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Black hole accretion versus star formation rate: theory confronts observations

Marta Volonteri, ^{1★} Pedro R. Capelo, ² Hagai Netzer, ³ Jillian Bellovary, ⁴ Massimo Dotti ⁵ and Fabio Governato ⁶

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ABSTRACT

We use a suite of hydrodynamical simulations of galaxy mergers to compare star formation rate (SFR) and black hole accretion rate (BHAR) for galaxies before the interaction ('stochastic' phase), during the 'merger' proper, lasting ~0.2–0.3 Gyr, and in the 'remnant' phase. We calculate the bivariate distribution of SFR and BHAR and define the regions in the SFR–BHAR plane that the three phases occupy. No strong correlation between BHAR and galaxy-wide SFR is found. A possible exception are galaxies with the highest SFR and the highest BHAR. We also bin the data in the same way used in several observational studies, by either measuring the mean SFR for AGN in different luminosity bins, or the mean BHAR for galaxies in bins of SFR. We find that the apparent contradiction or SFR versus BHAR for observed samples of AGN and star-forming galaxies is actually caused by binning effects. The two types of samples use different projections of the full bivariate distribution, and the full information would lead to unambiguous interpretation. We also find that a galaxy can be classified as AGN-dominated up to 1.5 Gyr after the merger-driven starburst took place. Our study is consistent with the suggestion that most low-luminosity AGN hosts do not show morphological disturbances.

Key words: galaxies: active – galaxies: interactions – galaxies: nuclei.

1 INTRODUCTION

The correlation (or lack thereof) between the black hole accretion rate (BHAR) and the star formation rate (SFR) of their host galaxies has been the subject of numerous investigations (e.g. Netzer 2009; Harrison et al. 2012; Mullaney et al. 2012; Rosario et al. 2012, 2013a; Hickox et al. 2014; Netzer et al. 2014, and references therein). One of the drivers behind these studies is the empirical correlation between BH mass and bulge mass in the local Universe (e.g. Marconi & Hunt 2003; Häring & Rix 2004; Gültekin et al. 2009) which suggests that the BH and stellar bulge have assembled in tandem (co-evolution). In the strictest view of co-evolution, to obtain a BH mass proportional to the bulge mass, BHAR should be proportional to SFR (or, at least, the SFR that builds up the bulge). Indeed, the cosmic total SFR and BHAR (i.e. the rate per unit comoving volume) seem to track each other at least to $z \sim 3$ (Heckman

et al. 2004; Merloni, Rudnick & Di Matteo 2004; Silverman et al. 2008, 2009; Madau & Dickinson 2014).

When SFR and BHAR are compared source by source, the connection appears to be weak, unless only the SFR in the central region is taken into account (<1 kpc; Diamond-Stanic & Rieke 2012, and references therein). A comprehensive work of this type by Rosario et al. (2012), shows that at low AGN luminosities, BHAR and SFR are uncorrelated, while at high AGN luminosities a significant correlation emerges. On the other hand, Mullaney et al. (2012) and Chen et al. (2013) study a sample of mass-selected and starforming galaxies, respectively, and suggest that the average BHAR correlates well with the average SFR, once the shorter variability time-scales of BHAR with respect to SFR are taken into account (see also Hickox et al. 2014).

In a companion paper (Volonteri et al. 2015), we have used new simulations of galaxy mergers and investigated the temporal correlation between SFR and BHAR, and their respective variability. We found that BHAR and nuclear (<100 pc) SFR are well correlated and vary on similar time-scales. However, we found that galaxy-wide (<5 kpc) SFR, which is used in statistical studies

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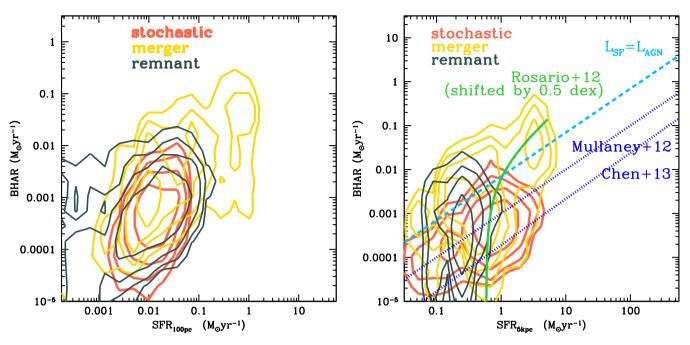


Figure 1. BHAR versus SFR within 100 pc (left) and within 5 kpc (right). Contours are based on 10 equally-spaced logarithmic levels, dropping the six lowest levels for clarity. The SFR is an average over the 100 Myr before the time-step used for the BHAR calculation. We distinguish stochastic (red), merger (gold) and remnant (dark grey) phases. We show curves from Rosario et al. (2012; AGN, green, see section 3.2 for details), Mullaney et al. 2012 (mass-selected galaxies, upper blue dotted line) and Chen et al. (2013; star-forming galaxies, lower blue dotted line). We also mark the line separating AGN and SF dominated regions (light blue dashed line).

(Mullaney et al. 2012; Rosario et al. 2012; Chen et al. 2013; Delvecchio et al. 2015), and BHAR are typically temporally uncorrelated, and have different variability time-scales, except during the short-lived merger proper (\sim 0.2–0.3 Gyr). We here compare SFR and BHAR by referring to the luminosity resulting from these processes ($L_{\rm SFR}$ and $L_{\rm AGN}$, respectively), and aim at providing a framework for comparison with observations.

In this Letter, we show how very detailed numerical simulation of galaxy mergers confirm the recent observational finding, provided the simulation results are treated in the same way as observations. In particular, we show that accounting for the different durations of the various stages of the mergers, and binning the results in a way similar to the observational methods, can bring the results very close to the observed correlations and reconciles seemingly conflicting results. The results presented here are based on our previous detailed work published in Capelo et al. (2015) and Volonteri et al. (2015). We refer the reader to those papers for a complete description of the simulations, of the physical implementation and of the analysis.

2 COMPARISON WITH OBSERVATIONS

The comparison of the numerical simulations with observed correlations has some caveats. As discussed by Neistein & Netzer (2014), measurements of SFRs are based on indicators that last $\gtrsim 100$ Myr, while the AGN luminosity, hence the BHAR, is measured instantaneously. Because of this, all our results concerning SFR represent averages over 100 Myr. The SFR is converted to SF luminosity, $L_{\rm SFR}$, by assuming that one solar mass per year corresponds to 10^{10} solar luminosities. BHAR is averaged instead over 1 Myr (see Volonteri et al. 2015, for a case where BHAR and SFR are measured over the same time interval). We distinguish between SF-dominated galaxies $(L_{\rm SF} > L_{\rm AGN})$ and AGN-dominated galaxies $(L_{\rm SF} < L_{\rm AGN})$.

Dense gas in the host galaxy may cause some obscuration, especially during the merger phase. We estimated the column density for a 1:2 merger, in cylinders of radius r=100 pc centred on the BHs, and found that it varies between $\sim 10^{22}$ (face-on) and $\sim 10^{23}$ cm⁻² (edge-on), with an increase by a factor of 2 during the merger proper. Using the WebPIMMS interface (with a photon index $\Gamma=1.4$), the average (maximal, for edge-on galaxies) dimming in the 2–10 keV band varies between 20 per cent (40 per cent) in the stochastic and remnant phase, to 40 per cent (55 per cent) during the merger if the source were at z=0, and it is almost negligible for z=3. We discuss this further in Section 2.2.

An additional difference between theory and observations is related to the fact that observed samples do not contain information about the fraction of sources in the stochastic, merger and remnant stages, and how many represent a merger with mass ratio of 1:2, 1:6 or 1:10. As shown in Fig. 1, the three stages cover differently the SFR–BHAR plane and a proper comparison with the observations requires an adjustment of the various numbers in the three groups.

The approach we adopt is to specify the boundaries of a region in the BHAR–SFR plane relevant to a population of relatively low-mass galaxies that includes quiescent galaxies and mergers with mass ratio from 1:1 to 1:10. Real correlations depend on the distribution of points within these boundaries, and will include galaxies of different masses and structural parameters, and this we cannot do quantitatively. One can, however, estimate the fraction of objects in each of the merger phases that is included in most observed samples, e.g. the remnant phase in low-redshift samples, and remnants+mergers in higher redshift samples. The duration of each phase must also be taken into account, since this determines the number of objects from this phase in a specific sample and hence the nature of the derived correlations that depend on such numbers. For example, the merger phase, being short-lived,

 \sim 0.2–0.3 Gyr, is sub-dominant in the general distribution, with respect to the stochastic and remnant phases, which last 0.8–1.3 Gyr each

2.1 Bivariate distribution

In Fig. 1, we show the bivariate distribution of BHAR and SFR within 100 pc, and within 5 kpc, distinguishing galaxies in the stochastic (red), merger (gold) and remnant (dark grey) phases. In the case of the central SFR, SFR_{100 pc} (left-hand panel), the stochastic and remnant phases occupy similar regions, while the merger phase, beyond a primary peak coincident with the other two phases, presents a secondary peak at higher BHAR and SFR_{100 pc}. The secondary peak is limited in its extent in SFR because averaging over 100 Myr removes the highest peaks of SFR, which are associated with the highest peaks in BHAR (cf. fig. 14 in Volonteri et al. 2015). As noted previously (Hopkins & Quataert 2010; Thacker et al. 2014; Diamond-Stanic & Rieke 2012; LaMassa et al. 2013; Esquej et al. 2014) SFR_{100 pc} is a better tracer of BHAR than SFR_{5 kpc}, i.e. they are better correlated. As discussed by Capelo et al. 2015 (see also Hopkins & Quataert 2010), the main drivers of the BHAR are the gas content (setting an overall 'normalization'; see also Rosario et al. 2013b; Vito et al. 2014) and local losses of angular momentum, and these became more clearly linked to SFR in the nuclear region. We have compared the results of our standard run to a low-resolution run (see Appendix of Volonteri et al. 2015). The degree of correlation is degraded at lower resolution, showing that our simulations are accurately representing inflows due to angular momentum loss. Additionally, because the nuclear region contains both cold inflowing gas and gas affected by thermal feedback, we calculate the accretion for each individual particle separately, rather than averaging over the gas broadly. The hot gas contribution is fairly negligible, because of its lower density and higher temperature.

The link between galaxy-wide SFR (SFR $_{5\,\mathrm{kpc}}$, right) and BHAR seems much weaker. Galaxies in the stochastic phase occupy a region in the BHAR–SFR $_{5\,\mathrm{kpc}}$ plane (red contours in Fig. 1) roughly tracing the region between the relations suggested by Mullaney et al. 2012 and Chen et al. 2013 to characterize star-forming galaxies (dotted blue lines). The merging systems (gold contours) approach the upper part of the curves shown in Rosario et al. (2012), who suggest that the highest luminosity AGN are merger driven, and they exhibit a tighter correlation between BHAR and SFR $_{5\,\mathrm{kpc}}$. In fact, in the merger phase, losses of angular momentum are driven by global dynamics and BHAR and SFR become somewhat better correlated even on large scales. In the remnant phase (grey contours) a wide range of BHAR can be associated to a given SFR $_{5\,\mathrm{kpc}}$.

The remnant phase is, perhaps, the most interesting when comparing to AGN observations, as 80 per cent of the AGN do not show any hint of a companion (e.g. Rosario et al. 2012) and are ordinary (massive) star-forming galaxies. For most of the remnant phase the simulated galaxies do not show strong morphological disturbances, however, the BHAR is sufficiently high at times that the galaxies enter the AGN dominated region of the BHAR-SFR_{5 kpc} diagram. Observationally, this has been a matter of some debate. Rosario et al. (2013b) find a strong connection between AGN activity and SFR, i.e. most AGN hosts are on the main sequence of star-forming galaxies (Elbaz et al. 2007; Noeske et al. 2007). It has been suggested that there is an excess of AGN in post-merger (Ellison et al. 2013) or post-starburst galaxies (Wild, Heckman & Charlot 2010), and that galaxies hosting moderate-luminosity AGN are transitioning from SF to quiescence (Schawinski et al. 2009). Our broad interpretation is that the AGN is more likely to be observed in the (much longer) remnant phase, where it enters the AGN-dominated region. In fact, in our calculations (fig. 13 in Volonteri et al. 2015), the AGN zigzags continuously in and out of this region.

We can now estimate the fraction of galaxies in different stages of the merger that are likely to be found in large observed samples. Galaxies in the merger phase would be considered AGN-dominated $(L_{\rm AGN} > L_{\rm SFR})$ for about 40 per cent of that phase, galaxies in the remnant phase are AGN-dominated for 25 per cent of their phase, and galaxies in the stochastic phase 7 per cent. However, the merger phase is much shorter. Therefore the overall probability, defined as the ratio between the time when $L_{\rm AGN} > L_{\rm SFR}$ and the total simulation time (stochastic + merger + remnant), of finding a galaxy in the remnant phase in the AGN-dominated region is almost twice as large than for a merging galaxy (13 per cent versus 7 per cent, respectively), and six times higher than for a galaxy in the stochastic phase.

2.2 BHAR versus SFR using binned data

Most observational papers cannot consider the full bivariate distribution of SFR and BHAR. Instead, they focus on the mean SFR for AGN in different luminosity bins (i.e. stacking in bins of BHAR; e.g. Rosario et al. 2012) or the mean BHAR in SFR bins (i.e. stacking in bins of SFR; e.g. Mullaney et al. 2012; Chen et al. 2013; Delvecchio et al. 2015). To compare with those studies, we binned our simulations in the same way as used by them: i.e. we calculate the mean SFR in bins of BHAR and the mean BHAR in bins of SFR. Our points are not uniform on the SFR and BHAR axes across the different evolutionary stages, while observational samples use relatively well distributed bins across their range. We therefore discard bins with few points (<50 for sub-samples, <300 for the complete sample), and draw only 50 or 300 random points, respectively, for the remaining bins.

The results are shown in Fig. 2 where the first way of binning is shown in the top panel and the second in the bottom panel. Remarkably, in the former case we recover the trends of Rosario et al. (2012), while in the latter we recover the trend found by Mullaney et al. (2012), Chen et al. (2013) and Delvecchio et al. (2015). We suggest that the different trends found for AGN and SF galaxies are in part caused by the different projections of the full bivariate distribution, and that the intrinsic distribution of properties in those samples is similar to the one shown in Fig. 1.

In light blue, we show a bona fide global sample which takes into account the time spent in each stage, and where we have enforced high statistical significance in each bin. This is the sample to be compared to observations (SFR $> 0.5~{\rm M}_{\odot}~{\rm yr}^{-1}$ seems also a plausible lower limit for the SFR). The process of averaging, and the statistical significance, has an obvious effect in removing the highest BHAR/SFR sources. For instance, when we require a large number of points per bin, the less populated high BHAR (SFR) bins disappear. Also, the mean in a BHAR (SFR) bin is driven by the more numerous sources at low SFR (BHAR).

In the top panel, except for the two bins at the highest BHAR, the SFR is consistent with being uncorrelated with BHAR. The wiggles show that the weight is dominated by the stochastic phase at low BHAR, then the remnant, and finally the merger phase take over. The bend at high BHAR is driven only by the last two points, dominated by the merger stage, as proposed by Rosario et al. 2012. However, the increase in SFR is only a factor of 2–3, similar to the enhancement seen by Stanley et al. (2015). An increase in the average masses of hosts of the most luminous AGN, or an enforced

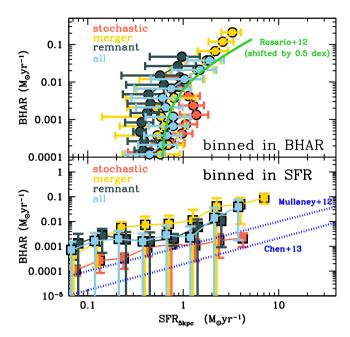


Figure 2. Mean SFR in bins of BHAR (top) and mean BHAR in bins of SFR (bottom); merger and remnant points have been slightly shifted to enable appreciating the size of the 1σ error bar. For the comparison, note that Rosario et al. (2012; green) measured the mean SFR for AGN in different luminosity bins, while Mullaney et al. (2012; blue) Chen et al. (2013; blue) measured the mean BHAR by stacking galaxies in bins of SFR.

correlation between SFR and the long term BHAR may also be responsible for the trend.

Note, again, that our simulations apply only to small-to-medium stellar mass systems. We can try to extrapolate the behaviour to more massive galaxies using the arguments outlined in Volonteri et al. 2015, which are based on the assumption that both BHAR and SFR would increase approximately linearly with stellar mass. In this case, larger galaxies would move by about one order of magnitude of BHAR per one order of magnitude of SFR in the plane shown in Fig. 1. This leads us to suggest that if our stellar masses and BHs were two orders of magnitude larger (a few times $\sim 10^{12}$ and 10^8 solar masses, respectively), then they would occupy a region surrounding the z=2 curve of Rosario et al. (2012) when we bin in BHAR. In Fig. 2, we show that the z=0 curve shifted by 0.5 dex in both BHAR and SFR (the shapes of all these curves in Rosario et al. 2012 are basically identical).

When binning in SFR (lower panel of Fig. 2), galaxies in the merger and remnant phases are \sim 0.5–1 dex above the extrapolation of the relation suggested by Mullaney et al. (2012) and Chen et al. (2013). The stochastic phase is closer to their fit in normalization, but has a shallower slope, caused by a drop in BHAR at the highest SFR, when, during a SF burst, SN feedback becomes strong enough to affect the gas properties. The full population has a best-fitting $\log(\text{BHAR}) = (1.05 \pm 0.29) \log(\text{SFR}) - 2.27 \pm 0.12$. If we included obscuration due to gas in the galaxy, the normalization decreases to -2.43. Furthermore, SF-dominated galaxies $(L_{\text{SF}} > L_{\text{AGN}})$, light squares in Fig. 3) have $\log(\text{BHAR}) = (0.93 \pm 0.17) \log(\text{SFR}) - 2.89 \pm 0.07$, slightly shallower and highly normalized than, but very close to, the curve proposed by Mullaney et al. (2012). Obscuration in the host would further decrease the normalization to -2.96.

Beyond binning, selection also matters. As highlighted by Stanley et al. (2015), their study and Rosario et al. (2012) start with

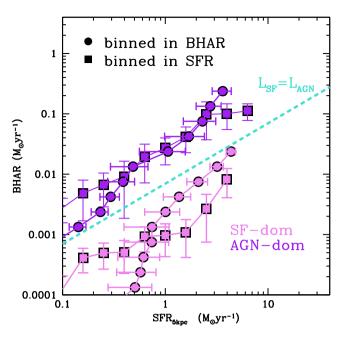


Figure 3. Trends of SFR and BHAR for binned data, dividing the sample in AGN-dominated (darker points) and SF-dominated (lighter points). The BHAR at a given SFR for AGN-dominated galaxies is about 20 times higher than for SF-dominated ones.

a population of AGN. Mass- or SF-selected samples start from a parent population of galaxies, which may or not host an AGN, and the average is likely dominated by sources which do not have AGN. Starting from an AGN population one studies the direct link between BHAR and SFR for galaxies in an active phase. A galaxy parent population instead highlights the broad BHAR-SFR co-evolution over cosmic time and the whole population. In Fig. 3, we propose a new way of looking at the SFR-BHAR connection. Observations often have only upper limits for either SFR or BHAR. In simulations, instead, we know when the BH is accreting or not, and what the SFR is, down to very low values (e.g. when SF is suppressed by SN feedback after a burst). We divide our sample in AGN- and SFdominated galaxies (L_{SF} larger or smaller than L_{AGN} , respectively) and estimate the increase in BHAR versus SFR in galaxies during the phase where BHs are actively growing. AGN-dominated galaxies have a BHAR-SFR relation with a normalization about 20 times higher than SF-dominated ones, i.e. their BHs grow relatively more than the galaxy stellar mass by this factor. As the BHAR, i.e. the AGN luminosity limit, is lowered, more of the SF-dominated region is sampled, and at low luminosity AGN-selected samples cross from AGN-dominated to SF-dominated, thus lessening, or erasing, the significance of a correlation between SFR and BHAR (see also Hickox et al. 2014).

3 CONCLUSIONS

We analyse a suite of high-resolution galaxy merger simulations to study BH and galaxy properties during various phases of the mergers. We calculate the BHAR and the SFR during the 'stochastic' phase (galaxies in isolation or in the early phases of an encounter), the 'merger' proper (when the merger dynamics dominates), and the 'remnant' phase (from the end of the merger to the return to quiescence). We find that in the remnant phase the BHAR can be sufficiently high at times to move the galaxy into the AGN-dominated region. The probability of finding a galaxy in the remnant phase,

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even long after the starburst took place, in the AGN-dominated region is twice as large than for a merging galaxy. The main goal in this Letter is to compare the relationships between BHAR and SFR with observations.

We find that different projections of the bivariate distribution recover different trends of the population. If the observations can be extrapolated to the simulated mass range, we are able to reconcile seemingly contradictory observational results when the mean SFR for AGN in different luminosity bins (e.g. Rosario et al. 2012) or the mean BHAR stacking galaxies in bins of SFR are measured (e.g. Mullaney et al. 2012; Chen et al. 2013; Delvecchio et al. 2015). Hickox et al. (2014), with an ansatz that over long time-scales BHAR and SFR are perfectly correlated, reach similar conclusions, also suggesting that future progress requires a direct measurement of the bivariate SFR/BHAR distribution.

The bivariate distribution derived from our simulations suggests that galaxies in the stochastic phase would not be considered AGN-dominated ($L_{\rm SF}\gg L_{\rm AGN}$). Galaxies would be AGN-dominated chiefly during the merger and remnant phases. The BHAR at a given SFR for AGN-dominated galaxies is about 20 times higher than for SF-dominated ones. During most of the merger proper and in the remnant phase a given SFR can be associated with a large range of BHAR. A possible exception is the groups of galaxies with the highest SFR and the highest BHAR, characterized by having $L_{\rm SF}\simeq L_{\rm AGN}$.

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