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Kev Points:

- Chorus and hiss amplitude statistics is presented based on 11 years of **THEMIS VLF** measurements
- Spatial extents of chorus and hiss waves have been determined in the statistical pattern
- Relation of chorus and hiss waves is shown in the statistical pattern

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Spatial Extent and Temporal Correlation of Chorus and Hiss: Statistical Results From Multipoint THEMIS Observations

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Abstract The spatial localization of whistler mode waves, determined by their generation process, propagation, and damping, is important for assessing the efficiency of wave-particle interactions and the dynamics of the radiation belts. We use 11 years of multipoint wave measurements in 2007–2017 from five Time History of Events and Macroscale Interactions during Substorms (THEMIS) spacecraft covering L = 2-10at low magnetic latitudes over all magnetic local times (MLTs) to characterize both the instantaneous spatial extent of chorus and hiss active regions and their wave amplitude distribution in the vicinity of their respective generation regions, making use of the time domain correlation technique. The correlation of lower band chorus wave power dynamics from two spacecraft measurements remains significant up to 250–800 km transverse to the background magnetic field, in agreement with the typical scale of a chorus element, with a shorter correlation length for higher peak wave power. Additionally, a coordinated analysis of average chorus amplitudes at separate locations shows that the radial extent of the chorus active region varies from 0.2 R_E on the nightside to 1 R_E on the dayside, with an azimuthal width of 1 MLT. The wave power correlation scale of hiss has no fine structure and is generally larger (1,500–3,000 km); its correlation remains nonnegligible up to 1.5 R_E. Such observations are consistent with chorus generation by localized electron injections and subsequent hiss amplification at lower L shells from seed chorus waves having propagated and spread over a wider region inside the plasmasphere. The connection between hiss and chorus is further explored statistically through the correlation between their respective wave power variations, suggesting a significant link between them.

1. Introduction

The dynamics of particles population in the Van Allen radiation belts is controlled by losses (mostly due to precipitation of the particles into the upper atmosphere) and refilling due to external injections, transport, and local acceleration processes. Understanding the dynamics of electron fluxes is of particular importance, since these high-energy particles may pose a threat to satellites whose orbits cross the radiation belts. Favored mechanisms for driving in situ acceleration and pitch angle scattering loss are the interaction of electrons with naturally generated plasma waves in the whistler mode: chorus and hiss (e.g., Horne et al., 2005; Lyons & Thorne, 1973; Mourenas et al., 2017). Chorus waves are observed as highly structured wave bursts with frequency rising or falling in time (called rising and falling tones, respectively) usually in two frequency bands separated by a wave power gap around $0.5 f_{cer}$ corresponding to lower band chorus (0.1–0.45 of the equatorial f_{ce}) and upper band chorus (0.5–0.9 of the equatorial f_{ce}). They propagate from the geomagnetic equator roughly along the magnetic field (Agapitov et al., 2013; Burtis & Helliwell, 1969; Helliwell, 1965; LeDocg et al., 1998; Santolík et al., 2005). Chorus emissions are observed in the inner magnetosphere outside the plasmapause throughout the dayside magnetosphere (Aryan et al., 2016; Meredith et al., 2001; Tsurutani & Smith, 1974). Intensification of chorus generation in the inner and middle magnetosphere is associated with substorm electron injections (Goldstein & Tsurutani, 1984; Meredith et al., 2001) and supposed to be carried out via the cyclotron instability caused by the temperature anisotropy of 10- to 50-keV electrons (e.g., Fu et al., 2014; Li et al., 2010; Sazhin & Hayakawa, 1992). Generation of oblique chorus waves has also been observed in relation with field-aligned beams (or plateau distributions) of ~1-keV electrons and a temperature anisotropy of electrons below 10 keV (Artemyev et al., 2016; Li et al., 2016; Mourenas et al., 2015). In order to quantify the parameters of wave-particle interactions with chorus, one must figure out and characterize

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their spatial extent. This spatial extent can also provide important additional information about the eventual presence of inhomogeneities, such as a possible fragmentation of hot electron injections during their propagation into the inner magnetosphere, their localization, or substructure (Gabrielse et al., 2014, 2017; Turner et al., 2017), and help to determine the most probable generation mechanisms responsible for the observed emissions.

This kind of information can be obtained making use of simultaneous very low frequency (VLF) measurements from multispacecraft missions through a correlation analysis. The first correlation study was performed based on a comparison of spectrograms from ISEE1 and ISEE2 for separation distances of about 400 km. The comparison of chorus electric field spectrograms showed that on a timescale of a few minutes or longer, similar chorus emissions could be identified aboard both spacecraft, but on shorter timescales the detailed correlation tended to be poor, and the correlation scale was reported to be of the order of 100 km (Gurnett et al., 1979). This showed a need to distinguish the properties of a single chorus element (the spatial extent, amplitude distribution, and coherence extent) from the global scale of the chorus active region-the total area with similar wave amplitudes and frequency band. The scales of the generation region of a single chorus burst were estimated theoretically (Helliwell, 1965; Omura et al., 2013; Trakhtengerts, 1999) and experimentally based on the data from coordinated VLF measurements aboard Cluster (Agapitov et al., 2011; Artemyev et al., 2016; A. Breneman et al., 2007; Santolik et al., 2003, 2004; Santolík & Gurnett, 2003), Time History of Events and Macroscale Interactions during Substorms (THEMIS; Agapitov et al., 2010), and the Van Allen Probes (Agapitov et al., 2017). Santolík and Gurnett (2003), Santolik et al. (2003, 2004), and Breneman et al. (2009) determined the correlation lengths of chorus-type whistler mode waves based on multipoint Cluster Wide-Band (WBD) Plasma Wave investigation in the vicinity of the chorus generation region during geomagnetic storms to be between approximately 60 and 120 km, which was suggested to correspond to the size of the chorus generation region transverse to the background magnetic field (Santolik et al., 2003). Later, observations of similar chorus wave packets at separations larger than 500 km (Agapitov et al., 2011) indicated that the transverse size of the chorus source region is often larger in the outer radiation belt, where the typical source region scale was found to be around 600 km at L ~ 4.5 (Agapitov et al., 2011). Besides, Agapitov et al. (2010) used simultaneous observations of intense chorus by the four THEMIS spacecraft in the outer magnetosphere (L > 9). They estimated the transverse dimensions of chorus sources to be about 3,000 km there (wave frequency was about 0.15 of the equatorial f_{cer} corresponding to a wavelength of about 300 km). Recently, an analysis of long-lasting VLF waveforms (9 hr of magnetic and electric field waveforms during 1 day) from the two Van Allen Probes have shown that a typical chorus burst can extend up to 600-800 km (for both lower and upper band chorus), while the amplitude of a given chorus element decays over a typical transverse scale of about 300 km (Agapitov et al., 2017). The statistical analysis of wave amplitude correlation based on continuous spectral data (one spectrum every 6 s calculated from 0.5-s interval; Kletzing et al., 2013) recorded onboard the two Van Allen Probes during almost 3 years, and thus, related not to a single element but to the size of the global active region, found scales about 1,000-2,000 km for chorus in the outer radiation belt (Aryan et al., 2016). However, the spatial extent of the chorus wave active region has not yet been provided separately in both L and magnetic local time (MLT).

Statistical information on the source size and active region extent is also lacking for hiss waves, which are intense broadband whistler mode waves observed in the region L = 2-5 between about 100 and 2–3 kHz at frequencies smaller than $0.1f_{ce}$ (Thorne et al., 1973; Meredith et al., 2004; Li et al., 2015). Hiss waves are thought to be amplified locally (Chen et al., 2012; Nakamura et al., 2016), probably from an embryonic source consisting of chorus waves having propagated inside the plasmasphere (Bortnik et al., 2008; Chen et al., 2012; Meredith et al., 2013). Hiss waves are mostly observed inside the plasmasphere but sometimes also at higher L shells inside high-density plasmaspheric plumes within the 14–24 MLT sector (e.g., Moldwin et al., 2004; Su et al., 2018, and references therein). Hiss waves strongly contribute to electron losses inside the plasmasphere and the formation of a slot region between the inner and outer radiation belt, at least at energies below 1–2 MeV (Lyons & Thorne, 1973; Ma et al., 2016; Mourenas et al., 2017). The present paper explores the fine structure of the chorus source scale, the distribution of the spatial scales of chorus and hiss, the extent of the wave active regions, and the connection between hiss and chorus in a statistical way, based on 11 years of THEMIS wave measurements over 2007–2017, using all five THEMIS spacecraft in 2008–2011 and three spacecraft since 2011.



Figure 1. The distribution of lower-band chorus $(0.1f_{ce} < f < 0.5f_{ce})$ and hiss $(f_{LH} < f < 0.1f_{ce})$ amplitude in (a) MLT/*L* shell, and λ/L shell domains for the (b) day 06–18 MLT and (c) night 18–06 MLT sectors of the magnetosphere during quiet (Kp < 3) geomagnetic conditions. The coverage by the THEMIS VLF measurements in 2007–2017 in the (d) MLT/*L* shell and (e) λ/L shell domains. (f–j) The distributions and coverage for perturbed (Kp > 3) geomagnetic conditions in the same format as in panels (a)–(e). MLT = magnetic local time; THEMIS = Time History of Events and Macroscale Interactions during Substorms; VLF = very low frequency.

2. Data Description

The THEMIS mission began in February 2007 and consists of the five identically equipped National Aeronautics and Space Administration satellites (THA, THB, THC, THD, and THE). The main goal of this mission is to carry out multipoint investigations of substorm phenomena in the tail of the terrestrial magnetosphere (Sibeck & Angelopoulos, 2008). For the current study, search coil magnetometer (SCM) observations (Roux et al., 2009) and electric field instrument measurements (Bonnell et al., 2008) in a bandwidth from 0.1 to 4 kHz have been analyzed. Spin-averaged background magnetic field measurements with 4 s⁻¹ sampling by the fluxgate magnetometers (FGMs; Auster et al., 2008) were used to determine the local electron cyclotron frequency f_{ce} and lower hybrid frequency f_{LH} . Continuous measurements of the magnetic and electric field mean wave power in six logarithmically spaced frequency ranges from 1 Hz to 4 kHz of filter bank (FBK) data were used to evaluate chorus and hiss wave amplitudes. One-component measurements of magnetic and electric field in the spacecraft rotation plane were used. Such components are mostly directed transverse to the background magnetic field. This is due to the spacecraft orbit close to the equatorial plane and the spacecraft rotation axis transverse to the ecliptic plane. Whistler mode waves with a field-aligned propagation direction have an almost circular polarization of their magnetic field, with one component transverse to the background magnetic field direction. This allows to estimate whistler mode waves' magnetic field amplitudes from the FBK data (Li et al., 2009). In the present study, the FBK data obtained in 2007–2017 from all five probes are analyzed to obtain the parameters of the distribution of chorus waves over a radial range from 3 to 10 $R_{\rm F}$ at magnetic latitudes (λ) less than 20°. The magnetic latitude of the spacecraft is estimated based on a pure dipolar magnetic field model. Although this model may significantly underestimate the actual magnetic latitude in the presence of field-line stretching during strong geomagnetic disturbances, it is worth noting that there were almost no measured chorus waves at L > 7 around local midnight (see Figures 1a and 1c). During



Figure 2. Distribution of chorus wave amplitude (at *L* shell range from 5 to 10) in the (a) night-morning (2200–0600 MLT), (b) day (0600–1400 MLT), and (c) evening (1400–2200 MLT) sectors on geomagnetic latitude for the range of the geomagnetic activity level. MLT = magnetic local time.

these years, the gradually modified spacecraft orbits provided a comprehensive coverage of all MLTs. Simultaneous measurements were available from identical THEMIS spacecraft with an interspacecraft distance $|\Delta r|$ varying from tens of $R_{\rm E}$ down to about 100 km. In the present study, chorus waves are categorized as waves with a frequency between ~0.1 and 0.5 times the equatorial $f_{\rm ce}$ observed at L > 5 outside the plasmasphere, whereas hiss waves must have a frequency strictly smaller than 0.1 times the equatorial $f_{\rm ce}$ (but larger than the lower hybrid frequency $f_{\rm LH}$) and be observed at L < 5 inside the plasmasphere. The spacecraft was considered to be located inside the plasmasphere when the density derived from spacecraft potential measurements was above a threshold density ~30 cm⁻³ (Li et al., 2009).

The distribution of chorus wave mean amplitude presented in Figure 1 agrees well with the distributions reported by Li et al. (2009), Li, Bortnik, et al. (2011), and Meredith et al. (2012) based on FBK as well as waveform (Li, Bortnik, et al., 2011) THEMIS measurements, showing a morning sector (00-13 MLT) maximum in agreement with that by Meredith et al. (2012), which was absent in the study by Li et al. (2009) simply due to a lack of THEMIS measurements in this sector during active time intervals in 2007–2008. The observed enhancement of wave amplitude with geomagnetic activity level from 5-7 to 20 pT (shown in Figure 2 for three MLT sectors where the distribution remains roughly similar when performing Kolmogorov-Smirnov comparisons; see Agapitov et al., 2015) is also in agreement with results based on measurements from Cluster satellites (Agapitov et al., 2013, 2015) and the Van Allen Probes (Agapitov et al., 2018; Aryan et al., 2014). The location of maximum chorus wave power shifts from $L \sim 8-10$ during quiet periods to $L \sim 6-8$ during high geomagnetic activity. The wave amplitude is localized in the vicinity of the geomagnetic equator at $|\lambda| < 10^{\circ}$ in the night-morning MLT sector (especially during intermediate and high geomagnetic activity) and distributed over the whole latitudinal range from -20° to $+20^{\circ}$ covered by THEMIS measurements in the day sector. On the dayside (06-18 MLT), the chorus wave amplitude distribution has a local minimum near the equator (see Figures 2b and 2c) in agreement with results based on Cluster and Van Allen Probes measurements (Agapitov et al., 2013, 2015, 2017).

The distribution of hiss wave power in Figure 1 is also consistent with previous statistics from Meredith et al. (2013) mainly based on a much smaller 1.5-year THEMIS data set combined with 10 years of Cluster data. Hiss wave power is stronger in the day sector, with a gap of the order of R_E between its upper *L* limit and lower *L* limit of the chorus active region, as expected if inward-propagating chorus waves constitute an embryonic source for plasmaspheric hiss as ray-tracing simulations have suggested (Bortnik et al., 2008; Chen et al., 2012). One also finds that the region of strong hiss gets reduced from L < 5 to L < 4 during high geomagnetic activity as the plasmapause gets closer to the Earth, while the peak hiss power shifts from 10 to 12 MLT.

3. Spatial Scales of the Generation Regions of Chorus and Hiss

One important advantage of the THEMIS mission is the availability of simultaneous measurements from five identical spacecraft with an interspacecraft distance $|\Delta r|$ varying from tens of R_E down to less than





Figure 3. Dynamic spectra of magnetic field waveforms recorded aboard (a) THC and (b) THD on 17 June 2007. (c, d) The corresponding FBK wave power with 3 s⁻¹ sampling. The time interval shown in (a) and (b) is marked by red lines in (c)–(f). (e, f) The 1-min-averaged wave power over 11–15 UT, with the interval considered in (c, d) marked by blue lines. FBK = filter bank.

about 100 km, mostly transverse to the background magnetic field in the inner magnetosphere. This gives the opportunity to study simultaneous and coordinated observations to determine the instantaneous distribution of wave parameters with no reliance on time-averaged data and, thus, to determine the instantaneous in situ radial profile of wave power as well as the typical scales of wave amplitude correlations.

Actually, chorus waves have a multiscale structure in time and space. The fine structure of the temporal as well as the spatial distribution of chorus is determined by the temporal and spatial extent of a single isolated chorus wave packet—a wave burst localized in time (0.1–0.3 s) with rising or falling frequency. Studies based on the measurements by Cluster (Agapitov et al., 2011; Santolik et al., 2003; Santolík et al., 2005), THEMIS (Agapitov et al., 2010), and the Van Allen Probes (Agapitov et al., 2017) have shown that the same chorus burst can be observed on different spacecraft as far as 1,000 km apart. One example of similarities in the waveforms from THA, THC, THD, and THE recorded on 17 July 2007 is presented in Figure 3. At that time, the THEMIS spacecraft were crossing a chorus active region (see details in Agapitov et al., 2010) with interspacecraft distances ranging from 100 to 2,000 km. The SCM instrument and electric field instrument aboard the spacecraft provide several burst intervals with waveforms recorded over about 32 s during each orbit, resulting in a relatively poor coverage of the equatorial region by simultaneous and coordinated waveform measurements. However, one can take advantage of the usual stability of average chorus and hiss wave properties during tens of seconds in most observations to use the FBK data for a correlation study of wave properties. Accordingly, we used data stored in three frequency channels of FBK that cover a frequency range 80 Hz < f < 4,000 Hz. Since f_{ce} drops from 13 kHz at L = 4 to 1,100 Hz at L = 9, we use data only from channels in the frequency range $0.1f_{ce} < f < 0.5f_{ce}$ (the frequency range of lower band chorus) and $f_{LH} < f < 0.1f_{ce}$ (corresponding to hiss).

The correlation coefficient was calculated making use of 2-min intervals of FBK signal in the chorus and hiss frequency ranges captured aboard two of the THEMIS spacecraft (normalized over each 50-km interval of interspacecraft distance). The statistical distribution of this correlation coefficient is presented in Figures 4b and 4d for chorus and in Figure 4f for hiss. The mean values of the distribution and the 95% confidence





Figure 4. Total number of 2-min time series registered simultaneously by pairs of THEMIS probes separated by $|\Delta r|$ interspacecraft distance and stored in FBK during 2007–2017 for (a, c) lower-band chorus and (e) hiss. The distribution function of the cross-correlation coefficient of (b, d) chorus (for wave amplitudes <10 and >10 pT in panels (b) and (d), respectively) and (f) hiss wave power signals captured on board each pair of THEMIS spacecraft as a function of interspacecraft distance $|\Delta r|$. The 95% confidence intervals of the mean values are shown by error bars, corresponding to twice the standard deviation σ of the distribution. THEMIS = Time History of Events and Macroscale Interactions during Substorms; FBK = filter bank.

interval are shown for each interval of interspacecraft distance. The corresponding numbers of processed intervals as a function of interspacecraft distance $|\Delta r|$ are presented in Figures 4a, 4c, and 4d, respectively. The fast decay of the correlation as the $|\Delta r|$ increases up to 500–800 km agrees well with the characteristic transverse scale of individual chorus bursts reported by Agapitov et al. (2011, 2017). Measurements are available from identical THEMIS spacecraft for an interspacecraft distance $|\Delta r|$ varying from tens of $R_{\rm F}$ down to about 70 km. Although the number of data points is maximal for $|\Delta r| > 2,000-3,000$ km, it remains sufficient up to $|\Delta r| = 2 R_{\rm E}$ and for $|\Delta r| > 150$ km at least for a reasonable statistical study (see Figures 4a–4c). Moreover, Figures 4b and 4d show that the radial scale of the correlation is significantly smaller (~250-400 km) for chorus bursts with higher-amplitude (>10 pT) than for lower-amplitude bursts. One possible explanation is that chorus source regions of relatively moderate temperature anisotropy of injected hot (5-30 keV) electrons (leading to weak linear wave growth) may naturally extend over larger distances than source regions of stronger temperature anisotropy (leading to fast wave growth). However, the smaller spatial extent of the source region producing chorus with larger wave amplitudes could also be related to the nonlinear stage of chorus wave generation and, in particular, the wave amplitude saturation process (e.g., see Omura et al., 2013; Trakhtengerts, 1999, and references therein). Besides, no significant variation of the scale of the chorus source region with L shell or MLT has been found.

The scale of the mean correlation decay of hiss at *L* shells <5 is much larger than for chorus in the radial direction, around 2,000–3,000 km in Figure 4f, with some significant correlation existing up to about 8,000 km. Let us define here the size of the hiss active region as the maximum distance reached half of the time (i.e., for 50% of the observations) between simultaneous but distant observations of >3-pT hiss waves. Figures 5 and 6 show the instantaneous spatial extent of hiss and chorus in the radial and azimuthal directions, respectively. The comparison shows that they are rather similar, reaching about 1 R_E in the radial direction and about 1.2 hr of MLT in the azimuthal direction. Moreover, hiss waves are generally observed (at different universal times) over a twice larger MLT domain than are chorus waves—and mainly at slightly later MLTs, as expected for a generation of hiss from an embryonic source made of chorus waves mainly entering the plasmasphere at high latitudes around noon (Bortnik et al., 2008; Chen et al., 2012; Meredith et al., 2013).

The full distribution (in number of minutes of measurements) of the relative separation between two THEMIS spacecraft in azimuthal and radial directions is shown in Figures 7a and 7b, respectively, as a function of the *L*





Figure 5. The cross-distributions of (a–d) hiss and (e–h) lower-band chorus wave amplitude captured simultaneously aboard each pair of two THEMIS spacecraft for different interspacecraft distance ranges. THEMIS = Time History of Events and Macroscale Interactions during Substorms.

shell of the first considered spacecraft in 2007–2017. The combined coverage of interspacecraft relative separations in the Δ MLT- Δ L domain (with Δ L = $L_i - L_j$ and Δ MLT = MLT_i - MLT_j, where *i* and *j* subscripts denote the two THEMIS spacecraft) is presented in Figures 7c and 7d for chorus (L > 5 and 22 < MLT < 14) and hiss (L < 5 and MLT > 10) separately. The coverage of the measurements allows to provide significant results in the processed range of interspacecraft distances $|\Delta$ MLT| < 2 and $|\Delta L| < 2$.



Figure 6. Spatial extent in the (a) radial and (b) azimuthal directions of the whistler mode waves, determined as the maximal distance at which the waves were observed simultaneously in at least 50% of the observations with more than 3-pT amplitudes. The spatial extents (the Δ MLT and ΔL , at which the probability to observe the wave activity simultaneously aboard two spacecraft decays to 0.5) for (c) hiss ($f_{LH} < f < 0.1f_{ce}$) and (d) lower-band chorus ($0.1f_{ce} < f < 0.5f_{ce}$) active regions. MLT = magnetic local time.





Figure 7. Distribution (in number *N* of minutes) of relative separations between two THEMIS spacecraft in the (a) azimuthal (Δ MLT) and (b) radial (Δ L) directions as a function of the *L* shell of the first spacecraft (determined by alphabetical order, e.g., THEMIS C is before THEMIS D), for the available measurements in 2007–2017. The distribution of measurements in the Δ MLT- Δ L domain for (c) chorus and (d) hiss. THEMIS = Time History of Events and Macroscale Interactions during Substorms.

3.1. Correlations Between Chorus and Hiss Observations

Next, we examine statistically the reliability of the model of hiss generation by local amplification from an embryonic source of chorus waves having propagated inside the plasmasphere (Bortnik et al., 2008, 2009; Chen et al., 2012 ; Meredith et al., 2013). Toward this goal, we provide in Figure 8a the probability over 2007–2017 to observe well-correlated chorus and hiss waves (i.e., with a correlation >0.7 between their wave power dynamics over 2 min) onboard two separate THEMIS spacecraft, in the same whistler mode frequency band above 100 Hz corresponding to hiss waves (Li et al., 2015), when one spacecraft was at L = 2-5 inside the plasmasphere (corresponding to widespread hiss; see Figure 6) and the other spacecraft was at L = 6-9 outside the plasmasphere (corresponding to widespread chorus; see Figure 6). We mostly consider chorus waves (with $0.1f_{ce} < f < 0.5f_{ce}$) observed at L > 6 and plasmaspheric hiss waves (with $f < 0.1f_{ce}$) observed at L < 5. The plasmapause position was further checked using THEMIS spacecraft potential measurements, which provide an estimate of the electron density (Bonnell et al., 2008; Li et al., 2010; Li, Bortnik, et al., 2011). Since the proposed "seed" chorus waves must propagate inward and then get amplified inside the plasmasphere over 5–10 s to produce hiss (Chen et al., 2012; Meredith et al., 2013), we only considered as significant the chorus waves occurring within 10 s before hiss and the hiss waves occurring during the 10 s following chorus. Moreover, we only considered waves with amplitudes >1 pT.

We processed similar frequency chorus and hiss waves with amplitudes >1 pT observed simultaneously inside and outside the plasmasphere at an interspacecraft distance larger than 2 $R_{\rm E}$ and Δ MLT smaller than 3 hr. A little more than 1,000,000 of such 2-min time intervals have been detected during 2007–2017, making use of different pairs of THEMIS spacecraft. Among them, 71,000 time intervals were found with a correlation coefficient greater than 0.5 between chorus and hiss wave power dynamics. Therefore, the global average probability of finding such a high correlation between chorus and hiss waves, calculated as the number of intervals of high correlation divided by the total number of intervals of chorus and hiss observations, is ~7%—but it reaches ~20% over the wide (L, MLT) region of maximal occurrence rates in Figure 8a. The percentage of chorus waves highly correlated with hiss waves observed slightly later (called chorus-hiss) and of hiss waves highly correlated with chorus waves detected just before (called hiss-chorus) are presented in Figure 8a. About 20% of chorus waves at L = 6.5-8.5 and 12–16 MLT are well correlated with hiss waves observed immediately afterward inside the plasmasphere, whereas there is a slightly lower probability (~12%) for hiss waves at L = 3-5 and 10-20 MLT to be well correlated with chorus waves detected just before. The probability of detecting hiss when chorus is observed is larger than the probability of detecting chorus when hiss is observed, presumably due to the smaller spatial extent of the chorus active region than the hiss active region already discussed in section 3 (see Figures 5, 6). Besides, the L-MLT regions of significant probability of well-correlated chorus and hiss observations are in good agreement with the active regions of chorus and hiss found by Meredith et al. (2013) during active periods (AE > 100 nT) at low latitudes $< 30^{\circ}$ (see Figure 1b and Agapitov et al., 2013, 2018).







Figure 8. (a) Percentage of chorus waves (highlighted by red) such that their wave power dynamics is highly correlated with the wave power dynamics of hiss observed less than 10 s afterward inside the plasmasphere, and percentage of plasmaspheric hiss waves (highlighted by blue) similarly well correlated with chorus waves outside the plasmasphere (i.e., hiss occurring less than 10 s after chorus, with a high correlation >0.7 between their respective wave power dynamics). (b) The corresponding distribution of time delays between wave amplitude peaks of highly correlated chorus and plasmaspheric hiss. (c) The distribution of radial separation ΔL between well-correlated chorus and plasmaspheric hiss and (d) the corresponding distribution of azimuthal separation $\Delta MLT = MLT$ (chorus) – MLT (hiss), where the blue curve indicates measurements in the 00–12 MLT sector and the red curve shows measurements in the 12–24 MLT sector. An example of correlated chorus and plasmaspheric hiss observed during about an hour by THE (outside the plasmasphere; shown in red) and THD (inside the plasmasphere; marked in blue) is shown in panels (e)–(i). (e) The plasma density estimated from THD and THE spacecraft potential measurements. (f) The background magnetic field magnitude. The dynamics spectra of magnetic field perturbations recorded aboard (g) THD and (h) THE. The red dashed curves indicate local values of $0.5f_{ce}$ in panel (g) and the local $0.05f_{ce}$ in panel (h). The white dashed curves in panel (g) show $0.5f_{ce}$ and $0.05f_{ce}$ values from panel (h) corresponding to the background magnetic field at the location of THE. (i) The wave power of chorus and hiss in the same frequency range of 100-1,000 Hz.

In addition, Figure 8b shows the distribution of time delay between each peak of chorus wave power and the temporally closest peak of hiss wave power recorded at lower *L* shells during events of highly correlated chorus and hiss. This time delay is mostly between 0 and 6 s, with a finite tail up to 20 s. These results agree well with the time delay between chorus and hiss intensity peaks expected in the frame of a model of hiss generation through local amplification of seed chorus waves coming from outside the plasmasphere (Bortnik et al., 2008; Chen et al., 2012; Meredith et al., 2013). Figures 8c and 8d show that such well-correlated chorus and hiss waves are generally separated by ~2–3 R_E in the radial direction, hiss being observed mainly 1–2 hr later in MLT than chorus being in the 0–12 MLT sector but slightly more often at earlier MLT in the 12–24 MLT sector. Hiss waves observed ~1–2 hr earlier in MLT than chorus in Figure 8d can correspond to chorus active regions extending ~1.0–1.5 hr earlier in MLT than the location of chorus measurements, since Figure 6 has shown the existence of such large chorus extensions around noon. Therefore, these results mostly agree again with ray-tracing simulations of hiss generation through local amplification of seed chorus waves (Bortnik et al., 2008, 2011; Chen et al., 2012; Meredith et al., 2013).

One selected event—far from unique in our data set but with prolonged and very clear characteristic features—is displayed in Figures 8e–8i. Strong modulations (over ~0.7- to 1.6-min time scales) of chorus wave power (in the frequency band $0.05f_{ce} < f < 0.5f_{ce}$) are observed aboard the THE spacecraft at high L = 7-8 outside the plasmasphere, well synchronized with modulations of hiss wave power (in the frequency band ~100–500 Hz below the local $0.05f_{ce}$) detected almost simultaneously at lower L = 2.4-4.5 inside the plasmasphere aboard the THD spacecraft (the plasmapause was then around L = 5.5 based on the density derived from spacecraft potential measurements and using a threshold density ~30 cm⁻³ as in Li et al., 2009). The similar frequency band of the observed chorus and hiss emissions (compare

100



Figures 8g and 8h) actually corresponds to the local $0.05f_{ce}$ - $0.5f_{ce}$ band estimated from background magnetic field measurements by THE outside the plasmasphere. The variation of the envelope of the wave power of plasmaspheric hiss over much longer time scales (~15–50 min) turns out to be similarly well synchronized with the long-term variation of chorus wave power. The difference in *L* shells between such highly correlated chorus and hiss varies from 3.5 to 4.5, and hiss waves are observed ~2 hr later in MLT than chorus. All these global characteristics agree with the results of ray-tracing codes that show a possible generation of hiss through inward propagation of chorus waves (Bortnik et al., 2008; Chen et al., 2012; Meredith et al., 2013). Similarly, close relationships between chorus and hiss have also been found by analyzing the global response of these waves to interplanetary disturbances in recent event (Liu et al., 2017; Liu et al., 2017; Su et al., 2015) and statistical studies (Yue et al., 2017).

However, one important question remains: What is the cause of the strong wave power modulation of chorus and hiss waves in Figures 8g-8i? A similar well-correlated modulation of hiss waves has been observed recently on two Van Allen Probes located far away from each other, and it was suggested that then-observed ultralow frequency (ULF) waves in the Pc4-Pc5 band could have produced it through a modulation of plasma density that in turn modulates the number of resonant hot (10-100 keV) electrons and wave growth rates (Li et al., 2017). Weak density modulations were also seen at times by THE (not shown), simultaneously with ULF waves in the event of Figures 8e-8i. However, chorus growth rates can also be modulated by ULF waves through their modulation of the total magnetic field, which can slightly modify the cyclotron-resonant parallel electron velocity, eventually leading to modulations of the temperature anisotropy of the resonant hot electrons when this anisotropy varies with parallel electron velocity (Li et al., 2011). During the event of Figures 8e-8i, if ULF waves in the 10to 25-mHz range (e.g., Takahashi et al., 2015) had been simultaneously present over a broad L range from ~2.5 to 9, the similar wave power modulations of chorus and hiss could have been produced separately by similar ULF waves modulating their respective growth rates, without any cause-and-effect relationship between these two different waves. During the event in Figures 8e-8i, however, strong ULF waves were measured aboard THE in the chorus source region at high L shells, but no significant ULF waves were detected aboard THD inside the plasmasphere (see Figure 8f) where hiss wave power nevertheless exhibited a similar temporal modulation as chorus wave power at much higher L (see Figure 8i). In the considered event, it is therefore likely that ULF waves present in the chorus source region first modulated chorus growth rates by varying the quantity and/or anisotropy of cyclotronresonant hot electrons (e.g., Li et al., 2017; Li, Thorne, et al., 2011; Xia et al., 2016), some of these power-modulated chorus waves later propagating inside the plasmasphere, where they provided a modulated seed of plasmaspheric hiss waves (Bortnik et al., 2008; Chen et al., 2012) accounting for the observed hiss modulation.

4. Discussion and Conclusions

In the present study, we examined the spatial parameters of chorus and hiss wave activity regions in a statistical way, making use of multispacecraft extremely low frequency/VLF measurements from the THEMIS project during 2007-2017. THEMIS extremely low frequency/VLF measurements in the FBK mode comprehensively covered all MLT and L shells from 2 to 10 for a wide range of geomagnetic activity. The present distributions of chorus wave occurrence rate and amplitudes extend the databases provided by Li et al. (2009), Meredith et al. (2012), and Meredith et al. (2013) and show that chorus wave activity is regularly observed in the dawn-day sector of the magnetosphere consistent with earlier results (Li et al., 2009; Meredith et al., 2001, 2004; Andre et al., 2002; Green et al., 2005; Cully et al., 2008; Pokhotelov et al., 2008; Agapitov et al., 2011, 2012, 2013, 2015, 2018; Aryan et al., 2016). An enhancement of chorus activity with a shift of maximum wave amplitudes toward lower L shells (from $L \sim 8-10$ to $L \sim 6-8$, partially following the plasmapause position) is observed during periods of higher geomagnetic activity. Lower-band chorus wave power on the nightside is localized near the equator below 10° of latitude and on the dayside lower-band chorus wave power is distributed at all latitudes ($|\lambda| < 20^{\circ}$) covered by THEMIS measurements with a local minimum of the wave amplitude at the geomagnetic equator and a maximum shifted to $\lambda \sim 5-10^{\circ}$ at L from 5 to 8. Hiss is observed mainly at L < 5 in the day sector of the plasmasphere (06–22 MLT) with the plasmapause position as its outer edge, and it is uniformly distributed over the latitude range covered by THEMIS measurements.



The correlation scales of the whistler mode waves were estimated, making use of multispacecraft analysis of the THEMIS FBK spectral power data collected during 2007–2017. The correlation scale of chorus wave power dynamics in the outer radiation belt (4 < L < 8) has been found to be around 500–800 km on average but only 250–400 km for the strongest bursts (see Figures 4a and 4b), providing an estimate of the distance over which the same chorus waves can be detected. Such results are in good agreement with previous studies of the average transverse dimension (Agapitov et al., 2011, 2017) and strong-burst (storm time) transverse dimension (Santolik et al., 2004) of the chorus source region, where the spatial extent of a single chorus element was estimated based on a finer chorus waveform correlation analysis. This inner scale of the chorus structure is important for nonlinear wave-particle interactions and probably determines the scale of bursty electron precipitations observed at ionospheric altitudes under the form of pulsating auroras (Nishimura et al., 2010, 2011) and microbursts (A. W. Breneman et al., 2017; Kersten et al., 2011; Mozer et al., 2018) depending on the energy of precipitating electrons. The significant level of correlation of the average chorus wave amplitude is observed at interspacecraft distances up to 2,000–3,000 km, consistent with the results of a correlation study based on Van Allen Probes VLF spectral measurements (Aryan et al., 2016). We estimated the instantaneous spatial extent of a chorus or hiss wave active region as a function of L shell and MLT, by evaluating the maximum interspacecraft distance where waves were detected simultaneously with amplitudes above 3 pT in more than 50% of observations. This spatial extent of the chorus active region, together with the corresponding average wave intensity, determines the actual effectiveness of wave-particle interactions at any given time in multidimensional guasi-linear diffusion codes of the radiation belts. The chorus wave active region was found to be $\sim 1 R_{\rm E}$ in radial and $\sim 1 \, \rm MLT$ in the azimuthal direction. It suggests that the effects of chorus waves on electrons (both precipitation and acceleration) are probably much more inhomogeneous (i.e., localized in L and MLT) than usually assumed in radiation belt codes, although the global effects (like electron lifetimes and acceleration rates) may end up to be correctly evaluated when considering the time- and MLT-averaged chorus wave intensity.

The chorus active regions are also much more localized in the radial direction on the nightside ($<0.3 R_F$) compared to the dayside (~0.9–1.2 $R_{\rm F}$), with a roughly constant MLT scale (~1 MLT) over the 6–13 MLT sector. Such scales presumably correspond to the spatial spread of hot (3-30 keV) electrons injected during substorms, which generate the observed chorus waves. Our results are therefore consistent with hot electron injections occurring in the form of initially narrow spatial bunches (of small MLT and radial spread) between local midnight and dawn, probably corresponding to incoming localized depolarizing flux bundles produced by bursty reconnections in the magnetotail (Gabrielse et al., 2014, 2017; Turner et al., 2017). At their arrival, such mainly localized successive injections cannot produce an MLT-continuous active region of chorus waves. However, once the injected electrons of different energies and azimuthal drift velocities have drifted toward sufficiently higher MLTs and have simultaneously spread to nearby L shells (due to convection electric fields, radial diffusion, or drift shell splitting), they can probably constitute relatively wide homogeneous spatial regions of chorus generation, explaining the observations of much wider (in MLT and L) chorus active regions after 6 MLT. Dispersed hot electron injections inside geosynchronous orbit may also be more frequent after 3 MLT than near local midnight, as suggested by statistical observations (Gabrielse et al., 2014). In the future, it would be interesting to examine the statistics of hot electron distributions in details to assess the reliability of such a global scenario.

Hiss, observed as a wideband signal often without clear structure, does show a slower decay of its crosscorrelation with distance than chorus, with a mean correlation scale of 2,000–3,000 km and some significant correlation up to about 8,000 km. The (instantaneous) spatial extent of the hiss active region at L < 5 was found to be larger or comparable to that of chorus, up to 3 hr of MLT in the azimuthal direction (with similar amplitudes), with a radial scale of about 1 R_E —although hiss waves can be observed at different times over a much larger MLT domain spanning the whole dayside, as also shown by Breneman et al., (2015) and Li et al. (2017). The nonuniformity of hiss wave power in MLT at any given time should generally have no significant impact on electron precipitation, because the time scale of hiss-induced precipitation is usually longer than 1 day for 0.02- to 5-MeV electrons (e.g., Mourenas et al., 2017), therefore much longer than their azimuthal drift period.

Significant temporal correlations have also been found between distant observations of chorus (outside the plasmasphere) and hiss (inside the plasmasphere). We indeed found that 20% of the observed chorus waves are well correlated with hiss waves detected less than 10 s afterward, with a correlation >0.7 between their



respective wave power dynamics. Such well-correlated chorus and hiss waves are generally separated by ~2- $3 R_{\rm E}$ in the radial direction, hiss being usually observed 1–2 hr later in MLT than chorus. The observed time delay between chorus and hiss wave power peaks recorded at distant L shells mostly varied between 0 and 6 s, with a finite tail up to 20 s. All these results are in good agreement with a model of hiss generation through local amplification of seed chorus waves coming from outside the plasmasphere (Bortnik et al., 2008; Chen et al., 2012; Meredith et al., 2013). Moreover, we have discussed one interesting event during which observed ULF waves likely modulated chorus wave power at high L shells, whereas no ULF waves were seen at lower L shells, where hiss waves nonetheless exhibited an intensity modulation well synchronized with the distant chorus intensity modulation. In the considered event, ULF waves present in the chorus source region likely modulated chorus wave growth rates by varying the quantity and/or anisotropy of cyclotron-resonant hot electrons, some of these power-modulated chorus waves later propagating inside the plasmasphere and providing a modulated seed of plasmaspheric hiss waves (Bortnik et al., 2008; Chen et al., 2012). Temporal variations of hot electron flux or temperature anisotropy at high L shells could also provide additional chorus power modulations—such variations could occur from a feedback of the generated chorus waves during and after the generation process, since chorus-electron interactions usually reduce the initial anisotropy (e.g., Fu et al., 2014) and can lead to a fast loss of hot electrons into the atmosphere (e.g., Artemyev et al., 2016; Breneman et al., 2017; Mozer et al., 2018; Nishimura et al., 2010, 2011). All these interesting conjectures will be worth further checking in the future with accurate multipoint measurements of waves and particles.

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