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▶ To cite this version:

P. Riviere-Marichalar, D. Barrado, B. Montesinos, G. Duchêne, H. Bouy, et al.. Gas and dust in the beta Pictoris moving group as seen by the Herschel Space Observatory. Astronomy and Astrophysics - A&A, 2014, 565, 10.1051/0004-6361/201322901. insu-03618582

HAL Id: insu-03618582 https://insu.hal.science/insu-03618582

Submitted on 24 Mar 2022

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Gas and dust in the beta Pictoris moving group as seen by the *Herschel* Space Observatory*

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Received 23 October 2013 / Accepted 27 March 2014

ABSTRACT

Context. Debris discs are thought to be formed through the collisional grinding of planetesimals, and then can be considered as the outcome of planet formation. Understanding the properties of gas and dust in debris discs can help us comprehend the architecture of extrasolar planetary systems. Herschel Space Observatory far-infrared (IR) photometry and spectroscopy have provided a valuable dataset for the study of debris discs gas and dust composition. This paper is part of a series of papers devoted to the study of Herschel-PACS observations of young stellar associations.

Aims. This work aims at studying the properties of discs in the beta Pictoris moving group (BPMG) through far-IR PACS observations of dust and gas.

Methods. We obtained Herschel-PACS far-IR photometric observations at 70, 100, and 160 μ m of 19 BPMG members, together with spectroscopic observations for four of them. These observations were centred at 63.18 μ m and 157 μ m, aiming to detect [OI] and [CII] emission. We incorporated the new far-IR observations in the SED of BPMG members and fitted modified blackbody models to better characterise the dust content.

Results. We have detected far-IR excess emission towards nine BPMG members, including the first detection of an IR excess towards HD 29391. The star HD 172555, shows [OI] emission, while HD 181296 shows [CII] emission, expanding the short list of debris discs with a gas detection. No debris disc in BPMG is detected in both [OI] and [CII]. The discs show dust temperatures in the range 55–264 K, with low dust masses ($<6.6 \times 10^{-5} M_{\oplus}$ to $0.2 M_{\oplus}$) and radii from blackbody models in the range 3 to <82 AU. All the objects with a gas detection are early spectral type stars with a hot dust component.

Key words. astrobiology – circumstellar matter – stars: evolution – Kuiper belt: general – astrochemistry

1. Introduction

Debris discs represent the last stage in the evolution of protoplanetary discs into planetary systems. The regions surrounding the star are replenished with dust through collisional cascades triggered by large bodies inside rings of circumstellar mater that resemble our own asteroid and Kuiper belts. Therefore, debris discs can be considered signposts to the likely existence of planets, or at least of comets (see e. g. Matthews et al. 2014, and references therein). Since the discovery of IR excess emission around the main sequence star Vega (Aumann et al. 1984), hundreds of stars with debris discs have been detected, mainly thanks to IRAS, ISO, and *Spitzer* Space Telescope observations.

Studies made with *Spitzer* Space Telescope data revealed a fraction (~15%) of debris discs around A to K stars (Su et al. 2006; Siegler et al. 2007; Trilling et al. 2007; Hillenbrand et al. 2008; Fujiwara et al. 2013).

The properties and evolution of debris discs have been extensively studied over the past decades (Wyatt 2008; Krivov 2010, and references therein). Most studies focussed on deriving dust disc properties by studying the shape of the spectral energy distribution (SED, see e.g. Bryden et al. 2006; Rebull et al. 2008). Nevertheless, resolved imaging has helped us understand the geometry of debris discs: about 100 of the brightest debris discs have been resolved at various wavelengths, from optical to the sub-mm (see e.g. Smith & Terrile 1984; Koerner et al. 1998; Schneider et al. 2006; Smith et al. 2009). They show different structures such as inner gaps, rings, clumps, spiral arms, and asymmetries, which can be signposts to the likely existence of

^{*} *Herschel* is an ESA space observatory with science instruments provided by European-led Principal Investigator consortia and with important participation from NASA.

planets (see e.g. Moro-Martin et al. 2007). Although dust has been extensively observed, studies have provided little information about its chemical composition. So far, only a few debris discs have shown solid-state features in their mid-IR spectra, indicating the presence of small and warm grains (Chen et al. 2006; Fujiwara et al. 2013). These debris discs are most probably in a transient state, and the small dust grains are produced during periods of intense collisional grinding (Lisse et al. 2008, 2009).

Gas in debris discs has only rarely been detected, although we need to know the total amount of gas present to understand the physical evolution of these systems. Only two debris-disc systems show traces of sub-millimetre CO emission, namely 49 Ceti and HD 21997, implying that the gas abundance in debris discs is low when compared with protoplanetary discs (Dent et al. 2005). Only the β Pic debris disc appears to have a full inventory of gaseous species, both atomic and molecular (see e.g. Lagrange et al. 1998; Thi et al. 2001; Roberge et al. 2006).

The Herschel Space Observatory (Pilbratt et al. 2010) has produced a valuable dataset for the study of circumstellar environments, allowing both the dust in the continuum and gas emission in the far-IR to be observed. The Photodetector Array Camera & Spectrometer-PACS (Poglitsch et al. 2010) can study debris discs with high sensitivity in the wavelength range (55–220 μ m) where cold dust emission peaks. PACS spectroscopy allows for detecting a few bright emission lines for molecular species such as CO, H2O, and OH and the strong cooling lines of [OI] and [CII] at 63 and 158 μ m, respectively. Some Herschel key time programmes, such as DUNES (Eiroa et al. 2010) and DEBRIS (Matthews et al. 2010), have focussed on studying debris discs in the far-IR. Dust around NEarby Stars (DUNES, Eiroa et al. 2013) studied a sample of 124 main sequence stars within 20 pc of the Sun and found an excess due to circumstellar dust in 25 of them (debris-disc rate ~20%), with a mean detection limit in the dust fractional luminosity of $L_{\rm d}/L_{*} \sim 2.0 \times 10^{-6}$. Some of the discs in the DUNES sample show a dust fractional luminosity that is only a few times higher than our Kuiper Belt, but these discs seem to be cooler and larger. Additionally, they discovered a new class of cool debris discs that show excess at 160 μ m but little or non excess at $100 \,\mu \text{m}$ (see Krivov et al. 2013).

Young stellar associations (Torres et al. 2006) most likely provide the best frame to study the evolution of circumstellar environments, because the ages of their members are well known. In this paper we present Herschel-PACS observations of the beta Pictoris moving group (BPMG), a young stellar association with a typical distance of ~30 pc. The most recent estimation of the age of the system by Binks & Jeffries (2014) $(21 \pm 4 \text{ Myr})$ is in good agreement with the $20 \pm 10 \text{ Myr}$ age derived by Barrado y Navascués et al. (1999). The star that names the group, β Pic, is an A6V star that harbours one of the best studied debris discs and that shares its space motion with a group of coeval stars that constitute the BPMG. Nineteen BPMG members were observed with PACS as part of the Herschel open time key programme GAS in Protoplanetary Systems (GASPS, P.I. W. Dent, Dent et al. 2013) that has observed ~250 stars in different associations with different ages, from Taurus (1–3 Myr) to Tucana Horologium (~30 Myr). Photometry at 70 and/or 100 and 160 μ m was obtained for the 19 stars. Spectroscopic observations were also obtained for four of them. We compared their SEDs with modified blackbody models and derived disc properties, such as inner radii and dust masses based on blackbody modeling results.

Table 1. Stellar parameters for BPMG members observed with *Herschel*-PACS.

Name	d	Sp. type	$T_{ m eff}$	L_*
_	(pc)	_	(K)	$L_{ m sun}$
AT Mic	10.7	M4V	3100	0.066
CD64-1208	29.2	K5V	4200	0.25
GJ 3305	29.8	M1V	3600	0.166
HD 203	39.1	F3V	6600	3.7
HD 29391	29.8	F0V	7400	5.43
HD 35850	26.8	F7V	6000	1.26
HD 45081	38.5	K4V	4200	0.329
HD 139084	39.8	K0V	5000	1.23
HD 146624	43.1	A0V	9750	22
HD 164249	46.9	F6V	6600	2.7
HD 172555	29.2	A6IV	7800	7.8
HD 174429	49.7	G9IV	5200	0.85
HD 181296	47.7	A0V	10000	23.1
HD 181327	50.6	F6V	6600	3.2
HD 199143	47.7	F7V	6000	1.93
HIP 10679	34.0	G2V	5800	0.87
HIP 10680	39.4	F5V	6200	2.4
HIP 11437	42.3	K6V	4400	0.242
HIP 12545	40.5	K6V	4000	0.238

Notes. Spectral types are taken from Zuckerman et al. (2001), Zuckerman & Song (2004), Torres et al. (2006). Stellar distances are taken from Zuckerman & Song (2004).

2. The sample

The BPMG members studied, together with their distances, spectral types, effective temperatures, and luminosities, are listed in Table 1. Archival photometry data were collected for each BPMG member in the sample, including Johnson, Stromgren, 2MASS, IRAC, WISE, AKARI, MIPS (Rebull et al. 2008), LABOCA (Nilsson et al. 2009, 2010), and SMA. We also included Spitzer-IRS data in the analysis, by collapsing the spectra taken from the archive to photometric points by using 1 μ m bins, including a 5% calibration uncertainty in the errors (see Chen et al. 2006). The different IRS spectral orders were scaled to be consistent with the photospheric flux at short wavelengths when possible. The inclusion of IRS spectra is a central interest, given that the onset of the excess is seen at mid-IR wavelengths. The lack of agreement between IRS data and WISE fluxes made us decide to exclude WISE data from the model analysis in Sect. 5.

The first step before modelling the whole SED is to derive the stellar parameters. For each star in the sample, we selected photometric points with a pure photospheric origin (typically in the $0.3-3.0\,\mu\mathrm{m}$ range) and used them to compare the non-excess SED with theoretical models using the Virtual Observatory SED Analyzer¹ (VOSA, Bayo et al. 2008). VOSA compares the photometric data with a grid of models to match the photosphere that produces the best fit, based on a χ^2 minimisation, and, if desired, on Bayesian analysis. The values for $T_{\rm eff}$ and L_* from the best fits are summarised in Table 1. In Fig. 1 we plot the Hertzsprung-Russell diagram for the stars in the sample.

3. Observations and data reduction

Nineteen BPMG members were observed in photometric mode and a subset of four BPMG members were also observed in spectral line $(63/189 \ \mu m)$ and range mode $(78/157 \ \mu m)$; see

http://svo2.cab.inta-csic.es/theory/vosa/

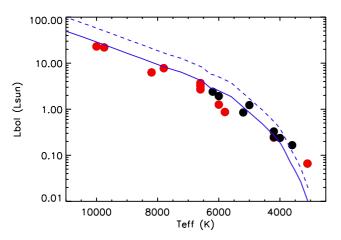


Fig. 1. Hertzsprung-Russell diagram for BPMG members in the sample. The blue solid line represents the 20 Myr isochrone, while the blue dashed one represents a 20 Myr isochrone with 2 times larger luminosity to include unresolved binary systems. The isochrones come from the models by Siess et al. (2000). Red filled circles are objects detected by PACS while black filled circles are objects not detected in PACS photometry.

Table 2 for an overview of the observations containing the PACS OBSIDs).

3.1. Photometric observations and data reduction

The PACS photometer simultaneously observes in either the 70 or 100 μm band, together with the 160 μm band, so when the source have been observed in both 70 and 100 μm , we have two images to combine in the 160 μm band. We observed 16 objects at 70 μm and another 16 at 100 μm , making a total of 19 observations in the 160 μm band (see Table 2). The exposure times range from 133 to 1122 s, based on the expected flux from the star. Each scan map was made with medium speed (20" s^-1), with scan legs of 3 arcmin and a separation of 4–5" between legs.

We reduced the photometric observations using HIPE 8, with the most recent version of the calibration files, in the same fashion as our TW Association data (Riviere-Marichalar et al. 2013). For bright IR-excess targets (i.e., those with flux greater than ~100 mJy), we used a version of the pipeline tuned for bright objects, while for faint objects and non-detected objects we used a different version optimised for noise-dominated maps. Both pipelines shared the following reduction steps: bad and saturated pixel flagging and removal, flat field correction, deglitching, high pass filtering, and map projection. We refer the reader to Riviere-Marichalar et al. (2013) for a detailed description of the differences between both pipelines. Photometric maps were projected into the final image with pixel scale 2"/pixel in the 70 and 100 μ m bands and with pixel scale 3"/pixel in the 160 μ m band. We also produced final maps with the native pixel scale of the detector that were used to perform the error calculation (3.2'')pixel for the 70/100 μ m bands and 6.4 for the 160 μ m band). When several images at the same wavelength are available for a single target, we combined all of them to improve the signal-to-noise ratio (S/N), averaging for each pixel and using the average sigma clipping algorithm to exclude bad pixels.

Aperture photometry was measured using an aperture of 6'' for the 70 and 100 μm bands and 12" for the 160 μm band. The annulus for sky subtraction was placed at 25–35" from the

star. We then applied an aperture correction for each band². Final fluxes are listed in Table 3. Noise errors consist of the standard deviation of the photometry obtained at several sky positions surrounding the target. PACS calibration uncertainties are 2.64, 2.75, and 4.15% for the 70, 100, and 160 μ m bands, respectively. Noise errors and calibration errors were added quadratically.

For non-detected sources, we derive upper limits as follows: we compute aperture photometry in the sky background in several pointings that surround the nominal position of the star on the detector. Then we compute the standard deviation, and we use it as the sky background value. The upper limits included in Table 3 are 3σ .

3.2. Spectroscopic observations and data reduction

Four BPMG members (HD 164249, HD 172555, HD 181296, and HD 181327) were observed with PACS in LineScan spectroscopic mode, targeting [OI] emission at 63.18 μ m. Riviere-Marichalar et al. (2012b) reported the detection of o-H₂O emission at 63.32 μm in the PACS spectra of eight T Tauri stars in Taurus, therefore we also searched for line emission at this wavelength. All the targets were also observed in RangeScan mode, aiming to detect [CII] emission at 157.74 μ m. PACS spectra were reduced using HIPE 9 with the latest version of the pipeline and the proper calibration files. Spectra were extracted from the central spaxel and aperture-corrected to account for flux spread in the surrounding spaxels. The line spectra from PACS typically show lower S/N near the spectrum edges. To account for that effect, we exclude from the spectra any wavelength shorter than 63.0 μ m or longer than 63.4 μ m in the 63 μ m range. In the RangeScan observations we exclude wavelengths shorter than 157 μ m and longer than 159 μ m.

Line fluxes were computed by applying a Gaussian fit to the line profile and calculating the integrated flux from that fit. Upper limits were computed by integrating a Gaussian with a width equal to the instrumental FWHM at the central wavelength, and maximum equal to three times the standard deviation of the continuum (upper limits are 3σ). HD 172555 was observed at two different epochs in LineScan mode, the emission being detected in both observations. When combining both epochs, small shifts in the line centres translate into an increase in noise. To avoid this effect, we computed line fluxes for the averaged spectrum after re-centring both spectra by artificially shifting them in a way that the line centre from the fit is exactly at the rest frame wavelength of the observed line. Line fluxes are shown in Table 4.

4. Results

4.1. Herschel/PACS photometry

Photometry results are shown in Table 3. Following Bryden et al. (2006), we used the parameters χ_{70} , χ_{100} , and χ_{160} to identify excess sources, defined as

$$\chi_{\text{band}} = \frac{F_{\text{obs,band}} - F_{*,\text{band}}}{\sigma_{\text{band}}} \tag{1}$$

where $F_{\rm obs,band}$ is the observed flux at any of the 70, 100, or 160 μ m bands, $F_{*,band}$ is the expected photospheric flux at the corresponding wavelength, and $\sigma_{\rm band}$ is the corresponding error,

² http://herschel.esac.esa.int/twiki/pub/Public/ PacsCalibrationWeb/

Table 2. Observation IDs for BPMG member observed with Herschel-PACS.

Name	Photometry OBSID	Wavelength ^a	Spectroscopy OBSID	Spectral ranges
_		(μm)		(μm)
AT Mic	1342196104, 1342209488, 1342209489	70, 70, 100	_	_
CD64-1208	1342209059, 1342209060, 1342209061	70,100, 70, 100	_	_
GJ 3305	1342224850, 1342224851, 1342224852, 1342224853	70, 70, 100, 100	_	_
HD 203	1342188366, 1342221118, 1342221119	70, 100, 100	_	_
HD 29391	1342190967, 1342216153, 1342216154	70, 100, 100	_	_
HD 35850	1342217746, 1342217747	70, 100	_	_
HD 45081	1342188506, 1342212832, 1342212833	100, 70, 70	_	_
HD 139084	1342216483, 1342216484, 1342216485, 1342216486	70, 70, 100, 100	_	_
HD 146624	1342215617, 1342215618, 1342215619, 1342215620	70, 70, 100, 100	_	_
HD 164249	1342183657, 1342215574, 1342215575	70, 100, 100	1342215648, 1342239388	63, 157
HD 172555	1342209059, 1342209060, 1342209061	70,100, 70, 100	1342215649, 1342228416, 1342228417	63, 157, 63
HD 174429	1342215576, 1342215577, 1342215578, 1342215579	70, 70, 100, 100	_	_
HD 181296	1342290955, 1342209056	100, 100,	1342209730, 1342239756	63, 157
HD 181327	1342183658, 1342209057, 1342209058	70, 100, 100	1342186311, 1342186810	63, 157
HD 199143	1342193550, 1342208861, 1342208862	70, 70, 70	_	_
HIP 10679	1342189193, 1342223862, 1342223863	70, 100, 100	_	_
HIP 10680	1342189193, 1342223862, 1342223863	70, 100, 100	_	_
HIP 11437	1342189210, 1342223864, 1342223865	70, 100, 100	_	_
HIP 12545	1342189150, 1342223574, 1342223575	70, 100, 100	_	_

Notes. (a) PACS photometer simultaneously observes in either the 70 or 100 μ m bands plus the 160 μ m band (see Sect. 3.1).

Table 3. *Herschel/*PACS photometry and estimated fluxes from the naked photosphere.

Name -	$F_{\nu,\text{obs}}(70 \mu\text{m})$ (mJy)	$F_{\nu, \text{phot}}(70 \mu\text{m})$ (mJy)	$F_{\nu,\text{obs}}(100\mu\text{m})$ (mJy)	$F_{\nu, \text{phot}}(100 \mu\text{m})$ (mJy)	$F_{\nu,\text{obs}}(160 \mu\text{m})$ (mJy)	$F_{\nu, \text{phot}}(160 \mu\text{m})$ (mJy)
AT Mic	13 ± 2	15 ± 3	_	7.3 ± 0.8	<14	2.8 ± 0.3
CD64-1208	<8	3.2 ± 0.4	<9	1.6 ± 0.2	<8	0.65 ± 0.08
GJ 3305	<4	3.1 ± 0.5	<4	1.5 ± 0.2	<8	0.6
HD 203	68 ± 3	6.5 ± 0.4	26 ± 2	3.2 ± 0.2	<14	1.24 ± 0.08
HD 139084	<7	4.2 ± 0.5	<10	2.0 ± 0.2	<29	0.81 ± 0.09
HD 146624	13 ± 2	9.1 ± 0.2	<7	4.47 ± 0.09	<13	1.75 ± 0.03
HD 164249	_	3.4 ± 0.2	493 ± 20	1.66 ± 0.07	249 ± 10	0.65 ± 0.03
HD 172555	191 ± 7	16.5 ± 0.5	79 ± 4	8.0 ± 0.3	31 ± 2	3.2 ± 0.1
HD 174429	<7	2.6 ± 0.4	<7	1.3 ± 0.2	<13	0.58 ± 0.07
HD 181296	_	7.8 ± 0.2	250 ± 8	3.81 ± 0.07	111 ± 6	1.49 ± 0.03
HD 181327	1453 ± 6	3.4 ± 0.2	1400 ± 40	1.62 ± 0.9	850 ± 40	0.63 ± 0.04
HD 199143	<5	3.6 ± 0.4	_	1.8 ± 0.2	<11	0.71 ± 0.07
HD 29391	22 ± 2	12.6 ± 0.8	17 ± 3	6.2 ± 0.4	< 16	2.4 ± 0.2
HD 35850	_	8.6 ± 0.6	42 ± 2	4.2 ± 0.3	18 ± 4	1.7 ± 0.1
HD 45081	<13	1.8 ± 0.2	<10	0.9 ± 0.1	<16	0.35 ± 0.05
HIP 10679	53 ± 3	2.7 ± 0.3	46 ± 2	1.3 ± 0.1	39 ± 5	0.51 ± 0.05
HIP 10680	<13	4.4 ± 0.3	<8	2.2 ± 0.02	<22	0.85 ± 0.06
HIP 11437	70 ± 4	1.34 ± 0.03	69 ± 3	0.65 ± 0.02	50 ± 4	0.255 ± 0.006
HIP 12545	<7	1.5 ± 0.2	_	0.7 ± 0.1	<19	0.29 ± 0.05

Table 4. Herschel/PACS spectroscopy.

Name	[OI] flux	S/N	[CII] flux	S/N
		$63 \mu m$		$157 \mu m$
_	(10^{-18} W/m^2)	_	(10^{-18} W/m^2)	_
HD 164249	<8.2	2.6	<3.6	0.3
HD 172555 ¹	9.7 ± 2.0	3.0	< 2.5	1.5
HD 181296	< 6.2	3.6	2.3 ± 0.6	3.2
HD 181327	<8.2	10.4	<7.6	5.8

Notes. Columns are target name, [OI] line flux at 63.18 μ m, S/N of the continuum at 63 μ m, [CII] line flux at 157.74 μ m, and S/N of the continuum at 157 μ m. ⁽¹⁾ Recentred and averaged spectrum.

which is the quadratic sum of the photometric error and the photosphere model error. To compute the error from the models, we

produced models with temperatures that are the best fit temperature $\pm 200~\rm K$ for each star and derived the errors as the mean difference between the best fit model and those models. We consider that excess sources at each band are those with $\chi_{band} > 3$.

The agreement between MIPS photometry (Rebull et al. 2008) and PACS photometry is generally good, with a mean difference smaller than 10%. We detected eight objects at 70 μ m out of 16 observed. Two of them, namely AT Mic and HD 146624, showed fluxes in agreement with photospheric emission ($\chi_{70} \simeq -0.7$ and 2.1, respectively). Plavchan et al. (2009) report the MIPS flux at 24 μ m towards AT Mic to be in excess over the photosphere; however, they predict a photospheric flux at 24 μ m of 114 mJy, while we predict a flux of 126 mJy ($\chi_{24} \sim 2.7$). Nevertheless, the authors did not include model errors when computing χ_{24} : if we include the model error at 24 μ m, we get $\chi_{24} \sim 0.4$. We conclude that the 24 μ m flux is compatible with

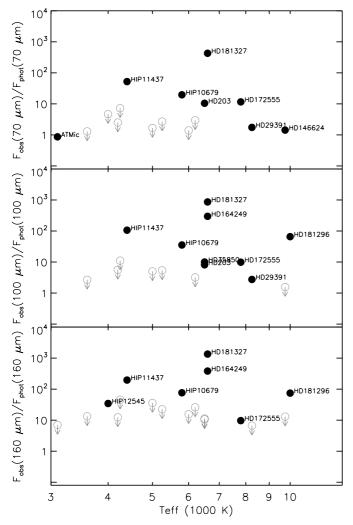


Fig. 2. Excess for the 70/100/160 μ m band versus $T_{\rm eff}$. Filled dots represent true detections, while empty circles with arrows represent 3σ upper limits.

pure photospheric emission. At 70 μ m, we measure a flux of (13 ± 2) mJy, compared to the value of 22 ± 5 mJy by Plavchan et al. (2009): we speculate that the larger MIPS flux is due to the larger MIPS beam size that results in a higher background pollution. HD 29391 showed a flux of 22 ± 2 mJy at 70 μ m and a flux of 17 ± 3 mJy at 100 μ m, while the expected photospheric fluxes are ~13 mJy and ~6.2 mJy, respectively. The difference between the expected and the observed flux is more than 3σ in both bands ($\chi_{70} \simeq 4.6$, $\chi_{100} \simeq 3.6$), and therefore HD 29391 has an infrared excess.

In Fig. 2 we show the ratio of the observed fluxes to the photospheric fluxes at $70/100/160~\mu m$ versus the effective temperature of the star. Objects not detected at $70~\mu m$ show upper limits that are one to ten times larger than the expected photospheric value. Upper limits at $100~\mu m$ also range from one to ten times the photospheric value. Upper limits at $160~\mu m$ are one to 100 times greater than the photospheric value. We can't exclude the presence of very cold, faint discs for some of the non-detected sources. We find no correlation between the strength of the IR excess at any of the PACS bands and the temperature of the central object, as other authors have reported (see Trilling et al. 2008; Eiroa et al. 2013).

To study whether we have spatially resolved the sources, we performed azimuthally averaged radial profiles in all the three PACS bands for every detected object and compared the results with the azimuthally averaged radial profile of a model point spread function source. Only HD 181327 seems to be more extended than the reference star α Boo in the 100 μ m band, with a FWHM of 7.58", compared to 6.94" for α Boo, a result already reported by Lebreton et al. (2012).

4.2. Herschel/PACS spectroscopy

Among the four BPMG members observed with PACS in LineScan mode, we detected the continuum level at 63 μ m for three of them, HD 164249 being the only exception. Continuum subtracted spectra at 63 μ m and 157 μ m are shown in Figs. 3 and 4 respectively.

The presence of atomic oxygen in HD 172555 was discussed in Riviere-Marichalar et al. (2012b), and we review the main results in Sect. 6.2. HD 181296 shows [CII] emission at 157.74 μ m with a S/N of 3.8. This is the first detection of gas emission towards HD 181296, adding HD 181296 to the short list of debris disc with a gas detection. We discuss the detection in detail in Sect. 6.2. None of the systems show water emission at 63.32 μ m.

5. Blackbody models

For each BPMG member in the sample, we have built the SED and included the new PACS photometric data, aiming to compare the observed photometry with synthetic models to study the BPMG disc properties. The simplest and most robust way of studying debris disc properties in BPMG members is to fit modified blackbody models to the dust emission. The photospheric contribution is estimated using photospheric models described by the stellar parameters listed in Table 1, with solar metallicity and $\log g = 4.5$. The modified blackbody is described by

$$F \propto F_{\nu}(T_{\text{dust}}) \times (\lambda_0/\lambda)^{\beta} \tag{2}$$

where F_{ν} is the flux emitted by a blackbody at a temperature $T_{\rm dust}$, $\beta=0$ if $\lambda<\lambda_0$. The parameters β , λ_0 and $T_{\rm dust}$ are free to vary. Here, λ_0 can be interpreted as a characteristic grain size, via $\lambda_0=2\pi a$ where a is the average grain size.

Such simple models provide estimates of the dust temperature and IR excess. The best fit model was determined through a reduced χ^2 minimisation, with

$$\chi_{\nu}^{2} = \frac{1}{\nu} \sum_{i=1}^{N} \frac{(F_{\text{obs},i} - F_{\text{mod},i})^{2}}{\sigma_{i}^{2}}$$
 (3)

where N is the total number of photometric points used in the fitting, v = N - n the degrees of freedom, n the number of free parameters, $F_{\text{obs},i}$ the observed flux of the ith photometric point, $F_{\text{mod},i}$ the i-th model flux, and σ is the error of the ith photometric point. Because the stellar properties are fixed, we only included excess points in the analysis. For the sources with IRS data available, we used the IRS spectra turning point as the starting point for the excess. When no IRS spectra were available, we just used the first photometric point showing excess.

We used a genetic algorithm $(GAbox^3)$ to explore the parameter space and find the set of free parameters that minimises Eq. (3). GAbox allowed us to rapidly identify the set of parameters that best fits our observations. Then, it conducts the estimation of the uncertainties by creating a small grid of parameters

³ This algorithm has already been used in Riviere-Marichalar et al. (2013) with similar purposes and is explained in Appendix A of Lillo-Box et al. (2014).

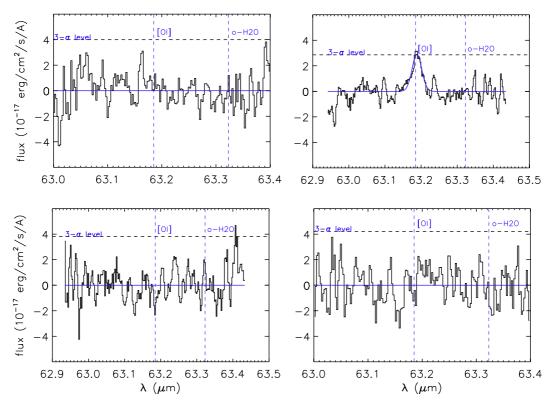


Fig. 3. Continuum-subtracted spectra of BPMG members at $63~\mu m$. The vertical blue dashed lines represent the position of the [OI] and o-H₂O. From left to right and top to bottom, targets are: HD 164249, HD 172555 (averaged and re-centred with respect to the rest frame wavelength), HD 181296, and HD 181327. We show 3σ limits as horizontal black dashed lines. For HD 172555 we show in blue a Gaussian fit to the data, while for the other sources the blue horizontal line depicts the position of the continuum.

around the best solution and computes the 3σ confidence levels for each parameter pair. The largest uncertainty among all pairs involving a particular parameter is then selected as the final uncertainty.

From the dust temperature of the blackbody models, the lower limit on the disc radius can be estimated by using

$$R_{\rm in} > \frac{1}{2} R_* \left(\frac{T_*}{T_{\rm dust}} \right)^2 \tag{4}$$

where T_* and R_* are the star effective temperature and radius. We note here that inner radii derived from blackbody dust temperatures are lower limits to the actual disc radii (see e.g. Rodriguez & Zuckerman 2012). The dust mass can be estimated using

$$M_{\text{dust}} = \frac{F_{\nu}(\lambda_{\text{Obs}})D^2}{\kappa_{\nu}B_{\nu}(T_{\text{dust}})}$$
 (5)

where D is the distance to the star, $B_{\nu}(T_{\rm dust})$ can be approximated by the Rayleigh-Jeans regime, $\kappa_{\nu} = 2 \times (1.3 \text{ mm}\lambda)^{\beta} \text{ cm}^2 \text{ g}^{-1}$ (Beckwith et al. 1990), and $F_{\nu}(\lambda_{\rm Obs})$ is the observed integrated flux density at a given wavelength emitting in the Rayleigh-Jeans regime. We use the flux at 160 μ m, since dust emission is always in the Rayleigh-Jeans regime for BPMG members at this wavelength (see Harvey et al. 2012; Riviere-Marichalar et al. 2013).

For the sources without detected excess, we computed upper limits on the infrared excess luminosity by using the upper limit with the smaller excess over the photosphere to scale the flux of a 70 K pure blackbody ($\beta=0$). The temperature was chosen to be similar to the median temperature of the objects with excess. Temperatures higher than 70 K can result in an excess over the photosphere at $\lambda \sim 30 \, \mu \mathrm{m}$ where the IRS spectra do not show any excess. For AT Mic, the detection of the photosphere

at 70 μ m makes the blackbody model fitting more restrictive: models with $T_{\rm dust} \geq 25$ K result in an excess flux at 70 μ m that is not observed, therefore we fixed the dust temperature to be 20 K. Results of the fitting are listed in Table 5, and model SEDs can be found in Fig. 5.

6. Discussion

In Fig. 6 we show an histogram of the infrared fractional excess $(L_{\rm IR}/L_*)$ distribution of BPMG members observed with PACS, including β Pic data from Rebull et al. (2008). The infrared fractional excess is computed from blackbody models as the fraction of dust luminosity over the stellar luminosity. The distribution of detected sources peaks at $\sim 10^{-4}$, while the distribution of non-detected sources upper limits peaks at $\sim 10^{-6}$. However, the faint tail of the distribution of detected objects coincides with the peak of the distribution for non-detected objects. The IR excess in BPMG members is due to the presence of a debris disc, with dust temperatures in the range 55 to 264 K, and a median temperature of 79 K. In the BPMG sample of objects observed with PACS, six out of nine objects have temperatures in the range 55 to 96 K and are therefore Kuiper belt-type discs. However, three out of nine, namely HD 203, HD 172555, and HD 181296, show temperatures that are much higher and may consist of a warm component or both a warm and a cold component, such as the one found in β Pic itself. Previous work suggests that a significant fraction of debris discs might have two-component architecture (see e.g. Morales et al. 2009; Su et al. 2013). Derived dust masses range from $<6.1 \times 10^{-5} M_{\oplus}$ to $0.20 M_{\oplus}$ and have an inner radii from 2.0 to 60 AU.

The detection ratio of objects with an IR-excess is \sim 50%, much higher than the standard value of 15% to 20% for older

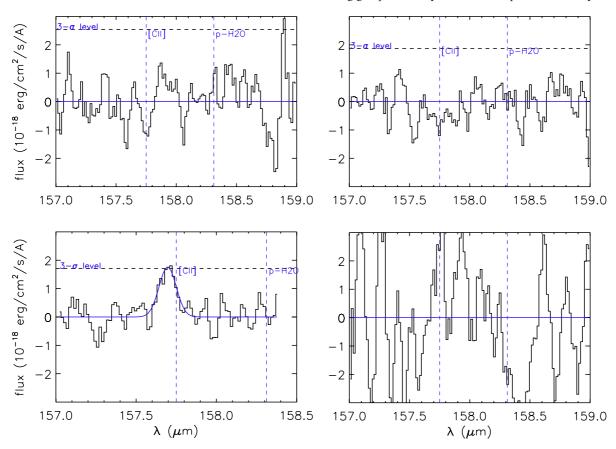


Fig. 4. Continuum-subtracted spectra of BPMG members at 157 μ m. The vertical blue dashed lines represent the position of the [OI] and o-H₂O. From left to right and top to bottom, targets are HD 164249, HD 172555 (averaged and re-centred with respect to the rest frame wavelength), HD 181296, and HD 181327. We show 3σ limits as horizontal black dashed lines. For HD 181296 we show in blue a Gaussian fit to the data, while for the other sources the blue horizontal line depicts the position of the continuum. The noise in HD 181327 spectrum is more than two times greater than in the other spectra.

Table 5. Black body models.

Name	T	β	λ_0	$L_{ m IR}/L_*$	$R_{ m dust}$	$M_{ m dust}$
_	(K)	_	(µm)	_	(AU)	(M_{\oplus})
HD 203	128^{+10}_{-11}	$1.6^{+0.56}_{-0.50}$	$58.4^{+8.0}_{-8.3}$	1.6×10^{-4}	9.1 ± 3.0	$<6.6 \times 10^{-5}$
HD 29391	55^{+21}_{-21}	$0.0^{+2.0}$	9.5 ⁺¹¹¹	2.3×10^{-6}	82^{+677}_{-75}	$<1.6 \times 10^{-3}$
HD 35850	$77.6^{+8.5}_{-15}$	$0.7^{+1.3}_{-0.7}$	100^{+60}_{-98}	3.7×10^{-5}	$16.7^{+10.3}_{-7.5}$	$(3.4 \pm 1.4) \times 10^{-4}$
HD 164249	$71.15^{+0.99}_{-0.54}$	$1.16^{+0.84}_{-0.90}$	119 ± 6	8.4×10^{-4}	25.2 ± 4.6	$(6.8^{+15.7}_{-5.3}) \times 10^{-3}$
HD 172555	$264.2^{+8.5}_{-8.4}$	$0.45^{+0.19}_{-0.11}$	$29.9^{+11}_{-8.9}$	7.5×10^{-4}	2.76 ± 0.77	$(2.7 \pm 1.2) \times 10^{-4}$
HD 181296	$161.7^{+2.9}_{-3.0}$	$0.0^{+2.0}$	(160_{-47})	2.4×10^{-4}	17.0 ± 1.8	$(1.3 \pm 0.08) \times 10^{-2}$
HD 181327	$79.04^{+0.07}_{-0.33}$	$0.0^{+0.20}$	(311^{+353}_{-290})	2.9×10^{-3}	22.6 ± 2.9	(0.20 ± 0.06)
HIP 10679	$96.2^{+2.9}_{-3.1}$	$0.0^{+2.0}$	(160_{-147})	3.0×10^{-4}	8.7 ± 2.5	$(3.7 \pm 0.6) \times 10^{-3}$
HIP 11437	$65.5^{+22}_{-2.9}$	$0.1^{+1.9}_{-0.1}$	(18^{+179}_{-15})	1.0×10^{-3}	8.3 ± 1.7	$(9.8 \pm 0.9) \times 10^{-3}$
AT Mic	20	0.0	_	$<2.5 \times 10^{-5}$	-	_
CD 64-1208	70	0.0	_	$<4.2 \times 10^{-5}$	_	_
GJ 3305	70	0.0	_	$<4.5 \times 10^{-5}$	_	_
HD 45081	70	0.0	_	$<9.4 \times 10^{-6}$	_	_
HD 139084	70	0.0	_	$< 8.4 \times 10^{-6}$	_	_
HD 146624	70	0.0	_	$< 5.0 \times 10^{-7}$	_	_
HD 174429	70	0.0	_	$< 2.3 \times 10^{-5}$	_	=
HD 199143	70	0.0	_	$< 2.6 \times 10^{-6}$	_	=
HIP 10680	70	0.0	_	$< 8.1 \times 10^{-6}$	_	-
HIP 12545	70	0.0	_	$<4.2 \times 10^{-5}$	_	-

Notes. For objects without a detected excess, we fixed $\beta = 0.0$.

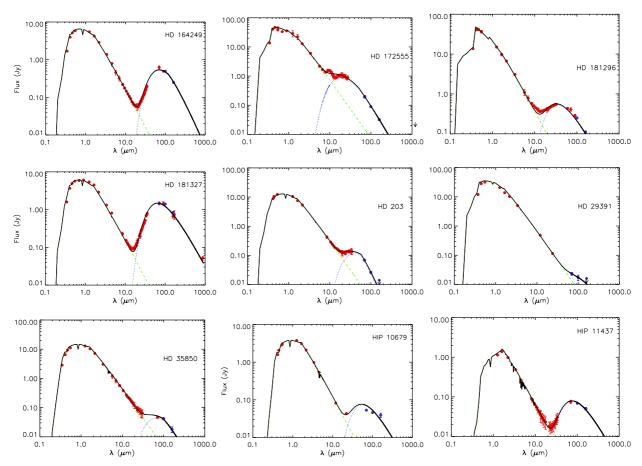


Fig. 5. Blackbody models for BPMG members detected with PACS. Red dots are photometric points from the catalogues and literature (see Sect. 2). Blue dots are PACS observations from this work. The black solid line represents the photosphere plus blackbody model. The doted blue line represents the blackbody that better fits the dust. The green dashed line represents the photospheric contribution.

systems (Beichman et al. 2006; Bryden et al. 2006). A trend towards a higher detection rate of debris discs among young stars is noted by Rieke et al. (2005). However, our ratio is computed using a small sample that is not robust against selection effects.

6.1. Notes on the individual SEDs

HD 203 was not detected at 160 μ m. The best fit temperature was $T_{\rm dust} = 128$ K, in good agreement with Rebull et al. (2008). The way far-IR emission plummets around 160 μ m suggests an under-abundance of large grains. An unusual, steep decrease in the SED in the PACS wavelength range of three debris discs in the DUNES survey is reported by Ertel et al. (2012), who define

$$\Delta_{\nu_1,\nu_2} = \frac{\log F_{\nu_2} - \log F_{\nu_1}}{\log \nu_2 - \log \nu_1}$$
 (6)

and find $\Delta_{70,100}$ values in the range 1.94 to 2.66, while we get $\Delta_{70,100} = 2.70$ for HD 203. The authors consider that the steep decay of the SED is the result of a deviation from the standard equilibrium in a collisional cascade, with a significative underabundance of large grains. A similar SED decay in the PACS wavelength range is also reported by Donaldson et al. (2012) for HD 3003, where a value of –4.4 was used for the exponent of the grain size distribution. Overall, it seems that we need a grain size distribution that is different from the expected in a collisional cascade in equilibrium to explain steep SEDs.

Our observations of HD 29391 show for the first time that the star is surrounded by a debris disc that produces a very tiny excess $(L_{\rm IR}/L_*=2.0\times10^{-6})$, the faintest in the sample), at a temperature of only 55 K, the coldest disc detected in BPMG. The lower limit on the inner radius is 82 AU, the largest among the sample, but it is poorly constrained due to large errors in $T_{\rm dust}$. We computed an upper limit in dust mass of $M_{\rm dust}<1.6\times10^{-3}~M_{\oplus}$.

HD 172555 is an A7 star that harbours a warm debris disc with solid silicate feature emission in its IRS spectrum (Chen et al. 2006) and with an infrared fractional luminosity that is ~86 times higher than the maximum value expected from steadystate collisional evolution (Wyatt et al. 2007). Lisse et al. (2009) analysed its IRS spectrum and propose that both an SiO gas tentative detection and a silicate feature around 10 μ m could be explained as the outcome of a hypervelocity collision (>10 km s⁻¹) between two planetary mass objects. The best modified blackbody model for HD 172555 has $T_{\text{dust}} = 264 \text{ K} \ (\beta = 0.5)$, the warmest in the sample. Although it is a poor fit for the IRS data, it demonstrates there is warm dust in the disc. After re-analyzing TReCs observations (Moerchen et al. 2010), Smith et al. (2012) show that the disc emission is resolved in the Q band (18.3 μ m), arising from r > 8 AU, but not in the N band (11.66 μ m), with the emission coming from 1.0 to 7.9 AU from the central star. These results suggest that there may be two dust populations in the system.

The SED of HD 181327, including the PACS data from GASPS, has recently been modelled by Lebreton et al. (2012), taking the imaging constraints by Schneider et al. (2006) into account, with a dust mass of $0.05~M_{\oplus}$ in grains from 1 to $1000~\mu m$ composed of a mixture of silicates, carbonaceous material and

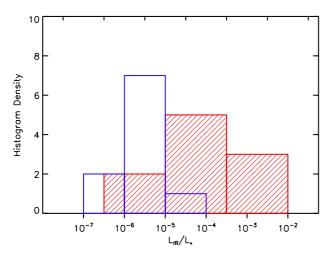


Fig. 6. Distribution of the infrared fractional excess for objects detected (red histogram filled with lines) and non-detected (blue histogram) with PACS.

amorphous ice with significant porosity. The best fit temperature for HD 181327 is 79 K, which results in a lower limit for the disc radius of ~23 AU, with a dust mass of 0.20 M_{\oplus} . We note here that the minimum radius of the disc is well below the value of 86 AU by Schneider et al. (2006), obtained by means of NICMOS coronographic observations of scattered light. That blackbody radii are a lower limit to actual debris disc radii is a well known fact (see e.g. Booth et al. 2013).

The IRS spectrum for HIP 10679 was too noisy to be included in the analysis of the SED. In Fig. 5 we can see that with a single blackbody fit we can not reproduce its IR emission. The best fit model, with $T=(92\pm4)$ K and $\beta=0.0$, produces a poor fit that clearly underestimates the flux at 160 μ m, while overestimating the flux at 70 μ m, arguing for a flatter slope. Our best fit temperature results in an inner radius of (8.7 ± 2.5) AU, much smaller than the value of 35 AU proposed by Rebull et al. (2008).

In Fig. 5 we can see that a single blackbody also fails to reproduce the IR emission towards HD 181296. The best fit model, with $T=161~\rm K$ and $\beta=0.0$ underestimates the flux at far-IR wavelengths. The poor fit quality might be explained by the need for a second, warm blackbody. The system was imaged by Smith et al. (2009), who derived $R_{\rm out}=45~\rm AU$ and propose that two different grain populations should be present.

6.2. Gas in BPMG debris discs

There are three circumstellar systems in the BPMG with detected gas emission: β Pic, HD 172555, and HD 181296. Given the age of the BPMG, a primordial origin for the gas is unlikely. The disc around β Pic has long been known to possess large amounts of gas (Hobbs et al. 1985; Lagrange et al. 1986; Roberge et al. 2000; Olofsson et al. 2001; Thi et al. 2001). Lagrange et al. (1987) showed that UV absorption lines profiles in β Pic show variations on time scales of months, signature of infalling gas (Lagrange-Henri et al. 1989). Olofsson et al. (2001) show Na I gas in Keplerian rotation. A braking mechanism is needed to keep Na stable against radiation pressure; Fernández et al. (2006) proposed that an overabundance of C by a factor of 10 with respect to solar abundance is enough to brake the gas. Observations with the Far-UV Space Explorer (FUSE) have shown that C is indeed overabundant by a factor of 20 with respect to other species (Roberge et al. 2006).

The detection of atomic oxygen emission at 63.18 μ m towards HD 172555 has previously been reported in Riviere-Marichalar et al. (2012a), and therefore we refer the reader to this paper, where a possible origin for the emission in a hypervelocity collision is discussed (see also Lisse et al. 2009; Johnson et al. 2012). Lisse et al. (2009) proposed a tentative detection of SiO gas, later on confirmed by Johnson et al. (2012), who concluded that the amount of oxygen gas proposed by Riviere-Marichalar et al. (2012a) is enough to keep SiO vapour from being destroyed by photo-dissociating photons. More recently, Kiefer et al. (2014) reported the detection of variable absorption signatures in the Ca II H and K lines, a signature of falling evaporating bodies (FEB), highlighting the similarities between β Pic and HD 172555.

We report the first detection of [CII] emission at 157 μ m towards HD 181296, with a flux of $(1.7 \pm 0.4)10^{-18}$ W/m². This is the third case of a debris disc where [CII] emission is present but no [OI] emission is observed, the other two systems being HD 32297 (Donaldson et al. 2013) and 49 Cet (Roberge et al. 2013). Following Roberge et al. (2013), and assuming the [CII] emission is optically thin, we can use

$$M_{\rm [CII]} = \frac{4\pi\lambda_{\rm [CII]}}{hc} \frac{F_{\rm int}md^2}{A_{10}\chi_{\rm u}} \tag{7}$$

to estimate a lower limit for the [CII] gas mass, where $\lambda_{\rm [CII]}$ is the rest frame wavelength of the transition, $F_{\rm int}$ the integrated line flux, m the mass of an atom, d the distance to the source, $A_{10} = 2.4 \times 10^{-6} \ {\rm s}^{-1}$ the spontaneous transition probability, and $\chi_{\rm u}$ the fraction of atoms in the upper level. Assuming local thermal equilibrium (LTE), we can compute $\chi_{\rm u}$ using

$$\chi_{\rm u} = \frac{(2J_{\rm u} + 1)}{Q(T_{\rm ex})} e^{-E_{\rm ul}/KT_{\rm ex}}$$
 (8)

where $J_{\rm u}$ is the angular momentum quantum number of the upper level, $g_{\rm u}$ the statistical weight of the upper level, $E_{\rm ul}$ the energy difference between the upper and the lower levelS ($E_{\rm l}/k=91.21$ K), $T_{\rm ex}$ the excitation temperature and $Q(T_{\rm ex})$ the partition function, which depends on the excitation temperature. Assuming a two-level atom, $Q(T_{\rm ex})$ is

$$Q(T_{\rm ex}) = g_{\rm l} + g_{\rm u} e^{-E_{\rm ul}/kT_{\rm ex}}$$
(9)

where g_1 is the statistical weight of the lower level. In Fig. 7 we show the distribution of $M_{\rm [CII]}$ as a function of $T_{\rm ex}$. Different excitation temperatures have led to very different gas masses, but we computed a lower limit of $8.1 \times 10^{-5}~M_{\oplus}$ valid in the range 1–2000 K. To translate this [CII] mass lower limit into a total C mass lower limit, we need knowledge about the ionization fraction. Following Roberge et al. (2013), we assume that HD 181296 has the same ionization fraction as β Pic (50%), given that they have a similar spectral type, both have optically thin debris discs, and assume that the ionization fraction is similar for both infrared fractional luminosities. A lower limit to the total carbon mass is then $1.6 \times 10^{-4}~M_{\oplus}$. If we assume that the carbon has a solar abundance (solar carbon mass fraction is 0.288%), we get $M_{\rm gas} \geq 0.056~M_{\oplus}$.

The origin of the gas in HD 181296 is difficult to assess without more species detected. The most probable mechanisms generating secondary gas in debris discs are grain-grain collisions, sublimation of dust grains, evaporation of comets, and photodesorption from dust grain surfaces. Photo-desorption of CO from grain surfaces, followed by photo-dissociation, should leave a similar amount of C and O atomic gas in the system, but we do not detect any [OI] emission. Therefore, it is unlikely

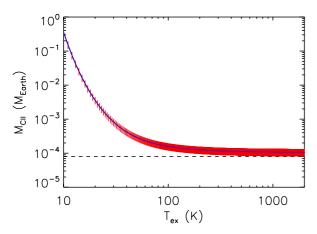


Fig. 7. Gas mass versus excitation temperature for [CII] in HD 181296. The red, vertical bars show the error in the [CII] gas mass determination. The doted, horizontal bar depicts the position of the lower limit to the [CII] gas mass.

Table 6. Overview of gas detections in debris discs.

Name -	Sp. type	Age (Myr)	T _{dust} (K)	Reference
βPic	A61	20 ± 10	130 ²	3
σ Her	$\mathrm{B9^4}$	200^{5}	200^{6}	4
HD 21997	$A3^7$	20 ± 10^{8}	52 ⁹	10
HD 32297	$A0^{11}$	30^{12}	240^{13}	13
HD 172555	$A6^{11}$	20 ± 10	264	14
HD 181296	$A0^{11}$	20 ± 10	162	15
49 Cet	$A1^{16}$	$15-180^{17}$	175^{16}	16
51 Oph	$B9.5^{18}$	$0.2-1.1^{17,*}$	$400-1000^{19}$	20

Notes. $T_{\rm dust}$ refers to the temperature of the warm dust when two components are present. ⁽¹⁾ Gray et al. (2006). ⁽²⁾ Rebull et al. (2008). ⁽³⁾ Vidal-Madjar et al. (1998). ⁽⁴⁾ Chen & Jura (2003). ⁽⁵⁾ Grosbol (1978). ⁽⁶⁾ Fajardo-Acosta et al. (1998). ⁽⁷⁾ Houk & Smith-Moore (1988). ⁽⁸⁾ Moór et al. (2006). ⁽⁹⁾ Nilsson et al. (2010). ⁽¹⁰⁾ Moór et al. (2011). ⁽¹¹⁾ Torres et al. (2006). ⁽¹²⁾ Kalas (2005). ⁽¹³⁾ Donaldson et al. (2013). ⁽¹⁴⁾ Riviere-Marichalar et al. (2012b). ⁽¹⁵⁾ Present work. ⁽¹⁶⁾ Roberge et al. (2013). ⁽¹⁷⁾ Montesinos et al. (2009). ⁽¹⁸⁾ Thi et al. (2013). ⁽¹⁹⁾ Fajardo-Acosta et al. (1993). ⁽²⁰⁾ Lecavelier Des Etangs et al. (1997). ^(*) Thi et al. (2013) consider it is a debris disc despite its young age.

that photo-desorption of CO is the main production mechanism. However, Roberge et al. (2000, 2006) show that in the case of β Pictoris only 2% of carbon gas could come from dissociation of CO, implying that many other materials could be acting as [CII] suppliers, such as amorphous carbon.

It is interesting to note that the presence of gas might be connected with the presence of a warm disc: three out of three BPMG members with a gaseous debris discs also have a hot excess component. In Table 6 we list debris discs with confirmed gas detections. We do not include the four shell stars proposed by Roberge & Weinberger (2008) in the list nor TWA 04B (Riviere-Marichalar et al. 2013) or HD 141569A (Thi et al. 2013), because their evolutionary stages need to be clarified. Seven out of eight debris discs with a solid gas detection have a warm disc component, the only exception being HD 21997, with a dust temperature of 52 K (Nilsson et al. 2010). However, Kóspál et al. (2013) argue that the gas surrounding HD 21997 could be a remnant of primordial gas. While the connection between the presence of hot dust and that of gas requires further

research, the significance of this trend is tantalizing. Considering this, HD 203 is a good candidate for future gas observations.

Another interesting fact is that the three stars share similar spectral spectral type (A0, for HD 181296, to A6 for HD 172555 and β Pic), which could imply that photo-desorption of gas from grain surfaces plays a major role. The relation is also patent when all debris discs with a solid gas detection are included: all eight stars in Table 6 are early spectral type stars (B9 to A6).

The presence of gas does not seem to be influenced by the strength of the infrared excess: HD 181327 shows an infrared excess that is more than three times that of HD 172555, yet it is not detected; HD 1624249 also shows a prominent excess, larger than the infrared excess for both HD 172555 and HD 181296; but again it shows no gas emission within the sensitivity of Herschel-PACS. The continuum emission at 160 μ m towards HD 172555 and HD 181296 is 0.031 and 0.111 Jy, respectively, compared to 0.850 Jy towards HD 181327 and 0.249 Jy towards HD 164249, so if continuum plays a major role in gas emission both HD 164249 and HD 181327 would have been easily detected.

Considering the variety of gas properties in BPMG, it seems that different mechanisms can operate at the same age, and stochastic events may play an important role by producing huge amounts of gas in a short period of time, while a more quiescent gas production results from grain dust collisions, evaporation, and cometary activity. Future research is needed to understand the gas production mechanisms in debris discs, and the identification of observables useful for distinguishing the various mechanisms is a main goal for future work. The detection of other gaseous species towards HD 172555 and HD 181296 will help us to understand the origin of the gas. Such observations can be made with ALMA, with the aim of detecting species like SiO, CO, and C I.

7. Summary and conclusions

We observed 19 BPMG members with the *Herschel* Space Observatory instrument PACS. Our main conclusions follow.

- 1. We have detected infrared emission in excess at $70/100/160~\mu m$ in 8/9/7 out of the 16/16/19 objects observed at each band. The detection ratio is greater than the typical value of 15-20% for older systems, but we highlight that our sample is small and not robust against selection effects.
- 2. We detected emission at 70 μ m towards AT Mic, HD 146624, and HD 29391 for the first time, HD 29391 was also detected at 100 μ m. The 70 μ m excess flux in HD 29391 results in a very low infrared fractional luminosity of $L_{\rm IR}/L_* \sim 2.0 \times 10^{-6}$.
- 3. We modelled the dust emission with modified blackbody models. Blackbody temperatures range from 58 K (HD 129391) to 229 K (HD 172555). Dust masses range from $6.6 \times 10^{-5}~M_{\oplus}$ to 0.20 M_{\oplus} and black body estimates for the inner radii range from 2 to ~82 AU.
- 4. The object HD 203 shows an unusually steep SED in the $70-100~\mu m$ range that can be a signpost of an underabundance of large grains.
- 5. One of the systems studied, HD 172555, shows [OI] emission at 63 μ m. Another of the systems studied, HD 181296, shows [CII] emission at 157 μ m.

HD 172555 and HD 181296 are both warm debris discs, and we note that the presence of gas could be linked to hot dust, in such a

way that the same phenomenon producing the gas also produces a population of small grains near the star.

Acknowledgements. We acknowledge the anonymous referee for a very interesting discussion that helped improve the quality of the paper. This research has been funded by Spanish grants BES-2008-003863, AYA 2010-21161-C02-02, AYA 2012-38897-C02-01, AYA 2011-26202, and PRICIT-S2009/ESP-1496. I. Kamp and P. Rivié-Marichalar acknowledge funding from an NWO MEERVOUD grant. We also acknowledge support from ANR (contract ANR-07-BLAN-0221) and PNPS of CNRS/INSU, France. C. Pinte acknowledges funding from the European Commission's 7th Framework Programme (contract PERG06-GA-2009-256513) and from Agence Nationale pour la Recherche (ANR) of France under contract ANR-2010-JCJC-0504-01.

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