_{Fe^{2+}-Mg}$, after taking account of olivine (Fo_{89}) non-ideality is 0.31; deviation of Kd_{Fe} r_{-Mg} from this value enables us to calculate Fe³⁺/ΣFe and thence $fO₂$. It is clear that addition (or subtraction) of just a few percent olivine can lead to dramatic reduction (or increase) in $\text{Kd}_{\text{Fe}^T-\text{Me}}$ and hence in calculated fO_2 . Using Eq. (10), alongside the reported MnO contents of olivine and melt inclusion, the required correction is a subtraction of just 0.5 wt% olivine, i.e., due to dissolution of the host olivine. The calculated $fO_2(\triangle NNO=2.2)$ accords well with that estimated for St. Vincent high-Mg basalt (∆NNO+1.5 to + 1.8) using spinel oxybarometry (Heath et al. 1998). Within the uncertainty of the measurement of MnO in olivine and glass, a range of olivine corrections is permissible. In Fig. 13, we show that a relative uncertainty of $\pm 10\%$ in Kd_{Mn−Mg} translates to uncertainty of $\pm 1.3\%$ in the olivine correction, or \pm 0.7 log units in $fO₂$. Evidently, having precise measurements of olivine and melt MnO contents is critical for this method, yet is often not a priority in melt inclusion studies. For comparison, using the method of Danyushevsky et al. (2000, 2002), Bouvier et al. (2008) estimate post-entrapment crystallisation of 6.5 wt% olivine, well outside the permissible bounds based on Mn–Mg exchange (Fig. 13).

Having established what is possible with our melt inclusion oxybarometer we have tested it against five olivine-hosted melt inclusions studies from the literature: Brounce et al (2014), Kelley and Cottrell (2009, 2012) and Moussallam et al (2014, 2016). In all of these cases melt Fe³⁺/ΣFe was determined by μ XANES, but using a lower fux than that used in our analyses. For these reasons the measured Fe³⁺/ΣFe is considered more reliable; however, the presence of various amounts sulphur in all of the analysed glasses raises the prospect of Fe redox during melt inclusion quenching through Fe–S redox couples. This is a perennial problem with microbeam analysis of $Fe³⁺/\Sigma Fe$ in sulphur-bearing glasses, but not for our methodology, as olivine-melt Fe–Mg exchange by difusion is unlikely to have taken place on quenching timescales; if it did, there would be evidence in the form of thin overgrowth rims on the walls of the melt inclusion.

To calculate $\text{Fe}^{3+}/\text{2Fe}$ in each of the melt inclusion sets we frst correct for post-entrapment modifcation using Eq. (10) with the measured $\text{Kd}_{\text{Mn}-\text{Mg}}$ and an equilibrium temperature from olivine-melt thermometry (or other petrological method) provided in the original sources. An iterative procedure was used. The corrected melt composition for each inclusion, following olivine addition or subtraction, was then used to calculate $Fe³⁺/\Sigma Fe$, using Eq. (8). For each volcano studied we present the average value of $\text{Fe}^{3+}/\text{\textstyle \Sigma \text{Fe}}$ and its standard deviation in Fig. 14. For five of the studied volcanoes our method reproduces the mean μ XANES Fe³⁺/ΣFe values within 1 s.d. The only exception is Sarigan, where our method consistently underestimates $Fe^{3+}/\Sigma Fe$. Notably these were the melt inclusions that required the greatest post-entrapment calculation (2–7 wt% olivine addition). Nonetheless, our method appears to have an accuracy of about \pm 0.04 in Fe³⁺/ΣFe and captures the broad variation from one sample suite to the next. It is a much simpler analytical method than µXANES (or SMS), requiring much less sample preparation, and can be performed easily on large numbers of melt inclusions, as a means to explore broad-scale variations in Fe³⁺/ΣFe (and fO_2) in space and in time. Moreover, as noted above, it is not susceptible to Fe–S redox exchange during quenching. However, our method does not take account of any difusive Fe–Mg exchange between melt inclusion and host olivine during cooling (e.g., Danyushevsky et al. 2002), a problem that also afflicts in situ measurements of Fe³⁺/ΣFe. As diffusive re-equilibration with the host olivine lowers the $Fe²⁺$ content of the melt inclusion, but not its $Fe³⁺$ content, such a process will lead to erroneously high estimated Fe³⁺/ΣFe (and hence *f*O₂) unless it is explicitly accounted for.

Acknowledgements This project has been carried out over 7 years, during which time many people have provided support and feedback. We thank S. Fawcett for meticulously preparing SMS wafers, C. Stamper for assistance at Diamond in the beginning of the project, A. Matzen for providing his olivine-melt experimental database, E. Cottrell and O. Shorttle for highlighting the issues of µXANES beam damage, and H. O'Neill for helpful discussions. K. Putirka and R. Lange provided exceptionally helpful reviews. We acknowledge the ESRF for provision of beam time at ID18. This work relied on the assistance of beamline scientists A. Chumakov, K. Ignatyev and T. Geraki. JB is grateful for visiting scientist support at University of Western Australia (Robert and Maude Gledden Fellowship), Caltech (Gordon and Betty Moore Scholarship) and Australian National University (ARC Discovery Project DP170103140). JB and EM acknowledge support from NERC project FAMOS (NE/P017371/1) and RB from NERC project DisEqm (NE/N018575/1).

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