



HAL
open science

Constraints on the richness-mass relation and the optical-SZE positional offset distribution for SZE-selected clusters

A. Saro, S. Bocquet, E. Rozo, B. A. Benson, J. Mohr, E. S. Rykoff, M. Soares-Santos, L. Bleem, S. Dodelson, P. Melchior, et al.

► **To cite this version:**

A. Saro, S. Bocquet, E. Rozo, B. A. Benson, J. Mohr, et al.. Constraints on the richness-mass relation and the optical-SZE positional offset distribution for SZE-selected clusters. Monthly Notices of the Royal Astronomical Society, 2015, 454, pp.2305-2319. 10.1093/mnras/stv2141 . insu-03644672

HAL Id: insu-03644672

<https://insu.hal.science/insu-03644672>

Submitted on 25 Apr 2022

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Constraints on the richness–mass relation and the optical-SZE positional offset distribution for SZE-selected clusters

A. Saro,^{1,2★} S. Bocquet,^{1,2} E. Rozo,³ B. A. Benson,^{4,5,6} J. Mohr,^{1,2,7} E. S. Rykoff,^{8,9} M. Soares-Santos,⁴ L. Bleem,^{5,10} S. Dodelson,^{4,5,6} P. Melchior,^{11,12} F. Sobreira,^{4,13} V. Upadhyay,¹⁴ J. Weller,^{2,7,15} T. Abbott,¹⁶ F. B. Abdalla,¹⁷ S. Allam,⁴ R. Armstrong,¹⁸ M. Banerji,¹⁹ A. H. Bauer,²⁰ M. Bayliss,^{21,22} A. Benoit-Lévy,¹⁷ G. M. Bernstein,¹⁸ E. Bertin,²³ M. Brodwin,²⁴ D. Brooks,¹⁷ E. Buckley-Geer,⁴ D. L. Burke,^{8,9} J. E. Carlstrom,^{5,6,25} R. Capasso,^{1,2} D. Capozzi,²⁶ A. Carnero Rosell,^{13,27} M. Carrasco Kind,^{28,29} I. Chiu,^{1,2} R. Covarrubias,²⁹ T. M. Crawford,^{5,6} M. Crocce,²⁰ C. B. D’Andrea,²⁶ L. N. da Costa,^{13,27} D. L. DePoy,³⁰ S. Desai,¹ T. de Haan,^{31,32} H. T. Diehl,⁴ J. P. Dietrich,^{1,2} P. Doel,¹⁷ C. E Cunha,⁸ T. F. Eifler,^{18,33} A. E. Evrard,^{34,35} A. Fausti Neto,¹³ E. Fernandez,³⁶ B. Flaugher,⁴ P. Fosalba,²⁰ J. Frieman,^{4,5} C. Gangkofner,^{1,2} E. Gaztanaga,²⁰ D. Gerdes,³⁵ D. Gruen,^{7,15} R. A. Gruendl,^{28,29} N. Gupta,^{1,2} C. Hennig,^{1,2} W. L. Holzappel,³² K. Honscheid,^{11,12} B. Jain,¹⁸ D. James,¹⁶ K. Kuehn,³⁷ N. Kuropatkin,⁴ O. Lahav,¹⁷ T. S. Li,³⁰ H. Lin,⁴ M. A. G. Maia,^{13,27} M. March,¹⁸ J. L. Marshall,³⁰ Paul Martini,^{11,38} M. McDonald,³⁹ C.J. Miller,^{34,35} R. Miquel,³⁶ B. Nord,⁴ R. Ogando,^{13,27} A. A. Plazas,^{33,40} C. L. Reichardt,^{32,41} A. K. Romer,⁴² A. Roodman,^{8,9} M. Sako,¹⁸ E. Sanchez,⁴³ M. Schubnell,³⁵ I. Sevilla,^{28,43} R. C. Smith,¹⁶ B. Stalder,^{22,44} A. A. Stark,²² V. Strazzullo,^{1,2} E. Suchyta,^{11,12} M. E. C. Swanson,²⁹ G. Tarle,³⁵ J. Thaler,⁴⁵ D. Thomas,²⁶ D. Tucker,⁴ V. Vikram,¹⁰ A. von der Linden,^{8,46,47} A. R. Walker,¹⁶ R. H. Wechsler,^{8,9,46} W. Wester,⁴ A. Zenteno⁴⁸ and K. E. Ziegler⁵

Affiliations are listed at the end of the paper

Accepted 2015 September 14. Received 2015 September 9; in original form 2015 June 25

ABSTRACT

We cross-match galaxy cluster candidates selected via their Sunyaev–Zel’dovich effect (SZE) signatures in 129.1 deg² of the South Pole Telescope 2500d SPT-SZ survey with optically identified clusters selected from the Dark Energy Survey science verification data. We identify 25 clusters between $0.1 \lesssim z \lesssim 0.8$ in the union of the SPT-SZ and redMaPPer (RM) samples. RM is an optical cluster finding algorithm that also returns a richness estimate for each cluster. We model the richness λ -mass relation with the following function $\langle \ln \lambda | M_{500} \rangle \propto B_\lambda \ln M_{500} + C_\lambda \ln E(z)$ and use SPT-SZ cluster masses and RM richnesses λ to constrain the parameters. We find $B_\lambda = 1.14^{+0.21}_{-0.18}$ and $C_\lambda = 0.73^{+0.77}_{-0.75}$. The associated scatter in mass at fixed richness is

* E-mail: saro@usm.lmu.de

$\sigma_{\ln M|\lambda} = 0.18^{+0.08}_{-0.05}$ at a characteristic richness $\lambda = 70$. We demonstrate that our model provides an adequate description of the matched sample, showing that the fraction of SPT-SZ-selected clusters with RM counterparts is consistent with expectations and that the fraction of RM-selected clusters with SPT-SZ counterparts is in mild tension with expectation. We model the optical-SZE cluster positional offset distribution with the sum of two Gaussians, showing that it is consistent with a dominant, centrally peaked population and a subdominant population characterized by larger offsets. We also cross-match the RM catalogue with SPT-SZ candidates below the official catalogue threshold significance $\xi = 4.5$, using the RM catalogue to provide optical confirmation and redshifts for 15 additional clusters with $\xi \in [4, 4.5]$.

Key words: catalogues – methods: data analysis – galaxies: abundances – galaxies: clusters: general – galaxies: haloes – galaxies: statistics – cosmology: miscellaneous – large-scale structure of Universe.

1 INTRODUCTION

Clusters of galaxies were first identified as overdense regions in the projected number counts of galaxies (e.g. Abell 1958; Zwicky, Herzog & Wild 1968). Nowadays, clusters are also regularly identified through their X-ray emission (e.g. Gioia et al. 1990; Vikhlinin et al. 1998; Böhringer et al. 2000; Pacaud et al. 2007; Šuhada et al. 2012), and at millimetre wavelengths through their Sunyaev–Zel’dovich effect (SZE) signatures (Sunyaev & Zel’dovich 1972). Large, homogeneously selected samples of clusters are useful for both cosmological and astrophysical studies, and such samples have recently begun to be produced using SZE selection (Staniszewski et al. 2009; Hasselfield et al. 2013; Bleem et al. 2015b; Planck Collaboration 2015). There is a longer history of large cluster samples selected from optical and near-infrared photometric surveys (e.g. Gladders & Yee 2000; Koester et al. 2007; Eisenhardt et al. 2008; Hao et al. 2010; Menanteau et al. 2010; Wen, Han & Liu 2012; Ascaso, Wittman & Dawson 2014; Bleem et al. 2015a; Rykoff et al. 2014, and references therein), and even larger samples will soon be available from ongoing and future surveys like the Dark Energy Survey (DES; The Dark Energy Survey Collaboration 2005),¹ KiDS (de Jong et al. 2013), *Euclid* (Laureijs et al. 2011) and LSST (LSST Dark Energy Science Collaboration 2012).

Reliable estimates of galaxy cluster masses play a key role in both cosmological and astrophysical cluster studies. First, the abundance of galaxy clusters as a function of mass is a well-known cosmological probe (White, Efstathiou & Frenk 1993; Bartlett & Silk 1994; Eke et al. 1998; Viana & Liddle 1999; Borgani et al. 2001; Vikhlinin et al. 2009; Mantz et al. 2010; Rozo et al. 2010; Allen, Evrard & Mantz 2011; Benson et al. 2013; Bocquet et al. 2015; Mantz et al. 2015, and many others). Secondly, accurate estimates of cluster masses are crucial in disentangling environmental effects from the secular evolution processes shaping galaxy formation (Mei et al. 2009; Zenteno et al. 2011; Muzzin et al. 2012).

In this paper, we calibrate the richness–mass relation for SZE-selected galaxy clusters detected in the DES science verification data (SVA1) using the red-sequence Matched-Filter Probabilistic Percolation (redMaPPer; Rykoff et al. 2014) cluster-finding algorithm. Specifically, we study the clusters detected via their SZE signatures in the South Pole Telescope SPT-SZ cluster survey (Bleem et al. 2015b, hereafter B15) that are also present in the redMaPPer catalogue. We also study the distribution of offsets between the SZE-derived centres and the associated optical centres, properly including the SZE positional uncertainties. Finally, we demonstrate

our ability to push to even lower candidate significance within the SPT-SZ candidate catalogue by taking advantage of the contiguous, deep, multiband imaging available through DES. In this respect, our study points towards the combined use of DES and SPT data sets to provide highly reliable extended SZE-selected cluster samples. We note that historically the optical follow-up of SPT-selected clusters was the original motivation for proposing DES.

The plan of the paper is as follows. In Section 2, we describe the galaxy cluster catalogues and the matching metric we use in this work. Section 3 describes the method we adopt to calibrate the SZE–mass and richness–mass relations. Our results are presented in Section 4. Section 5 contains a discussion of our findings and our conclusions. In the appendix, we provide a preliminary analysis of a cluster sample created using an independent cluster finding algorithm – the Voronoi Tessellation (VT) cluster finder – which helps to highlight areas where the VT algorithm can be improved. Throughout this work, we adopt $\Omega_M = 0.3$, $\Omega_\Lambda = 0.7$, $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$ and $\sigma_8 = 0.8$. Cluster masses are defined within R_{500} , the radius within which the density is 500 times the critical density of the Universe. Future analyses will include the dependence of the derived scaling–relation parameters on the adopted cosmology by simultaneously fitting for cosmological and scaling–relation parameters (e.g. Mantz et al. 2010; Rozo et al. 2010; Bocquet et al. 2015).

2 CLUSTER SAMPLE

2.1 SPT-SZ cluster catalogue

The SPT-SZ galaxy cluster sample used in this analysis has been selected via the cluster thermal SZE signatures in a point-source masked-region of 2365 deg^2 of the 2540 deg^2 (2500d) SPT-SZ survey using 95 and 150 GHz data. Typical instrumental noise is approximately 40 (18) μK_{CMB} arcmin and the beam FWHM is 1.6 (1.2) arcmin for the 95 (150) GHz maps. A multi-frequency matched filter is used to extract the cluster SZE signal in a manner designed to optimally measure the cluster signal given knowledge of the cluster profile, the SZE spectrum and the noise in the maps (Haehnelt & Tegmark 1996; Melin, Bartlett & Delabrouille 2006). The cluster gas profiles are assumed to be described by a projected isothermal β model (Cavaliere & Fusco-Femiano 1976) with $\beta = 1$. Note that, as discussed in Vanderlinde et al. (2010), the resulting SPT-SZ candidate catalogues are not sensitive to this assumption. The adopted model provides an SZE temperature decrement that is

¹ <http://www.darkenergysurvey.org>

maximum at the cluster centre and weakens with separation θ from the cluster centre as

$$\Delta T(\theta) = \Delta T_0 [1 + (\theta/\theta_c)^2]^{-1}, \quad (1)$$

where ΔT_0 is the central value and θ_c is the core radius. We adopt 12 different cluster profiles linearly spaced from $\theta_c = 0.25$ to 3 arcmin (Vanderlinde et al. 2010; Reichardt et al. 2013, B15). For each cluster, the maximum signal to noise across the 12 filtered maps is denoted as ξ . The SPT-SZ cluster candidates with $\xi > 4.5$ have been previously published in B15.

2.2 DES optical cluster catalogues

The DES Science Verification Data (DES-SVA1) that overlap SPT have been used to produce optically selected catalogues of clusters. In Section 2.2.1, we describe the acquisition and preparation of the DES-SVA1 data, and in Section 2.2.2, we describe the production of the redMaPPer cluster catalogue used in the primary analysis. We remind the reader that in Appendix A we present results of a preliminary analysis of the VT cluster catalogue.

2.2.1 DES-SVA1 data

The DES-SVA1 data include imaging of ~ 300 deg² over multiple disconnected fields (Sánchez et al. 2014; Banerji et al. 2015; Melchior et al. 2015), most of which overlap with the SPT-SZ survey. The DES-SVA1 data were acquired with the Dark Energy Camera (Diehl et al. 2012; Flaugher et al. 2012, 2015) over 78 nights, starting in Fall 2012 and ending early in 2013 with depth comparable to the nominal depth of the full DES survey (Rykoff et al., in preparation).

Data have been processed through the DES Data Management (DESDM; Desai et al. 2012) pipeline that is an advanced version of development versions described in several publications (Ngeow et al. 2006; Mohr et al. 2008, 2012). The data were calibrated in several stages leading to a *Gold* catalogue of DES-SVA1 galaxies (Rykoff et al., in preparation). The *Gold* catalogue covers ~ 250 deg² and is optimized for extragalactic science. In particular, it masks regions south of declination $\delta = -61^\circ$, avoiding the Large Magellanic Cloud and its high stellar densities. Furthermore, the footprint is restricted to the regions where we have coverage in all four bands.

2.2.2 redMaPPer cluster catalogue

The red-sequence Matched-Filter Probabilistic Percolation (redMaPPer; hereafter RM) algorithm is a cluster-finding algorithm based on the richness estimator of Rykoff et al. (2012). RM has been applied to photometric data from the Eighth Data Release (DR8) of the Sloan Digital Sky Survey (SDSS; Aihara et al. 2011) and to the SDSS Stripe 82 co-add data (Annis et al. 2014), and has been shown to provide excellent photometric redshifts, richness estimates that tightly correlate with external mass proxies, and very good completeness and purity (Rozo & Rykoff 2014; Rozo et al. 2014a, 2015b). We refer the reader to the paper by Rykoff et al. (2014) for a detailed description of the algorithm. Here, we briefly summarize the most salient features.

We employ an updated version of the algorithm (v6.3.3), with improvements summarized in Rozo & Rykoff (2014), Rozo et al. (2015a) and Rykoff et al. (in preparation). RM calibrates the colour of red-sequence galaxies using galaxy clusters with spectroscopic redshifts. RM uses this information to estimate the membership

probability of every galaxy in the vicinity of a galaxy cluster. The richness λ is thus defined as the sum of the membership probabilities (P_{RM}) over all galaxies (see section 4 of Rykoff et al. 2014 for more details):

$$\lambda = \sum P_{\text{RM}}. \quad (2)$$

In addition to the estimate of membership probabilities, the RM centring algorithm is also probabilistic. The centring probability P_{cen} is a likelihood-based estimate of the probability that the galaxy under consideration is a central galaxy. The centring likelihood includes the fact that the photometric redshift of the central galaxy must be consistent with the cluster redshift, that the central galaxy luminosity must be consistent with the expected luminosity of the central galaxy of a cluster of the observed richness, and that the galaxy density on a 300 kpc scale consistent with the galaxy density of central galaxies. The centring probability further accounts for the fact that every cluster has one and only one central galaxy, properly accounting for the relevant combinatoric factors. These probabilities have been tested on SDSS DR8 data using X-ray-selected galaxy clusters, and have been shown to produce cluster centres that are consistent with the X-ray centres (Rozo & Rykoff 2014).

The DES-SVA1 RM catalog was produced by running on a smaller footprint than that for the full SVA1 *Gold* sample. In particular, we restrict the catalogue to the regions where the z -band 10σ galaxy limiting magnitude is $z > 22$. In total, we use 148 deg² of DES-SVA1 imaging, with 129.1 deg² overlapping the SPT-SZ footprint. In this area, the largest fraction (124.6 deg²) is included in the so called DES-SVA1 SPT-E field. The final catalogue used in this work consists of 9281 clusters with $\lambda > 5$ and redshifts in the range $0.1 < z < 0.9$. Due to the varying depth of the DES-SVA1 catalog, RM produces a mask that determines the maximum redshift of the cluster search at any given location in the survey. As an example, the effective area in the SPT-E region at the highest redshift ($z > 0.85$) is only ~ 30 deg². In addition to the cluster catalogue, the RM algorithm also uses the survey mask to produce a set of random points with the same richness and redshift distribution as the clusters in the catalogue. The random points take into account the survey geometry and the physical extent of the clusters, and as with the clusters, only includes points that have < 20 per cent of the local region masked (see Rykoff et al. 2014).

2.3 Catalogue matching

We cross-match the SPT-SZ catalogue with the RM optical cluster catalogue following the method of Rozo et al. (2015b). First, we sort the SPT-SZ clusters to produce a list with decreasing SZE observable ξ , and we sort the RM catalogue to produce a list with decreasing richness. Secondly, we go down the SPT-SZ sorted list, associating each SPT cluster candidate with the richest RM cluster candidate whose centre lies within $1.5 R_{500}$ of the SZE centre. Thirdly, we remove the associated RM cluster from the list of possible counterparts when matching the remaining SPT-selected clusters.

R_{500} is first computed assuming the redshift of the optical counterpart and using the SZE-mass scaling-relation parameters adopted in B15. We subsequently check that our sample does not change when adopting our best-fitting scaling-relation parameters (see Section 3.1).

To test the robustness of our matching algorithm against chance associations, we first perform the above described procedure on a sample of randomly generated RM clusters as described in the previous section. Positions of clusters in this randomly generated sample

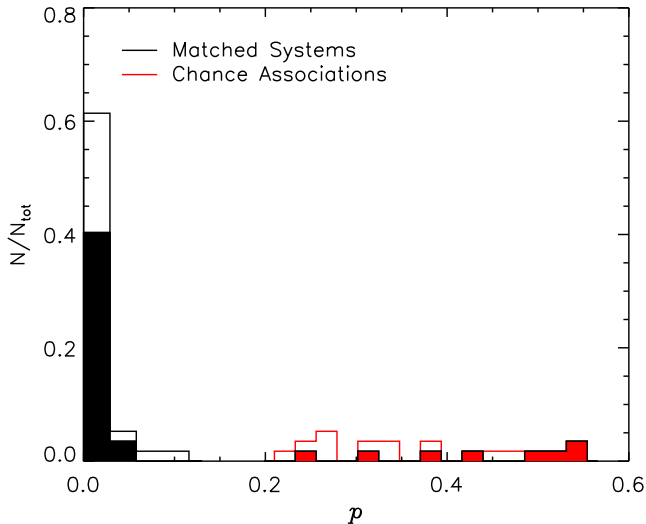


Figure 1. Distribution of probability p of chance associations for the $\xi > 4.5$ (filled) and $\xi > 4$ (open) samples. Black lines show the sample used in this work, red lines show the sample rejected (Section 2.3).

do not correlate with the positions of the SPT-SZ clusters. Using an ensemble of 10^4 random catalogues, we measure the distribution of richness in chance associations for each SPT-SZ candidate. We then apply the algorithm to match the real RM cluster catalogue with the SPT-SZ candidate list where $\xi > 4.5$. The distribution of probabilities p of chance associations estimated for each cluster candidate using the randomly generated samples is shown in Fig. 1 (filled histogram).

In Fig. 2 (left-hand panel) we show the resulting 84 and 95 per cent confidence limits (solid and dotted lines, respectively) in the richness distribution of the chance associations as a function of the SPT-SZ observable ξ . This test allows us to estimate the probability of chance superposition for each SPT-SZ cluster candidate. As detailed below, we use this information to determine whether or not to include particular matches for further analysis. As Fig. 2 shows, this filtering of the matched sample then ensures that chance superpositions are playing no more than a minor role even at $4 < \xi < 4.5$.

Within the DES-SVA1 region explored by RM there are 36 such SPT-SZ cluster candidates. Using information on the contamination fraction at $\xi > 4.5$ of the SPT-SZ candidate list and on the redshift distribution of the confirmed cluster candidates (Song et al. 2012, B15), we expect approximately nine of these candidates to be noise fluctuations and ~ 80 per cent of the real clusters to lie at $z < 0.8$. Therefore, assuming the optical catalogue is complete, the expected number of real cluster matches is ~ 22 . Similarly, one can estimate an expected number of real cluster matches (23.6) by scaling the total number of confirmed clusters in B15 below $z = 0.8$ (433) in the 2365 deg² of the SPT-SZ survey by the DES-SVA1 area overlapping SPT-SZ that has been processed with the RM cluster finder (129.1 deg²).

The actual number of matches to the SPT-SZ candidates is 33. Eight systems are foreground, low-richness RM clusters that have been erroneously associated with SPT-SZ candidates with either previously measured redshifts (four systems) or lower limits estimated in B15 (the remaining four candidates) that are at $z \gtrsim 0.8$; these systems are either noise fluctuations or real clusters that are at redshifts too high for them to be detected by RM. In fact, all of these systems have a probability p of chance associations estimated

from the randomly generated sample that is $p > 16$ per cent. Therefore, we remove these matches from the sample (red histograms in Fig. 1). This leaves 25 SPT-SZ candidates at $\xi > 4.5$ that have RM counterparts (black histograms in Fig. 1). We expect less than one false associations within this sample of 25 candidates. The associated optical richness as a function of the SPT-SZ significance is shown in the left-hand panel of Fig. 2. This number is somewhat larger but statistically consistent with the expected number of matches presented above. All 22 of the SPT-SZ confirmed clusters presented in B15 that lie at redshifts where they could be detected by RM are in this matched sample. According to our matching metric (which differs from the approach in B15), there are also three unconfirmed SPT-SZ candidates (i.e. candidates without identified optical counterparts in the B15 analysis) that have RM counterparts: SPT-CL J0502-6048, SPT-CL J0437-5307 and SPT-CL J0500-4551. The newly confirmed clusters are highlighted with large circles in Fig. 2.

Of the 25 SPT-SZ candidates with robust RM counterparts, we use 19 of them to calibrate the RM richness–mass relation. Six clusters are excluded from the analysis for the following reasons. Clusters with estimated redshift $z < 0.25$ in the SPT-SZ catalogue from B15 are highlighted in cyan in the left-hand panel of Fig. 2. Because the ξ –mass relation is robust only above this redshift (Vanderlinde et al. 2010), these systems are not used in the following analysis. Two systems (SPT-CL J0440-4744 and SPT-CL J0441-4502) are excluded from this analysis as they are detected in SPT-SZ regions that have been masked due to their proximity to point sources, which can compromise the SZE signal-to-noise measurement. In addition, we exclude the three clusters highlighted in magenta: SPT-CL J0417-4748, SPT-CL J0456-5116 and SPT-CL J0502-6048. These systems are strongly masked in the DES-SVA1 data; based on the SZE position, the masks cover 40 per cent of the total cluster region. As a result, the associated optical counterparts are highly miscentred, and the corresponding richness is severely biased. We note that the average centring failure rate caused by the detection mask is 12 per cent (3 clusters out of 25), in comparison to the corresponding rate in the SDSS RM catalogue, which is ≈ 1 –2 per cent. The difference reflects the fact that SDSS has a much larger contiguous area, while SVA1 has a more aggressive star mask. We expect this failure rate will decrease as the DES coverage increases, and object masking improves. Furthermore, improvements will be made to the RM algorithm to estimate the masked area not only at the putative centre of the cluster, but at all possible centres. In this way, clusters at high risk of mask-induced miscentring will be properly removed from the sample.

The B15 catalogue contains only SPT-SZ candidates with $\xi \geq 4.5$. In this work, we also apply the matching algorithm to SPT-SZ candidates at $4 < \xi < 4.5$. We identify 26 matches in this signal-to-noise range. The resulting probabilities of chance associations of the $\xi > 4$ sample are shown as open histograms in Fig. 1. Similarly to the $\xi > 4.5$ case, we exclude 11 of these systems, which have estimated probabilities p of chance associations $p > 16$ per cent. For the 15 matched systems, the expected number of false associations is also smaller than one. The remaining *cleaned* sample is shown as red points on the left-hand panel of Fig. 2. The resulting total number of SPT-SZ and RM associations at $\xi > 4$ is 40. This number is in good agreement with the expectation (~ 36) obtained using the number of SPT candidates above $\xi > 4$ in the DES-SVA1 region explored by RM (88) and correcting it by the expected number of noise fluctuations (~ 45) and the number of clusters above $z > 0.8$ (~ 7). We find that two $\xi < 4.5$ SPT-SZ candidates, SPT-CL J0501-4717 and SPT-CL J0439-5611, have probabilities of random associations

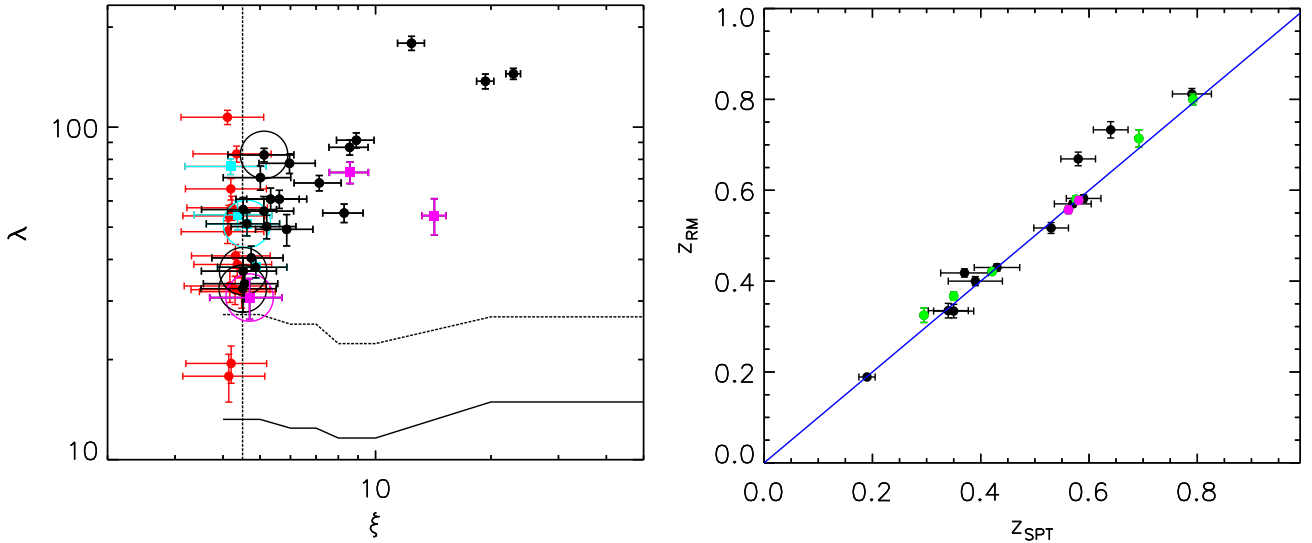


Figure 2. Left-hand panel: richness as a function of the SZE significance ξ for the matched cluster sample. SPT-SZ candidates with $\xi < 4.5$ (vertical line) are shown in red. Clusters at $z < 0.25$ (cyan) and clusters with miscentering due to a high masked fraction (magenta) are not used in the richness analysis. Large circles indicate the newly confirmed SPT-SZ candidates with $\xi > 4.5$. Solid and dashed lines represent the upper 84 and 95 percentiles in richness of chance associations of SPT-SZ candidates and clusters from the randomly generated RM catalogue. Right-hand panel: the estimated redshift for the RM sample as a function of SPT redshifts as presented in B15 from independent optical follow-up data. SPT-SZ candidates with spectroscopic redshift are shown in green. Magenta symbols are the same as in the left-hand panel.

larger than 5 percent, and therefore it is not clear whether these low-richness associations are correct (see Fig. 2).

The right-hand panel of Fig. 2 contains a comparison of the redshifts from the RM catalogues with the redshifts published in B15 (z_{SPT}) for the same clusters (obtained through dedicated optical/NIR followup by the SPT team or taken from the literature). Clusters with spectroscopic redshifts are highlighted in green. We note that the redshift estimates are not biased for the clusters affected by masking (magenta points). For SPT-SZ candidates with $4 < \xi < 4.5$ the SPT collaboration did not complete followup optical imaging, and therefore we adopt the redshifts of the RM optical cluster counterpart.

Table 1 contains all SPT candidates with RM counterparts used in this work. For newly confirmed SPT-SZ clusters, the associated z_{SPT} redshift is not given. We caution that the masses for low-redshift clusters ($z < 0.25$) may be underestimated due to filtering that is done to remove the noise component associated with the primary CMB.

3 MASS CALIBRATION METHOD

We apply the method described in Bocquet et al. (2015) to characterize the λ -mass relation of SPT-selected clusters. We refer the reader to the original paper (and to the appendix in particular) for a detailed description of the method. A similar approach has been adopted by Liu et al. (2015) for studying the SZE properties of an X-ray-selected cluster sample from the XMM-BCS survey (Desai et al. 2012; Šuhada et al. 2012). In this analysis, we consider the RM richness as a follow-up observable to the SZE-selected cluster sample. This choice is adequate as there are no SPT-SZ candidates with $\xi > 4.5$ missing RM counterparts in the redshift and spatial regime explored by the RM catalogue, so that the cross-sample can indeed be thought of as solely SPT selected. We note that this is not the case for SPT-SZ candidates with $4 < \xi \leq 4.5$ that do not have RM counterparts. However, the adopted method is also accurate under the assumption that cross-matching the SPT-SZ candidate list

with the RM cluster catalogue *cleans* the SPT-SZ candidate list, removing the expected noise fluctuations. Within this context the resulting cluster sample is therefore drawn from the halo mass function through the SPT-SZ selection in the redshift range explored by the RM catalogue.

In the following subsections, we describe the model we use to simultaneously constrain the SZE-mass relation (Section 3.1) and the richness-mass relation (Section 3.2).

3.1 The SZE-mass relation

Following previous SPT papers (Vanderlinde et al. 2010; Benson et al. 2013; Reichardt et al. 2013; Bocquet et al. 2015, B15), we define the unbiased SZE significance ζ as the average signal to noise a cluster would produce over many realizations of SPT data, if the cluster position and core radius were perfectly known. This quantity is related to the expectation value of ξ over many realizations of the SPT data by

$$\zeta = \sqrt{\langle \xi \rangle^2 - 3}, \tag{3}$$

where the bias in $\langle \xi \rangle$ is due to maximizing the signal to noise over three variables (cluster right ascension, declination and core radius). The scatter of the actual observable ξ with respect to $\langle \xi \rangle$ is characterized by a Gaussian of unit width. The SPT observable-mass relation $P(\zeta | M_{500}, z)$ is modelled as a lognormal distribution of mean

$$\begin{aligned} \langle \ln \zeta | M_{500}, z \rangle = & \ln A_{\text{SZE}} + B_{\text{SZE}} \ln \left(\frac{M_{500}}{3 \times 10^{14} h^{-1} M_{\odot}} \right) \\ & + C_{\text{SZE}} \ln \left(\frac{E(z)}{E(z = 0.6)} \right) \end{aligned} \tag{4}$$

and scatter D_{SZE} , and where $E(z) \equiv H(z)/H_0$. At low significance $\zeta \lesssim 2$, there is a non-negligible chance of multiple low-mass clusters overlapping within the same resolution element of the SPT beam. We account for this by only considering the brightest of these objects per approximate resolution element and we compute $P(\zeta_{\text{max}} | \zeta)$

Table 1. SPT-SZ cluster candidates with RM counterpart. We report the SPT-ID (1), right ascension (2) and declination (3), SPT peak detection significance ξ (4), corresponding core radius (5), richness λ (6), associated redshift from the RM catalogue (7) and SPT catalogue (8), and SPT derived masses (9). Coordinates are J2000.

SPT ID	RA	Dec.	ξ	θ_c (arcmin)	λ	z_{RM}	z_{SPT}	$M_{500}[10^{14}h_{70}^{-1}M_{\odot}]$
SPT-CL J0438-5419	69.574	-54.319	22.88	0.50	144.76 \pm 5.52	0.42 \pm 0.01	0.42	10.19 \pm 1.33
SPT-CL J0040-4407	10.199	-44.133	19.34	0.50	137.45 \pm 7.03	0.37 \pm 0.01	0.35	9.71 \pm 1.28
SPT-CL J0417-4748	64.344	-47.812	14.24	0.25	54.22 \pm 6.75	0.58 \pm 0.01	0.58	7.41 \pm 1.00
SPT-CL J0516-5430	79.149	-54.510	12.41	1.50	178.93 \pm 8.71	0.33 \pm 0.02	0.29	7.05 \pm 0.97
SPT-CL J0449-4901	72.273	-49.023	8.91	0.50	91.37 \pm 4.75	0.80 \pm 0.01	0.79	5.24 \pm 0.78
SPT-CL J0456-5116	74.115	-51.275	8.58	1.00	73.08 \pm 5.39	0.56 \pm 0.01	0.56	5.39 \pm 0.81
SPT-CL J0441-4855	70.450	-48.917	8.56	0.50	86.96 \pm 4.55	0.81 \pm 0.01	0.79 \pm 0.04	5.10 \pm 0.77
SPT-CL J0439-4600	69.807	-46.012	8.28	0.25	55.18 \pm 3.52	0.34 \pm 0.02	0.34 \pm 0.04	5.52 \pm 0.84
SPT-CL J0440-4657	70.229	-46.964	7.13	1.25	67.95 \pm 3.62	0.33 \pm 0.02	0.35 \pm 0.04	4.95 \pm 0.81
SPT-CL J0447-5055	71.843	-50.921	5.97	0.25	77.84 \pm 5.26	0.40 \pm 0.01	0.39 \pm 0.05	4.24 \pm 0.81
SPT-CL J0422-5140	65.591	-51.674	5.86	1.00	49.28 \pm 5.32	0.58 \pm 0.01	0.59 \pm 0.03	3.98 \pm 0.78
SPT-CL J0439-5330	69.928	-53.502	5.61	0.75	60.77 \pm 3.81	0.43 \pm 0.01	0.43 \pm 0.04	3.97 \pm 0.81
SPT-CL J0433-5630	68.249	-56.502	5.32	1.75	60.75 \pm 4.82	0.71 \pm 0.02	0.69	3.56 \pm 0.78
SPT-CL J0535-5956	83.791	-59.939	5.20	0.25	50.25 \pm 4.09	0.67 \pm 0.02	0.58 \pm 0.03	3.46 \pm 0.77
SPT-CL J0440-4744	70.242	-47.736	5.12	1.25	82.55 \pm 3.80	0.30 \pm 0.02	—	3.75 \pm 0.83
SPT-CL J0428-6049	67.026	-60.828	5.11	1.25	55.91 \pm 5.95	0.73 \pm 0.02	0.64 \pm 0.03	3.46 \pm 0.79
SPT-CL J0444-4352	71.162	-43.872	5.01	1.50	70.53 \pm 5.91	0.57 \pm 0.01	0.57 \pm 0.03	3.53 \pm 0.82
SPT-CL J0458-5741	74.598	-57.695	4.87	2.50	37.90 \pm 2.68	0.19 \pm 0.01	0.19 \pm 0.02	3.69 \pm 0.85
SPT-CL J0534-5937	83.606	-59.625	4.74	0.25	40.43 \pm 3.42	0.58 \pm 0.01	0.58	3.15 \pm 0.76
SPT-CL J0502-6048	75.724	-60.810	4.69	0.25	30.73 \pm 4.32	0.79 \pm 0.02	—	3.03 \pm 0.76
SPT-CL J0441-4502	70.345	-45.040	4.62	2.50	51.22 \pm 4.11	0.15 \pm 0.01	—	3.49 \pm 0.85
SPT-CL J0429-5233	67.430	-52.559	4.56	0.75	33.84 \pm 3.97	0.52 \pm 0.01	0.53 \pm 0.03	3.15 \pm 0.79
SPT-CL J0452-4806	73.002	-48.108	4.52	0.50	56.54 \pm 4.89	0.42 \pm 0.01	0.37 \pm 0.04	3.26 \pm 0.81
SPT-CL J0437-5307	69.259	-53.119	4.51	0.25	36.89 \pm 3.56	0.29 \pm 0.02	—	3.20 \pm 0.80
SPT-CL J0500-4551	75.209	-45.856	4.51	0.75	32.68 \pm 4.82	0.26 \pm 0.01	—	3.66 \pm 0.91
SPT-CL J0453-5027	73.307	-50.451	4.47	0.25	31.99 \pm 3.44	0.77 \pm 0.02	—	2.89 \pm 0.74
SPT-CL J0449-4440	72.473	-44.672	4.37	0.75	54.50 \pm 5.43	0.15 \pm 0.00	—	3.42 \pm 0.86
SPT-CL J0423-5506 ²	65.809	-55.104	4.36	1.25	38.65 \pm 3.44	0.27 \pm 0.02	—	3.26 \pm 0.83
SPT-CL J0451-5057	72.937	-50.965	4.34	0.50	83.11 \pm 4.62	0.76 \pm 0.01	—	2.81 \pm 0.74
SPT-CL J0438-4629	69.564	-46.488	4.31	0.50	41.01 \pm 3.37	0.43 \pm 0.01	—	3.07 \pm 0.79
SPT-CL J0456-4531	74.099	-45.523	4.30	0.25	32.45 \pm 3.18	0.29 \pm 0.02	—	3.17 \pm 0.81
SPT-CL J0431-5353	67.970	-53.896	4.22	0.50	57.23 \pm 4.76	0.75 \pm 0.02	—	2.74 \pm 0.73
SPT-CL J0501-4717	75.274	-47.294	4.20	3.00	19.47 \pm 2.50	0.35 \pm 0.02	—	3.34 \pm 0.88
SPT-CL J0518-5740	79.507	-57.670	4.19	0.25	65.22 \pm 4.82	0.82 \pm 0.01	—	2.60 \pm 0.70
SPT-CL J0438-4907	69.655	-49.117	4.19	1.75	76.20 \pm 4.16	0.24 \pm 0.01	—	3.13 \pm 0.81
SPT-CL J0513-5901	78.273	-59.029	4.17	0.25	33.29 \pm 3.58	0.61 \pm 0.01	—	2.75 \pm 0.73
SPT-CL J0451-4910	72.888	-49.178	4.14	0.25	54.10 \pm 4.10	0.73 \pm 0.02	—	2.71 \pm 0.73
SPT-CL J0439-5611	69.978	-56.192	4.14	0.50	17.82 \pm 2.93	0.28 \pm 0.02	—	3.10 \pm 0.81
SPT-CL J0532-5752	83.237	-57.877	4.11	0.50	48.49 \pm 3.81	0.77 \pm 0.02	—	2.59 \pm 0.71
SPT-CL J0449-5908	72.472	-59.142	4.11	1.25	107.14 \pm 5.29	0.77 \pm 0.01	—	2.68 \pm 0.73

following Crawford et al. (2010). The SPT observable–mass relation is therefore expanded to $P(\zeta_{\text{max}}|\xi)P(\zeta|M_{500}, z)$ and ζ_{max} is then converted to the observable ξ as in equation (3).

To calibrate the ξ – M relation we use the subsample of clusters with $\xi > 5$ and $z > 0.25$ from the 2500d SPT-SZ catalogue (B15). We determine the parameter values by abundance-matching the catalogue against our fixed reference cosmology. We predict the expected number of clusters as a function of mass and redshift using the halo mass function (Tinker et al. 2008). We convolve this mass function with the observable–mass relation accounting for its associated uncertainties, and compare the prediction with the data. Our approach here is effectively the opposite of the typical analysis, where cosmological parameters are deduced from the cluster sample using both priors and calibrating information to constrain the scaling-relation parameters (e.g. Benson et al. 2013; Bocquet et al. 2015); here, we assume perfect knowledge of cosmology to calibrate the scaling relation. Note that this method does not depend on any assumptions about hydrostatic equilibrium.

We assume flat priors on A_{SZE} , B_{SZE} , C_{SZE} and a Gaussian prior on $D_{\text{SZE}} = 0.18 \pm 0.07$; the latter corresponds to the posterior distribution derived from the cosmological analysis of the full SPT sample (de Haan et al., in preparation). We obtain the following parameters for the ξ –mass relation by maximizing the likelihood of obtaining the observed sample in ξ and redshift under the model derived from equations (3) and (4). The results are

$$\begin{aligned} A_{\text{SZE}} &= 4.02 \pm 0.16, & B_{\text{SZE}} &= 1.71 \pm 0.09, \\ C_{\text{SZE}} &= 0.49 \pm 0.16, & D_{\text{SZE}} &= 0.20 \pm 0.07. \end{aligned} \quad (5)$$

For every cluster in the sample we also calculate the associated mass distribution, accounting for selection effects:

$$P(M_{500}|\xi, z, \mathbf{p}) \propto P(\xi|M_{500}, z, \mathbf{p}) P(M_{500}|z, \mathbf{p}), \quad (6)$$

where the vector \mathbf{p} encapsulates cosmological and scaling-relation parameters and $P(\xi|M_{500}, z, \mathbf{p})$ is obtained from the ξ –mass scaling relation as described by equations (3) and (4). The halo mass

function $P(M_{500}|z, \mathbf{p})$ is the prior on the mass distribution at redshift z .

Masses derived for the matched cluster sample are shown in Table 1. We note that both these masses and the SZE scaling-relation parameters quoted here are different from the ones reported in B15. We adopt the same fixed cosmology as in B15, but in this analysis we consider data from the full SPT-SZ cluster-survey as opposed to just the sample from the initial 720 deg² (Reichardt et al. 2013).

3.2 Richness–mass relation

As for the SZE–mass relation (equation 4), we assume a power-law form for the λ –mass relation:

$$\langle \ln \lambda | M_{500}, z \rangle = \ln A_\lambda + B_\lambda \ln \left(\frac{M_{500}}{3 \times 10^{14} h^{-1} M_\odot} \right) + C_\lambda \ln \left(\frac{E(z)}{E(z=0.6)} \right), \quad (7)$$

where A_λ is the normalization, B_λ characterizes the mass dependence and C_λ characterizes the redshift evolution. An additional parameter D_λ describes the intrinsic scatter in λ , which is assumed to be lognormal and uncorrelated with the SZE scatter, with variance given by

$$\text{Var}(\ln \lambda | M_{500}) = \exp(-\langle \ln \lambda | M_{500} \rangle) + D_\lambda^2. \quad (8)$$

The first term above represents the Poisson noise associated with the number of galaxies in a halo at fixed mass, and therefore we define the intrinsic scatter D_λ as lognormal scatter in addition to Poisson noise. We assume flat priors on the distributions of A_λ , B_λ , C_λ and a positive flat prior for D_λ .

The probability that a cluster with SPT-SZ signal-to-noise ξ is observed to have a richness λ is

$$P(\lambda | \xi, z, \mathbf{p}) = \int dM_{500} P(\lambda | M_{500}, z, \mathbf{p}) P(M_{500} | \xi, z, \mathbf{p}). \quad (9)$$

The term $P(\lambda | M_{500}, z, \mathbf{p})$ contains the lognormal intrinsic scatter and normal measurement uncertainties in the observable λ . We use the above distribution to evaluate the likelihood of the matched cluster sample defined through our cross-matching procedure. Note that we simultaneously vary both the optical and SZE scaling-relation parameters, further including the SZE data set from B15 with $\xi > 5$, $z > 0.25$ for constraining the SZE–mass relation.

4 RESULTS

We present here the constraints on the richness–mass relation (Section 4.1) and then use these best-fitting parameters to explore whether the cumulative distribution of the matched samples are consistent with the expectations from the model (Section 4.2). Finally, in Section 4.3, we analysis the optical-SZE positional offset distribution.

4.1 redMaPPer richness–mass relation

We marginalize over the SZE-mass scaling-relation parameters and constrain the posterior distributions for the RM λ –mass scaling

² SPT-CL J0423-5506 was previously identified in Song et al. (2012) and Reichardt et al. (2013) at a redshift $z = 0.21 \pm 0.04$ with a signal to noise $\xi = 4.51$. The associated redshift estimated with the same analysis presented in B15 is $z = 0.25 \pm 0.036$.

Table 2. Best-fitting parameters and 68 per cent confidence level of the richness–mass scaling-relation parameters described by equations (7) and (8).

Catalogue	A_λ	B_λ	C_λ	D_λ
SPT-SZ+RM $\xi > 4.5$	$66.1^{+6.3}_{-5.9}$	$1.14^{+0.21}_{-0.18}$	$0.73^{+0.77}_{-0.75}$	$0.15^{+0.10}_{-0.07}$
SPT-SZ+RM $\xi > 4$	$69.8^{+6.0}_{-4.9}$	$1.17^{+0.19}_{-0.17}$	$1.71^{+0.63}_{-0.57}$	$0.20^{+0.09}_{-0.08}$

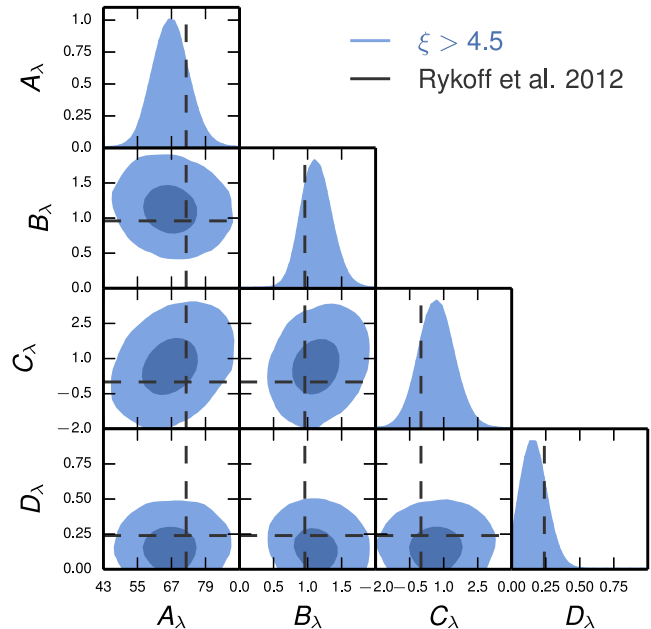


Figure 3. Posterior distribution for the four parameters of the λ –mass scaling relation (equations 7 and 8). Predictions from the scaling relation of Rykoff et al. (2012) are shown as dashed black lines. Best-fitting parameters and associated 1σ uncertainties appear in Table 2.

relation. Our best-fitting parameters and 68 per cent confidence level intervals are reported in Table 2 and shown in Fig. 3. We note that the slope of the λ –mass relation is consistent with 1 within 1σ (consistent, therefore, with the richness being proportional to the mass), and the model is consistent with no redshift evolution within 1σ (Andreon & Congdon 2014). Furthermore the resulting λ –mass relation is characterized by a remarkably low asymptotic intrinsic scatter, with $\sigma_{\ln \lambda | M_{500}} \rightarrow 0.15^{+0.10}_{-0.07}$ as $\langle \ln \lambda | M_{500} \rangle \rightarrow \infty$ (equation 8). Following Evrard et al. (2014), we estimate the characteristic scatter in mass at fixed richness to arrive at $\sigma_{\ln M} = 0.18^{+0.08}_{-0.05}$ at $\lambda = 70$, ~ 25 per cent larger than the corresponding characteristic scatter in mass at fixed ξ . We present in Fig. 4 the RM richness as a function of the SPT derived masses. Blue lines describe the best-fitting model and intrinsic scatter as derived from this analysis (Table 2) at a pivot point of $z = 0.6$.

We have verified that our results are not dominated by uncertainties in the SZE–mass scaling relation by fixing these parameters to their best-fitting values. Our results are only marginally improved in this case. Consequently, future analyses with larger samples are expected to considerably reduce the uncertainties of the recovered λ –mass scaling-relation parameters.

In Fig. 4, one cluster appears to be an obvious outlier: SPT-CL J0516-5435 ($\xi = 12.4$, $\lambda = 178.9$, $M_{500} = 7.05 \times 10^{14} h_{70}^{-1} M_\odot$). SPT-CL J0516-5345 is a well-known merger that is elongated in a north–south direction in the plane of the sky with an X-ray mass estimate nearly a factor of two larger than the SZE mass estimate.

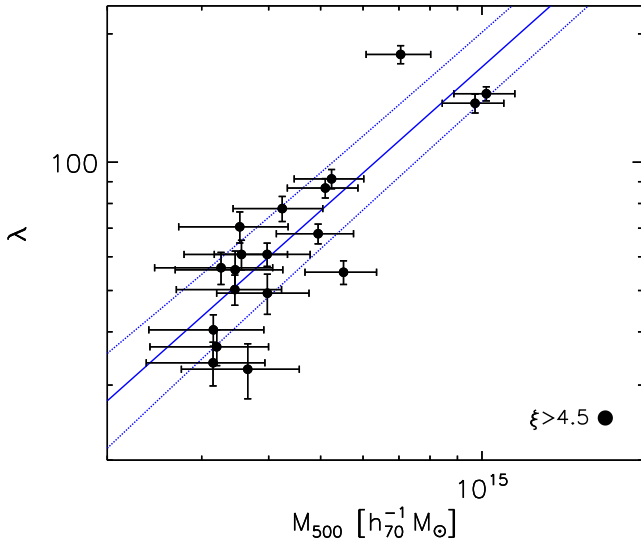


Figure 4. Richness as a function of the SPT derived masses for the calibration sample used in this analysis (Section 2.3). Blue lines show the best-fitting richness–mass relation and 1σ intrinsic scatter.

This cluster was in fact the strongest outlier in the sample of 14 clusters in Andersson et al. (2011). High et al. (2012) made a weak lensing measurement of SPT-CL J0516-5345, and found that there was a significant offset between the brightest central galaxy (BCG) and the weak-lensing centre, consistent with the merger hypothesis. Additionally, High et al. (2012) found that the weak-lensing mass was in better agreement with the SZE-mass estimate than the X-ray mass estimate, at a level consistent with elongation observed in the plane of the sky. Therefore, SPT-CL J0516-5345 appears to be an outlier due to true intrinsic scatter in the observable–mass relations, so we leave it in our analysis. We note, however, that whether or not we include SPT-CL J0516-5345 in the fit has a significant impact on our results. Our best-fitting parameters shift from $B_\lambda = 1.14^{+0.21}_{-0.18}$ and $D_\lambda = 0.15^{+0.10}_{-0.07}$ with this cluster, to $B_\lambda = 1.00^{+0.17}_{-0.15}$ and $D_\lambda = 0.05^{+0.07}_{-0.03}$ when SPT-CL J0516-5345 is not included in the fit. Whether SPT-CL J0516-5345 represents a rare event in a non-Gaussian tail in the distribution of richness of galaxy clusters or the recovered lognormal scatter obtained when cluster SPT-CL J0516-5349 is included is more correct will thus need to await future analyses with larger samples.

We convert the $P(M_{500}|\lambda)$ scaling relation derived by Rykoff et al. (2012) using abundance matching and the SDSS maxBCG cluster-catalogue (Koester et al. 2007) to a richness–mass relation so that we can compare it to our results (see also Evrard et al. 2014). The predictions from Rykoff et al. (2012) are shown as dashed lines in Fig. 3 under the assumption of no redshift evolution in the richness–mass relation. We note that all parameters of our derived RM λ –mass scaling relation for SPT-selected clusters are consistent with the Rykoff et al. (2012) values.

We repeat these analyses extending the sample to include those with $4 < \xi < 4.5$ and find similar results (Table 2). The largest difference is in the redshift evolution term which now has a best-fitting value $C_\lambda = 1.71^{+0.63}_{-0.57}$. While formally this difference does not have large statistical significance (1.3σ), it is coming from a sample that includes a large fraction of the same clusters, so it is likely statistically significant. A larger redshift evolution term would imply that higher redshift RM clusters are less massive at fixed λ . At the same time the derived scatter is also larger.

Two of the clusters in the $4 < \xi < 4.5$ range are compatible with false associations. Excluding the two matched clusters with the highest probability of random associations results in an $\sim 1\sigma$ shift in B_λ (from $B_\lambda = 1.17$ to $B_\lambda = 1.04$) and in an $\sim 0.5\sigma$ shift in C_λ (from $C_\lambda = 1.71$ to $C_\lambda = 1.42$), while the other parameters (A_λ and D_λ) are almost unchanged. Interestingly, even though we have increased the number of clusters in the SPT-SZ+RM sample by 40 per cent, the constraints on the scaling-relation parameters are only mildly tighter. There are two reasons for this. First these lower signal-to-noise SPT-SZ clusters have larger fractional mass uncertainties in ξ ($\langle \xi \rangle^{-1} \sim 0.23$ and $\langle \xi \rangle^{-1} \sim 0.13$, respectively for the $4 < \xi < 4.5$ and $\xi > 4.5$ samples). Secondly, the richnesses are also systematically lower, leading to a larger Poisson variance. Thus, each low- ξ cluster has less constraining power than a high- ξ cluster, reducing the impact of extending the sample to include the lower mass systems.

4.2 Consistency test of model

We also test the consistency of the adopted scaling relation model with the data by examining whether we are finding the expected number of matches with the correct distribution in richness. To do this, we focus on the SPT-E field, which at ~ 124.6 deg² is the largest contiguous region covered by the DES-SVA1 data. We carry out two different tests.

In the first test, we examine whether we are finding the expected number of SZE-selected clusters and whether these clusters have the expected number of optical matches with the correct λ distribution. We generate 10^6 Monte Carlo realizations of cluster samples extracted from the halo mass function above $M_{500} > 10^{13.5} h_{70}^{-1} M_\odot$ and assign richnesses λ and SPT-SZ significance ξ using the parameters we extract from our analysis of the real matched catalogue. These Monte Carlo mocks are generated taking into account the survey area as a function of redshift sampled by RM in the SPT-E field. We then apply the SZE selection – either $\xi > 4.5$ or $\xi > 4$ – and measure the cumulative distribution in λ of the SZE-selected samples. We then compare this to the same distribution in the real matched catalogue.

Shaded regions in the upper panels of Fig. 5 show 1σ , 2σ and 3σ confidence regions obtained from the mocks after marginalizing over the scaling-relation parameters. The solid black line shows the distribution from the real catalogue, which is in good agreement with the mocks. The largest observed difference is smaller than 1σ indicating that the adopted model provides a consistent description of the observed number and richness distribution of the SZE-selected sample. Similarly, a Kolmogorov–Smirnov (KS) test for the observed cumulative distribution of matched systems as a function of λ and the corresponding median distribution from the Monte Carlo simulations, shows that the null hypothesis of data being drawn from same distributions cannot be excluded and returns p –values 0.76 and 0.96 for $\xi > 4.5$ and $\xi > 4$, respectively.

The second test is focused on whether the RM cluster catalogue (which is significantly larger than the SPT-SZ catalogue) has the expected number of SZE matches with the correct λ distribution. Essentially, we take the observed RM catalogue as a starting point, and calculate the expected number of systems with SPT-SZ counterparts given the model and parameter constraints from our λ –mass likelihood analysis. This test differs from the first in that it takes the *observed* RM-selected cluster sample as a starting point, and such a test should be more sensitive to, for example, contamination in the RM catalogue. The formally correct way to evaluate a statistical difference between the RM-selected sample and the

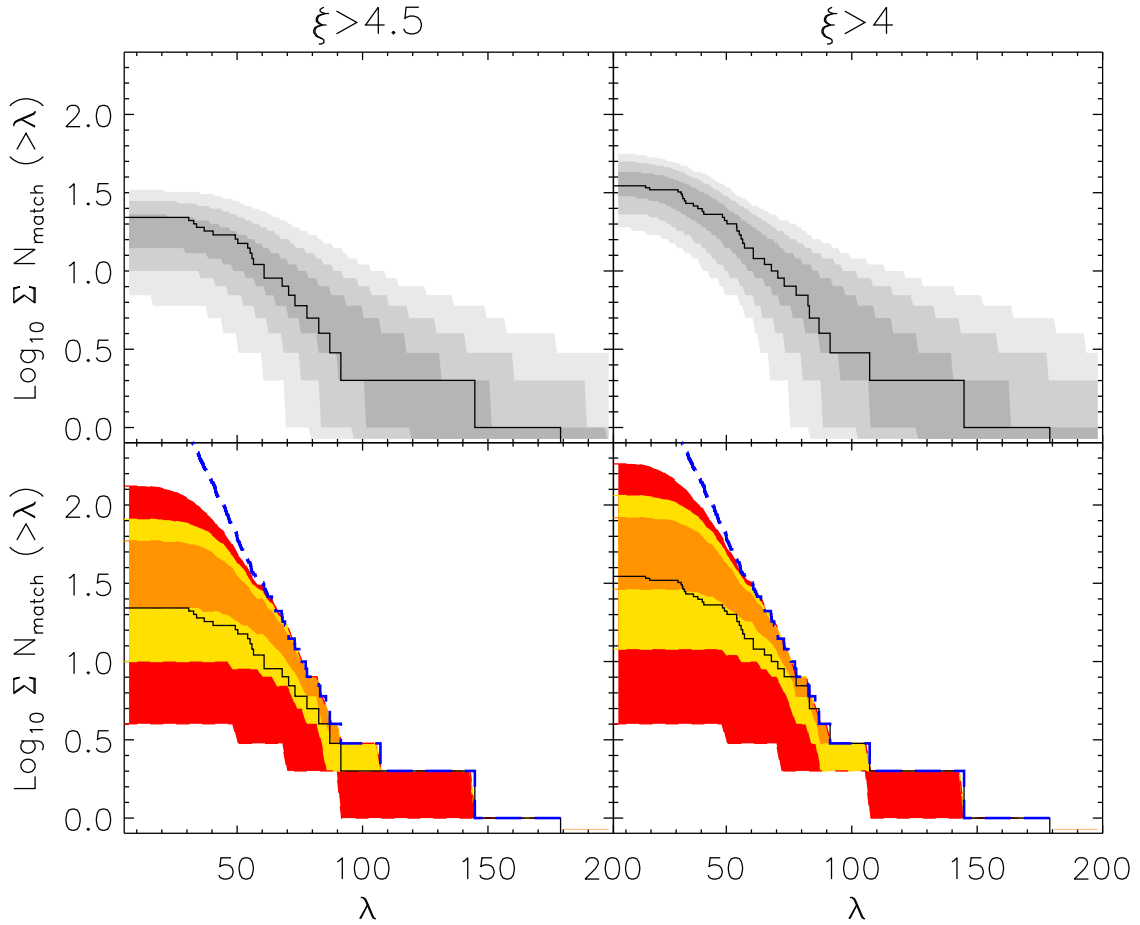


Figure 5. Consistency tests of our model (Section 3) and the SPT-SZ+RM catalogue from the $\sim 124 \text{ deg}^2$ SPT-E field. The solid black lines show the observed cumulative distribution in richness λ of the $\xi > 4.5$ (left-hand) and $\xi > 4$ (right-hand) SZE-selected samples. Upper panels: grey-scale regions show 1, 2 and 3σ regions predicted by drawing 10^6 SZE-selected samples from the mass function and assigning λ according to our scaling relation constraints. There is good agreement with the data. Lower panels: dashed blue lines show the cumulative distribution in λ of the full RM sample. Orange, yellow and red areas define 1, 2 and 3σ regions representing the predicted cumulative distribution of the SPT-SZ+RM catalogue using as input the full RM sample and the probability (equation 10) that each RM cluster will have an SPT-SZ counterpart, given our scaling relation constraints.

SPT-SZ+RM matched sample used in this analysis would be to calibrate the λ -mass and SZE-mass relations starting from the RM-selected sample. Such a study goes, however, beyond the scope of the current work and will be addressed in a future project. Therefore, the following analysis is only intended to be a consistency check.

For this purpose, we proceed by first computing, for each real RM-selected cluster in the SPT-E field, the probability P_m of that cluster also having $\xi > 4.5$ and therefore being in the matched sample. We define this probability as

$$\begin{aligned}
 P_m &= \int_{4.5}^{\infty} P(\xi|\lambda, z) d\xi \\
 &= \int_{4.5}^{\infty} d\xi \int dM_{500} P(\xi|M_{500})P(M_{500}|\lambda, z), \quad (10)
 \end{aligned}$$

where $P(M_{500}|\lambda, z) \propto P(\lambda|M_{500}, z)P(M_{500}, z)$ and $P(M_{500}, z)$ is the halo mass function (Tinker et al. 2008).

To predict the expected number of RM clusters with SZE counterparts, we then randomly sample the scaling-relation parameters, determining P_m for all the RM clusters in each case. We then Monte Carlo sample each P_m to produce randomly sampled matched cluster catalogues. We use the results from the ensemble

of random matched catalogues to produce the expected cumulative distributions in λ . Orange, yellow and red regions in the bottom panels of in Fig. 5 show the 1σ , 2σ and 3σ confidence regions, respectively, of the expected cumulative distribution in λ for $\xi > 4.5$ (left-hand) and $\xi > 4$ (right-hand), given the RM catalogue and scaling-relation parameter constraints as input. The dashed blue line shows the cumulative distribution of the entire sample of RM clusters in the area, while the solid black line shows the observed cumulative distribution of the real matched catalogue.

We note that the predicted number of SPT-SZ+RM matches in this case tends to be higher than that observed for $\lambda > 35$, but the tension is weak. Of some concern is the high λ end of the sample ($\lambda > 70$), where only 9 of the 17 RM selected clusters have SPT-SZ counterparts at $\xi > 4.5$ despite their having large probabilities indicating they should be in the SZE-selected sample. However, a KS test for the observed cumulative distribution of matched systems as a function of λ and the corresponding median distribution from the mocks shows that the null hypothesis of data being drawn from the same parent distribution cannot be excluded and returns p -values 0.90 and 0.33, respectively, for $\xi > 4.5$ and $\xi > 4$.

While the KS test is showing no evidence for the two distributions to differ, it is also known to be not very sensitive to the tails of the

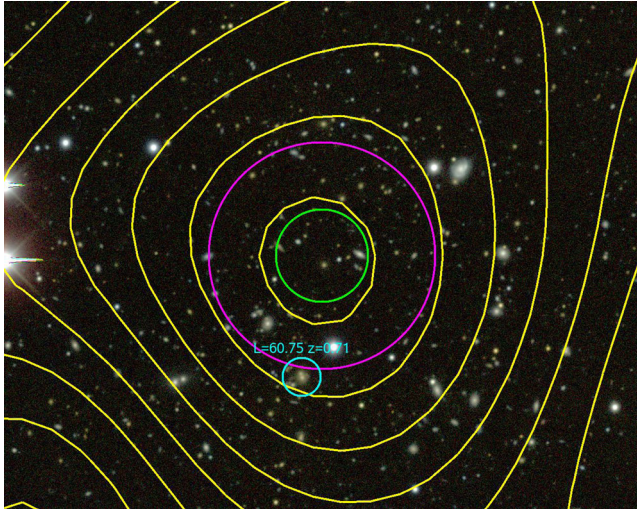


Figure 6. SPT-CL J0433-5630: DES-SVA1 *gri* pseudo colour image overlaid with SPT-SZ signal-to-noise contours (in steps $\Delta\xi = 1$). The magenta circle shows the projected $R_{500}/2$ radius at $z = 0.69$, while the green circle describes the 1σ SPT positional uncertainty (equation 11). The cyan label marks the associated $\lambda \sim 60$ RM cluster.

distributions (Moscovich-Eiger, Nadler & Spiegelman 2013). With this respect, it is also therefore interesting to directly compare the observed cumulative distribution with the cumulative distribution from the predicted number of matches. In this case, as for the classical KS test, we focus on the largest difference between the two distributions. Lower-left panel of Fig. 5 shows that at most, the observed cumulative distribution is in tension with expectations at the $\sim 2\sigma$ level, providing some indication of tension between the observed sample and our model. For example, when restricting ourselves to clusters with $\lambda > 70$, we find that our simulation results in a larger number of SPT+RM associations in 94.6 per cent of our Monte Carlo realizations. This tendency to observe fewer matches than expected given the size of the RM-selected sample could be explained by either contamination within the RM sample, additional incompleteness within the SPT-SZ sample beyond that caused by scatter in the SZE-mass relation, or simply a statistical fluctuation. Future work exploring the SZE properties of the lower mass systems along with an extension of the current analysis to the full overlap between DES and the SPT-SZ survey will sharpen this test.

4.3 Optical-SZE positional offset distribution

It has been shown that optical-miscentering can have a significant impact on the derived SZE signature for an optically selected sample (Biesiadzinski et al. 2012; Sehgal et al. 2013; Rozo & Rykoff 2014; Rozo et al. 2014a,b). We note, however, that in our case the SZE signal has been estimated at the SZE-determined position, so our results are not affected by optical miscentering. In fact, we can now use our data to constrain the distribution of offsets between the SZE-determined and the optically determined cluster centres.

As an example, we show in Fig. 6 the DES-SVA1 *gri* pseudo colour image of SPT-CL J0433-5630, an SPT-SZ selected cluster with $\xi \sim 5.3$ at redshift $z = 0.69$ (B15). Yellow contours show the SPT-SZ signal to noise in steps of $\Delta\xi = 1$, while the magenta circle describes the projected radius $R_{500}/2$. The cyan label refers to the associated RM cluster centre. This RM cluster has richness $\lambda \sim 60$. We note that the most probable central galaxy selected by RM is significantly offset from the SZE defined centre. As a result, the

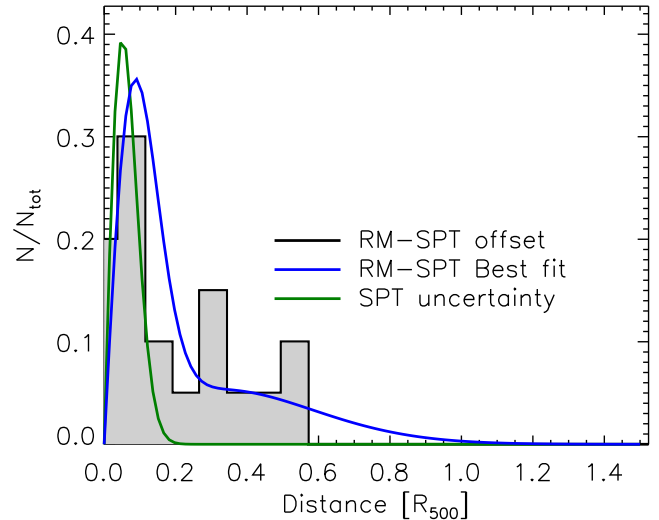


Figure 7. Filled histogram shows the measured fraction of SPT-SZ+RM clusters as a function of the optical-SZE positional offset in units of R_{500} . The green curve shows the SPT-SZ positional uncertainty, and the blue curves shows the best-fitting SZE-optical positional offset model.

measured SZE signature at the optical position ($\xi = 4.1$) would be significantly underestimated with respect to the derived unbiased quantity $\zeta = 5$ obtained through equation (3). We stress that this effect is not important for the scaling relation results reported in Section 4.1, as the sample analysed here is SZE selected.

Fig. 7 contains a normalized histogram of the distribution of cluster positional offsets in units of R_{500} for the $\xi > 4.5$ analysed SPT-SZ sample. Under the assumption that the measurement uncertainty from the optical side is negligible, we model this distribution as an underlying intrinsic positional offset distribution convolved with the SPT-SZ positional uncertainty.

The 1σ SPT-SZ positional uncertainty for a cluster with a pressure profile given by a spherical β model with $\beta = 1$ and core radius θ_c , detected with SPT-SZ significance ξ is described by

$$\Delta\theta = \xi^{-1} \sqrt{\theta_{\text{beam}}^2 + \theta_c^2}, \quad (11)$$

where $\theta_{\text{beam}} = 1.2$ arcmin is the beam FWHM (see Story et al. 2011 and equation 4 in Song et al. 2012 for more details). As a result, the expected distribution of positional offsets in the case in which the intrinsic one is a δ -function is shown (arbitrarily rescaled) as a green line.

Song et al. (2012) have shown that the intrinsic optical-SZE positional offset distribution for an SPT-SZ selected sample is consistent with the optical-X-ray positional offset distribution of X-ray selected clusters (Lin, Mohr & Stanford 2004). The offset distribution can be characterized by a large population of central galaxies with small offsets from the SZE centres and a less populated tail of central galaxies with large offsets (e.g. Lin et al. 2004; Rozo & Rykoff 2014; Lauer et al. 2014). We therefore parametrize the distribution of positional offsets between the RM centre and the SZE centre for x as

$$P(x) = 2\pi x \left(\frac{\rho_0}{2\pi\sigma_0^2} e^{-\frac{x^2}{2\sigma_0^2}} + \frac{1-\rho_0}{2\pi\sigma_1^2} e^{-\frac{x^2}{2\sigma_1^2}} \right) \quad (12)$$

where $x = r/R_{500}$, ρ_0 is the fraction of the population with small offsets and is characterized by variance σ_0^2 , while the population with larger offsets is characterized by variance σ_1^2 . While this model for the distribution was motivated by the expected intrinsic positional

Table 3. Best-fitting parameters and 68 per cent confidence level of the optical-SZE positional offset distribution.

Catalogue	ρ_0	$\sigma_0[R_{500}]$	$\sigma_1[R_{500}]$
RM- $\xi > 4.5$	$0.63^{+0.15}_{-0.25}$	$0.07^{+0.03}_{-0.02}$	$0.25^{+0.07}_{-0.06}$

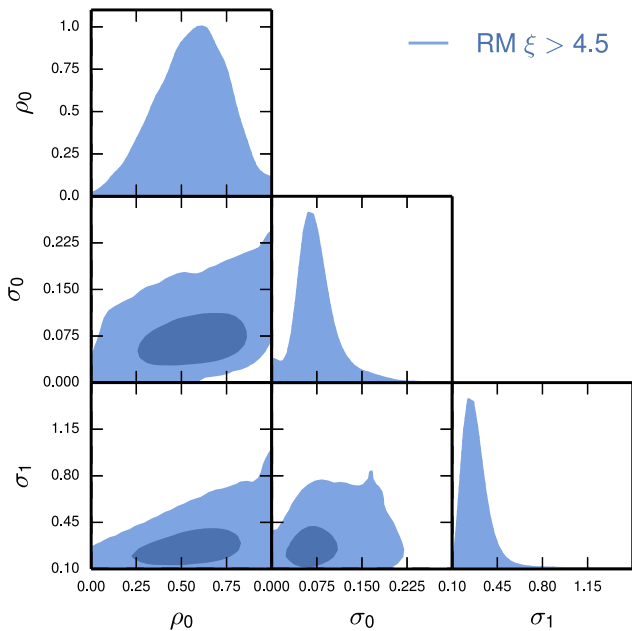


Figure 8. Posterior distribution for the 1σ and 2σ level of the three parameter model describing the positional offset distribution of equation (12). Best-fitting parameters are shown in Table 3.

offset distribution, the measured distribution will include both the actual physical SZE-central galaxy offset distribution and the systematics due to failures in identifying the correct cluster centre with the RM algorithm. For every cluster and parameter $\rho_0 \in [0, 1]$, $\sigma_0 \in [0, 1]$, and $\sigma_1 \in [\sigma_0, 1]$, we then compare the predicted offset distribution obtained by convolving the model with the SPT-SZ positional uncertainty of equation (11) to extract the associated likelihood. Best-fitting parameters and 68 per cent confidence intervals are shown in Table 3 and joint and fully marginalized parameter constraints are shown in Fig. 8.

We note that the positional offset distribution for the RM sample is consistent with a concentrated dominant population ($\rho_0 = 0.63^{+0.15}_{-0.25}$) of smaller offsets ($\sigma_0 = 0.07^{+0.03}_{-0.02} R_{500}$) and a subdominant population characterized by larger offsets ($\sigma_1 = 0.25^{+0.07}_{-0.06} R_{500}$). In the limit where the SZ centre and BCG are coincident for every cluster, the RM code (Rykoff et al. 2012, 2014) predicts the SZ-RM offset distribution for this sample to contain a centrally peaked component with normalization $\langle P_{\text{cen}} \rangle$ to be 0.79. While this is formally consistent with the measured SZ-RM offset distribution, intrinsic optical-SZE positional offsets, such as those presented by (e.g. Song et al. 2012) likely contribute to large separation component of the observed distribution shown in Fig. 7. Larger samples will be necessary to disentangle the impact of systematics due to failures in identifying the correct cluster centre with the RM algorithm from the offset distribution due to cluster morphology.

5 CONCLUSIONS

In this paper, we cross-match SZE-selected cluster candidates with $\xi > 4.5$ from the 2500d SPT-SZ survey (B15) with the optically selected cluster catalogue extracted from the DES science verification data DES-SVA1. The optically selected catalogue is created using the RM cluster-finding algorithm. We study the robustness of our matching algorithm by applying it to randomly generated RM catalogues.

Using the adopted matching algorithm in the 129.1 deg^2 of overlap between the two data sets, we create a matched catalogue of 33 clusters. Eight of these clusters are removed as likely chance superpositions that are identified using the randomly generated catalogues. The resulting 25 cluster sample includes all previously known $z < 0.8$ and $\xi > 4.5$ SPT-SZ clusters in this area (Song et al. 2012; Reichardt et al. 2013, B15) in addition to three previously unconfirmed SPT-SZ clusters.

We then study three characteristics of this cross-matched SPT-SZ+RM cluster sample.

(1) *The richness–mass relation of SPT-SZ selected clusters.* We calibrate the λ -mass relation from SZE measurements by applying the method described in Bocquet et al. (2015). In this analysis we assume a fixed fiducial cosmology and marginalize over the simultaneously calibrated SPT-SZ ξ -mass relation. We adopt flat priors on the richness–mass relation parameters. We find that the RM λ -mass relation for SPT-SZ selected clusters is characterized by a small asymptotic intrinsic scatter $D_\lambda = 0.15^{+0.10}_{-0.07}$ and by a slope $B_\lambda = 1.14^{+0.21}_{-0.18}$ that is consistent with unity. Our constraints are in good agreement with those of Rykoff et al. (2012) and show that the scatter in mass at fixed richness $\lambda = 70$ for this sample is only 25 per cent larger than the scatter in mass at fixed SPT-SZ observable ξ .

(2) *Consistency test of model and matched catalogue.* We carry out two consistency tests to determine whether there is tension between the observed matched sample and the expectations given the scaling relation model we have adopted. Both tests involve creating Monte Carlo generated cluster catalogues with associated richness and SPT-SZ significance derived from the fitted scaling relations. The first test checks whether the correct number of SZE-selected clusters is found and whether those clusters exhibit the correct number of optical matches with the expected λ distribution. As is clear from Fig. 5, the observations are perfectly consistent with the expectations from the model. Thus, our analysis shows that the data in our matched SPT-SZ+RM sample are well described by our adopted model. In the second test, we take the much larger *observed* RM catalogue as a starting point and use the model to test whether the expected number of SZE matches with the expected λ distribution is found. Unlike the first test, this one would in principle be sensitive to contamination within the RM sample. Here, the agreement is not as good because there is a tendency for there to be fewer observed matches than expected. However, the tension reaches the $\sim 2\sigma$ level at worst, and so there is no convincing evidence that our observed sample is inconsistent with the model.

(3) *The SZE-optical positional offset distribution.* We identify optical positional biases associated with 12 per cent of the sample due to the masking in the DES-SVA1 data. We remove these clusters and study the optical-SZE positional offset distribution for the rest of the matched sample. We model the underlying positional offset distribution as the sum of two Gaussians, while accounting for the SPT-SZ positional uncertainty. We show that the resulting distribution is consistent with being described by a dominant (63^{+15}_{-25} per cent) centrally peaked distribution with ($\sigma_0 = 0.07^{+0.03}_{-0.02}$

R_{500}) and a sub-dominant (~ 37 per cent) population characterized by larger separations ($\sigma_1 = 0.25_{-0.06}^{+0.07} R_{500}$). For the same population, the RM algorithm assumes that 79 per cent of the clusters will belong to a small-offset population, consistent with our observations.

We also match the SPT-SZ cluster candidates with $4 < \xi < 4.5$ to the RM optical cluster catalogues from DES-SVA1 to extend the mass range of the SZE-selected clusters. Including the SPT-SZ candidates between $\xi = 4$ and 4.5 increases the sample of matched clusters by ~ 40 per cent compared to the $\xi > 4.5$ sample, highlighting the potential synergies of SPT and DES in producing lower mass extensions of SZE-selected cluster samples. We show that this larger sample produces results that are broadly consistent with the $\xi > 4.5$ results, but only marginally tighter. This is due to the fact that mass constraints from lower signal-to-noise SPT clusters are somewhat weaker on a per cluster basis compared to the higher ξ sample. Future work benefiting from the larger region of overlap between the DES and SPT surveys will improve our derived constraints and help to better characterize the optical and SZE properties of cluster samples in terms of positional offsets, purity, and completeness. Moreover, the multiwavelength data sets available through DES and SPT enable characterization of the galaxy populations of large SZE-selected cluster samples, calibration of the SZE-selected cluster masses using weak lensing constraints, and many other promising studies.

ACKNOWLEDGEMENTS

We acknowledge the support by the DFG Cluster of Excellence ‘Origin and Structure of the Universe’, the Transregio program TR33 ‘The Dark Universe’ and the Ludwig-Maximilians University. The South Pole Telescope is supported by the National Science Foundation through grant PLR-1248097. Partial support is also provided by the NSF Physics Frontier Center grant PHY-1125897 to the Kavli Institute of Cosmological Physics at the University of Chicago, the Kavli Foundation and the Gordon and Betty Moore Foundation grant GBMF 947. AAS acknowledges a Pell grant from the Smithsonian Institution. TDH is supported by a Miller Research Fellowship. This work was partially completed at Fermilab, operated by Fermi Research Alliance, LLC under contract no. De-AC02-07CH11359 with the United States Department of Energy. The Dark Cosmology Centre is funded by the Danish National Research Foundation. CR acknowledges support from the Australian Research Council’s Discovery Projects scheme (DP150103208). Funding for the DES Projects has been provided by the US Department of Energy, the US National Science Foundation, the Ministry of Science and Education of Spain, the Science and Technology Facilities Council of the United Kingdom, the Higher Education Funding Council for England, the National Center for Supercomputing Applications at the University of Illinois at Urbana-Champaign, the Kavli Institute of Cosmological Physics at the University of Chicago, the Center for Cosmology and Astro-Particle Physics at the Ohio State University, the Mitchell Institute for Fundamental Physics and Astronomy at Texas A&M University, Financiadora de Estudos e Projetos, Fundação Carlos Chagas Filho de Amparo à Pesquisa do Estado do Rio de Janeiro, Conselho Nacional de Desenvolvimento Científico e Tecnológico and the Ministério da Ciência e Tecnologia, the Deutsche Forschungsgemeinschaft and the Collaborating Institutions in the Dark Energy Survey. The DES participants from Spanish institutions are partially supported by MINECO under grants AYA2012-39559,

ESP2013-48274, FPA2013-47986 and Centro de Excelencia Severo Ochoa SEV-2012-0234, some of which include ERDF funds from the European Union. The Collaborating Institutions are Argonne National Laboratory, the University of California at Santa Cruz, the University of Cambridge, Centro de Investigaciones Energeticas, Medioambientales y Tecnologicas-Madrid, the University of Chicago, University College London, the DES-Brazil Consortium, the Eidgenössische Technische Hochschule (ETH) Zürich, Fermi National Accelerator Laboratory, the University of Edinburgh, the University of Illinois at Urbana-Champaign, the Institut de Ciències de l’Espai (IEEC/CSIC), the Institut de Física d’Altes Energies, Lawrence Berkeley National Laboratory, the Ludwig-Maximilians Universität and the associated Excellence Cluster Universe, the University of Michigan, the National Optical Astronomy Observatory, the University of Nottingham, The Ohio State University, the University of Pennsylvania, the University of Portsmouth, SLAC National Accelerator Laboratory, Stanford University, the University of Sussex and Texas A&M University.

REFERENCES

- Abell G. O., 1958, *ApJS*, 3, 211
 Aihara H. et al., 2011, *ApJS*, 193, 29
 Allen S. W., Evrard A. E., Mantz A. B., 2011, *ARA&A*, 49, 409
 Andersson K. et al., 2011, *ApJ*, 738, 48
 Andreon S., Congdon P., 2014, *A&A*, 568, A23
 Annis J. et al., 2014, *ApJ*, 794, 120
 Ascaso B., Wittman D., Dawson W., 2014, *MNRAS*, 439, 1980
 Banerji M. et al., 2015, *MNRAS*, 446, 2523
 Bartlett J. G., Silk J., 1994, *ApJ*, 423, 12
 Benson B. A. et al., 2013, *ApJ*, 763, 147
 Biesiadzinski T., McMahon J., Miller C. J., Nord B., Shaw L., 2012, *ApJ*, 757, 1
 Bleem L. E., Stalder B., Brodwin M., Busha M. T., Gladders M. D., High F. W., Rest A., Wechsler R. H., 2015a, *ApJS*, 216, 20
 Bleem L. E. et al., 2015b, *ApJS*, 216, 27 (B15)
 Bocquet S. et al., 2015, *ApJ*, 799, 214
 Böhringer H. et al., 2000, *ApJS*, 129, 435
 Borgani S. et al., 2001, *ApJ*, 561, 13
 Cavaliere A., Fusco-Femiano R., 1976, *A&A*, 49, 137
 Crawford T. M., Switzer E. R., Holzappel W. L., Reichardt C. L., Marrone D. P., Vieira J. D., 2010, *ApJ*, 718, 513
 de Jong J. T. A., Verdoes Kleijn G. A., Kuijken K. H., Valentijn E. A., 2013, *Exp. Astron.*, 35, 25
 Desai S. et al., 2012, *ApJ*, 757, 83
 Diehl T. et al., 2012, *Phys. Procedia*, 37, 1332
 Eisenhardt P. R. M. et al., 2008, *ApJ*, 684, 905
 Eke V. R., Cole S., Frenk C. S., Patrick Henry J., 1998, *MNRAS*, 298, 1145
 Evrard A. E., Arnault P., Huterer D., Farahi A., 2014, *MNRAS*, 441, 3562
 Flaugher B. L. et al., 2012, in McLean I. S., Ramsay S. K., Takami H., eds, *Proc. SPIE Conf. Ser. Vol. 8446, Ground-based and Airborne Instrumentation for Astronomy IV*. SPIE, Bellingham, p. 844611
 Flaugher B. et al., 2015, preprint ([arXiv:1504.02900](https://arxiv.org/abs/1504.02900))
 Gioia I. M., Maccacaro T., Schild R. E., Wolter A., Stocke J. T., Morris S. L., Henry J. P., 1990, *ApJS*, 72, 567
 Gladders M. D., Yee H. K. C., 2000, *AJ*, 120, 2148
 Haehnelt M. G., Tegmark M., 1996, *MNRAS*, 279, 545
 Hao J. et al., 2010, *ApJS*, 191, 254
 Hasselfield M. et al., 2013, *J. Cosmol. Astropart. Phys.*, 7, 8
 High F. W. et al., 2012, *ApJ*, 758, 68
 Koester B. P. et al., 2007, *ApJ*, 660, 239
 Lauer T. R., Postman M., Strauss M. A., Graves G. J., Chisari N. E., 2014, *ApJ*, 797, 82
 Laureijs R. et al., 2011, preprint ([arXiv:1110.3193](https://arxiv.org/abs/1110.3193))
 Lin Y., Mohr J. J., Stanford S. A., 2004, *ApJ*, 610, 745
 Liu J. et al., 2015, *MNRAS*, 448, 2085

- LSST Dark Energy Science Collaboration 2012, preprint (arXiv:1211.0310)
- Mantz A., Allen S. W., Ebeling H., Rapetti D., Drica-Wagner A., 2010, MNRAS, 406, 1773
- Mantz A. B. et al., 2015, MNRAS, 446, 2205
- Mei S. et al., 2009, ApJ, 690, 42
- Melchior P. et al., 2015, MNRAS, 449, 2219
- Melin J.-B., Bartlett J. G., Delabrouille J., 2006, A&A, 459, 341
- Menanteau F. et al., 2010, ApJS, 191, 340
- Mohr J. J. et al., 2008, in Brissenden R. J., Silva D. R., eds, Proc. SPIE Conf. Ser. Vol. 7016, Observatory Operations: Strategies, Processes, and Systems II. SPIE, Bellingham, p. 70160L
- Mohr J. J. et al., 2012, in Radziwill N. M., Chiozzi G., eds, Proc. SPIE Conf. Ser. Vol. 8451, Software and Cyberinfrastructure for Astronomy II. SPIE, Bellingham, p. 84510D
- Moscovich-Eiger A., Nadler B., Spiegelman C., 2013, preprint (arXiv:e-prints)
- Muzzin A. et al., 2012, ApJ, 746, 188
- Ngeow C. et al., 2006, in Silva D. R., Doxsey R. E., eds, Proc. SPIE Conf. Ser. Vol. 6270, Observatory Operations: Strategies, Processes, and Systems. SPIE, Bellingham, p. 627023
- Pacaud F. et al., 2007, MNRAS, 382, 1289
- Planck Collaboration 2015, preprint (arXiv:e-prints)
- Reichardt C. L. et al., 2013, ApJ, 763, 127
- Rozo E., Rykoff E. S., 2014, ApJ, 783, 80
- Rozo E. et al., 2010, ApJ, 708, 645
- Rozo E. et al. 2015a, preprint (arXiv:1507.05460)
- Rozo E., Evrard A. E., Rykoff E. S., Bartlett J. G., 2014a, MNRAS, 438, 62
- Rozo E., Bartlett J. G., Evrard A. E., Rykoff E. S., 2014b, MNRAS, 438, 78
- Rozo E., Rykoff E. S., Bartlett J. G., Melin J. B., 2015b, MNRAS, 450, 592
- Rykoff E. S. et al., 2012, ApJ, 746, 178
- Rykoff E. S. et al., 2014, ApJ, 785, 104
- Sánchez C. et al., 2014, MNRAS, 445, 1482
- Sehgal N. et al., 2013, ApJ, 767, 38
- Soares-Santos M. et al., 2011, ApJ, 727, 45
- Song J. et al., 2012, ApJ, 761, 22
- Staniszewski Z. et al., 2009, ApJ, 701, 32
- Story K. et al., 2011, ApJ, 735, L36
- Šuhada R. et al., 2012, A&A, 537, A39
- Sunyaev R. A., Zel'dovich Y. B., 1972, Comments Astrophys. Space Phys., 4, 173
- The Dark Energy Survey Collaboration 2005, preprint (arXiv:e-prints)
- Tinker J., Kravtsov A. V., Klypin A., Abazajian K., Warren M., Yepes G., Gottlöber S., Holz D. E., 2008, ApJ, 688, 709
- Vanderlinde K. et al., 2010, ApJ, 722, 1180
- Viana P., Liddle A., 1999, MNRAS, 303, 535
- Vikhlinin A., McNamara B., Forman W., Jones C., Quintana H., Hornstrup A., 1998, ApJ, 502, 558
- Vikhlinin A. et al., 2009, ApJ, 692, 1060
- Wen Z. L., Han J. L., Liu F. S., 2012, ApJS, 199, 34
- White S. D. M., Efstathiou G., Frenk C. S., 1993, MNRAS, 262, 1023
- Wiesner M. P., Lin H., Soares-Santos M., 2015, MNRAS, 452, 701
- Zenteno A. et al., 2011, ApJ, 734, 3
- Zwicky F., Herzog E., Wild P., 1968, Catalogue of Galaxies and of Clusters of Galaxies. California Institute of Technology, Pasadena

APPENDIX A: STUDY OF MASS-RICHNESS SCATTER FOR CLUSTERS FOUND WITH THE VT METHOD

The scatter of the mass–richness relation is a primary source of systematic uncertainties in cosmological measurements using galaxy clusters. To ensure that this and other astrophysical systematics are kept under control in future cosmology analyses, the DES collaboration has pursued the development of multiple cluster-finding algorithms. Here, we present initial results obtained using the analysis framework described throughout this paper on a DES-SVA1

cluster catalogue created using the VT cluster finder (Soares-Santos et al. 2011). The VT method is fundamentally different from RM. Specifically, VT uses photometric redshifts to detect clusters in 2+1 dimensions, and is designed to produce a cluster catalogue up to $z \sim 1$ and down to mass $M \sim 10^{13.5} M_{\odot}$ without any assumptions about the colours of galaxies in cluster environments. VT has been tested on DES simulations (Soares-Santos et al. 2011) and on SDSS data. The mean mass–richness relation has been calibrated using a stacked weak-lensing analysis of the SDSS VT clusters (Wiesner, Lin & Soares-Santos 2015). In this appendix, we describe the first study of the scatter of the mass–richness relation using our analysis framework. This study demonstrates that VT, although not as mature a cluster finder as RM, is a promising algorithm. Our report provides an assessment of the current VT performance and helps to identify areas where improvements should be made. Some of these improvements can be applied to other cluster-finding algorithms, too.

A1 VT method

To detect clusters with the VT method, we build 2D tessellations in each photometric redshift shell and flag galaxies that lie in high-density cells as cluster members. The density threshold is set in a non-parametric way from the 2-point correlation function of that given shell. This takes advantage of the fact that the distribution of VT cell densities can be uniquely predicted for any given point process. The 2-point function is a good description of the point process of the background galaxies on the sky. Clusters cause a small deviation from the predicted distribution, and we take the point where that deviation is maximized as the threshold for detection.

A2 VT catalogue for DES SVA1 data

For the DES SVA1 data, the final catalogue consists of 12 948 clusters with richness $N_{\text{vt}} > 5$ and redshifts in the range $0.15 < z < 1$. N_{vt} is defined as the number of member galaxies. The catalogue covers the SPT-E and SPT-W regions of the SVA1 total footprint. We use DESDM data products as inputs, namely the *Gold* galaxy catalogue, plus photometric redshifts and mask information. We used a mask to apply magnitude cuts $10 < \text{mag_auto_i} < 23.5$. The photometric redshift information was obtained using a neural network method (Sánchez et al. 2014).

We match the VT and SPT-SZ catalogues using the same method as described in Section 2.3. We sort the SPT-SZ cluster sample according to decreasing SPT observable ξ and sort the VT catalogues according to decreasing richness. Then, we associate the SPT cluster candidate with the richest cluster candidate whose centre lies within $1.5 R_{500}$ of the SPT-SZ centre. We finally remove the associated optical cluster from the list of possible counterparts when matching the remaining SPT-selected clusters. This procedure results in 42 VT clusters matched to $\xi > 4.5$ SPT-SZ clusters.

A3 Results

In Fig. A1 we show the optical richness N_{vt} as a function of ξ (left-hand panel) for the 42 matched clusters. The scatter in the richness– ξ plot indicates that the VT richness performs poorly as a mass indicator. Improvements to the method are being developed based on these findings. Specifically, future work will explore using the galaxy magnitudes to calculate total stellar masses, which can then be used as a mass proxy.

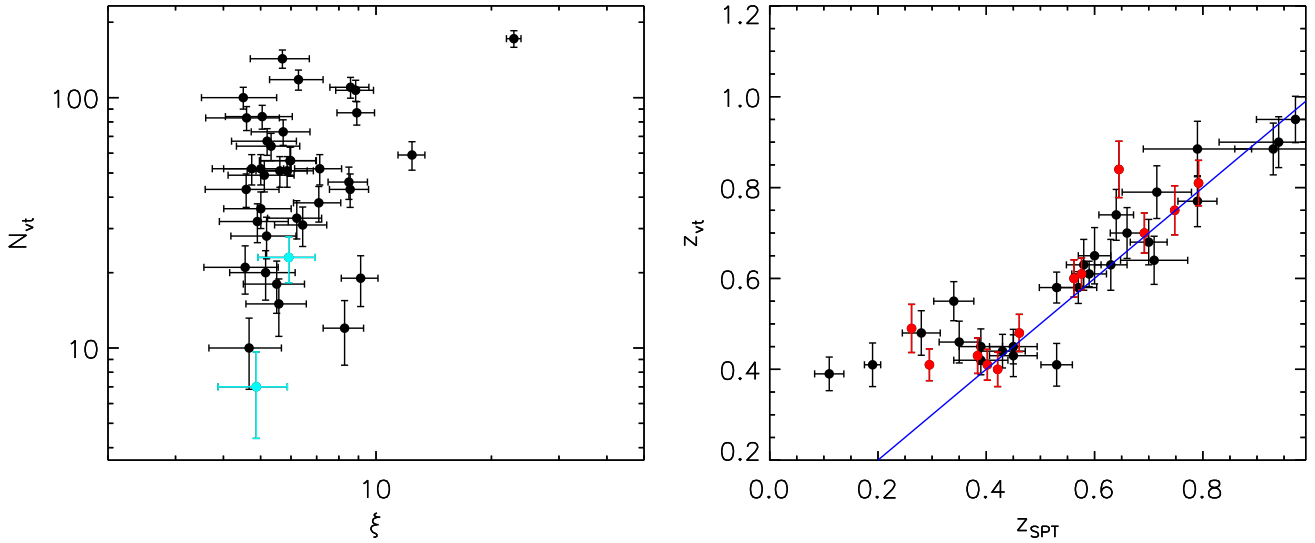


Figure A1. Left-hand panel: richness as a function of the SPT-SZ significance ξ for the SPT-SZ+VT cluster sample. Clusters at $z < 0.25$ are not used for the richness–mass fit and are shown in cyan. Right-hand panel: the estimated redshift for the VT sample as a function of SPT-SZ redshifts as presented in B15. SPT-SZ candidates with spectroscopic redshift are shown in red.

Fig. A1 also shows the VT estimated redshifts versus the redshifts determined in SPT follow-up observations (right-hand panel). There is good agreement for $z > 0.35$. Deviations at lower z are found to arise from problems in the calculation of the 2-point function predictions at $z < 0.35$, and a fix is underway.

We also obtained the best-fitting parameters, and the corresponding 68 per cent confidence level uncertainties, for the richness–mass scaling relation described by equations (7) and (8):

$$\begin{aligned} A_{VT} &= 48.1^{+6.9}_{-6.3}, & B_{VT} &= 0.56^{+0.51}_{-0.25}, \\ C_{VT} &= -0.51^{+1.82}_{-1.48}, & D_{VT} &= 0.64^{+0.29}_{-0.14}. \end{aligned} \quad (\text{A1})$$

This result is consistent with the calibration for the mean relation obtained with weak lensing (Wiesner et al. 2015). The uncertainties, however, are larger than those obtained for RM. We expect that improvements to the richness estimator will result in better performance in future applications of the VT method. VT is not as mature a cluster finder method as RM but the complementarity of the two techniques argues for further development of this alternative. This study allowed us to identify a key area for improvement and establishes a framework for future assessment of the mass–richness scatter for VT clusters.

¹Department of Physics, Ludwig-Maximilians-Universitaet, Scheinerstr. 1, D-81679 Muenchen, Germany

²Excellence Cluster Universe, Boltzmannstr. 2, D-85748 Garching, Germany

³Department of Physics, University of Arizona, 1118 E 4th St, Tucson, AZ 85721, USA

⁴Fermi National Accelerator Laboratory, P. O. Box 500, Batavia, IL 60510, USA

⁵Kavli Institute for Cosmological Physics, University of Chicago, Chicago, IL 60637, USA

⁶Department of Astronomy and Astrophysics, University of Chicago, 5640 South Ellis Avenue, Chicago, IL 60637, USA

⁷Max Planck Institute for Extraterrestrial Physics, Giessenbachstrasse, D-85748 Garching, Germany

⁸Kavli Institute for Particle Astrophysics and Cosmology, P. O. Box 2450, Stanford University, Stanford, CA 94305, USA

⁹SLAC National Accelerator Laboratory, Menlo Park, CA 94025, USA

¹⁰Argonne National Laboratory, 9700 South Cass Avenue, Lemont, IL 60439, USA

¹¹Center for Cosmology and Astro-Particle Physics, The Ohio State University, Columbus, OH 43210, USA

¹²Department of Physics, The Ohio State University, Columbus, OH 43210, USA

¹³Laboratório Interinstitucional de e-Astronomia - LIneA, Rua Gal. José Cristino 77, Rio de Janeiro, RJ 20921-400, Brazil

¹⁴CERN, CH-1211 Geneva 23, Switzerland

¹⁵Universitäts-Sternwarte, Fakultät für Physik, Ludwig-Maximilians Universität München, Scheinerstr. 1, D-81679 München, Germany

¹⁶Cerro Tololo Inter-American Observatory, National Optical Astronomy Observatory, Casilla 603, La Serena, Chile

¹⁷Department of Physics and Astronomy, University College London, Gower Street, London WC1E 6BT, UK

¹⁸Department of Physics and Astronomy, University of Pennsylvania, Philadelphia, PA 19104, USA

¹⁹Institute of Astronomy, University of Cambridge, Madingley Road, Cambridge CB3 0HA, UK

²⁰Facultat de Ciències, Institut de Ciències de l’Espai, IEEC-CSIC, Campus UAB, Torre C5 par-2, E-08193 Bellaterra, Barcelona, Spain

²¹Department of Physics, Harvard University, 17 Oxford Street, Cambridge, MA 02138, USA

²²Harvard-Smithsonian Center for Astrophysics, 60 Garden Street, Cambridge, MA 02138, USA

²³Institut d’Astrophysique de Paris, Univ. Pierre et Marie Curie and CNRS UMR7095, F-75014 Paris, France

²⁴Department of Physics and Astronomy, University of Missouri, 5110 Rockhill Road, Kansas City, MO 64110, USA

²⁵Department of Physics and Enrico Fermi Institute, University of Chicago, 5640 South Ellis Avenue, Chicago, IL 60637, USA

²⁶Institute of Cosmology and Gravitation, University of Portsmouth, Portsmouth PO1 3FX, UK

²⁷Observatório Nacional, Rua Gal. José Cristino 77, Rio de Janeiro, RJ 20921-400, Brazil

²⁸Department of Astronomy, University of Illinois, 1002 W. Green Street, Urbana, IL 61801, USA

²⁹National Center for Supercomputing Applications, 1205 West Clark St, Urbana, IL 61801, USA

³⁰George P. and Cynthia Woods Mitchell Institute for Fundamental Physics and Astronomy, and Department of Physics and Astronomy, Texas A&M University, College Station, TX 77843, USA

³¹Department of Physics, McGill University, 3600 Rue University, Montreal, QC H3A 2T8, Canada

³²*Department of Physics, University of California, Berkeley, CA 94720, USA*

³³*Jet Propulsion Laboratory, California Institute of Technology, 4800 Oak Grove Dr., Pasadena, CA 91109, USA*

³⁴*Department of Astronomy, University of Michigan, Ann Arbor, MI 48109, USA*

³⁵*Department of Physics, University of Michigan, Ann Arbor, MI 48109, USA*

³⁶*Institut de Física d'Altes Energies, Universitat Autònoma de Barcelona, E-08193 Bellaterra, Barcelona, Spain*

³⁷*Australian Astronomical Observatory, North Ryde, NSW 2113, Australia*

³⁸*Department of Astronomy, The Ohio State University, Columbus, OH 43210, USA*

³⁹*Kavli Institute for Astrophysics and Space Research, Massachusetts Institute of Technology, 77 Massachusetts Avenue, Cambridge, MA 02139, USA*

⁴⁰*Brookhaven National Laboratory, Bldg 510, Upton, NY 11973, USA*

⁴¹*School of Physics, University of Melbourne, Parkville, VIC 3010, Australia*

⁴²*Department of Physics and Astronomy, Pevensey Building, University of Sussex, Brighton BN1 9QH, UK*

⁴³*Centro de Investigaciones Energéticas, Medioambientales y Tecnológicas (CIEMAT), E-28040 Madrid, Spain*

⁴⁴*Institute for Astronomy, University of Hawaii at Manoa, Honolulu, HI 96822, USA*

⁴⁵*Department of Physics, University of Illinois, 1110 W. Green St, Urbana, IL 61801, USA*

⁴⁶*Department of Physics, Stanford University, 382 Via Pueblo Mall, Stanford, CA 94305, USA*

⁴⁷*Dark Cosmology Centre, Niels Bohr Institute, University of Copenhagen, Juliane Maries Vej 30, DK-2100 Copenhagen Ø, Denmark*

⁴⁸*Cerro Tololo Inter-American Observatory, Casilla 603, La Serena, Chile*

This paper has been typeset from a $\text{\TeX}/\text{\LaTeX}$ file prepared by the author.