

# Invited review: Infrared spectroscopy of planetary atmospheres: Searching for insights into their past and present histories

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#### Invited review: Infrared spectroscopy of planetary atmospheres : 1 Searching for insights into their past and present histories 2 3 4 Thérèse Encrenaz 5 LESIA, Observatoire de Paris - PSL, CNRS, France 6 7 Gerard P. Kuiper Prize Lecture, AAS-DPS General Assembly, October 2021 8 Submitted to Icarus, November 30<sup>th</sup>, 2021 9 Revised, December 30<sup>th</sup>, 2021 10 11 12 13 14 Corresponding author: 15 16 Thérèse Encrenaz LESIA, Paris Observatory, 5 place Janssen 92195 Meudon, France 17 Therese.encrenaz@obspm.fr 18 19

#### Abstract

22 This article reflects my personal experience and illustrates some developments in the field of 23 planetary spectroscopy, achieved within the Planetology Group of Paris Observatory over the 24 past fifty years. Over these decades, planetary spectroscopy has led to the identification of 25 minor atmospheric species with mixing ratios as low as a few parts per billion (ppbv). Since 26 the early 1970s, we have been measuring infrared spectra of giant planets and their satellites, 27 in order to search for new minor species and to determine elemental and isotopic ratios, for 28 better constraining their formation and evolution processes. In particular, observations of 29 Jupiter at the time of the Shoemaker-Levy 9 collision, in 1994, allowed us to monitor the 30 thermal sequence and the formation of stratospheric water vapor; a few years later, spectra taken by the Infrared Space Observatory confirmed the presence of an external oxygen source 31 32 in all giant planets and Titan. More recently, the advent of bi-dimensional infrared arrays has 33 allowed us to map the distribution of these components over planetary disks, and to use them 34 as tracers of dynamical and photochemical processes. In complement to in-orbit observations 35 recorded by space missions, high-resolution spectral mapping from the ground allows us to 36 obtain instantaneous global maps of the planets, and thus to trace transient phenomena or 37 temporal variations of minor atmospheric species over short and long timescales. Over the 38 past twenty years, we have been monitoring the behavior of minor atmospheric species on 39 Mars and Venus, using the TEXES (Texas Echelon Cross Echelle Spectrograph) at the NASA IRTF 40 (InfraRed Telescope Facility) at Maunakea Observatory, and the EXES (Echelon Cross Echelle 41 Spectrograph) aboard SOFIA (Stratospheric Observatory For Infrared Astronomy): H<sub>2</sub>O<sub>2</sub> and 42 H<sub>2</sub>O on Mars, D/H on Mars, SO<sub>2</sub> and H<sub>2</sub>O at the cloud top of Venus. In addition, in an attempt 43 to build a 3-D image of the sulfur and water cycles on Venus, we have obtained maps of SO, 44 SO<sub>2</sub> and HDO in its upper mesosphere (about 20 km above the cloud top) using the ALMA 45 (Atacama Large Millimeter/submillimeter Array) facility in Chile. These observations are 46 presented and discussed in the light of global dynamical and photochemical models, and 47 interpreted in the context of the past and present history of these planets. In the conclusion, 48 I present the perspectives of this work for the forthcoming development of exoplanetary 49 spectroscopy. 50

51 **Key words:** Atmospheres, evolution; Jupiter, atmosphere; Mars, atmosphere; Venus, 32 atmosphere

#### 54 1. Introduction

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Infrared spectroscopy of planets is ideally suited for studying the composition of their neutral atmospheres. Indeed, the infrared range is the region where the strongest spectral signatures of neutral molecules are found, in the form of rotational bands or vibration-rotation bands, while transitions associated with radicals, ions or atoms, corresponding to higher energies, are found in the visible and ultraviolet range. In addition, the thermal emission of solar-system planets is maximum in the infrared range.

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63 Considerable progress has been achieved in the domain of planetary spectroscopy over the 64 five past decades. This is illustrated, in the case of Jupiter, by a comparison between an 65 infrared spectrum recorded by F. C. Gillett and his team in 1969, and the spectrum recorded 66 by the Short Wavelength Spectrograph (SWS) of the Infrared Space Observatory (ISO: 67 www.cosmos.esa.int/web/iso) in 1996 (Figure 1). In the 1960s, only ground-based 68 spectrophotometric measurements could be obtained in the thermal infrared range, with the 69 use of circular variable filters providing a spectral resolving power of about 50.

70

71 The Jupiter spectrum recorded in 1969 already reveals important features. As for all solar 72 system objects, the spectrum shows two components, the solar reflected radiation which 73 peaks at 0.5 µm, and the thermal emission of the planet that dominates in the mid-infrared 74 range. In the first case, the signatures of atmospheric components appear in absorption in 75 front of the solar continuum. In the second case, they appear either in emission or in 76 absorption, depending on the region probed: emission lines appear in the stratosphere (where 77 the temperature increases with altitude) while tropospheric lines (formed in a region below 78 where the temperature decreases as the altitude increases) are in absorption (see e.g. Hanel 79 et al. 2003). The spectrum recorded by Gillett et al. (1969) shows a strong thermal continuum 80 in the 5- $\mu$ m window where the lower troposphere of Jupiter can be probed, and an emission 81 at 7.7  $\mu$ m attributed to methane, showing evidence for a temperature inversion above the 82 tropopause. However, no specific identification of individual molecule could be made, as such 83 an assessment requires a higher spectral resolution. This analysis became possible with the 84 development of Fourier-Transform spectrometers and grating spectrometers. The Jupiter 85 spectrum recorded in 1996 by ISO-SWS corresponds to a spectral resolving power of about 86 2000, sufficient to resolve the molecular bands, either in absorption (in the troposphere, 87 around 5  $\mu$ m and 10  $\mu$ m) or in emission (in the stratosphere, around 7.7  $\mu$ m and 12  $\mu$ m).



Figure 1. Top: The spectrum of Jupiter recorded from the ground by Gillett et al. (1969), with a
spectral resolving power of 50. The black curves indicate the reflected solar component (below
8 μm), and blackbodies at 230 K (around 5 μm) and. 125 K (above 8 μm). Bottom: the same
spectrum recorded by the Short Wavelength Spectrometer of the Infrared Space Observatory
in 1996 with a spectral resolving power of about 2000. Sources: Gillett et al. 1969 (top),
Encrenaz et al. 2003 (bottom).

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97 In addition to the intrinsic information about Jupiter, the spectrum recorded by Gillett et al. 98 (1969) is interesting, as it more or less corresponds to what we can observe today on the 99 brightest exoplanets, mostly with the use of the Spitzer telescope. Its comparison with higher-100 resolution spectra from ISO thus allows us to imagine the enhanced quality of spectra that we 101 may hope to get for exoplanets in the thermal infrared in the future.

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103 The purpose of this article is to present a personal and biased illustration of some 104 developments that at Paris have been achieved Observatory 105 (www.observatoiredeparis.psl.eu) in the field of planetary spectroscopy over the past fifty 106 years, following the progress realized in the development of infrared detectors, spectrometers 107 and telescopes. The first part (Section 2) describes giant planet observations, starting with the 108 spectrum of Jupiter in the 10-µm window (2.1), the monitoring of the SL9-Jupiter collision in 109 1994 (2.2), and the ISO observations in the late 1990s (2.3), followed by ground-based 110 observations of Uranus and Neptune in the 2000s at the InfraRed Telescope Facility (IRTF) at 111 the Maunakea Observatories (www.maunakeaobservatories.org) and the Very Large 112 Telescope (VLT: www.eso.org/public/teles-instr/paranal-observatory/) in Cerro Paranal, Chile 113 (2.4 and 2.5), leading us from disk-integrated spectra to spectral mapping. The second part of 114 this article (Section 3) is mostly devoted to the spectral mapping of Mars (3.1) and Venus (3.2), 115 achieved with the TEXES thermal imaging spectrometer at the IRTF over the past twenty years. 116 Section 4 presents a brief review of planetary spectroscopy's main contributions, a discussion 117 about some remaining open questions, and finally some perspectives about exoplanetary 118 spectroscopy. 119

120

#### 121 2. Giant planets

Until the beginning of the XXI<sup>st</sup> century, in most cases, infrared planetary spectroscopy was 122 achieved with single detectors, and disk-integrated spectra of planets were recorded. In the case 123 124 of the giant planets, infrared spectroscopy had two main objectives: (1) identify minor species 125 to better understand the physico-chemical processes at work in these atmospheres, and (2) determine elemental and isotopic abundance ratios in the giant planets to better constrain 126 127 their formation scenario. Indeed, in the 1970s and 1980s, two possible formation models 128 were proposed: (1) the direct collapse model, in which the giant planets are formed by 129 gravitational instability from the direct collapse of matter within the protosolar disk (Cameron 130 1973, 1988; Cameron & Pine 1973), and (2) the core-accretion model, in which the protosolar 131 matter is accreted around a solid core, mostly made of ice, whose mass is in the range of 10-132 15 terrestrial masses (Mizuno, 1980; Pollack et al. 1996). Elemental and isotopic ratios are an excellent diagnostic to discriminate between these two models; indeed, no enrichment versus 133 134 hydrogen with respect to the solar nebula is expected in the first case, while an enrichment is expected in the second case, due to the high content of the heavy elements in the ices. The larger 135 the mass fraction of the icy core in the giant planet, the stronger the enrichment is (in the case 136 137 of exoplanets, we now use the term of "metallicity", by analogy with the stars). Assuming a 138 core mass of 10 terrestrial masses, a simple calculation (Owen & Encrenaz 2006) shows that the expected enrichment is about 3-4 for Jupiter (core mass fraction of 3%), 7-9 for Saturn (core 139 140 mass fraction of 10%) and above 30 for Uranus and Neptune (core mass fraction > 50%); the uncertainty reflects a variation with time in the choice of the elemental abundance values in the 141 142 Sun. However, it must be kept in mind that there are underlying assumptions in this calculation; in particular, the volatile elements in the ices are assumed to be homogeneously mixed with the 143 144 external layers of the planet during the collapse and post-collapse phases of the planet's 145 formation. As shown by the results of the Juno probe regarding the internal structure of 146 Jupiter (Guillot 2019), the assumption of an initial icy core of ten terrestrial masses is clearly 147 oversimplified. Nevertheless, possibly as a result of a fortuitous coincidence, the enrichments 148 measured in the giant planets happen to be consistent with the predictions of this simple 149 model (Owen & Encrenaz 2006).

150 Among the elemental ratios, the easiest one to be measured with infrared spectroscopy is 151 C/H, which can be directly inferred from the CH<sub>4</sub>/H<sub>2</sub> ratio. In the case of Jupiter and Saturn, methane is uniformly distributed with altitude and over the disks, and its spectral signatures 152 153 are present everywhere over the infrared spectrum. Another key element is the D/H ratio. 154 Deuterium was initially formed in the Big Bang and is continuously destroyed in stars through 155 stellar nucleosynthesis. In the case of the gaseous giants, mostly made of their protosolar gas, 156 the D/H ratio should reflect its value in the protosolar nebula (expected to be intermediate between the primordial value and its present value, measured in the local interstellar 157 158 medium). In contrast, D/H is expected to be enriched in the icy giants, because, as observed 159 in the interstellar medium as well as in the laboratory, deuterated species are known to be 160 enriched in ices, as an effect of ion-molecule reactions. Thus, the D/H ratio can be used as an 161 indicator of the formation temperature: the larger the ratio, the colder is the formation 162 environment; in particular, the measurement of D/H in water has been widely used for 163 constraining the origin of comets (Encrenaz 2008). In the giant planets, D/H can be measured 164 either from HD/H<sub>2</sub> or CH<sub>3</sub>D/CH<sub>4</sub> and we attempted both methods in the case of Jupiter.

165

In the 1970s, HD transitions had been detected in the reflected solar component of the Jupiter
 spectrum, but their interpretation was difficult, due to the effect of scattering in the radiative
 transfer. With Toby Owen (SUNY, Stony Brook) and Michel Combes (Paris Observatory), we

- 169 used an indirect method to minimize as much as possible the scattering effects. We selected 170 weak transitions of HD and CH<sub>4</sub> of comparable intensities, close in wavelength, and we directly 171 estimated the HD/CH<sub>4</sub> ratio (hence the D/C ratio) from the intensity ratio of these two 172 transitions. Using an independent measurement of C/H, we inferred a D/H ratio of about 2.3 173 +/-1.1 10<sup>-5</sup> (Combes et al. 1978; Combes and Encrenaz 1979).
- 174

Twenty years later, the use of ISO made possible the direct measurement of HD in the far infrared range on the four giant planets (see below, 2.3), confirming our earlier estimate. In the late 1990s, the mass spectrometry in-situ measurements of the Galileo probe definitely favored the core-accretion model (Owen et al. 1999).

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180 <u>2.1 Jupiter in the 10µm window</u>

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182 In the early 1970s, there was no existing instrument capable of measuring the mid-infrared 183 spectrum of Jupiter with enough spectral resolution to separate the molecular components 184 of the planet's atmosphere. In contrast, Pierre Connes (Institut d'Optique, Orsay) and his team 185 had been very successful in building a Fourier Transform interferometer for measuring the 186 near-infrared spectrum of all bright planets (Connes et al. 1969) with a spectral resolution of 187 0.25 cm<sup>-1</sup> (corresponding to a resolving power R higher than 15000). This may have influenced 188 the choice that Michel Combes made to build a Fourier Transform interferometer to observe 189 Jupiter in the 10- $\mu$ m atmospheric window, previously observed by Gillett et al. (1969) at low 190 resolution (Figure 1). The instrument built at Paris Observatory, however, was much less 191 ambitious than the one built by Connes et al. (1969), as its highest spectral resolving power 192 was only about 1000. The spectrum of Jupiter (Figure 2) between 10 and 13 µm was recorded 193 with this instrument at the 2.2-m telescope of the University of Hawaii in 1977. It allowed the 194 observation of stratospheric C<sub>2</sub>H<sub>2</sub> and C<sub>2</sub>H<sub>6</sub> - already detected by Ridgway et al (1976) - in 195 emission, as well as tropospheric NH<sub>3</sub> and <sup>15</sup>NH<sub>3</sub> in absorption. Figure 2 also shows, for 196 comparison, the same spectrum of Jupiter as observed by ISO (also shown in Figure 1 with a 197 different scale). The advantage of ISO was, of course, to cover the entire infrared range, as 198 will be discussed below (Section 2.3). Using the 8-9 µm part of the spectrum, we were able to 199 set constraints on the D/H ratio, in agreement with our estimate derived from HD/CH<sub>4</sub> 200 (Encrenaz et al. 1980).



Figure 2. Left: the spectrum of Jupiter, recorded between 10 and 13 μm at Maunakea
 Observatory, using a Fourier Transform interferometer (Encrenaz et al. 1978). Right: the
 spectrum of Jupiter in the same spectral range, recorded by ISO-SWS (Data from Encrenaz et
 al. 1996).

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208 The spectra shown in Figures 1 and 2 were acquired with single detectors which recorded the 209 whole disk of Jupiter, or (in the case of ISO-SWS) the central part of the disk. Starting in the 210 1990s, spectral mapping became possible with the use of detector arrays, and later with bidimensional detectors. Some were used on ground-based facilities, like the 3.5-m New 211 212 Technology Telescope (NTT) at the European Southern Observatory (ESO: www.eso.org) in La 213 Silla, Chile; some others were flown on space missions, like NIMS (Near Infrared Mapping 214 Spectrometer) aboard the Galileo mission. These two means were used at the time of the 215 collision of comet Shoemaker-Levy 9 with Jupiter in 1994; these observations are described 216 hereafter.

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219

218 <u>2.2 The collision of Shoemaker-Levy 9 with Jupiter</u>

220 Discovered in April 1993, the comet Shoemaker-Levy 9 immediately appeared to be an 221 exceptional object, having disrupted during its previous close flyby of Jupiter in July 1992. The 222 "ring of pearls", made of about twenty individual fragments (named, in chronological order, 223 from A to W), was orbiting around Jupiter and its collision with the planet was foreseen in July 224 1994 (Chodas & Yeomans 1996). This discovery led to an exceptional observing campaign 225 including many observatories, from ground and space, covering the whole spectral range, 226 from high energies to radio wavelengths. The comet crashed in the atmosphere of Jupiter at 227 the 44°S latitude, in a series of consecutive impacts which spread in longitude, as a result of 228 Jupiter's rotation, between July 16 and July 23, 1994. This rare event provided us with a unique 229 opportunity to monitor remotely in real time the behavior of a planetary atmosphere 230 subjected to a major meteoritic impact (Noll et al. 1996).

- 232 2.2.1 Temperature monitoring at ESO/NTT
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234 Our observing program used the NTT telescope at ESO (La Silla, Chile), equipped with a bi-235 dimensional detector operating in the near infrared. Our plan was to monitor the behavior of a transition of  $H_{3^+}$ , at 3.53  $\mu$ m; this ion had been detected in the upper Jovian stratosphere a 236 237 few years earlier (Drossart et al. 1989), and is a tracer of the physical conditions of this region. 238 At the time of the H impact (1994 July 18, 19:25:53 UT), the weather was clear and we were 239 ready for monitoring the line emission on the screen. Because of the geometry of the event, 240 the crash was occurring on the hidden side of the planet, and ground-based observers had to 241 wait for fifteen minutes for the impact site to appear at the limb and become visible. At the 242 exact expected time, a strong sinusoidal signal appeared on the screen, about twenty times 243 stronger than the  $H_{3^+}$  emission (Figure 3); it faded slowly and disappeared after about 30 244 minutes. This unexpected signal was indeed from Jupiter: it came from the thermal emission of the CH<sub>4</sub> molecules at high J-levels, usually too faint to be detected. From the observations, 245 246 the temperature of the emitting region could be monitored (from 750 K at the beginning to 247 300 K thirty minutes later; Encrenaz et al. 1995). In addition, enhanced H<sub>3</sub><sup>+</sup> emissions were 248 observed at 44N and 44S latitude, similarly distributed in longitude, showing evidence for a 249 northern counterpart of the impact phenomenon (Schulz et al. 1995). This story tells us that, 250 in any case, we have to be prepared for the unexpected...

251



**Figure 3.** Emission spectrum of methane around 3.53  $\mu$ m at the time of the H impact. The  $v_3$ band peaks at 3.3  $\mu$ m and its high-J components are usually not observed in the spectrum of Jupiter. At the time of the H impact, the temperature increased enough for these components to be observable. Monitoring them with time gives us a record of the temperature decrease in the Jovian stratosphere after the H impact. Left: first spectrum observed (solid line) compared with a model (dashed line; T = 750 K, P = 0.01 mb); right: spectrum observed 6 minutes later (solid line) compared with a model (T = 640 K, P = 0.01 mb). Source: Encrenaz et al. 1995.

#### 2.2.2 Water monitoring with Galileo/NIMS

263 At the time of the SL9 collision, the Galileo probe (www.jpl.nasa.gov/missions/galileo) was on 264 its way to Jupiter. It was in a unique position to observe the crash, as its instruments were in 265 direct view of the impact. Aboard Galileo, NIMS was an imaging spectrometer operating 266 between 0.7 and 5.2 µm. In the fixed grating mode used for the SL9 observations, 267 measurements are recorded simultaneously at 17 fixed wavelengths (Carlson et al. 1995). 268 Data of the impact were limited to the  $1.8 - 4.4 \,\mu\text{m}$  range, due to contamination by the 269 reflected solar light and the thermal emission outside this spectral range. Most of the Jovian 270 disk was included in the NIMS pixel.

271

272 Observations took place at the time of the splash phases following the G and R impacts, 273 between 400 and 600 seconds after the fireball event (Encrenaz et al. 1997). In spite of the 274 discrete wavelength sampling, emissions due to water and methane were easily identified, 275 around 2.7 and 3.3  $\mu$ m respectively. The CH<sub>4</sub> emission was used to infer the temperature 276 evolution of the sequence, and the water content was inferred from the 2.7-µm emission. 277 Figure 4 summarizes the results for the G and R events. In both cases, a maximum temperature 278 of 1000 K was inferred, nine minutes after the fireball event, at pressure levels of 15 µbar and 279 10 µbar for G and R respectively. The inferred H<sub>2</sub>O mass was (0.6 – 2.9)  $10^{12}$  g for the G impact, 280 and about ten times less for the R impact; these results were found in agreement with KAO 281 measurements obtained at 6.2 and 7.7  $\mu$ m (Sprague et al. 1996). The data also show that the 282 water formation started in the splash phase, about 7 and 9 minutes after the explosion, for G 283 and R respectively; it is also interesting to see that no carbon dioxide was detected. This 284 observation would not have been possible from ground-based telescopes. 285



286

Figure 4. Time sequence of the NIMS spectra during the splash phase following the G impact
(left) and the R impact (right) at the time of the SL9 collision with Jupiter. The NIMS scans are
summed over 1-minute intervals between 6 and 10 minutes after the fireball. The signal is
increasing with time from bottom to top. Source: Encrenaz et al. 1997.

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The abundance of water inferred from our measurements has to be compared with the one of CO, the most abundant molecule, which was monitored in the millimeter range (2.5 10<sup>14</sup> g, Lellouch et al. 1995). The relatively low water content suggests that the penetration level of

- the impactors was above the tropospheric water cloud of Jupiter. The diameter of the comet,
  before its fractionation, was estimated to be about two kilometers (Solem, 1995).
- 297
- Many other new species were discovered at the level of  $10^{11} 10^{12}$  g, in particular HCN which remained visible for many years. In 1996, it appeared that it was also the case of H<sub>2</sub>O, which was detected in the stratosphere of Jupiter by the ISO satellite (see below, Section 2.3).
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#### 302 <u>2.3 The giant planets as seen by ISO</u>

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304 The Infrared Space Observatory (ISO: www.cosmos.esa.int/web/iso) was a European satellite, 305 launched by the European Space Agency (ESA: www.esa.int) in 1995, in operation until 1998. 306 It consisted of a helium-cooled 60-cm telescope and four infrared instruments (a camera, a 307 photometer and two spectrometers). It was designed for the study of the cold Universe, in 308 particular the interstellar medium and protoplanetary disks. As a Mission Scientist aboard this 309 mission, I had the opportunity of coordinating a solar system observing program which was 310 especially fruitful in the case of giant planets and comets, in particular thanks to the Short 311 Wavelength Spectrometer (SWS), whose Principal Investigator was Thijs de Graauw (SRON, 312 Netherlands; de Graauw et al. 1996).

313

Figure 5 shows the spectra of Jupiter, Saturn and Neptune, as seen by ISO, between 2.3 and 16 μm (Encrenaz et al. 2003). The reflected component prevails below 4 μm, and the thermal

emission dominates around 5  $\mu$ m and above 6  $\mu$ m. In the case of Saturn, the 5- $\mu$ m emission 317 is a mixture of reflected sunlight and thermal emission.



318

**Figure 5.** The spectra of Jupiter, Saturn and Neptune (log scale) between 2 and 16  $\mu$ m, as

320 observed by ISO-SWS. The spectra of Saturn and Neptune are disk-integrated, while the Jupiter 321 spectrum covers the central part of the disk. The spectrum of Uranus is not shown; it would

322 appear only in the  $C_2H_2$  band at 13.7  $\mu$ m, with an intensity equal to a third of the Neptune

323 value. Source: Encrenaz et al. (2003a).

325 Regarding the giant planets, the main discoveries of ISO were the evidence for an external 326 source of oxygen, and the measurement of deuterated hydrogen HD. In addition, new minor 327 species were detected: C<sub>3</sub>H<sub>4</sub> and C<sub>4</sub>H<sub>2</sub> in Jupiter (Fouchet et al. 2000) and Saturn (de Graauw 328 et al. 1997) with SWS, CH<sub>3</sub> in Saturn (Bézard et al. 1998) and Neptune (Bézard et al. 1999), and 329 C<sub>6</sub>H<sub>6</sub> outside the auroral regions of Jupiter and Saturn (Bézard et al. 2001). C<sub>2</sub>H<sub>4</sub> was detected 330 on Neptune, with the ISOPHOT instrument (Schulz et al. 1999). ISO also allowed for 331 significantly improved constraints on the abundances of all the hydrocarbon photochemical 332 products on the giant planets, including information on their vertical distributions (e.g. C<sub>2</sub>H<sub>2</sub>, 333 see Moses et al. 2000a, 2005).

334

#### 335 2.3.1 An external oxygen source in the stratospheres of the giant planets

336 337 In spite of the detection of water in Jupiter's stratosphere at the time of the SL9 collision, 338 nobody expected to detect this molecule two years later. ISO detected H<sub>2</sub>O in Jupiter, but also 339 in the stratospheres of the other giant planets and Titan (Feuchtgruber et al. 1997; Coustenis 340 et al. 1998; Figure 6). Several transitions appeared in emission between 20 and 40  $\mu$ m, 341 corresponding, within a factor of 10, to mixing ratios in the range of a few 10<sup>-9</sup>. In addition, 342  $CO_2$  was detected at 15  $\mu$ m – also by ISO-SWS – on Jupiter, Saturn and Neptune (de Graauw 343 et al. 1997; Bézard, 2000; Lellouch et al. 2002), confirming the existence of an oxygen reservoir 344 in the stratospheres of these planets. Because of the presence of the tropopause which acts, 345 for all objects, as a cold trap for H<sub>2</sub>O, the origin of this reservoir had to be external. Two origins 346 were proposed: (1) an interplanetary flux of meteoritic particles, as comets or smaller debris 347 (Moses et al. 2000b) and (2) a contribution coming from the rings and satellites. In the case of 348 Jupiter, a subsequent map obtained with the PACS imaging spectrometer of the Herschel 349 submillimeter satellite has shown that the SL9 collision was indeed responsible for the 350 distribution of stratospheric water as a function of latitude (Cavalié et al. 2013). What could 351 be the origin of H<sub>2</sub>O in the other giant planets? In the case of Saturn, one would expect the 352 influence of rings and satellites to be more important, while the cometary source could prevail 353 on Uranus and Neptune.



**Figure 6.** Detection of the  $H_2O$  emission line at 39.4  $\mu$ m on the four giant planets and Titan. In the case of Jupiter, the observation was made using the Fabry-Perot mode of SWS, with a spectral resolving power of 31000; in all other cases, the grating mode was used (R = 2000). Source: Feuchtgruber et al. (1999).

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#### 361 *2.3.2 Deuterium in the giant planets*

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363 Before the launch of ISO, it was not possible to obtain a measurement of D/H in the four giant 364 planets using a single instrumental method. This became possible with ISO, using the R2 and R3 transitions of HD, at 37.7 µm and 28.5 µm respectively. From these transitions, D/H ratios 365 of 2.25 +/- 0.35 x  $10^{-5}$  and 1.7 (+0.75, -0.45) x  $10^{-5}$  were retrieved for Jupiter and Saturn 366 respectively (Lellouch et al. 2001), and values of 5.5 x 10<sup>-5</sup> and 6.5 x 10<sup>-5</sup> were inferred for 367 Uranus and Neptune respectively (Feuchtgruber et al. 1999). The values of Uranus and 368 369 Neptune were later reduced to  $4.4 + - 0.4 \times 10^{-5}$  and  $4.1 + - 0.4 \times 10^{-5}$ , using the R1 HD line at 370 56 µm with the PACS instrument aboard the Herschel submillimeter satellite (Feuchtgruber et 371 al 2013).

372

373 Two comments can be made about these results. First, the D/H ratio in Saturn seems to be 374 slightly lower than the Jupiter value, which is not expected by the models, since the mass 375 fraction of the initial core, in the case of Saturn, is expected to be larger than in the case of 376 Jupiter. We have presently no explanation for this discrepancy. The second comment is that 377 D/H in the giant planets follows the trend expected in the case of the core-accretion model, 378 with larger values in Uranus and Neptune, as compared with Jupiter and Saturn. This 379 measurement, as the one of C/H in the four giants, comes as a support to the core-accretion 380 model of the giant planets (Owen and Encrenaz 2006).

382 2.4 Uranus

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381

As shown in Figure 5, the 5-μm region, free of methane absorption, is an atmospheric window
 where the deep troposphere of Jupiter and Saturn can be probed. In the case of Jupiter, the

386 radiation is fully of thermal origin while, in the case of Saturn, there is a significant contribution 387 of solar reflected radiation. In the case of Uranus and Neptune, this atmospheric window 388 should be fully of solar reflected origin, since the temperature of the probed region is 389 expected to be around 120 K, and the blackbody radiation corresponding to such a low 390 temperature is too weak to be detectable. As shown in Figure 5, ISO was not able to detect 391 this component in Uranus nor in Neptune, so, in 2000, we began an observing program at 392 IRTF, using the SpeX imaging spectrometer at the IRTF (Rayner et al. 2003). Our objective was 393 to search for tropospheric CO and (possibly) PH<sub>3</sub> in Uranus (easier to observe than Neptune); 394 both species were expected to be in absorption, as in the case of the two other giant planets, 395 and their detection would bring an important diagnostic for constraining the planet's origin 396 and evolution.

397

398 Our observations led to a surprise. We did not detect the continuum at 4.5 µm, but, instead, 399 we observed a set of  $H_3^+$  emission lines, between 3.5 and 4.7  $\mu$ m. The  $H_3^+$  ion had been 400 discovered in Jupiter ten years earlier (Drossart et al. 1990), and later on Saturn and Uranus (Trafton et al. 1993), so it was not a new discovery, but this observation allowed us to infer 401 402 the rotational temperature of the emitting region, in the upper stratosphere, and an estimate 403 of the H<sub>3</sub><sup>+</sup> abundance in Uranus. Combined with other data, our results are consistent with a 404 correlation between the H<sub>3</sub><sup>+</sup> emission and the solar radiation, modulated by the solar cycle 405 (Encrenaz et al. 2003b). However, the 4.5- $\mu$ m continuum signal was too weak to be analyzed. 406 In 2003, we used a more sensitive instrument, the ISAAC imaging spectrometer at the Very 407 Large Telescope. Again, the result was surprising. The weak continuum at 4.5 µm was indeed 408 detected, but the spectrum appeared to be dominated by a series of CO fluorescence lines (Encrenaz et al. 2004a; Figure 7). Such a fluorescence phenomenon is not common in giant 409 410 planets' atmospheres; it has been observed only in CH<sub>4</sub> emission lines at 3.3 µm in the case of 411 Jupiter and Saturn, at high stratospheric levels (Drossart et al. 1999). From this observation, a mixing ratio of 3 x 10<sup>-8</sup> was tentatively inferred in the stratosphere of Uranus. From the 412 413 absence of CO and PH<sub>3</sub> features in the weak continuum, upper limits of 2 x 10<sup>-8</sup> and 10<sup>-6</sup> were 414 respectively derived for the tropospheric mixing ratios of these species, suggesting an external 415 origin for the stratospheric CO. The detection of CO was later confirmed with the Herschel 416 submillimeter satellite, but the CO abundance was found to be about 8 x  $10^{-9}$ , i.e. significantly 417 lower than the earlier estimate (Teanby and Irwin, 2013; Cavalié et al; 2014).

418



420 Figure 7. Left: The spectrum of Uranus between 3.6 and 4.1 µm, observed with SpeX at IRTF 421 (Maunakea Observatories). All emission lines are due to  $H_3^+$ . Solid line: September 2000; 422 dashed line: September 2001. The inferred rotational temperature is 560 K in 2000, and 640 K 423 in 2001. Source: Encrenaz et al. (2003b). Right: the spectrum of Uranus (black line) recorded 424 with ISAAC at the VLT (ESO, Paranal) between 4.6 and 5.0  $\mu$ m. H<sub>3</sub><sup>+</sup> emission lines are present 425 at 4.67 and 4.87 µm. CO emission appears between 4.6 and 4.8 µm. Red line: best fit model 426 including CO fluorescence. Green line: model including CO thermal emission. Source: Encrenaz 427 et al. (2004a).

428

430

#### 429 <u>2.5 Neptune</u>

Although SpeX at the IRTF and ISAAC at the VLT were using bi-dimensional detectors, our
observations of Uranus were not spatially resolved and integrated most of the planetary disk.
A few years later, another imaging spectrophotometer, also at the VLT, allowed us to obtain
the first thermal images of Uranus and Neptune.

435

436 In September 2006, we used VISIR (VLT Imager and Spectrometer for mid Infrared; Lagage et 437 al. 2004) in the imaging mode to map the disks of Uranus and Neptune in different filters to 438 probe different atmospheric levels. In particular, the 18.7 µm filter probes the tropopause 439 where the temperature is minimum. At this wavelength, the spectrum is dominated by 440 hydrogen collision-induced absorption; since H<sub>2</sub> is uniformly distributed over the planet, a 441 map at 18.7  $\mu$ m shows the temperature distribution at the tropopause. The result was very 442 surprising (Figure 8). While the map of Uranus showed moderate variations as a function of 443 latitude (Orton et al. 2009), the map of Neptune exhibited a strong maximum near the south 444 pole, corresponding to a maximum temperature of 62 K, 7 K above the disk-averaged 445 temperature (Orton et al. 2007). This temperature increase has been tentatively interpreted 446 as the consequence of the maximum solar irradiation received at the south pole over several 447 decades, as southern solstice occurred in July 2005. A consequence of this anomaly is the 448 release of methane, elsewhere cold-trapped in the troposphere, to escape in the 449 stratosphere; this result provides us with an explanation for the high abundances of 450 stratospheric methane reported previously (Baines & Hammel 1994; Baines et al. 1995), 451 significantly higher than those corresponding to the saturation vapor pressure at the average 452 minimum temperature (55 K). Another possible explanation is adiabatic heating from 453 downwelling or other dynamical phenomena (Karkoschka and Tomasko, 2011; de Pater et al. 454 2014; Lellouch et al. 2015).



Figure 8. Left: images of Uranus (top) and Neptune (bottom) recorded at 18.7 μm with VISIR
at the VLT (ESO, Paranal) in September 2006. In both cases, the north pole is at the top of the
figure. Source: Orton et al. (2007, 2015). Right: the disk-integrated spectrum of Neptune,
recorded with the TEXES imaging spectrometer at IRTF (Maunakea Observatories) in June
2003, with a spectral resolving power of 80000. Top: the observed spectrum; bottom: the
synthetic spectrum. Source: personal archives.

463

464 In the early 2000s, another thermal imaging spectrometer became available. TEXES (Texas 465 Echelon Cross-Echelle Spectrograph), designed and built by John Lacy (University of Texas), 466 and operated by John Lacy, Matt Richter and Tommy Greathouse, was mounted at the IRTF 467 (Lacy et al. 2002). This instrument, operating between 5 and 25 µm, combines imaging 468 capabilities with different spectroscopic modes, including a high spectral resolution mode (R 469 > 80000). With such a high resolving power, this instrument appeared to be perfectly suited 470 for planetary mapping of minor atmospheric species. In the case of Neptune, we used TEXES 471 in 2003 to observe the stratospheric emission band of ethane  $(C_2H_2)$ . Figure 8 (right) shows 472 the disk-integrated spectrum of Neptune and illustrates the exceptional quality of the data. 473 These data also provided us with the resolved spectrum of acetylene  $C_2H_4$  on Neptune. In 474 2007, TEXES was used at the 8-m GEMINI telescope, also at the Maunakea Observatories, to 475 map the disk of Neptune and to retrieve information on the distribution of temperature and 476 various hydrocarbons (Greathouse et al. 2011).

477 478

#### 479 **3. Terrestrial planets**

480

481 It is well established that the atmospheres of the terrestrial planets – Venus, the Earth, and 482 Mars – exhibit, at their surfaces, extremely different conditions. At 0.7 au from the Sun, the 483 surface pressure of Venus (93 bars) is almost a hundred times the terrestrial value, and its 484 temperature is close to 730 K; in contrast, on the other side of the Earth, at 1.52 au from the 485 Sun, the surface pressure of Mars, averaged over the season, is less than one hundredth of a 486 bar, and its mean temperature is around 220 K. What is less well known are the reasons for 487 such extreme conditions. It is generally believed that the three planets were formed in a 488 common environment, relatively close to the Sun, from the accretion of solid refractory 489 particles that led to the formation of dense, rocky planets. This scenario, known as the core 490 accretion model of planetary formation, is also supported by the fact that both Venus and 491 Mars have a  $CO_2$ -dominated atmosphere with a few percent of  $N_2$ , which was also probably 492 the case of the primordial atmosphere of the Earth. The presence of massive amounts of water 493 on Earth today suggests that, originally, water was equally present on Venus and Mars, at their 494 surface and in their atmospheres; indeed, we find signatures of past water on both planets, 495 especially through the D/H ratio, as will be discussed below. Why and how did the three 496 planets evolve, from relatively comparable initial conditions, to the diverging conditions that 497 we observe today? Answering this question is a major challenge in today's planetary research. 498

499 In this section, I will present observations of Mars and Venus that have been performed from 500 the ground over the past two decades, using the TEXES instrument at IRTF (Maunakea 501 Observatories). As described above (Section 2.4), TEXES is an imaging spectrometer operating 502 in the thermal range, which allows us to obtain instantaneous maps of minor species over 503 planetary disks. Thanks to the high resolving power of TEXES (up to 80000), it is possible to 504 monitor trace species with vertical mixing ratios in the range of a few ppbs, and to follow their 505 day/night, latitudinal or seasonal variations, as well as transient phenomena. These 506 observations can provide a precious complement to in-orbit space observations, which are 507 often limited to specific local times and restricted planetary areas.

508

509 In the next sub-sections, I will describe long-term observing campaigns that have been 510 performed to monitor the seasonal behavior of hydrogen peroxide and water vapor on Mars 511 (Section 3.1) and to study the sulfur and water cycles on Venus (Section 3.2). In addition, TEXES 512 observations have been used to derive upper limits of SO<sub>2</sub> on Mars (Encrenaz et al. 2011) and 513 PH<sub>3</sub> on Venus (Encrenaz et al. 2020a).

- 514 515
- 515 <u>3.1 Mars</u> 516

517 With its very tenuous atmosphere, Mars has a "simple" spectrum, fully dominated by CO<sub>2</sub>, 518 both in its reflected sunlight and its thermal emission. Weak signatures of water vapor and 519 carbon monoxide are also present. Until 2004, the only other identified trace component was 520 ozone O<sub>3</sub>. For over two decades, the climate of Mars has been successfully modelled through 521 the use of Global Climate Models (GCMs) derived from those developed for the terrestrial 522 climate (Forget et al. 1999). Generally speaking, these models provide a good fit of the 523 observations, illustrating that the main mechanisms (circulation, condensation, 524 photochemistry) are globally well understood.

525

526 Still, there are pending open questions. One of them raised after the Viking exploration: what 527 was the oxidant responsible for the destruction of organics at the surface of Mars? Hydrogen 528 peroxide (H<sub>2</sub>O<sub>2</sub>) was soon suspected, as its presence was predicted by photochemical models, 529 and was searched for repeatedly until Todd Clancy and his team identified it in the 530 submillimeter range (Clancy et al. 2004). In parallel, using TEXES at IRTF, we have searched for 531 H<sub>2</sub>O<sub>2</sub>, finally detected it, and mapped it on Mars to monitor its seasonal behavior (see below, 532 Section 3.1.1). Another pending question was the water cycle on Mars today, and, even more important for understanding the water evolution, the water content of the early Mars. Bothquestions are addressed below (Sections 3.1.2 and 3.1.3).

535

537

#### 536 3.1.1 Hydrogen peroxide mapping and seasonal behavior

538 At the time of the Viking exploration, the negative results of the biological experiments 539 suggested the presence of a powerful oxidizer, and hydrogen peroxide (H<sub>2</sub>O<sub>2</sub>) was suggested 540 as a possible candidate (Oyama and Berdahl 1977). The presence of small amounts of  $H_2O_2$ 541 (less than 100 ppbv) was actually expected on the basis of photochemical models (Atreya and 542 Gu 1994; Nair et al. 1994; Krasnopolsky 1995). After unsuccessful attempts (Krasnopolsky et 543 al 1997, Encrenaz et al. 2002), the detection of H<sub>2</sub>O<sub>2</sub> on Mars was first announced by Clancy 544 et al. (2004), using the JCMT sub-millimeter telescope at Maunakea Observatories). In parallel 545 we searched again for  $H_2O_2$  and were able to detect and map the molecule (Encrenaz et al. 546 2004b), using the TEXES imaging spectrometer at IRTF.

547 As mentioned above (Section 2.4), TEXES has the advantage of combining a high spectral 548 resolving power and very good imaging capabilities, which makes it a perfect tool for mapping 549 minor atmospheric species. We selected the  $H_2O_2 v_6$  band centered at 1266 cm<sup>-1</sup> (7.9  $\mu$ m), in 550 a region relatively free from telluric contamination, and we chose a small spectral interval, 551 free of telluric absorption, where weak neighboring transitions of both H<sub>2</sub>O<sub>2</sub> and CO<sub>2</sub> are 552 present. In order to estimate the H<sub>2</sub>O<sub>2</sub> mixing ratio over the Martian disk, we simply ratio the  $H_2O_2$  line depth by the  $CO_2$  line depth. This first-order method (actually derived from our 553 554 earlier analyses: Combes and Encrenaz 1979) has the advantage of minimizing the 555 uncertainties associated with surface and atmospheric properties, including thermal profile, 556 dust opacity and airmass factor. This procedure best applies in the case of optically thin lines, 557 but can also be used if the depths of the lines (convolved at the TEXES spectral resolution) are 558 less than a few percent. Indeed, synthetic calculations have shown that in this case, within a few percent, the H<sub>2</sub>O<sub>2</sub> mixing ratio varies linearly with the H<sub>2</sub>O<sub>2</sub>/CO<sub>2</sub> line depth ratio, over a 559 560 wide range of atmospheric parameters (Encrenaz et al. 2008).

561 We started searching for H<sub>2</sub>O<sub>2</sub> in February 2001, when Mars was near northern summer 562 solstice (Ls = 112°). This situation looked favorable as it corresponds to a maximum of water 563 vapor near the north pole and, according to photochemical models, H<sub>2</sub>O and H<sub>2</sub>O<sub>2</sub> are 564 expected to be correlated. Surprisingly,  $H_2O_2$  was not detected and the inferred upper limit (6 565 ppbv) was difficult to reconcile with the models (Encrenaz et al. 2002). Two years later, we 566 detected H<sub>2</sub>O<sub>2</sub> with a maximum mixing ratio of 40 ppbv (Encrenaz et al. 2004). Then, we 567 regularly monitored the H<sub>2</sub>O<sub>2</sub> distribution as a function of latitude and season until 2019. 568 Figure 9 shows an example of the spectrum used to infer the  $H_2O_2$  mixing ratio and two 569 examples of maps corresponding to different seasons, compared with GCM models (Forget et 570 al. 1999). In both cases, the agreement is satisfactory.



**Figure 9.** Left panel, top: the TEXES spectrum (black) recorded in the region of maximum  $H_2O_2$ on June 2003. A model of  $H_2O_2$  absorption is shown for a mixing ratio of 100 ppbv. Left, bottom: an enlargement showing the  $H_2O_2$  and  $CO_2$  lines used to derive the map. Right panel, top: the TEXES map corresponding to the spectrum shown on the left (June 2003), compared with a GCM-IPSL model. Source: Encrenaz et al. (2004). Note that the absolute value of the  $H_2O_2$  vmr was rescaled after 2015 (see text and Figure 10). Right, bottom: a  $H_2O_2$  map obtained for another season (May 2016, Ls = 148°). Source: Encrenaz et al. (2019a).

579 Figure 10 shows the seasonal variations of  $H_2O_2$  recorded between 2001 and 2019. Most of 580 the data are from TEXES, with the exception of two sub-millimeter measurements, the first 581 detection by Clancy et al. (2004) for Ls = 250°, and an upper limit by Hartogh et al. (2010) with 582 the Herschel space observatory (www.esa.int/Science Exploration/Space Science/Herschel) 583 for Ls = 77°. In 2015, a problem arose when it was realized that the  $H_2O_2$  spectroscopic 584 intensities used before, taken from the GEISA-2003 database, were lower than the GEISA-585 2015 database by a factor close to 2. As a result, all TEXES values published prior to 2015 (open squares in Figure 10) had to be recalibrated (black and green squares). The data were 586 587 compared to different models, involving homogeneous chemistry (Lefèvre et al. 2008, 588 Daerden et al. 2018), and heterogeneous chemistry on water ice particles (Lefèvre et al. 2008). 589 Our first conclusion (Encrenaz et al. 2012a) was that the comparison between the data and 590 the models favored heterogeneous chemistry; however, the recalibration of the H<sub>2</sub>O<sub>2</sub> data 591 showed that there is a discrepancy between the data and the models: in most of the cases, 592 the measurements are lower than expected by a factor close to 2, whichever homogeneous 593 or heterogeneous chemistry is considered. We have presently no explanation for this 594 discrepancy.

595 Another surprising result is shown in Figure 10. For several values of the solar longitude (Ls 596 around 80°, 100°, 150° and 207°), the  $H_2O_2$  abundances show different values, depending on 597 the Martian year. In the last case, a possible explanation is that the upper limit of  $H_2O_2$  was 598 obtained in July 2018, at the time of the global dust storm. Models developed by Daerden et 599 al. (2018) were able to account for the absence of  $H_2O_2$  under these conditions (Encrenaz et al. 2019a). But, in the three other cases, we have no explanation for the observed interannualvariations.



602

603 *Figure 10.* Seasonal variations of the  $H_2O_2$  mixing ratio as a function of the solar longitude Ls. 604 Open squares: TEXES data prior to 2015 (using the GEISA-2003 database); black squares: final 605 dataset, including the TEXES recalibrated data and the two sub-millimeter measurements at 606 77° by Hartogh et al. (2010) and 250° by Clancy et al. (2004). Green squares: TEXES data from 607 2016 and 2018. The 2018 upper limit corresponds to the global dust storm; other discrepancies 608 are unexplained. Models: GCM-IPSL, with homogeneous chemistry (red) and heterogeneous 609 chemistry (purple), by Lefèvre et al. (2008); GEM model with homogeneous chemistry (blue) 610 by Daerden et al. (2019). Source: Encrenaz et al. (2019a).

611 In summary, our analysis of the behavior of hydrogen peroxide on Mars reveals that its 612 observed abundances are below the predictions. This discrepancy might suggest the possible 613 presence of another oxidizer, still to be identified, or could be due to another loss process for 614 H<sub>2</sub>O<sub>2</sub>, not taken into account in the current photochemical models.

- 615
- 616 3.1.2 Water vapor on Mars
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618 In June 2003, as we were searching for  $H_2O_2$  around 8  $\mu$ m, we realized that a few weak HDO 619 transitions were present in the spectral interval that we had selected for  $H_2O_2$  (1227–1243 cm<sup>-</sup> 620 <sup>1</sup>). We took advantage of these transitions to obtain water vapor maps, simultaneously with 621 H<sub>2</sub>O<sub>2</sub> maps. It was possible to use HDO as a proxy for H<sub>2</sub>O, because, as shown by Montmessin 622 et al. (2005), the variations of  $H_2O$  over the Martian disk are expected to be much stronger 623 than the variations of D/H (derived from the HDO/H<sub>2</sub>O ratio). Indeed, the water seasonal cycle 624 of Mars has been well known since the Viking experiments (Jakosky and Farmer 1982) and 625 later with the TES instrument aboard the Mars Global Surveyor orbiter (Smith 2002, 2004). 626 Water vapor was found to exhibit a strong maximum at high northern latitudes around the 627 northern summer solstice, and a weaker maximum at the south pole near southern summer 628 solstice, with latitudinal variations by a factor as high as 10. The D/H ratio in the Martian 629 atmosphere was found to be enriched by a factor 5-6 with respect to the VSMOW (Vienna 630 Standard Mean Ocean Water, i.e.  $1.56 \times 10^{-4}$ ) terrestrial value (Owen et al. 1988, Krasnopolsky 631 et al. 1997; see Section 3.1.3), and no large variations over the disk were expected (Fouchet 632 and Lellouch 2000); this was later confirmed by more sophisticated models, showing that 633 seasonal variations of D/H are expected to be less than 10 - 15% for latitudes between 60°N 634 and 60°S (Montmessin et al. 2005).

635

636 Figure 11 shows examples of water vapor maps obtained using the line depth ratio method 637 described above. The first example corresponds to the same dataset as shown in Figure 10 638 (top right). We used a HDO transition at 1240.0 cm<sup>-1</sup>, very close to the H<sub>2</sub>O<sub>2</sub> doublet shown in 639 Figure 10, and we used the CO<sub>2</sub> line at 1241.6 cm<sup>-1</sup> (already used for the  $H_2O_2/CO_2$  line depth 640 ratio) for comparison. The second example is extracted from the May 2008 dataset; in this case, we used another HDO transition at 1236.3 cm<sup>-1</sup>. In both cases, the observed H<sub>2</sub>O maps 641 642 were found to be in agreement with the GCM predictions (Forget et al. 1999), both for the 643 absolute water vapor abundance and for the H<sub>2</sub>O spatial distribution. In both cases, the water 644 vapor content is maximum in the northern hemisphere, as expected during northern spring 645 and early summer (Smith, 2002, 2004).

646

647



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Figure 11. Left: TEXES spectra (black) of two HDO transitions used to retrieve the water vapor 649 650 maps shown on the middle of the figure. Top: Observations of June 2003 (Ls = 112°); the best 651 fit is obtained for a  $H_2O$  mixing ratio of 300 ppmv (red curve; the green and blue curves 652 correspond to 200 and 400 ppmv respectively). Bottom: Observations of May 2008 ( $Ls = 80^{\circ}$ ); 653 the best fit is obtained for  $H_2O = 250 \text{ ppmv}$  (blue curve; the green and red curves correspond 654 to 100 and 500 ppmv respectively). Middle: HDO maps inferred from the spectra (top: June 655 2003; bottom: May 2008). The unit on the maps is the HDO/CO<sub>2</sub> line depth ratio. Right: GCM 656 maps calculated for the conditions of the observations (top: June 2003; bottom: May 2008).

The unit is the H<sub>2</sub>O /CO<sub>2</sub> mixing ratio. Sources: Encrenaz et al. (2005) (top) and Encrenaz et al.
(2010) (bottom).

659

660 3.1.3 D/H mapping on Mars with EXES on SOFIA

661

662 While the seasonal cycle of water vapor on Mars is well understood by the climate models, 663 the very low abundance of water on Mars (as well as on Venus) today remains a puzzle. It is 664 now generally accepted that, as in the case of the Earth, the water abundance in the primitive 665 atmospheres of Mars and Venus was much higher than today; however, its abundance is 666 poorly constrained presently. In the case of Mars, there are many signs (including channel 667 incision, putative northern ocean, deuterium enrichment etc (see e.g. Catling and Kasting 668 2017), leading to estimates of the early water inventory between a hundred and a thousand 669 meters of GEL (global equivalent layer).

670

671 The HDO/H<sub>2</sub>O ratio on Mars is a diagnostic of the early water content of the planet, because, 672 in the upper atmosphere, the H atoms escape more easily than the D atoms. Such a 673 mechanism is especially efficient in the case of Venus, where a D/H enrichment by a factor 674 120 has been inferred from the study of HDO and  $H_2O$  in the lower atmosphere (Bézard et al. 675 1990). In the case of Mars, the enrichment (factor of about 5) is more modest, but still very 676 significant for the very tenuous atmosphere of Mars. An early estimate of D/H ( < 1.2 VSMOW) 677 was derived from SNC meteorites (Usui et al. 2012, Hallis et al. 2012), and a D/H enrichment 678 by a factor 3 was inferred by the Curiosity rover in ancient clays, dated from 3.5 Gy ago 679 (Mahaffy et al. 2015). Still, the measurement of D/H on Mars has led to very different 680 estimates of the early water reservoir (Krasnopolsky 2002, Villanueva et al. 2015), due to the 681 uncertainty in the fractionation factor indicating the efficiency of D escape relative to H 682 escape. Another uncertainty is the determination of the current D/H on Mars, globally 683 averaged over latitudes, longitudes and seasons.

684

685 As a contribution to this study, we have set up an observing program to globally map the D/H 686 ratio over the Martian disk and for different seasons. Another objective was to study the 687 present water seasonal cycle through fractionation due to differential condensation 688 processes. Indeed, deuterium is known to be enriched in the ice phase, through the so-called 689 Vapor Pressure Isotopic Effect (VPIE; Merlivat & Lief, 1967). On Mars, the D/H ratio is expected 690 to be maximum at the north pole at northern summer solstice, just after the content of the 691 northern water ice cap is released in the atmosphere (Montmessin et al. 2005). Our objective 692 was thus to test the VPIE mechanism.

693

694 In order to determine the Martian D/H ratio in water using our line depth ratio method, we 695 needed to use weak neighboring transitions of HDO and H<sub>2</sub>O. which requires that the 696 observations take place above the terrestrial troposphere, so as to eliminate the terrestrial 697 water vapor absorption lines. The EXES (Echelle Cross Echelon Spectrograph) aboard the 698 Stratospheric for Observatory Infrared Astronomy (SOFIA: 699 www.nasa.gov/mission\_pages/SOFIA/overview/index.html), jointly developed by NASA and 700 DLR, gave us this opportunity. EXES, under the PIship of Matt Richter (UC Davis, CA), is directly 701 derived from TEXES, with a slightly enlarged spectral range (4.5 – 28.3  $\mu$ m) and increased 702 spectral resolving power (close to 10<sup>5</sup>). Several flights occurred between 2014 (when EXES 703 became operational) and 2018. Observations were performed around 1360 cm<sup>-1</sup> (7.2  $\mu$ m) 704 where weak transitions of HDO,  $H_2O$  and  $CO_2$  occur. The advantage of the line depth ratio 705 method is that it allows us to derive the HDO/ $H_2O$  directly, without using the  $CO_2$  transitions, 706 which removes uncertainties associated with the thermal structure. Still, we also used the  $CO_2$ 707 transition to infer the HDO and  $H_2O$  mixing ratios and to check they were in agreement with 708 the GCMs (Encrenaz et al. 2016a, 2018, 2019b).

709

710 Figure 12 shows three maps of D/H obtained with EXES in 2014, 2016 and 2018 respectively. 711 In all cases, the D/H maps appear to be remarkably uniform, with a disk-integrated D/H value 712 between 4.0 and 4.4 times the (VSMOW) terrestrial value. The figure also shows an example 713 of disk-integrated spectrum, taken in April 2014; the best fit corresponds to a D/H value of 4.4 714 (+ 1.0, - 0.6) x VSMOW. The values inferred from the two other maps are very similar. Our 715 conclusion is thus that we see no significant change of the disk-integrated value of D/H over 716 the two observed seasons (early northern summer and southern solstice), nor any strong 717 variation of D/H over the Martian disk. Regarding the spatial distribution of D/H, however, we 718 have to keep in mind that the spatial resolution of SOFIA (3 arcsec) is, unfortunately, very 719 poor, so that our images are significantly blurred, as compared, for example, with those of 720 Villanueva et al (2015). Another conclusion from our study is that the D/H variations over the 721 Martian disk are in agreement with the predictions of the VPIE model (Montmessin et al. 722 2005), and our disk-integrated value of D/H is in agreement with the measurements by 723 Krasnopolsky (2015).

724



725

Figure 12. Left panel: The EXES disk-integrated spectrum of Mars around 7.2 μm, recorded on
April 8, 2014. Right panel: 3 maps of D/H obtained in 2014, 2016 and 2018. The disk-integrated
D/H values are 4.4, 4.0 and 4.3 x VSMOW respectively. The minimum on the 2014 map
corresponds to the Tharsis plateau, in agreement the VPIE model predictions (Montmessin et
al. 2005) and with the observations by Villanueva et al. (2015). Sources: Encrenaz et al. 2016a,
2018, 2019b

- 732 733 <u>3.2 Venus</u>
- 734

In spite of their common composition, dominated by carbon dioxide, the atmosphere of Venus is by far less well understood than the Martian atmosphere. While global climate models of Mars have been developed for over twenty years (Forget et al. 1999), GCMs of Venus are still in progress (Lebonnois et al. 2016, Gilli et al. 2021, Navarro et al. 2021), and consider separately the troposphere below the clouds and the mesosphere above. The chemistry of the Venus atmosphere is dominated by the cycles of sulfur and water. Both SO<sub>2</sub> and H<sub>2</sub>O are abundant in the troposphere (150 ppmv and 30 ppmv respectively; Bézard & de Bergh 2007), 742 and they are transported above the main cloud deck by Hadley convection (Prinn & Fegley 743 1987). SO<sub>2</sub> is photodissociated above the clouds by the solar UV radiation and combines with 744 H<sub>2</sub>O to form H<sub>2</sub>SO<sub>4</sub> which condenses to form the clouds (Zhang et al. 2010, 2012). At the cloud 745 top and above the clouds, water is depleted by a factor of about 30 (Fedorova et al. 2008), 746 while  $SO_2$  is depleted by a factor of 300–1000 (with mixing ratios of 150-500 ppby; Marcq et 747 al. 2011, 2013; Belyaev et al. 2012); this factor is much larger than expected if  $SO_2$  were only 748 associated with  $H_2O$  to form  $H_2SO_4$ . Above the clouds,  $SO_2$  continues to be depleted as the 749 altitude increases, down to the level of a few ppbv. What happens to sulfur between the 750 troposphere and the mesosphere? It is probably trapped in the clouds in the form of aerosols, 751 still to be identified. But another surprise occurs in the upper mesosphere : submillimeter 752 observations have shown that SO<sub>2</sub> is again present at the level of a few tens or hundreds of 753 ppbv (Sandor et al. 2010), which implies the presence of another sulfur reservoir, which also 754 remains to be identified.

755

756 There are other open questions related to the atmosphere and the surface of Venus (Taylor 757 2014). What is the origin of the long-term variations of the SO<sub>2</sub> abundance at the cloud top 758 (Esposito, 2004; Marcq et al. 2013)? Is there active volcanism today, as suggested by infrared observations by Venus Express (Smrekar et al. 2010)? What is the origin of the super-rotation 759 760 and what are the mechanisms of the atmospheric circulation? What are the coupling 761 mechanisms between the troposphere and the mesosphere? These questions (and others) 762 will be addressed by the future space missions toward Venus, Veritas, Envision and Da Vinci, 763 in the next decade.

- 764
- 765 3.2.1 SO<sub>2</sub> and HDO mapping with TEXES

766

In the context of the successful exploration of Venus by the Venus Express mission, and then by the Akatsuki mission, we set up an observing program, using TEXES at IRTF, to monitor both SO<sub>2</sub> and H<sub>2</sub>O at the cloud top in support of these missions. Indeed, both molecules have strong infrared transitions around 7  $\mu$ m and can be simultaneously observed, together with CO<sub>2</sub>. As in the case of Mars, the simultaneous observation of weak transitions of these three molecules is needed to estimate their mixing ratios.

773

774 Our observing campaign started in January 2012 and is still ongoing, after the interruption due 775 to the Covid crisis in 2020. Most of the observations were recorded at 7.4  $\mu$ m (where the cloud 776 top is probed) and 19  $\mu$ m (where the radiation comes from within the clouds, a few km below 777 the cloud top). In addition, some observations were done at 8.6 µm where the radiation comes 778 from a few km above the cloud top. Indeed, this wavelength corresponds to the maximum of 779 the H<sub>2</sub>SO<sub>4</sub> extinction coefficient, responsible for the infrared opacity of the cloud (Zasova et 780 al. 1993). HDO (used as a proxy of H<sub>2</sub>O, as in the case of Mars) is observed at 7.4  $\mu$ m while SO<sub>2</sub> 781 is observed at both 7.4 and 19  $\mu$ m, i.e. at two slightly different levels, at the cloud top and 782 within the clouds. Following earlier measurements (Fedorova et al. 2008; Krasnopolsky 2010), 783 we assume a D/H value of 200 times the terrestrial value above the clouds. 784

Figure 13 shows the spectral interval used to retrieve the  $SO_2$  and HDO maps at 7.3  $\mu$ m, with examples of  $SO_2$  and HDO maps. It was soon realized that both molecules behave very differently (Encrenaz et al. 2012b, 2013). While HDO maps are homogeneous over the disk,  $SO_2$  maps show localized plumes which induce strong variations over the disk, by as much as a factor 5. Repeated observations over several consecutive hours show that these plumes follow the 4-day rotation of the cloud top, and that their lifetime is limited to a few hours. They were also visible in the 19- $\mu$ m maps (Encrenaz et al. 2016b). It is interesting to note that the TEXES maps of SO<sub>2</sub> in the infrared range are in very good agreement with the UV maps recorded by the Akatsuki spacecraft (Encrenaz et al. 2019c); this agreement illustrates that the TEXES data can be used to monitor SO<sub>2</sub> on the night side of Venus where UV observations cannot be done.

796

The existence of the SO<sub>2</sub> plumes was unexpected: since both SO<sub>2</sub> and H<sub>2</sub>O are associated with the formation of H<sub>2</sub>SO<sub>4</sub>, why do we see SO<sub>2</sub> plumes, but no H<sub>2</sub>O plumes? According to GCM simulations, the SO<sub>2</sub> plumes are the result of convective motions within the upper cloud (Lefèvre M. et al. 2018, 2020), combined with a strong SO<sub>2</sub> depletion with altitude just above the cloud top; the H<sub>2</sub>O mixing ratio, in contrast, is more or less constant in this region, which would explain the absence of H<sub>2</sub>O plumes (Lefèvre F., priv. comm.).

803



804

Figure 13. Left: Examples of disk-integrated TEXES spectra at 7.4 μm, showing weak transitions
of SO<sub>2</sub>, HDO and CO<sub>2</sub>. Right panel: Examples of SO<sub>2</sub> maps (middle column) and HDO maps (right
column). The two SO<sub>2</sub> maps are separated by 2 hours. It can be seen that in the second map,
the SO<sub>2</sub> plume is shifted by 7.5°, in agreement with the 4-days rotation of the clouds. Source:
Encrenaz et al. 2016b, 2019c.

810

811 We have used our whole data base (2012 – 2019) to study the long-term variations of SO<sub>2</sub> and 812 HDO at the cloud top, and we have found a clear anticorrelation (cc = -0.8; Encrenaz et al. 2020). As shown in Figure 14, this anticorrelation is actually predicted by the dynamical 813 814 photochemical models of Shao et al. (2021). Between 2014 and 2019, we have observed an 815 overall increase of  $SO_2$  and a slow decrease of  $H_2O$ ; in 2019, the mixing ratios of both 816 molecules were comparable, at the level of about 500 ppbv, while H<sub>2</sub>O was significantly more 817 abundant than SO<sub>2</sub> during the previous years. Unfortunately, no data were obtained in 2020; 818 however, new measurements performed in the summer and fall 2021 indicate that the SO<sub>2</sub> 819 abundance has decreased again. In addition, we have performed a statistical study of the 820 appearance of  $SO_2$  plumes as a function of the local hour. We have shown evidence for a 821 double maximum at the positions of the terminators, which might be the signature of a semidiurnal wave (Figure 14). This behavior was also observed, at a slightly higher atmospheric
level, by the SPICAV UV experiment (Marcq et al. 2020; Figure 14)

#### 824



#### 825

Figure 14. Left: The distribution of the SO<sub>2</sub> plumes as a function of the local hour, as observed
by (top) SPICAV on Venus Express (Marcq et al. 2020), and (bottom) TEXES. Right: The H<sub>2</sub>O/SO<sub>2</sub>
anticorrelation, as predicted by the models (top, Shao et al. 2021) and as observed by TEXES
(bottom; data from Encrenaz et al. 2020b).

830

#### 831 3.2.2 Mesospheric mapping of sulfur and water with ALMA

832

833 As mentioned earlier (3.2), the SO<sub>2</sub> vertical distribution is strongly depleted above the clouds, 834 as compared with its tropospheric value. However a second maximum has been detected in 835 the upper mesosphere, at an altitude of about 80 - 90 km by submillimeter measurements of 836 SO and SO<sub>2</sub> (Sandor et al., 2010, 2012). These observations, which exhibit a strong temporal 837 variability, suggest the presence of another sulfur reservoir in the upper mesosphere, possibly 838 in the form of aerosols. In 2011, just before the beginning of our TEXES campaign, we used 839 the ALMA (Atacama Large Millimeter/ submillimeter Array) to map these species; a HDO 840 transition was again used as a proxy to map the water vapor simultaneously. Observations 841 were performed on November 14 and 15, 2011, at frequencies close to 345 GHz (SO and SO<sub>2</sub>) 842 and 336 GHz (HDO).

843

844 Figure 15 shows the evolution of the SO, SO<sub>2</sub> and HDO maps over 24 hours (Encrenaz et al. 845 2015). It can be seen that their distributions over the disk are very inhomogeneous, and 846 strongly vary from day to day. The SO transition (also shown in Figure 15) is much stronger 847 than the SO<sub>2</sub> line located in the wing of the SO line, which explains the poor signal-to-noise of 848 the SO<sub>2</sub> map. Still, it can be seen that SO<sub>2</sub> is detected on November 14, but not on November 849 15. The behavior of water vapor is interesting. It is more or less homogeneous on November 850 15, but shows a strong localized maximum the day before. There is no apparent correlation 851 between the SO and SO<sub>2</sub> maps, nor between the SO and  $H_2O$  maps. As observed at the cloud 852 top, the behavior of SO<sub>2</sub> in the upper atmosphere shows strong local and short-term temporal 853 variations. In contrast, water vapor, which is homogeneous at the cloud top, also shows these 854 local temporal variations in the upper atmosphere. The explanation might be that, at higher 855 levels, the water vapor content is constrained by its saturation curve; it thus depends upon

- the mesospheric temperature which is known to exhibit strong daily variations.
- 857



858

**Figure 15.** Left: the SO transition at 346.528 GHz observed on Venus with ALMA on November 14, 2011. A weak transition of  $SO_2$  is visible in the wing of the SO line. The best fit (red curve) is obtained with SO = 8 ppbv and  $SO_2 = 12$  ppbv. Right: maps of SO (top),  $SO_2$  (middle) and HDO (bottom) recorded with ALMA on November 14 and 15, 2011. The unit is the depth of the transition. The blue color indicates a high abundance and the red color corresponds to zero. Source: Encrenaz et al. (2015).

865 866

# 867 4. Conclusions and perspectives

868

This last section reviews some achievements of planetary infrared spectroscopy and some remaining questions, as well as possible future developments, especially in the field of exoplanetary spectroscopy.

872

## 873 <u>4.1 Planetary infrared spectroscopy: a perennial tool</u>

874 Until the early 1970s, planetary spectroscopy, performed in the visible and near-infrared 875 range, led to the identification of some atmospheric species (CO<sub>2</sub> on Mars and Venus, CH<sub>4</sub> and 876 NH<sub>3</sub> in Jupiter). Absorption lines were observed in front of the reflected solar continuum, 877 leading to an estimate of the abundance of the observed species and (from the line width) the 878 effective pressure of the formation region. After 1973, spectroscopic observations of the 879 Jovian thermal emission became possible in the 4.6-µm and 10-µm windows, leading to the 880 identification of several minor tropospheric and stratospheric species (Figures 1 and 2). Later, 881 thanks to the IRIS instruments of the two Voyager spacecraft, the thermal profile and the 882  $H_2/He$  ratio were retrieved in the atmospheres of the giant planets (Hanel et al. 2003). Later, 883 the use of very high spectral resolution in the millimeter range, using heterodyne spectroscopy, led to the determination of velocity fields on Mars (Moreno et al. 2009). More 884 885 recently, wind velocities, inferred from high-resolution measurements in the visible range, 886 were measured at the cloud top of Venus (Machado et al. 2017). 887

888 Planetary spectroscopy can bring surprises. We have seen the case of the fast response of 889 Jupiter at the time of the SL9 impact (2.2). In several cases, the spectroscopic identification of 890 new species has shown evidence for unexpected physico-chemical processes in planetary atmospheres. It was the case for the detection of phosphine PH<sub>3</sub> on Jupiter (Ridgway et al. 891 892 1976), and later Saturn (Tokunaga et al. 1980); this molecule, whose presence is not expected 893 on the basis of thermochemical calculations, has been interpreted as the signature of vertical 894 transport, bringing PH<sub>3</sub> from deep tropospheric levels up to the upper troposphere in a time 895 shorter than the chemical destruction time of the molecule; because of the strong internal 896 source of Saturn (Hanel et al. 1983), probably associated with a stronger vertical transport, 897 this effect is stronger on Saturn than on Jupiter. Another example was the unexpected 898 discovery of a large abundance of CO in the stratosphere on Neptune (Rosenqvist et al. 1992; 899 Marten et al. 1993). In this case, the explanation was the existence of an external oxygen 900 source, later confirmed by the discovery of stratospheric water vapor and carbon dioxide in 901 the stratospheres of the giant planets (Feuchtgruber et al. 1997; 2.3.1). In the case of Neptune, 902 the detection of stratospheric CO and HCN (Marten et al. 1993) and more recently CS Moreno 903 et al. 2017) suggest that a large comet hit the planet within the past thousand years (Lellouch 904 et al. 2005).

905

907

#### 906 <u>4.2 A few open questions...</u>

Regarding the terrestrial planets, there are still open questions about the mechanisms
responsible for the diverging evolution of their atmospheres. Some of them are listed below
(see, in particular, Esposito et al. 2007, Taylor 2014 and Catling & Kasting 2017).

911

912 Did Venus have an early ocean? If yes, could life have appeared and developed before the 913 escape of water? If so, could life have escaped the hostile surface conditions and migrated to 914 the level of the clouds, where temperature and pressure conditions are temperate? Regarding 915 present conditions, is active volcanism responsible for the long-term variations of sulfur 916 dioxide at the cloud top? What are the interaction mechanisms between the surface and the 917 lower atmosphere? What are the sources and sinks of sulfur on Venus, and what is the 918 coupling between the troposphere and the mesosphere? What is the origin of the super-919 rotation at the cloud top?

920

921 Active volcanism at the surface of Venus cannot be traced by the monitoring of sulfur dioxide 922 and water vapor at the cloud top, as performed by SPICAV aboard Venus Express or by TEXES 923 at IRTF. Indeed, as mentioned above (3.2), the short-term SO<sub>2</sub> variability observed by TEXES 924 can be explained by convective motions within the clouds, associated to the steep decrease 925 of SO<sub>2</sub> from below to above the clouds (Lefèvre et al. 2018, 2020); its long-term variability 926 could also possibly be associated with intrinsic variability inside the clouds. Thermal mapping 927 of the surface on the night side, at specific wavelengths located outside the CO<sub>2</sub> absorption bands (0.9, 1.1, 1.7, 2.3 µm) can bring new information on active volcanism through changes 928 929 of emissivity, as previously reported by Venus Express observations (Smrekar et al. 2010) or 930 variability in the SO<sub>2</sub> and H<sub>2</sub>O abundances (Marcq et al. 2021). The best tool for this study, 931 however, will be the Synthetic Aperture Radar (SAR) which is planned aboard the Veritas and 932 EnVision orbiters, to be launched by NASA and ESA respectively for the coming decade.

934 In the case of Mars, there is a pending problem associated with the paradox of the early faint 935 Sun (see, e.g. Catling & Kasting 2017). When the Sun's radiation was only 70% of its present 936 value, how could the surface temperature be compatible with the presence of liquid water? 937 In contrast with the previous concept of a wet and warm early atmosphere, present models 938 tend to favor a cold and dry environment, with episodic outgassing due to volcanic and/or 939 meteoritic events. Could life have appeared on Mars in such an environment, and if so, could 940 we find extinct or extant traces of it? This quest is triggering the on-going exploration program 941 of Mars.

942

943 Also related to this question is the controversy about the possible presence of methane. Since 944 2003, several ground-based and space observations have reported the sporadic presence of 945 methane (Krasnopolsky et al. 2004, Formisano et al. 2004, Mumma et al. 2009). Because some 946 of these measurements were controversial, and because of their potential implication in 947 terms of astrobiology (the origin of methane on Mars could be either abiotic or biogenic; see 948 e.g. Atreya et al. 2007, 2011), these announcements raised a lot of discussion in the 949 community. A few years later, the Curiosity rover reported again the transient occurrence of 950 methane plumes, over a low persistent background of 0.2 - 0.6 ppbv (Webster et al. 2015, 2021); at the same time, the Trace Gas Orbiter reported a very stringent upper limit, ten times 951 952 below the background detection of Curiosity (Montmessin et al. 2021). As suggested by 953 Moores et al. (2019), this contradiction could be at least partly due to the day/night cycle of 954 the planetary boundary layer which may be reduced to a few meters during the night (when 955 the CH<sub>4</sub> detection is reported by Curiosity), inhibiting the mixing of methane at higher altitudes 956 (above a few kilometers) where TGO is observing (Webster et al. 2021). The outstanding 957 question regarding methane on Mars is not so much its source (as it could be trapped in 958 clathrates under the surface through abiotic processes) but its fast destruction at the surface. 959 It has been suggested that surface oxidizers, like H<sub>2</sub>O<sub>2</sub> or perchlorates (XClO<sub>4</sub>), could be 960 involved in the destruction of organics at the Martian surface; in particular, large quantities of 961 H<sub>2</sub>O<sub>2</sub> could be produced by triboelectric processes at the time of dust storms, and be stored 962 in the regolith (Encrenaz et al. 2012, Atreya et al. 2021). If released at the same time of CH<sub>4</sub> 963 plumes, this oxidizer could rapidly destroy methane at the surface. These coupled surface-964 atmosphere processes might also explain the discrepancy between the observed  $H_2O_2$ 965 abundance in the gas phase and the model predictions (Section 3.1). It is likely that, in the 966 future, the study of such surface-atmosphere interactions will be crucial for better 967 understanding the seasonal behavior of oxidized species on Mars.

968

#### 969 <u>4.3 The future of ground-based planetary spectral mapping</u>

970

Our observing campaigns on Mars and Venus were performed on a telescope of modest size,
the 3-m IRTF, in an excellent high-altitude site, Maunakea Observatory. With the advent of
10-m class telescopes, and the future perspective of ELTs, we could wonder if these spectral
mapping programs could benefit from these developments. Actually, the question is already
raised since, in a few occasions, the TEXES instrument has been mounted at the 8-m GeminiN telescope, at the Maunakea observatory.

977

In the case of large planetary disks (more than ten arcsec in diameter), like Venus, Mars,
Jupiter and Saturn, the answer appears to be negative. Indeed, a compromise has to be found
between the spectral resolving power (about 80000), the spatial resolution (about 1 arcsec),

the required signal-to-noise (over about 100 per pixel), and the time needed to build a map
(typically less than an hour). If a 10-m telescope is used, the number of pixels needed to scan
the planet becomes very large and the time needed to build a map becomes prohibitive.

984

985 In contrast, the use of TEXES (or similar thermal imaging spectrometers) on 10-m class 986 telescopes is perfectly adapted to the study of small objects, like Io, Titan, or Uranus and 987 Neptune. TEXES has been successfully used at IRTF on Io (Spencer et al. 2005, Tsang et al. 988 2012, 2013) and on Titan, where two new molecules, propane C<sub>3</sub>H<sub>8</sub> and propadiene CH<sub>2</sub>CCH<sub>2</sub> 989 were detected (Roe et al. 2003, Lombardo et al. 2019). The use of a larger telescope would 990 allow us to map these species. The spectral mapping of hydrocarbons emissions ( $CH_4$ ,  $C_2H_2$ 991 and C<sub>2</sub>H<sub>6</sub>) on Neptune by TEXES at the Gemini telescope has been used to infer the thermal 992 field and the spatial distribution of minor species in the stratosphere of Neptune (Greathouse 993 et al. 2011). Such studies could be developed on large telescopes in the future, and be applied 994 to other objects (Uranus, the icy Galilean satellites, Ceres...).

995

997

996 <u>4.4 From solar system planets to exoplanets</u>

998 Over the past twenty years, the development of transit spectroscopy of exoplanets has given 999 us some insight into the atmospheric composition of exoplanets. Two methods have been 1000 used. In the case of primary transits (when the planet passes in front of its host-star), the 1001 atmosphere is observed in transmission at terminator; when it passes behind the star, the 1002 dayside emission of the exoplanet is observed. In the second case, the signal is more difficult 1003 to observe in the near infrared, but becomes stronger at longer wavelengths as the 1004 [planet/star] flux contrast increases (Tinetti et al. 2013). Most of our knowledge on the 1005 atmospheric composition of the exoplanets comes from transit observations; however, near-1006 infrared spectroscopy is becoming performed also on young giant exoplanets detected by the 1007 imaging technique, and is expected to develop rapidly in the future.

1008

1009 The number of exoplanets where atmospheric species have been detected is close to a 1010 hundred. In addition to many atoms (Na, Ca, K, Mg, Fe, O, H, He...), we find several molecules 1011 or radicals (H<sub>2</sub>O, CH<sub>4</sub>, NH<sub>3</sub>, CO, CO<sub>2</sub>, HCN, C<sub>2</sub>H<sub>2</sub>, C<sub>2</sub>, CN, NH<sub>2</sub>...) and more exotic species (AlO, 1012 Al<sub>2</sub>O<sub>3</sub>, SiO, TiO<sub>2</sub>, VO...; see Schneider et al. 2011 and http://exoplanet.eu)). It must be noted 1013 that these detections can be tentative or questionable; reviews on this topic can be found, in 1014 particular, in Seager (2008), Deming and Seager (2017), and Madhusudhan (2019). In addition 1015 of the intrinsic weakness of the signal, their identification is made more difficult than in the 1016 case of solar system planets for two reasons: (1) hot Jupiters (the easiest targets), as an effect 1017 of their high temperature, exhibit broad molecular bands which tend to overlap; (2) in the 1018 case of secondary transits, the shapes of the bands (in emission or in absorption) depend upon 1019 the thermal profile of the exoplanet which must be simultaneously retrieved.

1020

1021 Can we anticipate the atmospheric composition of a given exoplanet? In the case of a 1022 transiting object, we know its mass and its distance to its host-star. Knowing the spectral type 1023 of this star, we can estimate the exoplanet's equilibrium temperature. We can then plot a (T, 1024 M) diagram and, using solar system objects as analogs, make an initial guess on their 1025 atmospheric composition, based on the general laws of thermochemical equilibrium (Tinetti 1026 et al. 2013). This diagram, shown in Figure 16, may be used for a very simplified classification 1027 of exoplanets. 1029 On the mass scale, the value of 10 terrestrial masses acts as a threshold, because for more 1030 massive nuclei, the gravity field is expected to capture the surrounding material (mostly 1031 hydrogen and helium) and lead to the formation of a giant planet. Below this threshold, as a 1032 first approximation, the planet can be considered as a solid planet. On the temperature scale, 1033 a threshold is given by the temperature of the ice line, which can be approximated to the 1034 temperature of water condensation in a dilute medium (i.e. about 180 – 200 K). Below this 1035 temperature, we expect to have small icy nuclei (Titan, Pluto) or giant planets formed by 1036 nuclear accretion. At higher temperature, we expect to have rocky planets (Earth, Mars, Venus 1037 or Mercury); the hot exoplanets have no equivalent in the solar system. If we assume the core-1038 accretion scenario for the planetary formation, we do not expect hot giants to be present in 1039 this diagram; the existence of hot giant exoplanets is attributed to migration processes.



1028



1041

Figure 16. A classification of exoplanets as a function of temperature and mass. The 1042 temperature is the equilibrium temperature, inferred from the distance to the host star and 1043 1044 the spectral type of the star. Detected exoplanets are indicated in light grey (Source: 1045 exoplanet.eu). The vertical dashed line around T = 200 K indicates the water ice line. The 1046 horizontal dotted line indicates the threshold at 10 terrestrial masses, separating solid planets 1047 from giant ones. Solar system planets are indicated with black symbols. The hot planets have no equivalent in the solar system. The young giant exoplanets in the process of cooling are 1048 1049 missing in this diagram.

1050

1051 On the basis of this simple diagram, we can estimate what could be the atmospheric 1052 composition of each class of exoplanet. Assuming thermochemical equilibrium, we expect 1053 carbon (and nitrogen) to be in the form of  $CH_4$  (and  $NH_3$ , respectively) at low temperature and 1054 high pressure, and in the form of CO or  $CO_2$  (and  $N_2$ , respectively) under opposite conditions. 1055 This is actually what we observe in the solar system (in the case of the small icy bodies, the 1056 presence of nitrogen is explained as the effect of  $NH_3$  photodissociation).

1058 There are several caveats to this very simple classification. The first one is that it does not take 1059 into account the effect of migration, which is known to be an efficient mechanism in 1060 exoplanetary systems. This is probably the reason why planets of intermediate masses (super-1061 Earths and sub-Neptunes) do not appear on this diagram, while they are known to be very 1062 common in the Galaxy (Figure 16). Migration, in particular, could be responsible for the 1063 presence of small, very dense exoplanets in the close vicinity of their stars; these objects could 1064 be the remnant nuclei of previous massive giant objects, initially formed at large distances 1065 (Mocquet et al. 2014). In addition, our diagram does not consider the case of warm young 1066 exoplanets, distant from their star, now revealed with imaging techniques. Finally, another 1067 caveat is a possible departure from thermochemical equilibrium within the exoplanet's 1068 atmospheres: as an example, both CH<sub>4</sub> and CO and have been detected on two giant 1069 exoplanets, HD208458b and HD189733b. Nevertheless, in the absence of other constraints, 1070 this simple diagram could be considered as a very first step toward the characterization of 1071 planetary atmospheres.

1072

1073 Two main mechanisms could be responsible for a departure from thermochemical 1074 equilibrium. The first one is photochemistry. In the case of terrestrial planets, it is responsible 1075 for the  $CO_2$  and  $H_2O$  photodissociation and, in the case of the Earth, for the formation of the ozone layer. In the case of giant planets, it is responsible for the formation of a variety of 1076 1077 hydrocarbons from methane photochemistry. The second mechanism is vertical transport, 1078 which carries up molecules from the deep troposphere up to observable levels (at pressure 1079 levels of a few bars or less) in a lifetime shorter than the destruction lifetime of the molecule. 1080 The best example is PH<sub>3</sub> in Jupiter and Saturn (see Section 4.1), and the same situation could 1081 occur also on gaseous exoplanets.

1082

In summary, in view of the high diversity of the planetary spectra observed in the solar system
 - a « simple » multiple system with quasi-coplanar and concentric orbits ! - we must be
 prepared to discover an extreme variety of exoplanetary spectra. In particular, an important
 means will be the JWST, starting in 2022. A few years later, the European Ariel mission is
 expected to provide us with a whole inventory of infrared exoplanetary spectra of all kinds of
 targets.

1089 1090

## 1091 **5. Acknowledgements**

1092

1100

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1110

Since the early years of my carrier, I have been observing at Maunakea Observatory. This is a wonderful place and I was always happy to spend a couple of weeks there. Since the early 2000s, I have been using the Texas Echelon Cross Echelle Spectrograph (TEXES) imaging spectrometer at the NASA Infrared Telescope Facility (IRTF), whose director was Alan Tokunaga. I am most grateful to Tommy Greathouse who has been operating TEXES for about twenty years; working with him at IRTF has been (and still is) a great experience.

1117

In 2009, at the occasion of a conference in Barcelona, I started being interested in exoplanets.
At this occasion I met Giovanna Tinetti, and I am very grateful to her for having introduced me
to the new field of exoplanetary spectroscopy.

1121

1122 I had many friendly and long-lasting collaborations during my career. The first one was with 1123 Jean-Pierre Bibring and his group at the Institut d'Astrophysique Spatiale (Orsay, France), as 1124 well as Vassili Moroz at the Space Research Institute (IKI) in Moscow. Another one was with 1125 the Planetary Group of Jet Propulsion Laboratory, in particular, Glenn Orton and Bob Carlson; 1126 I am especially grateful to Sam Gulkis who allowed us (my husband Pierre and me) to come to 1127 JPL frequently. I also want to mention my long-term collaboration with Sushil Atreya from the 1128 University of Michigan; it is always a pleasure for me to exchange with him and hear his 1129 innovative ideas. Over the years, I have also developed a friendly and fruitful collaboration 1130 with the team initiated by François Forget in Paris, which includes many bright scientists from 1131 LMD (Laboratoire de Météorologie Dynamique) and LATMOS (Laboratoire Atmosphères, 1132 Milieux, Observations Spatiales), and has developed Global Climatic Models for most of the 1133 planets and satellites with an atmosphere.

1134

Over the three past decades, I have spent quite some time writing text books and popular books. I have been most lucky to meet two very good friends and co-authors whom I want to thank, James Lequeux and Fabienne Casoli, both from Paris Observatory. And, last but not least, my profound thanks to my family, my husband Pierre Encrenaz (a radio-astronomer who turned out to be also a planetologist) and our children and grand-children.

1140

1 would like to conclude this paper with a comment. When I started working in astronomy,
there were very few women in this field, especially in the US. I remember a conference in the
US, in the early 1970s, attended by more than a hundred participants and only three women...
I am very happy to see so many young women attending the DPS conference these days, and

- 1145 I hope that this tendency will continue and grow in the future.
- 1146

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- 1150
- 1151

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