

# Permittivity and electrical resistivity measurements and estimations during the recovery of DNAPL in saturated porous media: 2D tank experiments

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1	Permittivity and electrical resistivity measurements and estimations during the
2	recovery of DNAPL in saturated porous media: 2D tank experiments
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# 13 Graphical abstract



# 15 Highlights

- 2D tank experiments were performed to estimate DNAPL recovery
- DNAPL recovery was monitored with permittivity and electrical resistivity measurements
- DNAPL saturation estimation with permittivity is very accurate
- Electrical resistivity monitoring allows less accurate estimation of DNAPL saturation
- 20

### 21 Abstract

22 Pumping experiments were performed in a 2D tank in order to estimate the recovery yield and rate 23 of pure heavy chlorinated organic compounds (DNAPL; dense non-aqueous phase liquids). Several 24 operating configurations were considered: permeability of the saturated zone, pumping flow rates, addition of surfactant (to reduce capillary effects), and thermal treatment (to reduce DNAPL 25 viscosity). The experiments were monitored with permittivity ( $\epsilon$ ) and electrical resistivity ( $\rho$ ) 26 27 measurements. The experiments were also monitored with photography allowing, based on image 28 interpretation, to accurately convert optical densities into water saturations  $(S_w)$  and to validate the 29 geophysical data collected. Average S<sub>w</sub> were determined in all detection areas from permittivities and 30 electrical resistivities. These average  $S_w$  were used to calculate theoretical permittivities using the complex refractive index model (CRIM) and theoretical electrical resistivities using Archie's law. We 31 32 found a good correlation between measured permittivities and the estimated permittivities with image. 33 Conversely, the correlation was less accurate between measured and estimated electrical resistivities except for low electrical resistivities (*i.e.* high  $S_w$ ). It is therefore not possible to accurately quantify 34 35 water saturations using electrical resistivity monitoring alone. However, the accuracy is sufficient (especially for high values of  $S_w$ ) to highlight differences between the three considered treatment 36 technologies (*i.e.* without enhancement, with chemical and/or thermal enhancements). A combination 37 38 of monitored electrical resistivity (which provides an integrated vision of the cone of depression) and 39 monitored permittivity (which provides accurate but spatially limited information), and imaging allows a more accurate calibration of the relationship between resistivities and Sw and quantification of 40 41 remediation rates and yields.

42

#### 43 Keywords

Dense non-aqueous phase liquids, two-phase flow, porous media, free product recovery, electrical
resistivity monitoring, permittivity monitoring.

### 46 1. Introduction

47 Chlorinated Organic Compounds (COCs) are particularly persistent and toxic pollutants that may 48 significantly contaminate different environmental media (soil, indoor air and groundwater [NIEHS (2015); IARC (2018); Arp and Hale (2019)], over very long periods of time. These compounds have 49 been produced in large quantities since the middle of the 20<sup>th</sup> century mostly as solvents, pesticides, 50 electrical insulators [Cohen and Mercer (1993); Kueper et al. (2003)]. COCs are very hydrophobic and 51 52 are denser than water. In the event of chronic or accidental contamination, COCs infiltrate into the 53 vadose and saturated zones down to the substratum to form DNAPL (dense non-aqueous phase liquid) pools [Cohen and Mercer (1993); Kueper et al. (2003)]. The recovery of DNAPL free-phase (mobile) 54 55 product is usually performed by implementing "pump and treat" methods. However, these methods are not very efficient and commonly last for many years [ITRC (2002); Falta et al. (2005b); McGuire 56 et al. (2006)]. Also, this remediation technology does not allow a homogeneous recovery of the 57 DNAPL. Thus, it is essential to better: i. spatialize the residual saturations of DNAPL; and ii. monitor 58 the applied remediation operation in order to implement additional remediation measures on "hot 59 60 spots" and thus improve recovery rates [Maire et al. (2018)]. Geophysical techniques can help to better 61 characterize these pollution sources [Power et al. (2014)].

62 Geophysical investigation, used to characterize hydrocarbon contaminated sites, mainly relies on 63 the electrical properties of hydrocarbons because their presence may change some of the medium's 64 physical properties, changes which can be detected using electrical measurements such as resistivity 65 measurements (classic or stationary and complex or transitory). These NAPLs have high electrical 66 resistivity and low dielectric constant (compared to water). They respectively decreases and increases 67 with increasing degradation, as does electric chargeability, *i.e.* the capacity to polarize, which is significant and depends on their degradation state [Brown et al. (2003); Brown et al. (2004); Schmutz 68 69 et al. (2010); Smallwood (2012)]. Since the early 2000s, geophysical methods have been the subject of 70 intensive research to study environmental problems associated with hydrocarbons and to determine the 71 relationship between the geophysical signal and the contamination. Also, several electrical resistivity monitoring experiments have been conducted to follow biogeochemical remediation in sites polluted 72 by chlorinated solvents and refined hydrocarbons [Nyquist et al. (1999); Power et al. (2014); Noel 73

et al. (2016a)]. Deng *et al.* (2017) have reported the importance of using laboratory tomography measurements for electrical resistivity (classic) combined when possible with imaging techniques, to improve the interpretation of resistivity measurements on sites polluted by NAPL and composed of materials with varying permeabilities [Deng et al. (2017)].

This research aims to connect DNAPL saturations and electrical resistivity measurements. Resistivity of soils depends of two different mechanisms: electrolytic and surface conduction (*i.e.* presence of clay) [Revil (2012)]. Both terms are saturation-dependent. It should be noted that there is no consensus regarding the interdependence of water saturation and the surface conductivity [Breede et al. (2011); Laloy et al. (2011)]. Archie's law has been shown to provide relatively accurate estimates of DNAPL [Kamona et al. (2003); Colombano et al. (2020); Iravani et al. (2020a)] in case of electrolytic conduction.

85 One of the most commonly-used physical models of dielectric permittivity  $(\varepsilon)$ , *i.e.*, the Complex 86 Refractive Index Model (CRIM), describes soil as a mixture of particles, water and air [Birchak et al. (1974); Roth et al. (1990); Endres and Knight (1992)]. In the 1990s, studies aimed at connecting 87 DNAPL saturation  $(S_n)$  in saturated soils to the  $\varepsilon$  measured by time domain reflectometry (TDR) 88 probes [Redman et al. (1991); Redman and DeRyck (1994); Brewster et al. (1995)]. Perchloroethylene 89 90 samples, spiked at controlled saturations, were monitored by TDR. The measured permittivities allow to estimate S<sub>n</sub> in saturated porous media with good accuracy using the mixing models [Redman and 91 92 DeRyck (1994)]. Ajo-Franklin et al. (2004) obtained degrees of correlation higher than 97% using the 93 CRIM model in the case of trichloroethylene (TCE) in the saturated zone [Ajo-Franklin et al. (2004)]. Finally, Persson and Berndtsson (2002) showed that the mixing model used led to S<sub>n</sub> estimation errors 94 95 up to 5% in the case of saturated sand "contaminated" by sunflower seed oil, synthetic motor oil and 96 paraffin [Persson and Berndtsson (2002)]. To minimize this error, Persson and Berndtsson coupled 97 TDR probes with the conductivity measurements.

98 Geophysical methods are complementary of intrusive investigation methods such as boreholes with 99 soil sample analyses, piezometers with water analyses, or soil-gas wells with gas analyses, and may 100 provide relatively spatially continuous information, where resolution is a function of electrodes 101 spacing and depth imaging [Chambers et al. (2004); Reynolds (2011)]. Several studies have documented the potential of using electrical methods to monitor remediation of soil and groundwater
at both laboratory and field experimental scales [Sogade et al. (2006); Cardarelli and Di Filippo
(2009); Noel et al. (2016a); Noel et al. (2016b)]. Combining both electrical conductivity and dielectric
permittivity has been proven to be very effective for saturated soil characterization [Linde et al.
(2006); Brovelli and Cassiani (2011)].

107 The aim of this study was to demonstrate the accuracy of electrical resistivity and permittivity 108 monitoring methods as a mean to locate and estimate the saturation of DNAPL during laboratory 109 pumping experiments in a water-saturated porous medium. DNAPL recovery was performed by pumping without enhancement as well as with chemical and thermal enhancements. Adding surfactant 110 aims to reduce capillary effects, whereas heating aims to reduce DNAPL viscosity, while both 111 methods aim to increase DNAPL recovery yield and rate during pumping. While the detailed results of 112 enhancement effects on DNAPL recovery are described elsewhere, this paper focuses on the 113 114 geophysical monitoring methods.

115

#### 116 2. Materials and methods

The DNAPL pumping experiments in the 2D tank were performed with glass beads (GB) packing 117 of two distinct diameters (0.1 and 0.5 mm), different flow rates (50, 150, 220 mL.min<sup>-1</sup>), without 118 119 thermal enhancement (at 20 °C), with chemical enhancements (with sodium dodecyl benzene 120 sulfonate, SDBS) at its critical micellar concentration (CMC, to avoid DNAPL dissolution) and/or thermal enhancement at 50 °C (to avoid DNAPL volatilization). The use of glass beads helped avoid 121 potential artefact effects associated with porous medium heterogeneity, while the two different 122 diameters enabled to reproduce distinct hydraulic conductivities  $(1.30 \times 10^{-10} \text{ m}^2 \text{ and } 6.73 \times 10^{-12} \text{ m}^2 \text{ for})$ 123 0.5 mm GB and 0.1 mm GB, respectively). 124

The DNAPL used for the experiments was sampled at the Tavaux site (June 25, 2014), i.e., a large chloralkali chemical plant located in the center-east of France. It was stored at 4 °C until use. DNAPL was filtered just before its use (EMD Millipore, 0.45  $\mu$ m). This DNAPL is a mixture of heavy chlorinated organic compounds (weight percent; wt%): hexachlorobutadiene-HCBD (58%), hexachloroethane-HCA (14%), perchloroethylene-PCE (8%), and pentachlorobenzene (3.5%), carbon tetrachloride (4%), trichloroethylene-TCE (2%), hexachlorobenzene (1%) [Cazaux et al. (2014)]. The
water used for all experiments was BRGM tap water. It was degassed using VWR Ultrasonic Cleaner
(USC500D: 60 °C, 45 Hz, 60 min).

133

## 134 2.1 DNAPL pumping experiments in the 2D tank and image interpretation

The dimensions of the 2D tank, made of polyvinylidene difluoride (PVDF), were: L=50.00 cm,
W=30.00 cm and H=7.00 cm (Figure 1).

137







Figure 1: Scheme of the experimental 2D tank, its main features and dimensions

140

Water was injected into the glass bead porous medium from the bottom of the 2D tank until it
reached a height of 30.0 cm (at a flow rate of 0.05 mL.min<sup>-1</sup>). Then, the DNAPL was injected to a
height of

144 15.0 cm (in order to reach the irreducible saturation,  $S_{rw}$ ) at a flow rate of 0.01 mL.min<sup>-1</sup>. The central 145 bottom access was used to pump the DNAPL during imbibition (red arrow in Figure 1). Several pipes 146 were inserted inside each counter-channel to regulate the static level of air-water interface (at a height 147 of 30.00 cm) and the static level of DNAPL-water interface (at a height of 15.00 cm) [Colombano 148 et al. (2021)]. Two types of peristaltic pumps were used: Watson Marlow 205 and Watson Marlow 149 530U; their flow rate ranges are  $6.00 \times 10^{-4}$  to 22.00 mL.min<sup>-1</sup> (± 0.5 %) and  $4.00 \times 10^{-3}$  to  $2.40 \times 10^{-3}$  150 mL.min<sup>-1</sup>

151  $(\pm 0.5 \%)$ , respectively.

The experiments were monitored with photography allowing, based on image interpretation to convert optical densities (OD) into  $S_w$ . The average  $S_w$  values were estimated in the detection areas of the dielectric permittivity and electrical resistivity sensors (assuming a homogeneous medium). The theoretical dielectric permittivity and electrical resistivity of the bulk ( $\varepsilon_{est}$  and  $\rho_{c,est}$ , respectively) were estimated according to water and DNAPL saturations (using suitable mixing models) in theses detection areas. They were then compared to the permittivity and electrical resistivity measured ( $\varepsilon_{meas}$ and  $\rho_{c,meas}$ , respectively).

159

# 160 2.2 *Permittivity monitoring*

Dielectric permittivities were monitored using capacitance/frequency domain technology (Decagon 161 Devices 5TE 40567) with a 70 MHz frequency. As well as permittivity (-), these dielectric probes 162 monitor temperature (°C), and bulk electrical conductivity-EC (dS.m<sup>-1</sup>). The probes were calibrated 163 164 and were connected to a Campbell Scientific CR1000 (4M) data logger and only used to acquire temperature and permittivity data. Data acquisition frequency was 1 signal measurement per 30 s. In 165 total, 15 dielectric probes were arranged in the 2D tank. On the basis of the experiments conducted in 166 167 1D cells and 1D columns [Colombano (2019)], the detection area for the dielectric probes is estimated 168 at around 2 mm of the probe (Figure 2) [Colombano (2019)].

169



171 Figure 2: Position of dielectric probes in the 2D tank with associated detection areas for DNAPL

173 The permittivity values were corrected relative to the reference values measured in air and water174 using the following equation (Eq. 1) [Kargas and Soulis (2012)]:

$$\varepsilon_{corrected} = \frac{\varepsilon_{w-\text{theoretical}}(\varepsilon_{meas} - \varepsilon_{air})}{(\varepsilon_w - \varepsilon_{air})} + \varepsilon_{air-\text{theoretical}} \qquad \text{Eq. 1}$$

where  $\varepsilon_{w-\text{theoretical}}$  (-) is the relative permittivity of pure water ( $\varepsilon_{w-\text{theoretical}}$ =80 at 20 °C),  $\varepsilon_{meas}$  (-) is the effective relative permittivity of the medium measured with the dielectric probe during the experiments,  $\varepsilon_{air}$  (-) is the relative permittivity of air measured at the beginning of the experiment,  $\varepsilon_w$  (-) is the relative permittivity of tap water measured at the beginning of the experiment, and  $\varepsilon_{air-\text{theoretical}}$ =1).

180 Permittivity measurements in these detection areas were compared with image interpretations181 (which allow to estimate the average water saturation in the same detection areas).

182

## 183 2.3 Electrical resistivity monitoring

Electrical resistivity was monitored using unpolarizable potential electrodes for voltage 184 185 measurement, metallic stainless electrodes for current injection, a digital resistivity meter and a data 186 acquisition software. Cu/CuSO<sub>4</sub> unpolarizable electrodes were used after preparation according to the 187 method developed by Noel (2014) and derived from Maineult (2004) [Maineult et al. (2004); Noel 188 (2014)]. The current injection electrodes were made of stainless Inox 316L. The resistivity meter was 189 SIP LAB IV and we injected sinusoidal current shape. We recorded data using a frequency of 1.4 Hz 190 because it is close to the one used in the field [Chambers et al. (2004); Constable and Srnka (2007); 191 Han et al. (2015); Deparis et al. (2019)]. In total, 30 unpolarizable potential Cu/CuSO<sub>4</sub> electrodes and 192 30 metallic rod stainless current electrodes were arranged on the back of the 2D tank. On the basis of experiments conducted in 1D cells and 1D columns [Colombano (2019)], the detection areas for 193 194 resistivity was estimated at approx. 2.50 mm around those electrodes (Figure 3) for porous media saturated with DNAPL. Two dipole-dipole injection and reception configurations were programmed 195 (vertical and crossed). 196

198



Figure 3: Position of unpolarizable potential and current injection electrodes in the 2D tank with associated detection areas for DNAPL: a) axial dipole-dipole and b) azimuthal dipole-dipole

In all, there were 47 resistivity monitoring zones. The electrode configurations were axial dipoledipole (for the vertical line detection areas) and azimuthal dipole-dipole (for square detection areas). The values measured were resistance values,  $R_e$ , which were transformed into apparent electrical resistivity values,  $\rho_c$ , using the geometric coefficient,  $K_g$  (Eq. 2 and Eq. 3) [Reynolds (2011)]:

$$K_g = \frac{\rho_0}{R_e}$$
 Eq. 2

$$\rho_c = R_e K_g \tag{Eq. 3}$$

where  $K_g$  (m) is the geometric coefficient,  $\rho_0$  ( $\Omega$ .m) is the measured resistivity value using a reference liquid,  $R_e$  ( $\Omega$ ) is the electrical resistance measured with SIP system using a reference liquid, and  $\rho_c$  ( $\Omega$ .m) is the electrical resistivity of the bulk. The conductivity of the tap water, measured at the beginning of each experiment, made it possible to determine  $\rho_0$  and to calculate  $K_g$  (using Eq. 2) with filling 2D tank with water only. The value of  $K_g$ was then used throughout the experiment to transform  $R_e$  into  $\rho_c$  (using Eq. 3).

213

#### 214 **3.** Theory

The relative permittivity,  $\varepsilon_r$ , is defined as the ratio of the real dielectric permittivity,  $\varepsilon'$ , and the free space permittivity,  $\varepsilon_0$  (Eq. 4):

$$\varepsilon_r = \frac{\varepsilon'}{\varepsilon_0}$$
 Eq. 4

217 where  $\varepsilon_r$  (= $\varepsilon$ ) (-) is the relative permittivity,  $\varepsilon'$  (F.m<sup>-1</sup>) is the real dielectric permittivity, and  $\varepsilon_0$ 218 (F.m<sup>-1</sup>) is the free space permittivity.

In this article, later, all dielectric permittivity values are expressed relatively. In addition, for the sake of simplification,  $\varepsilon_r$  will be called  $\varepsilon$ .

The dielectric probe instrument sends a high-frequency electromagnetic step pulse through a transmission line of known length L, and the pulse is reflected back at the end of the line. From the travel time of the pulse analysis, the soil's bulk dielectric constant is computed [Topp et al. (1980); Persson and Berndtsson (2002)]. The dielectric constant can be calculated as (Eq. 5):

$$\varepsilon = \left(\frac{c_{\varepsilon}}{v_{\varepsilon}}\right)^2$$
 Eq. 5

where:  $c_{\varepsilon}$  (m.s<sup>-1</sup>) is the speed of light (velocity of electromagnetic waves) in vacuum ( $c_{\varepsilon} = 299\ 792\ \text{km.s}^{-1}$ ),  $v_{\varepsilon}$  (m.s<sup>-1</sup>) is the function of the propagation velocity ( $v_{\varepsilon} = 2L_{\varepsilon}/t_{\varepsilon}$ ),  $L_{\varepsilon}$  (m) is the waveguide length, and  $t_{\varepsilon}$  (s) is the travel time for the pulse to traverse the length of the embedded waveguide (down and back: 2L).

The CRIM model for multi-phase dielectric permittivity is based on Eq. 6 [Birchak et al. (1974);
Roth et al. (1990); Endres and Knight (1992)]:

$$\varepsilon = \left[\sum_{i=1}^{N} v_i \varepsilon_i^{\alpha_{\varepsilon}}\right]^{1/\alpha_{\varepsilon}}$$
Eq. 6

where  $\varepsilon$  (-) is the effective relative permittivity of the multi-phase system,  $\varepsilon_i$  (-) is the relative permittivity of the *i* phase,  $v_i$  (-) is the volume fraction of the *i* phase, and  $\alpha_{\varepsilon}$  (-) is an empirical constant related to the geometry of the grains and their spatial distribution.

Ajo-Franklin *et al.* (2004) report that for a three-phase mineral/water/NAPL (TCE) mixture and assuming that  $\alpha_c$ =0.5, the CRIM equation becomes (Eq. 7) [Ajo-Franklin et al. (2004)]:

$$\varepsilon = \left[ \emptyset \left( S_w \sqrt{\varepsilon_w} + S_n \sqrt{\varepsilon_n} \right) + (1 - \emptyset) \sqrt{\varepsilon_s} \right]^2$$
 Eq. 7

where  $\emptyset$  (-) is the porosity,  $S_w$  (-) is the water (wetting fluid) saturation,  $\varepsilon_w$  (-) is the relative permittivity of water,  $S_n$  (-) is the DNAPL (non-wetting fluid) saturation,  $\varepsilon_n$  (-) is the relative permittivity of DNAPL, and  $\varepsilon_s$  (-) is the relative permittivity of soil particles.

The values measured were resistance values,  $R_e$ , which were transformed into apparent electrical resistivity values,  $\rho_c$ , using the geometric coefficient,  $K_g$ , (Eq. 2 and Eq. 3), as explained in the section 241 2.3.

Archie's law (1942), adapted by Winsauer et al. (1952), describes how resistivity depends on
porosity if ionic conduction in the pore fluid dominates other conduction mechanisms in the rocks (Eq.
8) [Archie (1942); Winsauer et al. (1952)]:

$$\phi \rho_c = \rho_{c,w} a_c \phi^{-m_c}$$
 Eq. 8

where  $\rho_c$  ( $\Omega$ .m) is the electrical resistivity of the bulk,  $\rho_{c,w}$  ( $\Omega$ .m) is the electrical resistivity of the water at temperature T,  $a_c$  (-) is an empirical parameter (typically equal to 1 but that can vary from <1 for intergranular porosity to > 1 for joint porosity), and  $m_c$  (-) is the cementing factor (also an empirical parameter usually approximately equal to 2, but that can vary from 1.2 for unconsolidated sediments to 3.5 for crystalline rocks; Friedman (2005) produced an exhaustive list of cementation exponent value [Friedman (2005)]).

Adapted Archie's law (referred to as Archie's law in the rest of the document) seems to be a good approximation when the conductivity is controlled by electrolytic conduction, which is often the case for non-consolidated porous media with a low clay content [Árnason et al. (2000)]. Glover (2010) has proposed, on the basis of Archie's law, the following equations for calculating S<sub>w</sub> for a two-phase system, where the conductivity of GB is neglected as it is very low ( $\sigma_{c,GB}=10^{-20}$  S.m<sup>-1</sup>) (Eq. 9 to Eq. 12) [Glover (2010)]:

$$\sigma_{c,bulk} = \sigma_{c,DNAPL} [\emptyset(1 - S_w)]^{m_1} + \sigma_{c,water} [\emptyset(S_w)]^{m_2}$$
 Eq. 9

$$\left(-\frac{\phi_1^2}{2}\right)m_2^2 + \left(\phi_1 + \frac{\phi_1^2}{2}\right)m_2 - \phi_1^{m_1} = 0$$
 Eq. 10

$$m_{2} = \frac{-\left(\emptyset_{1} + \frac{\emptyset_{1}^{2}}{2}\right) \pm \left[\left(\emptyset_{1} + \frac{\emptyset_{1}^{2}}{2}\right)^{2} - 4\left(-\frac{\emptyset_{1}^{2}}{2}\right)\left(-\emptyset_{1}^{m_{1}}\right)\right]^{\frac{1}{2}}}{-\emptyset_{1}^{2}}$$
Eq. 11

$$m_2 = \frac{-(4\phi_1 + 2\phi_1^2) \pm (4\phi_1^2 + 4\phi_1^3 + \phi_1^4 - 8\phi_1^2\phi_1^{m_1})^{\frac{1}{2}}}{-4\phi_1^2}$$
Eq. 12

where  $\sigma_{c,bulk} = \sigma_c$  (S.m<sup>-1</sup>) is the electrical conductivity of the bulk,  $\sigma_{c,DNAPL}$  (S.m<sup>-1</sup>) is the electrical conductivity of DNAPL, m<sub>1</sub> (-) is the cementation exponent of DNAPL phase,  $\sigma_{c,water}$ (S.m<sup>-1</sup>) is the electrical conductivity of water, m<sub>2</sub> (-) is the cementation exponent of water phase, and  $\phi_1 = \phi(1 - S_w); \phi_2 = \phi(S_w).$ 

The DNAPL-water interface charges (*i.e.* surface conductivity) have been neglected because i. the specific surface is negligible (surface tension is high) [Brusseau et al. (2006); Linde et al. (2007)] and, ii. the ionic exchanges between the two phases are weak (DNAPL is poorly soluble in water) [Revil (2017)].

265

# 266 4. Results and discussion

Previous work has shown that one can accurately determine water and DNAPL saturations by image interpretation during drainage-imbibition experiments in 1D cells and pumping experiments in 2D tank [Colombano (2019); Colombano et al. (2020); Iravani et al. (2020b); Philippe et al. (2020); Colombano et al. (2021)]. The purpose of the experiments presented here is to compare the measured permittivities and electrical resistivities with those estimated on the basis of the water saturations determined using image interpretation. Figure 4 shows raw photographs and estimations of DNAPL saturations with image interpretation during a pumping experiment conducted with 0.5 mm GB with a 274 flow

275 150 mL.min<sup>-1</sup> and without enhancement.









Figure 5: Evolution of the permittivity in the 2D tank with 0.5 mm GB with a flow rate of

150 mL.min<sup>-1</sup> (without enhancement)

DP 1 to 5 located in the upper part of the 2D tank (Figure 2) were only surrounded by water  $(S_w=1)$ . They had the same values when the experiments began (31.49 ± 0.51). These values were similar to the ones measured and corrected by Colombano *et al.* (2020) in the 1D cells.

DP 11 to 15 were in the lower portion of the 2D tank (Figure 2). As pumping began, the water 295 contents were equivalent to  $S_{rw}$ . The monitored permittivities for  $S_{rw}$  were around 11.39 ± 0.39 and 296 297 agreed with those measured and corrected by Colombano et al., 2020 in the 1D cells and columns. As 298 pumping ended, DP 13, located in the lower central position, detected the arrival of the cone of 299 depression (at t=23 min). DP 6 to 10 were in the middle of the 2D tank (Figure 2). Logically DP 8, 300 located in the tank's centre, was the first to measure the permittivity increase; then as pumping continued, DP 7 and 9, located on either side of this central position, showed increased permittivities 301 302 (and therefore decreasing DNAPL saturations). As pumping ended,  $\varepsilon_{meas}$  at DP 7 to 9 were on average 303  $27.27 \pm 1.10$ , which matches with measurements at the end of imbibition in 1D cells and 1D columns 304 [Colombano et al. (2020)]. The  $\varepsilon_{meas}$  at DP 6 and 10 (located at the cone of depression) were 305 respectively

306  $23.41 \pm 0.29$  and  $21.46 \pm 0.77$ .

Figure 6 shows variations in water saturation, measured permittivity and estimated permittivity during the pumping experiment. The water saturation was calculated from image interpretation (*i.e.* from optical densities in the DP's detection area). The permittivity estimate was calculated from  $S_w$ and

Eq. 13 for 0.5 mm GB (adaptation of the CRIM model from drainage-imbibition experiments
[Colombano et al. (2020)]).

$$\varepsilon = [\phi(S_w \varepsilon_w^{0.7} + S_n \varepsilon_n^{0.7}) + (1 - \phi) \varepsilon_s^{0.7}]^{\frac{1}{0.7}}$$
 Eq. 13





Figure 6: Average water saturation, measured and estimated permittivity during the pumping
experiment in the 2D tank with 0.5 mm GB and a flow rate of 150 mL.min<sup>-1</sup> (without enhancement)
317

318 The  $\varepsilon_{est}$  (from image interpretations) were very similar to the  $\varepsilon_{meas}$  at the start and end of pumping. 319 When the migration front moved to the DP's detection area, the  $\varepsilon_{meas}$  were systematically lower than the  $\varepsilon_{est}$ . The value assigned to  $S_w$  is the average of the  $S_w$  at the detection area. This value is 320 transformed into permittivity according to the CRIM model. This model is only valid for 321 322 homogeneous media. When the detection area is crossed by the migration front, this zone is a mixing 323 of contents similar to the residual saturation of DNAPL (Sm) and contents similar to the residual 324 saturation of water  $(S_{rw})$ . Therefore, for two heterogeneous zones, the weight of the portion characteristic of  $S_{rw}$  influences the global  $\varepsilon_{meas}$  more that the portion characteristic of  $S_m$ . Eq. 6, 325 326 combined with the permittivity values for DNAPL and water ( $\varepsilon_n=3.11$  and  $\varepsilon_w=80$ ), shows that the 327 propagation velocity function is 6 times faster for DNAPL than for water. It is therefore consistent 328 that, when the DP detection area is crossed by the migration front,  $\varepsilon_{meas}$  is systematically lower than 329 Eest.

This shift between the estimated and measured permittivities in the transition zones was studied for soils in the unsaturated zone (soil-air-water system). Some authors have shown that the CRIM model (based on the permittivity of free water, of air and solids) may not be suitable in all situations, especially in the transition phases (when  $S_w$  are between  $S_{rn}$  and  $S_{rw}$ ). These authors have proposed to consider a 4<sup>th</sup> constituent, the bound water (whose dielectric constant is very different from that of free water) [Dasberg and Hopmans (1992); Capparelli et al. (2018)]. An equation derived from the CRIM model incorporating the bound water has been proposed [Dobson et al. (1985); Dirksen and Dasberg (1993)]. Capparelli et al. (2018) carried out monitoring of  $S_w$  variations with dielectric probes in unsaturated soils. The experimental results were compared to the classical CRIM model, but also to the four-phase dielectric mixing model (with bound water), and the classical CRIM model with a variable exponent. The best results were obtained with the three-phase dielectric mixing model with a variable exponent ( $\alpha_e$ ) [Capparelli et al. (2018)].

- Figure 7 compares measured and estimated permittivity (with image interpretation) in the 2D tankfor 0.5 and 0.1 mm GB packings in the absence of any enhancement.
- 344





b) 0.1 mm GB



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Normalized root-mean-square error (NRMSE) were 0.12 and 0.14 for 0.5 and 0.1 mm GB, respectively. The measured permittivities agreed with the estimated permittivities (based on image interpretation). As expected, NRMSE for the 0.5 mm GB was lower than for the 0.1 mm GB, which shows superior correlation in the case of 0.5 mm GB (because of the fingering process which is more pronounced for 0.1 mm GB). In both cases, four distinct zones appear:



355	• A transition zone with values between the $S_{rw}$ zone and the $S_{rm}$ zone; here the majority of
356	points were below the linear regression line ( $\epsilon_{meas}$ are below $\epsilon_{est}$ ),
357	• A $S_{rn}$ zone with values around 27.00 to 28.00 for the 0.5 mm GB and 25.00 to 26.00 for the
358	0.1 mm GB (end of pumping),
359	• A water zone which corresponds to $\varepsilon_{meas}$ at DP 1 to 5 (where S <sub>w</sub> =1), with 31.00 to 32.00 for
360	0.5 mm GB and 32.00 to 33.00 for 0.1 mm GB.
361	Figure 8 displays a surface plot of the permittivity measured on an image at the end of pumping
362	experiment. The permittivity values were calculated using an interpolation technique (Kriging or

Gaussian process regression) with Surfer® software.



363



365

Figure 8: Surface plot of measured permittivity within an image with 0.5 mm GB and a flow rate of
150 mL.min<sup>-1</sup> (without enhancement) at t=18 min

368

369 One can see that the permittivity and the  $S_n$  do not overlap completely. This is explained by two 370 phenomena: i. the kriging method (under Surfer® software) is only based on 15 points; ii. the 371 transition zones cause an offset for the measurements. However, this does illustrate that the cone of 372 depression can be determined by permittivity monitoring.

# 374 4.1.2 Experiments with chemical or thermal enhancement

The modelling and experimental results show that chemical enhancement has a beneficial effect on 375 376 recovered DNAPL volumes. The use of chemical enhancement was proportionally more advantageous for lower flow rates than for higher ones. The cone of depression radius and height increased with 377 378 added surfactant. Thermal enhancement had no beneficial effect on DNAPL recovery rate or yield. 379 Heating the porous media had a negative effect on cone of depression radius and height [Colombano 380 et al. (2021)]. Figure 9 shows a comparison of measured permittivity in the 2D tank (with and without 381 enhancements) for 0.5 mm GB. Figure 9 also shows that by using DP it is possible to assess the 382 differences in remediation yields and rates between recovery of free product without enhancement, with chemical enhancement and with thermal enhancement. The DP located at the heart of the cone of 383 depression globally showed faster and larger variations in permittivities for chemical enhancement 384 than without enhancement (this is more visible for the 0.1 GB than for the 0.5 mm GB). Compared to 385 the reference (no enhancement), the DP monitoring data show that with chemical enhancement, the 386 radius of action of the cone of depression is higher, the  $S_n$  are lower and the  $S_n$  fall faster than without 387 388 enhancement. Moreover, the low NRMSE values, *i.e.* around 0.10 for 0.5 mm GB and lower than 0.14 for 0.1 mm GB (calculated based on graphs  $\varepsilon_{meas} = f(\varepsilon_{est})$ ), suggest that the S<sub>n</sub> can be quantified from the 389 390 dielectric probes (see Figure 1A and Figure 2A, Appendix).

391





Figure 9: Evolution of measured permittivity in the 2D tank for 0.5 mm GB with a flow rate of

394

150 mL.min<sup>-1</sup> (with and without enhancements)

 $S_n$  can be quantified accurately from the dielectric probes with and without enhancement, except for the transition zone. It is possible to highlight the differences in remediation yield and rate with permittivity measurements, but the detection area is spatially limited.

399

400 4.2 Electrical resistivity monitoring

# 401 *4.2.1 Experiments without enhancement*

Figure 10 shows changes in resistivity as a function of time during the pumping experiment conducted with 0.5 mm GB with a flow rate of 150 mL.min<sup>-1</sup> (without enhancement). The resistivity values were calculated from Eq. 2 and conductivity measurements obtained from blank test (reference) before starting the experiments.



b) R21 to R35: square detection areas (azimuthal dipole-dipole)

--- (blue): top line
--- (yellow): medium line
--- (red): bottom line

408 Figure 10: Evolution of the resistivity in the left and central parts of the 2D tank for 0.5 mm GB with a
409 flow rate of 150 mL.min<sup>-1</sup> (without enhancement) – a) R1 to R10 and b) R21 to R35

410

The resistivity variations are identical in the left and right parts of the 2D tank; therefore, only the left and middle parts of the 2D tank are shown in Figure 10. Detection areas R1 to R10 were vertical line. R with an even number (R2 to R10) group detection areas are in the upper section; they cover two thirds of the upper portion of the 2D tank (see Figure 3). At the start of the experiment, they were therefore submerged by more than one third of DNAPL (at a content corresponding to  $S_{rw}$ ) and by a little less than two thirds of water ( $S_w$ =1). From the start of pumping one saw that the resistivity values 417 ( $\rho_c$ ) decreased for R10, located in the central section of the cone of depression. Logically the other 418 detection areas located in more lateral positions of the cone of depression had smaller and delayed 419 reductions. We saw the same phenomena for the vertical line detection areas in the lower section (odd 420 numbers between R1 and R9).  $\rho_c$  value is divided by a factor of 4 in the first 5 minutes for R9 (which 421 is in the central part of the 2D tank).

The square detection areas were arranged along three horizontal lines: the top line only in contact 422 423 with water (in blue in Figure 10 b), the medium line at the upper portion of the DNAPL (in yellow), 424 the bottom line at the lower portion of the DNAPL (in red). The resistivities measured in the top line 425 and in the bottom line were relatively constant during the experiment. The resistivities measured at 426 R33 (located in the center of the bottom line) decreased slightly, which shows that the cone of 427 depression partially reached this area. As for the medium line, the resistivities at R34 (located at the 428 center of the 2D tank) decreased faster. At the end of pumping, this area was completely filled by 429 values around  $S_{\rm m}$  (according to the image interpretations).

Figure 11 shows how water saturation, measured and estimated resistivity varied during the pumping experiment. Transforming optical density into S<sub>w</sub> allows a mean S<sub>w</sub> value to be assigned by detection area. This value is transformed into resistivity according to Eq. 14 for 0.5 mm GB (adaptation of the Archie's law from drainage-imbibition experiments [Colombano et al. (2020)].

$$\rho_c = \frac{1}{\sigma_{c,bulk}} = \frac{1}{\sigma_{c,DNAPL} [\emptyset(1 - S_w)]^{1.752} + \sigma_{c,water} [\emptyset(S_w)]^{1.9}}$$
Eq. 14

434



b) R21 to R34: square detection areas (azimuthal dipole-dipole)



438 experiment in the left and central parts of the 2D tank with 0.5 mm GB with a flow rate of

150 mL.min<sup>-1</sup> (without enhancement) – a) R1 to R10 and b) R21 to R34

440

439

Globally, one see that the variations in  $\rho_{c,meas}$  match the variations in  $\rho_{c,est}$ . However, these variations show that the differences between the  $\rho_{c,meas}$  and  $\rho_{c,est}$  are not constant. The fitting for the experiments with Archie's law (during the experiments in 1D cells and 1D columns) was not as good as for the permittivities [Colombano et al. (2020)]. The  $\rho_{c,meas}/\rho_{c,est}$  ratios increase as the resistivity values increase. The differences between the estimated and measured resistivities may have several causes: In a porous medium filled with non-wetting and wetting fluids, considering a suitable
cementation factor that considers the saturation condition, a better estimation of resistivity
with Archie's law would be possible [Byun et al. (2019)].

At a pore scale, the resistivity indexes did not generally obey Archie's law in the non-homogeneous zones (for example with a transition zone constituted of values close to S<sub>m</sub> and S<sub>rw</sub>). The water films play a role not only on the pore space connectivity but also on resistivity; therefore, we can over or underestimate resistivity depending on the thickness of these water films around grains of sand or glass beads [Bernabé et al. (2011); Li et al. (2015); Bernabé et al. (2016)].

456 Many other parameters could be considered at different scales in electrical conductivity 457 models in porous media: tortuosity, pore size distribution, pore-conductance distributions, interconnectivity and universal power law of percolation [Glover (2010); Cai et al. (2017); 458 Ghanbarian and Sahimi (2017)]. Jougnot et al. (2018) reported that the distribution between 459 460 wetting and non-wetting fluids (with different electrical conductivity) in porous media and the 461 resulting heterogeneity lead to electrical current channelling with strong effects on bulk conductivities. They suggest a strong impact of pore-scale on upscaled bulk electrical 462 conductivity in terms of magnitude and anisotropy [Jougnot et al. (2018)]. 463

These parameters have not been considered in our experiments. However, Figure 12 shows that it is possible, as a first approach, to establish correlations between  $\rho_{c,meas}$  and  $\rho_{c,est}$ . Figure 12 compares the measured and estimated resistivity (with image interpretation) in the 2D tank for 0.5 and 0.1 mm GB.



Figure 12: Comparison of measured and estimated resistivity (with image interpretation) in the 2D
tank with a flow rate of 150 mL.min<sup>-1</sup> (without enhancement) for a) 0.5 and b) 0.1 mm GB

The NRMSE were 0.16 and 0.27 for 0.5 and 0.1 mm GB, respectively. The NRMSE values are quite high, which shows that the measured permittivities were not well correlated with the estimated permittivities (based on image interpretation). Moreover, the NRMSE was lower for the 0.5 mm GB than for the 0.1 mm GB because of the fingering process which is more pronounced for 0.1 mm GB. The data dispersivities were quite high, in particular for the 0.1 mm GB for which the fingerings and heterogeneities in DNAPL distribution were clearly visible.

For values corresponding to zones close to the  $S_{rw}$  (*i.e.* when all of the detection area had values close to the  $S_{rw}$ ), the  $\rho_{c,meas}/\rho_{c,est}$  ratios were close to those determined during the 1D cell and 1D column experiments [Colombano et al. (2020)]; these ratios were on average 2.97 for 0.5 mm GB and 4.32 for 0.1 mm GB. In large part, these ratios explain the slopes of above one (y = 1.15x and y = 1.24x for

482 0.5 mm GB and 0.1 mm GB, respectively).

For values close to the  $S_{rn}$  (when all of the detection area has values close to the  $S_{rn}$ ), the  $\rho_{c,meas}/\rho_{c,est}$ ratios were close to those determined during the 1D cell and 1D column experiments; these ratios were on average 1.05 for 0.5 mm GB et 2.3 and 0.1 mm GB (taking into account the correction factors). A transition zone was visible between the two previously mentioned zones. These were the detection areas for the DNAPL-water interface (and therefore, the values close to both  $S_{rw}$  and  $S_m$ ). In this zone, the  $\rho_{c,meas}/\rho_{c,est}$  ratio was variable but globally higher than those determined for the zone relative to  $S_{rm}$ . These ratios are the second explanation for the slopes of the curves. Therefore, we can distinguish, just like for permittivity, four distinct zones (Figure 13):

- 491  $S_{rw}$  zone: start of pumping
- 492 Transition zone: this zone corresponds to values between the S<sub>rw</sub> zone and the S<sub>rm</sub> zone; here,
   493 the majority of points were above the linear regression line (ρ<sub>c,meas</sub> are above ρ<sub>c,est</sub>),
- 494  $S_{rn}$  zone: end of pumping,
- Water zone: this zone corresponds to the detection areas situated in the upper zone of the 2D
  tank.



transition zone. The areas of the bottom line (red dots) are the furthest from the linear regression. The

areas of the top line (green dots), filled largely by water, are obviously closer to the linear regressionline (the calibrations are better for low resistivities).

As for the square detection areas (azimuthal dipole-dipole electrode configurations), the top line measurements (brown dots), corresponding to  $S_w=1$ , agree with the estimation perfectly. The bottomline measurements (pink dots) corresponding to  $S_{rw}$ , are overestimated by the factors previously described. The points outside the group of dots correspond to the arrival of the cone of depression. In that case, the dots get further from the linear regression line (this is the transition zone). Finally, the measurements at the medium line (black dots) correspond in large part to the transition zone and are further from the linear regression line.

For lower resistivity values (corresponding to the detection areas entirely concerned with  $S_{rm}$ ) and for the highest resistivity values (corresponding to zones entirely concerned with  $S_{rw}$ ), we see that the points of course fall closer to the linear regression line.

Figure 14 shows a surface plot of the resistivity measured on an image at the end of pumpingexperiment.

520



521

Figure 14: Surface plot of measured resistivity with an image with 0.5 mm GB and a flow rate of
150 mL.min<sup>-1</sup> (without enhancement) at t=18 min

This image was made using Surfer® software with kriging method only with the square detection areas (R21 to R47). We see that although they are not perfectly superimposed, the resistivity measurements do show the cone of depression.

528

### 529 4.2.2 *Experiments with chemical or thermal enhancement*

Figure 15 shows a comparison of measured resistivity in the 2D tank (with and without thermal and 530 531 chemical enhancements) for 0.5 mm GB with a flow rate of 150 mL.min<sup>-1</sup>. The electrical resistivity 532 monitoring had higher standard deviations (often in the order of 30%) than permittivity monitoring. Therefore, we cannot closely estimate the residual saturations from electrical resistivity alone. 533 However, the precision was sufficient (especially, for values close to the S<sub>m</sub>) to show differences 534 between the three treatment technologies. Figure 15 shows that chemical enhancement has higher 535 remediation yields and rates: the radius of action being higher, the S<sub>n</sub> being lower and S<sub>n</sub> falling more 536 quickly. In contrast, the thermal enhancement was less effective with resistivity drops being lower and 537 538 slower.

539





Figure 15: Evolution of the measured resistivity in the 2D tank for 0.5 mm GB with a flow rate of
150 mL.min<sup>-1</sup> (with and without enhancements)

Figures 3A and 4A (Appendix) compare the measured and estimated resistivity (with image interpretation) in the 2D tank for 0.5 and 0.1 mm GB with chemical and thermal enhancements,

546 respectively. The results are consistent with those obtained without enhancement: we distinguish the 547 same 4 areas and the NRMSE are close to those obtained previously.

548 By combining electrical resistivity monitoring (which gives an integrative view of the cone of 549 depression) with DP (which gives precise but spatially limited data), and image analysing one can 550 better closely calibrate resistivities with S<sub>w</sub> and better quantify remediation rates and yields.

551 As for applications at field conditions, the main methods for indirectly interpolating data between 552 boreholes (with soil analysis) are: piezometers (with water analysis) and soil-gas wells (with gas 553 analysis). Remediation monitoring by geophysical methods (including resistivities for DNAPLs) are 554 complementary and useful to better interpolate data between boreholes. The main monitoring 555 approach is based on differences in resistivities before, during and after treatment. However, quantifying the  $S_n$  remains difficult. Monitoring the permittivities, continuously in the field, would 556 557 contribute to better quantify the S<sub>n</sub> locally and to calibrate resistivity measurements on these points. This permittivity/resistivity coupling would: i. better quantify Sn (during the diagnostic and 558 559 remediation phases); ii. limit field investigations by direct measurements (boreholes, soil-gas wells 560 and piezometers); and iii. improve remediation operations (by optimizing pumping flow rates and radii 561 of cone of depression as a function of the flows and pressures applied).

562

## 563 5. Conclusions

564 The objective of the DNAPL pumping experiments in the 2D tank was to validate permittivity and 565 electrical resistivity measurements and estimations using image interpretation. To achieve this, the 566 average optical densities (corresponding to the detection area for the dielectric probes) were 567 determined throughout the experiments. These optical densities were converted into water saturation 568 then into permittivities using the CRIM model (whose parameters were validated during previous 1D cell and 1D column experiments). Comparison of  $\varepsilon_{meas}$  and  $\varepsilon_{est}$  between experiments without 569 570 enhancement, with chemical enhancement and with thermal enhancement shows very good 571 correlations. For instance, the NRMSE were 0.12 and 0.14 for 0.5 and 0.1 mm GB, respectively. This shows that the measured permittivities agreed well with the image interpretation. As expected, the 572 NRMSE for the 0.5 mm GB was higher than for the 0.1 mm GB which indicates a better correlation of 573

values for 0.5 mm GB experiments, which display less fingering. All experiments clearly displayed 574 575 four distinct zones: S<sub>rw</sub> zone (start of pumping), S<sub>rn</sub> zone (end of pumping), water zone (corresponding 576 to S<sub>w</sub>=1), and transition zone (this zone corresponds to values between the S<sub>rw</sub> and S<sub>rn</sub> zones). The correlations are very good for the first three zones. For the transition zone, we see that the majority of 577  $\varepsilon_{\text{meas}}$  were lower than  $\varepsilon_{\text{est}}$ . This is because when the detection area is crossed by the migration front, 578 579 this area displays both  $S_{m}$  and  $S_{rw}$  areas. Therefore, for two heterogeneous zones, the weight of the portion similar to  $S_{rw}$  impacts on the global  $\epsilon_{meas}$  more than the portion similar to  $S_{rn}$ . This confirms 580 581 that the CRIM model is only valid for homogeneous media. It will be necessary to quantify this 582 transition zone using mixing models.

583 The image interpretation results were compared with measured resistivities. The average optical densities (corresponding to the detection area of the dipole-dipole injection and reception 584 configurations) were determined throughout the experiments. These optical densities were transformed 585 by water saturation into resistivity using Archie's law. The comparison of results on  $\rho_{c,meas}$  and  $\rho_{c,est}$ 586 were not very promising for any of the experiments, with or without enhancements. The NRMSE 587 588 varied between 0.16 and 0.27 for 0.5 and 0.1 mm GB, respectively. The  $\rho_{c,meas}/\rho_{c,est}$  ratios were higher for low  $S_w$  (close to the  $S_{rw}$ ). Note that for low resistivities (*i.e.* high  $S_w$ ), the  $\rho_{c,meas}$ - $\rho_{c,est}$  correlations 589 590 are good. Therefore, it is not possible to quantify water saturations accurately only by measuring 591 electrical resistivities. However, the accuracy is sufficient (especially for values close to S<sub>w</sub>) to 592 highlight differences between the three treatment approaches (i.e. without and with enhancement). 593 Measurements allowed the distinction between four specific zones:  $S_{rw}$  zone, transition zone,  $S_{rn}$  zone, 594 water zone.

595

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609 Appendix

610

Figures 1A and 2A compare measured and estimated permittivity (with image interpretation) in the 2D tank for 0.5 and 0.1 mm GB in the presence of chemical enhancement and with thermal enhancement, respectively.

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Figure 1A: Comparison of measured and estimated permittivity (with image interpretation) in the 2D
tank with a flow rate of 150 mL.min<sup>-1</sup> (with chemical enhancement) for a) 0.5 and b) 0.1 mm GB



Figure 2A: Comparison of measured and estimated permittivity (with image interpretation) in the 2D
tank with a flow rate of 150 mL.min<sup>-1</sup> (with thermal enhancement) for a) 0.5 and b) 0.1 mm GB

Adding SDBS at its CMC has no effect on permittivity measurements; permittivities are therefore estimated using Eq. 13. The NRMSE are 0.10 and 0.13 for 0.5 and 0.1 mm GB, respectively (Figure 1A). The NRMSE for the 0.5 mm GB is logically lower than those of the 0.1 mm GB.

Figure 2A compares measured and estimated permittivity (with image interpretation) in the 2D tank for 0.5 and 0.1 mm GB and with thermal enhancement. The permittivity values were influenced by temperature variations. The variations in  $\varepsilon_n$  and  $\varepsilon_w$  as a function of temperature were quantified and integrated in the CRIM model (Eq. 13). These variations were estimated using Eq. 1A and Eq. 2A for 0.5 mm GB [Colombano et al. (2020)]:

$$\varepsilon_n = 0.0367.T + 4.51$$
 Eq. 1A

$$\varepsilon_w = 0.0231.T + 31.50$$
 Eq. 2A

629 where T (°C) is the temperature in Celsius.

As shown in Figure 1A and Figure 2A, the NRMSE are between 0.10 and 0.14. The NRMSE for the 0.5 mm GB are logically lower than those of the 0.1 mm GB (0.10 vs 0.14 for example for thermal enhancement) because the fingering is more important. In the case of thermal and chemical enhancements, four zones are clearly distinguished:  $S_{rw}$  zone, transition zone,  $S_{rm}$  zone and water zone. As stated above, in the transition zone the majority of the  $\varepsilon_{meas}$  were below  $\varepsilon_{est}$ .

635

Figures 3A and 4A compare the measured and estimated resistivity (with image interpretation) inthe 2D tank for 0.5 and 0.1 mm GB with chemical and thermal enhancements, respectively.





640 Figure 3A: Comparison of measured and estimated resistivity (with image interpretation) in the 2D

tank with a flow rate of 150 mL.min<sup>-1</sup> (with chemical enhancement) for a) 0.5 and b) 0.1 mm GB

642





Adding SDBS at its CMC has a negligible impact on the measured electrical resistivities
[Colombano et al. (2020)]. Eq. 14 was therefore considered to estimate the resistivity. The NRMSE

648 were 0.17 and 0.27 for 0.5 and 0.1 mm GB, respectively (Figure 3A), which means that the data 649 dispersivity is higher with the 0.1 mm GB. The  $\rho_{c,meas}/\rho_{c,est}$  ratios for values close to S<sub>rw</sub> are high: 3.20 650 and 5.99 for 0.5 and 0.1 mm GB, respectively.

Figure 4A compares the measured and estimated resistivity (with image interpretation) in the 2D tank for 0.5 and 0.1 mm GB. The resistivity estimates were calculated by considering Eq. 14 and the resistivity changes related to the increase in temperature. The increase in temperature results in an increase in ionic mobility and a decrease in resistivities [Dakhnov (1962); Hayashi (2004); Light et al. (2005); Grellier et al. (2008)]. The decrease in the resistivity of water and the DNAPL was determined on the basis of experiments and the Dakhnov equation (Eq. 3A) [Colombano et al. (2020)]:

$$\rho_{c,f} = \frac{\rho_{c,f_0}}{1 + \alpha_{c,f}(T - T_0)}$$
 Eq. 3A

657 where  $\rho_{c,f}$  ( $\Omega$ .m) is the electrical resistivity of the fluid at temperature *T* (°C),  $\rho_{c,f_0}$  ( $\Omega$ .m) is the 658 electrical resistivity of the fluid at temperature T<sub>0</sub>, and  $\alpha_{c,w}$  (°C<sup>-1</sup>) is the temperature coefficient of 659 resistivity ( $\alpha_{c,w}\approx 0.023$  °C<sup>-1</sup> for T<sub>0</sub>=23 °C, and 0.025 °C<sup>-1</sup> for T<sub>0</sub>=0 °C)

The NRMSE were 0.18 and 0.27 for 0.5 and 0.1 mm GB, respectively (Figure 3A). The dispersivity was higher with 0.1 mm GB than with 0.5 mm GB. Here also, we note that the  $\rho_{c,meas}/\rho_{c,est}$ ratios for values close to S<sub>rw</sub> were high: 4.46 and 9.37 for 0.5 and 0.1 mm GB, respectively. As stated above, we also distinguished four zones: S<sub>rw</sub> zone, transition zone, S<sub>rn</sub> zone, water zone.

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# 667 Notations

ac	Empirical parameter	-
$c_{\epsilon}$	Speed of light (velocity of electromagnetic waves) in vacuum	m.s <sup>-1</sup>
$K_g$	Geometric coefficient	m
$L_{arepsilon}$	Waveguide length	m
$m_{1,} m_{2}$	Cementation exponent of DNAPL and water phases	-
m <sub>c</sub>	Cementing factor	-
$R_e$	Electrical resistance measured with SIP system using a reference liquid	Ω
$S_{n,} S_{w}$	DNAPL (non-wetting fluid) and water (wetting fluid) saturations	-
Т	Temperature	°C
t	Travel time for the pulse to traverse the length of the embedded	S
$\iota_{\mathcal{E}}$	waveguide	3
$v_i$	Volume fraction of the <i>i</i> phase	-
$\alpha_{c,w}$	Temperature coefficient of resistivity	$^{\circ}C^{-1}$
a	Empirical constant related to the geometry of the grains and their	
$u_{\epsilon}$	spatial distribution	-
$\varepsilon = \varepsilon_r$	Relative permittivity	-
ε <sub>0</sub>	Free space permittivity	F.m <sup>-1</sup>
<u>C</u> .	Relative permittivity of air measured at the beginning of the	
Eair	experiment	-
$\epsilon_{air-theoretical}$	Theoretical relative permittivity of pure air ( $\varepsilon_{air-theoretical}=1$ )	-
ε'	Real dielectric permittivity	F.m <sup>-1</sup>
C	Effective relative permittivity of the medium estimated with the	
Eest	water saturation (based on image interpretation)	-
ε <sub>i</sub>	Relative permittivity of the <i>i</i> phase	-
ε <sub>meas</sub>	Effective relative permittivity of the medium measured	-
ε <sub>n</sub>	Relative permittivity of DNAPL	-
ε <sub>s</sub>	Relative permittivity of soil particles	-
$\epsilon_{\rm w}$	Relative permittivity of water	-
$\epsilon_{w-theoretical}$	Relative permittivity of pure water ( $\varepsilon_{w-theoretical}$ =80 at 20 °C)	-
$ ho_c$	Electrical resistivity of the bulk	$\Omega.m$
0	Electrical resistivity of the bulk estimated with the water saturation	0 m
$\rho_{c,est}$	(based on image interpretation)	\$2,111
$ ho_{c,f}$	Electrical resistivity of the fluid at temperature T	$\Omega.m$
$\rho_{c,f0}$	Electrical resistivity of the fluid at temperature T <sub>0</sub>	$\Omega.m$
$\rho_{c,meas}$	Electrical resistivity measured of the bulk	$\Omega.m$
$\rho_{c,w}$	Electrical resistivity of the water at temperature T	$\Omega.m$
$ ho_0$	Measured resistivity value using a reference liquid	$\Omega.m$
$\sigma_{c,bulk=}\sigma_c$	Electrical conductivity of the bulk	S.m <sup>-1</sup>
$\sigma_{c,DNAPL}$	Electrical conductivity of DNAPL	S.m <sup>-1</sup>
$\sigma_{c,water}$	Electrical conductivity of water	$S.m^{-1}$
$\upsilon_{\epsilon}$	Function of the propagation velocity	m.s <sup>-1</sup>
Ø	Porosity	-

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