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## To cite this version:

M. Trevisan, G. A. Mamon. A finer view of the conditional galaxy luminosity function and magnitude-gap statistics. Monthly Notices of the Royal Astronomical Society, 2017, 471, pp.20222038. 10.1093/mnras/stx1656 . insu-03747430

HAL Id: insu-03747430
https://insu.hal.science/insu-03747430
Submitted on 8 Aug 2022

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# A finer view of the conditional galaxy luminosity function and magnitude-gap statistics 

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Accepted 2017 June 28. Received 2017 June 26; in original form 2017 May 30


#### Abstract

The gap between first- and second-ranked galaxy magnitudes in groups is often considered a tracer of their merger histories, which in turn may affect galaxy properties, and also serves to test galaxy luminosity functions (LFs). We remeasure the conditional luminosity function (CLF) of the Main Galaxy Sample of the SDSS in an appropriately cleaned subsample of groups from the Yang catalogue. We find that, at low group masses, our best-fitting CLF has steeper satellite high ends, yet higher ratios of characteristic satellite to central luminosities in comparison with the CLF of Yang et al. The observed fractions of groups with large and small magnitude gaps as well as the Tremaine \& Richstone statistics are not compatible with either a single Schechter LF or with a Schechter-like satellite plus lognormal central LF. These gap statistics, which naturally depend on the size of the subsamples, and also on the maximum projected radius, $R_{\max }$, for defining the second brightest galaxy, can only be reproduced with two-component CLFs if we allow small gap groups to preferentially have two central galaxies, as expected when groups merge. Finally, we find that the trend of higher gap for higher group velocity dispersion, $\sigma_{\mathrm{v}}$, at a given richness, discovered by Hearin et al., is strongly reduced when we consider $\sigma_{\mathrm{v}}$ in bins of richness, and virtually disappears when we use group mass instead of $\sigma_{\mathrm{v}}$. This limits the applicability of gaps in refining cosmographic studies based on cluster counts.


Key words: galaxies: evolution - galaxies: formation - galaxies: groups: general-galaxies: luminosity function, mass function.

## 1 INTRODUCTION

One of the most fundamental properties of galaxy populations is their distributions of stellar masses and luminosities. Indeed, as galaxies form in dark matter haloes, their stellar mass function (SMF) and luminosity function (LF) combine information on the halo mass function. Moreover, galaxies evolve in many ways, through both internal and environmental processes, and these leave imprints in the SMF and LF. For example, the lower mass groups of galaxies (excluding the higher mass clusters) should be the primary sites of galaxy mergers (e.g. Mamon 1992).

In a series of important articles, Yang and co-authors have modified our view of the galaxy SMF and LF. Yang et al. $(2005,2007)$ first designed a new type of group finder that extracts groups from galaxy catalogues in redshift space. In their group finder, the most luminous galaxies or the most massive in stars (hereafter called central galaxies) have a special role in the group, as in most semi-analytical models (SAMs) of galaxy formation (Kauffmann et al. 1999; Croton et al. 2006; De Lucia \& Blaizot 2007; Guo et al. 2011; Henriques et al. 2015). Testing group finders on mock

[^0]catalogues, Duarte \& Mamon (2015) found that both their own group finder MAGGIE and the Yang et al. group finder are much more accurate than the popular Friends-of-Friends (FoF) method (Huchra \& Geller 1982; Berlind et al. 2006; Robotham et al. 2011). Indeed, the groups extracted with MAGGIE and the Yang et al. (2007) group algorithm are much less prone to be secondary fragments of more massive groups and lead to much more accurate and less biased group masses than one infers for FoF groups using the virial theorem.
In a subsequent work, Yang, Mo \& van den Bosch (2008, hereafter Y08) have shown that the galaxy LF strongly depends on group mass, in particular the characteristic luminosity at the knee of the LF increases with group mass. Y08 thus measured the conditional luminosity function (CLF). Moreover, Y08 established that in sufficiently narrow bins of group (halo) mass, the CLF was bimodal: the non-central (hereafter satellite) galaxies follow a luminosity distribution close to the gamma distribution known as the Schechter (1976) function that represents well the full LF, while the central galaxies are well fitted by a lognormal distribution. Thus, Y08 confirmed the idea that centrals and satellites are distinct populations (White \& Frenk 1991; Kravtsov et al. 2004). These results were subsequently generalized to the SMF by Yang, Mo \& van den Bosch (2009).

The Yang model leads to a stellar mass-halo mass relation (SMHM) where the upper envelope increases with halo mass, and with a gap between the centrals and the satellite galaxies. This particular SMHM has been reproduced quite accurately by Cattaneo et al. (2011) in an analytical model of galaxy formation used to predict stellar masses for galaxies living in groups of given mass and at a given redshift and run on the halo merger trees extracted from a high-resolution cosmological simulation.

The idea of a gap between centrals and satellites has been around for a long time. Using $N$-body simulations of isolated groups, various teams recognized early on that mergers will lead to runaway growth of the most massive galaxy (Carnevali, Cavaliere \& Santangelo 1981; Schneider \& Gunn 1982; Barnes 1989). This growth occurs independently of the merger mechanism (Mamon 1987), whether the group evolves by direct merging between galaxies or by orbital decay via dynamical friction that causes galaxies to lose energy and angular momentum against a diffuse background (e.g. Schneider \& Gunn 1982). In both scenarios, the growth of the central, brightest group galaxy (BGG) occurs at the expense of the second-brightest group galaxy (SBGG), because the merger cross-section of the SBGG is greater than that of the smaller, less luminous galaxies, and because the dynamical friction time scales as the inverse of the galaxy subhalo mass (Chandrasekhar 1943), leading to faster orbital decay of the SBGG, hence more rapid merging with the BGG. The early $N$-body simulations of Mamon (1987) showed that the gap grows as log time after a waiting period of 0.5 to 1 Gyr (see Farhang et al. 2017).

In the hope of finding groups that have been the sites of such rapid growth of the most massive (central) galaxy, people searched for groups with very large magnitude gaps between their first- and second-ranked members. Such studies were boosted by the discovery of diffuse X-ray sources, roughly as X-ray luminous as the most luminous groups, with a giant elliptical galaxy surrounded by low-luminosity satellites (Ponman et al. 1994; Mulchaey \& Zabludoff 1999). These systems were recognized as having evolved by early galaxy mergers, which ceased long ago without replenishment of the group's most luminous satellites by new infalling members. Jones et al. (2003) defined fossil groups (FGs; as coined by Ponman et al. 1994) to be X-ray luminous (bolometric $L_{\mathrm{X}}>10^{42} h_{50}^{-2} \mathrm{erg} \mathrm{s}^{-1}$ ) and with a large absolute magnitude gap: $\mathcal{M}_{2}-\mathcal{M}_{1}>2$, where the SBGG is the second most luminous galaxy within a maximum projected radius of $R_{\max }=0.5 r_{200}$. Assuming that high X-ray luminosity is a proxy for high total mass, one can study FGs with optical group catalogues, by seeking high-mass groups with such large magnitude gaps.

However, without information on the group mass, only a fraction of large-gap groups (LGGs) are in fact FGs (Raouf et al. 2014). The origin of LGGs and FGs is still a matter of debate. Studies of the evolution of FGs in cosmological simulations seem to indicate that the mass assembly histories of FG haloes differ from those of small-gap groups (SGGs; D'Onghia et al. 2005; Dariush et al. 2007; Raouf et al. 2014; Farhang et al. 2017).

The magnitude gap serves also as a diagnostic of the LF. Tremaine \& Richstone (1977) derived two simple statistics to test whether the observed magnitude gaps are consistent with one or several cumulative LFs that have power-law asymptotic behaviour at the faint and bright ends. They defined the quantities
$T_{1}=\frac{\sigma\left(\mathcal{M}_{1}\right)}{\left\langle\mathcal{M}_{2}-\mathcal{M}_{1}\right\rangle}$ and $T_{2}=\frac{1}{\sqrt{0.677}} \frac{\sigma\left(\mathcal{M}_{2}-\mathcal{M}_{1}\right)}{\left\langle\mathcal{M}_{2}-\mathcal{M}_{1}\right\rangle}$,
where $\sigma\left(\mathcal{M}_{1}\right)$ corresponds to the standard deviation of the absolute magnitude of the first-ranked galaxies, while $\left\langle\mathcal{M}_{2}-\mathcal{M}_{1}\right\rangle$ and
$\sigma\left(\mathcal{M}_{2}-\mathcal{M}_{1}\right)$ are the respective mean and standard deviation of the magnitude-gap distribution. Using Cauchy-Schwarz inequalities, Tremaine \& Richstone show that $T_{1}$ and $T_{2}>1$ for groups with galaxy LFs with power-law shapes for their cumulative forms in both their faint and bright ends. In group samples with $T_{1}$ and $T_{2}$ lower than unity, the first-ranked galaxies are abnormally bright compared to the second-ranked galaxies, unless the group samples are small $(\lesssim 40)$ or the group richnesses are low $(\lesssim 8)$ (see table 7 of Mamon 1987).

Values of $T_{1}$ and $T_{2}$ smaller than unity have been found in several studies: Tremaine \& Richstone measured $T_{1}=0.45$ for their full cluster sample and $T_{1}=0.72$ for their subsample of clusters with over 75 members. Loh \& Strauss (2006, hereafter LS06) found $T_{1}=0.75 \pm 0.1$ and $T_{2}=0.86 \pm 0.1$ in nearby rich clusters in the Sloan Digital Sky Survey (SDSS) dominated by Luminous Red Galaxies (LRGs). Lin, Ostriker \& Miller (2010, hereafter LOM10) determined $T_{1}=0.84 \pm 0.01$ and $T_{2}=0.77 \pm 0.01$ for low-luminosity SDSS clusters, but $T_{1}=0.70 \pm 0.1$ and $T_{2}=0.79 \pm 0.01$ in luminous ones. Furthermore, Díaz-Giménez et al. (2012, hereafter DG+12) obtained even smaller values for a complete sample of compact groups drawn from the 2MASS catalogue: $T_{1}=0.51 \pm 0.06$ and $T_{2}=0.70 \pm 0.06$, with slightly lower values for mock compact group samples extracted from SAMs of galaxy formation. These low values of $T_{1}$ and $T_{2}$ imply that the cumulative LF can depart from power-law behaviour at the faint or bright ends.

The simple $N$-body simulations of Mamon (1987) indicated that galaxy mergers in groups lead to low values of $T_{1}$ and $T_{2}$, as they tend to build the first-ranked galaxy at the expense of the second-ranked one, as confirmed by Smith et al. (2010) through observations.

It has been suggested that galaxy properties are not only a function of their halo mass, but also of their halo assembly history (Gao, Springel \& White 2005), which is termed galaxy assembly bias. Perhaps the first unequivocal proof of such galaxy assembly bias was the measurement of greater dark matter concentrations of haloes of red galaxies compared to blue galaxies of the same stellar or halo mass (Wojtak \& Mamon 2013), given that red galaxies have older stellar populations while, at a given mass, more concentrated haloes assemble earlier (Wechsler et al. 2002).

If large magnitude gaps are caused by an earlier merger history, one would then expect that the BGGs in LGGs should show older stellar populations. In a previous study (Trevisan, Mamon \& Khosroshahi 2017, hereafter Paper I), we investigated whether such a difference in the evolution of LGGs is imprinted in the global properties of the stellar content and in the star formation history (SFH) of the BGGs. Through a detailed reconstruction of the stellar assembly of galaxies, we found that, after removing the dependence with galaxy velocity dispersion or with stellar mass, there is no correlation with magnitude gap of BGG ages, metallicities and SFHs.

This lack of variation of SFH with $\Delta \mathcal{M}_{12}$ suggests that all BGGs are formed in a very similar way regardless of the magnitude gap. Therefore, groups with large $\Delta \mathcal{M}_{12}$ might merely represent an evolutionary phase of galactic systems. Indeed, analyses of SAMs indicate that most LGGs turn into regular groups by seeing the gap filled by infalling luminous galaxies (von Benda-Beckmann et al. 2008; Dariush et al. 2010). This suggests that the 'true' SBGGs may lie further than $0.5 r_{\text {vir }}$ away from the BGG (e.g. Gozaliasl et al. 2014).

It is natural to expect that the distribution of magnitude gaps should be consistent with the general LF or the CLF. Paranjape \& Sheth (2012) claim that the gap distributions are
consistent with the order statistics of luminosities sampled from a single LF independent of halo mass, as long as the LF is accurate at the bright end. This appears to be in contradiction with the two-component CLF of Y08. Hearin et al. (2013) analysed both SDSS groups and mock groups to find that the richness-mass relation of groups differs between LGGs and SGGs: LGGs have higher masses (measured by their velocity dispersions) than SGGs of same richness. They therefore suggest that group masses can be made more precise once the gap is factored in. More accurate group masses would be beneficial for cosmographic studies based on the group/cluster mass function. Motivated by this finding, More (2012) demonstrated that, at fixed richness, the CLF should indeed lead to a correlation of $\Delta \mathcal{M}_{12}$ with group mass.

In this article, we wish to address three questions related to magnitude gaps of groups. (1) Is the Y08 model for the CLF correct, in particular are the satellite bright-end shape and the ratio of satellite to central characteristic luminosities independent of group mass? (2) Is the CLF consistent with the statistics of magnitude gaps? (3) Is the richness-mass relation of groups a function of gap, allowing more accurate group mass determinations? We will answer these questions by analysing a doubly complete subsample of the SDSS, suitably cleaned for edges and nearby saturated stars, using the latest Yang group catalogue.

This paper is organized as follows. In Section 2, we describe the sample of groups and the data used in our analysis. In Section 3, we fit different CLF models, and in Section 4, we compute the statistics of the bright end of these CLFs. In Section 5, we discuss our results and in Section 6, we present the summary and the conclusions of our study. We adopt three-year Wilkinson Microwave Anisotropy Probe cosmological parameters of a flat $\Lambda$ Cold Dark Matter Universe with $\Omega_{\mathrm{m}}=0.275, \Omega_{\Lambda}=0.725, \Omega_{\mathrm{b}}=0.046$ and $H_{0}=70.2 \mathrm{kms}^{-1} \mathrm{Mpc}^{-1}$, to be consistent with parameters of the group catalogue used in this study.

## 2 SAMPLE AND DATA

The galaxy groups were selected from the updated version of the catalogue compiled by Yang et al. (2007). ${ }^{1}$ The new catalogue contains 473482 groups drawn from a sample of 601751 galaxies mostly from the Sloan Digital Sky Survey's Data Release 7 (SDSSDR7; Abazajian et al. 2009).

The radii $r_{200, \mathrm{~m}}$, i.e. the radii of spheres that are 200 times denser than the mean density of the Universe, are derived from the $M_{200, \mathrm{~m}}$ masses given in the Yang et al. catalogue, which are based on abundance matching with the group luminosities. We then calculated the virial radii ( $r_{\text {vir }}=r_{200, \mathrm{c}}$, where $r_{200, \mathrm{c}}$ are the radii of spheres that are 200 times denser than the critical $^{2}$ density of the Universe) and masses ( $\left.M_{\text {vir }} \equiv M_{200, \mathrm{c}}=200 H^{2}(z) r_{200, \mathrm{c}}^{3} / G\right)$ by assuming the Navarro, Frenk \& White (1996, hereafter NFW) profile and the concentration-mass relation given by Dutton \& Macciò (2014).
We used biweight and gapper scale estimators (see Beers, Flynn \& Gebhardt 1990, and references therein) to determine the group velocity dispersions, $\sigma_{\mathrm{v}}$, for groups with $N_{\text {vir }} \geq 10$ and $N_{\text {vir }}<10$, respectively, where $N_{\text {vir }}$ is the number of galaxies within $r_{\text {vir }}$. The $\sigma_{\mathrm{v}}$ values were corrected to the group rest frame by dividing by the factor ( $1+z_{\text {group }}$ ) Harrison \& Noonan (1979).

[^1]We selected groups that satisfy the following criteria:
(i) redshifts in the range from 0.015 to 0.07 ;
(ii) masses $\log \left(M_{\text {vir }} / \mathrm{M}_{\odot}\right) \geq 13.0$;
(iii) at least two member galaxies within $0.5 r_{\text {vir }}$ brighter than $\mathcal{M}_{r}^{\text {Petro }} \leq-19.57$, where $\mathcal{M}_{r}^{\text {Petro }}$ is the $k$-corrected SDSS Petrosian absolute magnitude in the $r$ band;
(iv) the magnitude gap, defined as the difference between the $k$-corrected SDSS $r$-band Petrosian absolute magnitudes of the BGG and SBGG galaxies within half the virial radius, i.e.

$$
\Delta \mathcal{M}_{12}=\mathcal{M}_{r, 2}^{\text {Petro }}-\mathcal{M}_{r, 1}^{\text {Petro }}
$$

is smaller than 2.47 mag.
The lower redshift limit was chosen to avoid selecting groups too close to the edge of the catalogue (the groups were defined using galaxies at $0.01<z_{\text {gal }}<0.2$.). The upper limit was optimized to obtain the largest possible number of groups with $\Delta \mathcal{M}_{12} \geq 2 \mathrm{mag}$, given the other criteria and taking into account the variation of $\mathcal{M}_{r}^{\text {Petro }}$ and $\Delta \mathcal{M}_{12}$ limits with $z$.
To establish the $\log M_{\text {vir }}$ completeness limit, we compared the halo mass function of our sample with the theoretical halo mass function computed using the HMFCALC tool ${ }^{3}$ (Murray, Power \& Robotham 2013). The adopted halo mass lower limit, $\log \left(M_{\text {vir }} / \mathrm{M}_{\odot}\right) \geq 13.0$, corresponds to the value above which the difference between the observed and theoretical mass functions is smaller than $\sim 0.1$ dex.
We did not use the galaxy absolute magnitudes of the Yang group catalogue, but instead retrieved the apparent galaxy magnitudes from the SDSS-DR12 data base (Alam et al. 2015). Since our sample is within a small range in redshift, we did not apply any evolution correction with redshift (see e.g. Blanton et al. 2003b; Yang et al. 2007), ${ }^{4}$ and the absolute magnitudes are simply given by
$\mathcal{M}_{r}^{\text {Petro }}=m_{r}^{\text {Petro }}-\operatorname{DM}(z)-k_{0.1}(z)$,
where $m_{r}^{\text {Petro }}$ is the Petrosian apparent magnitude in the $r$ band corrected for Galactic extinction and $\mathrm{DM}(z)$ is the distance modulus. The $k$-corrections, $k_{0.1}(z)$, were obtained with the ксогRECT code (version 4_2) of Blanton et al. (2003a), choosing as reference the median redshift of the SDSS main galaxy sample (MGS; $z=0.1$ ). ${ }^{5}$
We determine the 95 per cent limit in absolute magnitudes following the geometric approach similar to that described by Garilli, Maccagni \& Andreon (1999) and La Barbera et al. (2010). We first determine the 95 th percentile of the extinction-corrected apparent magnitude, $m_{r}^{\text {Petro }}$, in bins of $\mathcal{M}_{r}^{\text {Petro }}$ and then perform a linear fit to the 95 th-percentile points, so that the value of $\mathcal{M}_{r}^{\text {Petro }}$ where the best-fitting line intersects $m_{r}^{\text {Petro }}=17.77$ defines the absolute magnitude of 95 per cent completeness. This leads to a 95 per cent completeness limit of $\mathcal{M}_{r}^{\text {Petro }} \leq-19.57$ for our sample.

This absolute magnitude limit in turn leads to a sample complete up to $\Delta \mathcal{M}_{12}=2.47 \mathrm{mag}$, as illustrated in fig. 1 of Paper I.

The above criteria lead to a sample of 2319 groups.

[^2]
### 2.1 Spectroscopic incompleteness

The SDSS fibre collision limit prevents neighbouring fibres from being closer than 55 arcsec , therefore affecting the completeness of the SDSS spectroscopy in high-density regions. This spectroscopic incompleteness may lead to an incorrect identification of the BGGs and SBGGs. Following the approach adopted in Paper I to address this issue, we used the SDSS photometric catalogue to identify galaxies without SDSS-DR7 spectra that could be BGGs or SBGGs.

Concerning BGGs, we first selected all the photometric SDSSDR12 galaxies within one virial radius from the luminosityweighted centre of each group that are brighter than the group's BGG. We then retrieved the spectroscopic redshifts available in SDSS-DR12, and the NASA/IPAC Extragalactic Database (NED), to check if these galaxies lie within the redshift range of the groups. We adopt a maximum redshift separation given by
$\left|z-z_{\text {group }}\right| c<2.7 \sigma_{\mathrm{v}}$.
The factor 2.7 in equation (3) provides optimal rejection of outliers to recover the line-of-sight velocity dispersion profile expected for a single-component NFW model (Mamon, Biviano \& Murante 2010). If the group contains a galaxy brighter than its BGG, but the galaxy has no redshift available in either SDSS-DR12 or NED, then the group is discarded. These criteria led us to discard 228 groups that may have incorrect BGG identifications.

For SBGGs, we followed a similar approach, by retrieving from the photometric catalogue all galaxies within $0.5 r_{\text {vir }}$ from the BGGs that are brighter than the SBGG of that group. Groups are excluded if any of these galaxies have redshifts in the range specified by equation (3), or if there is no spectroscopic redshift. We thus discarded an additional 192 groups that may have incorrect SBGG identifications.

### 2.2 Bright stars and edges of the survey

The proper identification of BGGs and SBGGs might also be affected by the presence of very bright stars lying near the line of sight that may hide or prevent an accurate photometric measurement of BGGs or SBGGs. Moreover, BGGs and/or SBGGs of groups close to the edges of the survey may lie outside the area covered by SDSS. Therefore, we ensure that at least 95 percent of the region within $0.5 r_{\text {vir }}$ from the group centre lies within the SDSS coverage area and are not masked by bright stars.

To determine the fraction, $f_{\text {mask }}$, of the area within $0.5 r_{\text {vir }}$ that lies within the SDSS masks for bright stars or that are outside the boundaries of the survey, we adopted the SDSS-DR7 spectroscopic angular selection function mask ${ }^{6}$ provided by the NYU Value-Added Galaxy Catalog team (Blanton et al. 2005) and assembled with the package mangle 2.1 (Hamilton \& Tegmark 2004; Swanson et al. 2008). Only groups with $f_{\text {mask }}<5$ percent are selected, leading to a sample of 1597 (out of 1899) groups.

In summary, as displayed in Table 1 (column $R_{\max } / r_{\text {vir }}=0.5$ ), following the criteria (i) to (iv) listed in the beginning of Section 2, but discarding a total of $228+192=420$ groups with incomplete SDSS-DR7 spectroscopy (see Section 2.1), and 302 additional groups that are in the edge of the survey and/or affected by the presence of bright stars, we obtain our final sample of 1597 groups, among which 152 have $\Delta \mathcal{M}_{12}>2$ mag.

[^3]Table 1. Steps in cleaning the group sample for different maximum distance, $R_{\max }$, allowed for the SBGGs.

|  | $R_{\max } / r_{\text {vir }}$ |  |  |  |
| :--- | ---: | ---: | ---: | ---: |
| Criteria | 0.5 | 1.0 | 1.5 | 2.0 |
| (1) Initial number of groups | 2319 | 2900 | 3021 | 3027 |
| (2) Spectroscopic incompleteness |  |  |  |  |
| BGG (within $1 r_{\text {vir }}$ ) | 228 | 271 | 286 | 289 |
| SBGG (within $R_{\max }$ ) | 192 | 286 | 367 | 465 |
| (3) Near edges or bright stars | 302 | 349 | 346 | 353 |
| (4) Final group sample | $\mathbf{1 5 9 7}$ | $\mathbf{1 9 9 4}$ | $\mathbf{2 0 2 2}$ | $\mathbf{1 9 2 0}$ |
| LGGs $\left(\Delta \mathcal{M}_{12} \geq 2.0\right.$ mag) | 152 | 111 | 93 | 78 |
| SGGs $\left(\Delta \mathcal{M}_{12} \leq 0.3\right.$ mag) | 275 | 446 | 474 | 460 |
| (5) Satellites at $\boldsymbol{R} \leq \boldsymbol{R}_{\text {max }}$ | $\mathbf{4 9 6 1}$ | $\mathbf{9 2 3 4}$ | $\mathbf{1 1 1 5 2}$ | $\mathbf{1 1 0 9 5}$ |

### 2.3 Samples with $\boldsymbol{R}_{\max }>0.5 r_{\text {vir }}$

As mentioned in Section 1, the value of $\Delta \mathcal{M}_{12}$ might vary with the maximum distance allowed for obtaining the SBGG. To investigate how the gap varies with $R_{\max }$, we select samples with $\Delta \mathcal{M}_{12}$ defined within $0.5,1.0,1.5$ and $2.0 r_{\text {vir }}$, following the same criteria listed in Section 2. Groups with incomplete SDSS-DR7 spectroscopy, near the edges of the survey or with bright stars nearby, were excluded following the same approach described in Sections 2.1 and 2.2.

The group samples are summarized in Table 1. The samples defined within $0.5 r_{\text {vir }}$ and in one virial radius have 1425 groups in common, and 1237 groups are present in all four samples.

## 3 CONDITIONAL LUMINOSITY FUNCTIONS

The galaxy CLF can be estimated by directly counting galaxies in groups of our doubly complete sample of galaxies. The absolute Petrosian magnitudes are converted to luminosities using the solar absolute magnitude in the $r$ band, redshifted to $z=0.1, \mathcal{M}_{r} \odot=$ 4.76 (Blanton et al. 2003b).

As in Y08, we assume that the total CLF of haloes with $M_{\text {vir }}=M$ is the sum of the CLFs of the central and the satellite galaxies:

$$
\begin{equation*}
\Phi(L \mid M)=\frac{\mathrm{d} N}{\mathrm{~d} \log L}=\Phi_{\mathrm{c}}(L \mid M)+\Phi_{\mathrm{s}}(L \mid M) . \tag{4}
\end{equation*}
$$

The central galaxies follow a lognormal distribution,
$\Phi_{\mathrm{c}}(L \mid M)=\frac{1}{\sqrt{2 \pi} \sigma_{\mathrm{c}}} \exp \left[-\frac{\left(\log L-\log L_{\mathrm{c}}\right)^{2}}{2 \sigma_{\mathrm{c}}^{2}}\right]$,
while the satellite luminosities are distributed according to the probability distribution function

$$
\begin{align*}
\Phi_{\mathrm{s}}(L \mid M)= & \frac{\beta}{\left.\Gamma\left[(\alpha+1) / \beta,\left(L_{\min } / L_{\mathrm{s}}\right)^{\beta}\right)\right]} \\
& \times\left(\frac{L}{L_{\mathrm{s}}}\right)^{\alpha+1} \exp \left[-\left(\frac{L}{L_{\mathrm{s}}}\right)^{\beta}\right] \tag{6}
\end{align*}
$$

where $\Gamma(a, x)=\int_{x}^{\infty} t^{a-1} \exp (-t) \mathrm{d} t$ is the upper incomplete Gamma function. Equation (6) corresponds to the Schechter (1976) LF when $\beta=1$, while Y08 adopted $\beta=2$ and $L_{\mathrm{s}}=L_{\mathrm{c}} / 10^{0.25}$. Hereafter, a galaxy is referred to as central (satellite) if its luminosity is drawn from the CLF specific to the centrals (satellites; equations 5 and 6).

Instead of fitting the CLFs in bins of halo mass, we adopt a different approach. We assume that the parameters $\log L_{\mathrm{c}}, \sigma_{\mathrm{c}}, \log L_{\mathrm{s}}$, $\alpha$ and $\beta$ vary linearly with $\log M_{\mathrm{vir}}$, and determined the best-fitting
parameters through maximum likelihood estimation (MLE). We assume that the brightest galaxy is always the central galaxy and we fit $\Phi_{\mathrm{c}}(L \mid M)$ and $\Phi_{\mathrm{s}}(L \mid M)$ separately. The probability of observing a galaxy with luminosity $L_{i}$ in a halo with mass $M_{i}$ is given by
$p\left(L_{i} \mid M_{i}\right)=\Phi\left(L_{i} \mid M_{i}\right) p\left(M_{i}\right)$,
where $\Phi=\Phi_{\mathrm{c}}$ or $\Phi_{\mathrm{s}}$ for central and satellite galaxies, respectively, and the probability of observing a halo with mass $M_{i}, p\left(M_{i}\right)$ is given by the halo mass function. We estimate $p\left(M_{i}\right)$ by fitting a secondorder polynomial in $\log -\log$ to the distribution of halo masses from the Yang catalogue (converted from overdensities of 200 times the mean density of the Universe to 200 times the critical density of the Universe).

Given the observed luminosities, $\boldsymbol{L}_{\text {cen }}$, of a sample with $N_{\mathrm{c}}$ central galaxies, from equations (5) and (7) we obtain that the likelihood of a CLF described by the parameters $\boldsymbol{\theta}=\left\{\log L_{\mathrm{c}}, \sigma_{\mathrm{c}}\right\}$ is given by

$$
\begin{align*}
-\ln \mathcal{L}_{\mathrm{c}}\left(\boldsymbol{\theta} \mid \boldsymbol{L}_{\mathrm{cen}}\right)= & -\ln \prod_{i=1}^{N_{\mathrm{c}}} \Phi_{\mathrm{c}}\left(L_{i} \mid M_{i}\right) p\left(M_{i}\right) \\
= & \sum_{i=1}^{N_{\mathrm{c}}} \ln \left(2 \pi \sigma_{\mathrm{c}}\right)+\sum_{i=1}^{N_{\mathrm{c}}} \frac{\left(\log L_{i}-\log L_{\mathrm{c}}\right)^{2}}{2 \sigma_{\mathrm{c}}^{2}} \\
& -\sum_{i=1}^{N_{\mathrm{c}}} \ln p\left(M_{i}\right) \tag{8}
\end{align*}
$$

where $\sigma_{\mathrm{c}}$ and $L_{\mathrm{c}}$ are functions of the halo mass $M_{i}$. Similarly, for a sample of $N_{\mathrm{s}}$ satellite galaxies with luminosities $\boldsymbol{L}_{\text {sat }}$, the likelihood of a CLF described by the parameters $\boldsymbol{\theta}=\left\{\log L_{\mathrm{s}}, \alpha, \beta\right\}$ is given by

$$
\begin{align*}
-\ln \mathcal{L}_{\mathrm{s}}\left(\boldsymbol{\theta} \mid \boldsymbol{L}_{\mathrm{sat}}\right)= & -\ln \prod_{i=1}^{N_{\mathrm{s}}} \Phi_{\mathrm{s}}\left(L_{i} \mid M_{i}\right) p\left(M_{i}\right) \\
= & \sum_{i=1}^{N_{\mathrm{s}}} \ln \left\{\frac{1}{\beta} \Gamma\left[\frac{\alpha+1}{\beta},\left(\frac{L_{\mathrm{min}}}{L_{\mathrm{s}}}\right)^{\beta}\right]\right\} \\
& -\sum_{i=1}^{N_{\mathrm{s}}}(\alpha+1) \ln \left(\frac{L_{i}}{L_{\mathrm{s}}}\right)+\sum_{i=1}^{N_{\mathrm{s}}}\left(\frac{L_{i}}{L_{\mathrm{s}}}\right)^{\beta} \\
& -\sum_{i=1}^{N_{\mathrm{s}}} \ln p\left(M_{i}\right) \tag{9}
\end{align*}
$$

where $\alpha, \beta$ and $L_{\mathrm{s}}$ are functions of halo mass $M_{i}$.
We performed the MLE, first for the centrals, using equation (8), and then for the satellites, using equation (9). We performed the minimization using the function mle2 from the R package bbmle (Bolker \& R Development Core Team 2017), adopting the method L-BFGS-B (Byrd et al. 1995), which allows us to specify upper and lower bounds for each variable. The uncertainties in the parameters were estimated by bootstrapping the sample 200 times.

### 3.1 Central galaxies

Fig. 1 displays the LF of central galaxies (solid squares), computed as the best-fitting CLFs averaged over the halo masses of the groups in our sample, i.e. $\Phi(L)=\sum_{i=1}^{N_{\text {groups }}} \Phi\left(L \mid M_{i}\right) / N_{\text {groups }}$. Our fits are very close to those of Y08.


Figure 1. Total LFs of central and satellite galaxies, computed as the bestfitting CLFs integrated over the halo masses of the groups in our sample, i.e. $\Phi(L)=\sum_{i=1}^{N_{\text {groups }}} \Phi\left(L \mid M_{i}\right) / N_{\text {groups }}$. For the satellites, we show the bestfitting models with two free parameters, i.e. when we fit the linear relation with $\log M_{\text {vir }}$ of only one of the parameters $\alpha$ (red dashed), $\log L_{\mathrm{s}}$ (green dotted), or $\beta$ (blue dot-dashed lines). The orange long-dashed, cyan solid, and purple short-long dashed lines indicate the models with 4 degrees of freedom, and the model shown as the black solid line has six free parameters. The data are shown as open circles (satellite galaxies) and filled squares (central galaxies). The LFs by Y08 are indicated as black dashed lines. The lower panel shows the residuals of the satellite LF relative to the ' $L_{\mathrm{s}}, \beta$ ' model.

Fig. 2 shows our best-fitting model in bins of halo mass. The relations among $\log L_{\mathrm{c}}, \sigma_{\mathrm{c}}$ and halo mass are given by

$$
\begin{align*}
\log \left(\frac{L_{\mathrm{c}}}{\mathrm{~L}_{\odot}}\right) & =(10.90 \pm 0.01)+(0.28 \pm 0.01) \log \left(\frac{M_{\mathrm{vir}}}{10^{14} \mathrm{M}_{\odot}}\right) \\
\frac{\sigma_{\mathrm{c}}}{\mathrm{~L}_{\odot}} & =(0.15 \pm 0.01)+(0.02 \pm 0.01) \log \left(\frac{M_{\mathrm{vir}}}{10^{14} \mathrm{M}_{\odot}}\right) \tag{10}
\end{align*}
$$

In comparison with Y08, our centrals have lower luminosity at low group mass and higher luminosity at high group mass. This is illustrated in Fig. 3, where we compare our CLF fits with the CLFs of Y08 and their fits. Their results are given in bins of $M_{180, m}$, i.e. the mass within spheres that are 180 times denser than the mean density of the Universe, which we converted to $M_{200, c}$. Moreover, their luminosities are corrected for evolution. After estimating it from our sample, we subtracted the evolutionary correction from the $\mathrm{Y} 08 \log L_{\mathrm{c}}$ values and fitted the relations with $\log M_{\text {vir }}$. We find $\log \left(L_{\mathrm{c}} / \mathrm{L}_{\odot}\right)=10.88+0.21 \log \left[M_{\text {vir }} /\left(10^{14} \mathrm{M}_{\odot}\right)\right]$ and $\sigma_{\mathrm{c}} / \mathrm{L}_{\odot}=0.15+0.01 \log \left[M_{\text {vir }} /\left(10^{14} \mathrm{M}_{\odot}\right)\right]$, with errors in the coefficients of 0.01 .

Our relation between $\log L_{\mathrm{c}}$ and halo mass is indeed steeper than that of Y08, and the central galaxies in haloes with $\log \left(M_{\text {vir }} / \mathrm{M}_{\odot}\right) \sim$ 15 are 25 percent ( 0.1 dex) more luminous than predicted by the Y08 CLFs. It is not clear what is the source of this discrepancy,


Figure 2. CLFs in four bins of group (halo) mass. Same notations as in Fig. 1. Each bin contains the same number of groups ( 480 groups). The luminosity ranges are the same for the four panels.
but it might be due to differences in the definitions of the samples: Y08 include galaxies up to $z=0.2$, while our sample contains only groups at $z<0.07$. Besides, their analysis was based on SDSS-DR4, while ours is defined from DR7, and we use DR12 photometric measurements, which include an improved background subtraction in very luminous galaxies as the centrals in the most massive groups. We discuss this further in Section 5.1.

### 3.2 Satellite luminosity functions

Following the same approach adopted for the central galaxies, we performed MLE to obtain the parameters of the satellite CLFs. We compared models with different degrees of freedom. We assume that the three parameters $L_{\mathrm{s}}$ (characteristic luminosity), $\alpha$ (faint-end
slope) and $\beta$ (bright-end shape) vary linearly with $\log M_{\text {vir }}$, but allow ourselves to fix instead one or several of these three parameters to their Y08 values (thus independent of $M_{\text {vir }}$ ). The models are listed in Table 2, where we indicate which of these relations are fitted in each model. In models 3,4 and $6, \alpha$ is fixed to the Y08 relation. The Y08 linear relations with $\log M_{\text {vir }}$ were obtained by fitting their values of $\alpha$ in bins of halo mass (middle panel in Fig. 4). In models 2 , 4 and 7 , we fix $\log L_{\mathrm{s}}=\log L_{\mathrm{c}}-0.25$ (as Y08), and in models 2, 3 and 5, $\beta=2$ (as Y08).

Fig. 1 displays the best-fitting CLFs integrated over the halo masses of the groups in our sample and Fig. 2 shows our bestfitting models in bins of halo mass. The relations with halo mass of $\log \left(L_{\mathrm{c}} / L_{\mathrm{s}}\right), \alpha$ and $\beta$ are presented in Fig. 4 and in Table 3. We find a positive correlation between $\log \left(L_{\mathrm{c}} / L_{\mathrm{s}}\right)$ and $\log M_{\text {vir }}$ for all the


Figure 3. Relation between $L_{\mathrm{c}}$ (top) and $\sigma_{\mathrm{c}}$ (bottom) with the group halo mass. Our best-fitting relations and errors are indicated by the solid lines and shaded areas. The black diamonds represent the results from Y08 in bins of $M_{\mathrm{vir}}$, and the dashed lines show the best linear fits to their values.

Table 2. Comparison between satellite CLF models. (1) Free parameters of the model; (2) p-values of a KS and (3) AD tests; (4) $\Delta \mathrm{BIC}$ and (5) $\Delta \mathrm{AIC}$ relative to the model with the lowest BIC (model 6) and AIC (model 8) values. Since model 1 has no free parameters, the values of BIC and AIC correspond to $-2 \ln \mathcal{L}_{\mathrm{s}}$.

| Model | $p$-values |  |  | $\Delta$ BIC |
| :--- | :---: | :---: | :---: | :---: |
|  | KS test | AD test | $\Delta$ AIC |  |
| (1) | (2) | (3) | (4) | $(5)$ |
| 1. Yang et al. (2008) | $2 \times 10^{-4}$ | $9 \times 10^{-6}$ | 64 | 100 |
| 2. $\alpha$ | $6 \times 10^{-1}$ | $6 \times 10^{-1}$ | 4.7 | 25.7 |
| 3. $L_{\mathrm{s}}$ | $2 \times 10^{-2}$ | $7 \times 10^{-3}$ | 66.4 | 87.4 |
| 4. $\beta$ | $7 \times 10^{-5}$ | $4 \times 10^{-6}$ | 123.1 | 144.1 |
| 5. $L_{\mathrm{s}}, \alpha$ | $7 \times 10^{-1}$ | $5 \times 10^{-1}$ | 16.0 | 22.3 |
| 6. $L_{\mathrm{s}}, \beta$ | $1 \times 10^{-1}$ | $4 \times 10^{-2}$ | $\mathbf{0 . 0}$ | 6.4 |
| 7. $\alpha, \beta$ | $6 \times 10^{-1}$ | $6 \times 10^{-1}$ | 5.6 | 12.0 |
| 8. $L_{\mathrm{s}}, \alpha, \beta$ | $2 \times 10^{-1}$ | $7 \times 10^{-2}$ | 8.3 | $\mathbf{0 . 0}$ |

models where we allow $L_{\mathrm{s}}$ to vary, except for model 5, for which $\log \left(L_{\mathrm{c}} / L_{\mathrm{s}}\right) \sim 0.27$ (basically constant). As in Y08, we also find that $\alpha$ decreases with halo mass; however, our best-fitting relations are steeper than those of Y08. Finally, $\beta$ is roughly constant for models 4 and 5 , but in models 6 and 8 , it varies from $\beta \sim 3.4$ to $\sim 1.4$ between $\log \left(M_{\text {vir }} / \mathrm{M}_{\odot}\right)=13$ and 15 .

We used the Kolmogorov-Smirnov (KS) and Anderson-Darling (AD) tests to check the goodness of fit of each model. The resulting $p$-values are shown in Table 2. The $p$-values of the AD tests are typically smaller than those of the KS tests, since the former is more sensitive to the tails of the distributions. Both tests indicate that models 1 (Y08), 3 (free $\log L_{\mathrm{s}}$ ) and 4 (free $\beta$ ) do not provide a good description of the data $(p<0.05)$.

We also computed the Bayesian information criterion (BIC) and the Akaike information criterion (AIC). These criteria are a measure of the relative quality of statistical models for a given set of data, and the model with the lowest BIC/AIC is preferred. There is strong evidence against a model when $\Delta \mathrm{BIC}>6$, and it is decisive when $\Delta \mathrm{BIC}>10$ (Kass \& Raftery 1995; Raftery 1995).


Figure 4. Relations with group halo mass of $\log \left(L_{\mathrm{c}} / L_{\mathrm{s}}\right)$ (upper), $\alpha$ (middle), $\beta$ (lower panel) for CLFs with different degrees of freedom.

Table 2 shows the values of BIC and AIC relative to the best models according to each of these criteria (i.e. models 6 and 8 , respectively). There is a disagreement between these two criteria: BIC indicates that model 6 is the best model, while model 8 is the preferred one according to AIC (BIC penalizes model complexity more heavily, while AIC leads to more complex models.).

Therefore, models 1, 3 and 4 can be clearly rejected based on the $p$-values, $\Delta \mathrm{BIC}$ and $\Delta \mathrm{AIC}$. Models 2,5 and 7 can be discarded according to $\triangle \mathrm{BIC}$ or $\triangle \mathrm{AIC}$ (or both). Model 8 corresponds to the best model based on $\triangle \mathrm{AIC}$, but there is strong evidence against it based on $\triangle B I C$. Finally, model 6 is the best model based on $\triangle B I C$, and there is only marginal rejection based on the AD test and $\Delta \mathrm{AIC}$. Therefore, we adopt model 6 as our standard model throughout the paper. Moreover, model 6 has the advantage of adopting the value of $\alpha$ from Y08 who probe the CLFs to fainter luminosities with their flux-limited sample. However, as we discuss later, the main conclusions of our study do not depend on the particular choice of the model.

Table 3. Parameters of the satellite CLF models. (1) Free parameters of the model; (2), (4) and (6) Values of the parameters $L_{\mathrm{s}}, \alpha$ and $\beta$ for $M_{\mathrm{vir}}=10^{14} \mathrm{M}_{\odot}$, respectively. (3), (5) and (7) Slopes of the linear relations with $\log M_{\mathrm{vir}}$. Our preferred model is indicated in bold.

| Model | $\begin{gathered} \log \left(L_{\mathrm{s}} / \mathrm{L}_{\odot}\right) \\ {\left[M_{\text {vir }}=10^{14} \mathrm{M}_{\odot}\right]} \end{gathered}$ | $\mathrm{d} \log L_{\mathrm{s}} / \mathrm{d} \log M_{\mathrm{vir}}$ | $\left[\begin{array}{c} \alpha \\ {\left[M_{\mathrm{vir}}=10^{14} \mathrm{M}_{\odot}\right]} \end{array}\right.$ | $\mathrm{d} \alpha / \mathrm{d} \log M_{\mathrm{vir}}$ | $\begin{aligned} & \beta \\ & {\left[M_{\mathrm{vir}}=10^{14} \mathrm{M}_{\odot}\right]} \end{aligned}$ | $\mathrm{d} \beta / \mathrm{d} \log M_{\mathrm{vir}}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| (1) | (2) | (3) | (4) | (5) | (6) | (7) |
| 1. Yang et al. (2008) | $10.63 \pm 0.01$ | $0.21 \pm 0.01$ | $-1.25 \pm 0.01$ | $-0.33 \pm 0.02$ | 2.00 | 0.00 |
| 2. $\alpha$ | 10.65 | 0.28 | $-1.18 \pm 0.04$ | $-0.61 \pm 0.03$ | 2.00 | 0.00 |
| 3. $L_{\mathrm{s}}$ | $10.65 \pm 0.01$ | $0.21 \pm 0.01$ | -1.25 | -0.33 | 2.00 | 0.00 |
| 4. $\beta$ | 10.65 | 0.28 | -1.25 | -0.33 | $1.98 \pm 0.11$ | $-0.05 \pm 0.07$ |
| 5. $L_{\mathrm{s}}, \alpha$ | $10.63 \pm 0.02$ | $0.28 \pm 0.01$ | $-1.13 \pm 0.05$ | $-0.65 \pm 0.03$ | 2.00 | 0.00 |
| 6. $L_{\mathrm{S}}, \boldsymbol{\beta}$ | $10.67 \pm 0.01$ | $\mathbf{0 . 1 9} \pm \mathbf{0 . 0 1}$ | -1.25 | -0.33 | $\mathbf{2 . 3 8} \pm \mathbf{0 . 1 4}$ | $\mathbf{- 0 . 9 8} \pm 0.09$ |
| 7. $\alpha, \beta$ | 10.65 | 0.28 | $-1.17 \pm 0.04$ | $-0.64 \pm 0.03$ | $2.07 \pm 0.12$ | $-0.26 \pm 0.08$ |
| 8. $L_{\mathrm{S}}, \alpha, \beta$ | $10.68 \pm 0.02$ | $0.22 \pm 0.01$ | $-1.25 \pm 0.05$ | $-0.48 \pm 0.03$ | $2.35 \pm 0.20$ | $-0.76 \pm 0.10$ |

## 4 MAGNITUDE-GAP STATISTICS FROM THE CONDITIONAL LUMINOSITY FUNCTIONS

### 4.1 Gap statistics, predictions from the conditional luminosity function and previous results

We now test the CLF using the statistics of the magnitude gap $\Delta \mathcal{M}_{12}=\mathcal{M}_{2}-\mathcal{M}_{1}$. Fig. 5 illustrates five statistics of the gap as a function of group richness (number of galaxies, $N_{\text {gals }}$, with $\mathcal{M}_{r}^{\text {Petro }} \leq$ -19.57 , left) and halo mass $M_{\text {vir }}$ (right). We can see that $\Delta \mathcal{M}_{12}$ slightly decreases with $N_{\text {gals }}$ and $\log M_{\text {vir }}$ (Figs 5a and b). The median value of the magnitude gap varies from $\Delta \mathcal{M}_{12} \sim 1.0 \mathrm{mag}$ in groups with $N_{\text {gals }}=2$ to $\Delta \mathcal{M}_{12} \sim 0.9 \mathrm{mag}$ in groups with high multiplicity ( $N_{\text {gals }} \gtrsim 10$ ). A similar variation is observed in Fig. 5(b): $\Delta \mathcal{M}_{12}$ decreases from $\sim 1.0 \mathrm{mag}$ in low-mass haloes $\left[\log \left(M_{\text {vir }} / \mathrm{M}_{\odot}\right)<\right.$ 13.1] to $\Delta \mathcal{M}_{12} \sim 0.9 \mathrm{mag}$ in haloes with $\log \left(M_{\text {vir }} / \mathrm{M}_{\odot}\right) \gtrsim 14$. The Spearman and Kendall correlation tests confirm the correlation of gap with $N_{\text {gals }}\left(p=8 \times 10^{-4}\right)$ and $\log M_{\text {vir }}\left(p=7-8 \times 10^{-4}\right)$.

The fraction of LGGs decreases with $N_{\text {gals }}$ and halo mass (panels c and d in Fig. 5). Around 12 percent of low-multiplicity groups have large magnitude gaps, and this fraction decreases to $\sim 4$ percent among rich groups. The fraction of SGGs is roughly constant and varies between 10 and 20 percent.

We applied the Tremaine \& Richstone (1977) statistics (see Section 1 and equation 1) to test if the observed magnitude gaps are consistent with a cumulative LF that has power-law behaviour at the faint and bright ends. The values of $T_{1}$ and $T_{2}$ are presented in Figs $5(\mathrm{~g})$ to (j). They were computed in the bins of $N_{\text {gals }}$ and $\log M_{\text {vir }}$. We find very low values of $T_{1}(\lesssim 0.5)$ and $T_{2}(\lesssim 0.8)$. This suggests that the LFs of galaxies, conditional to the group richness or group mass, are not consistent with power-law behaviour of the CLF at its faint and bright ends.

These very low values of $T_{1}$ and $T_{2}$ may be caused by the low multiplicity of most of the considered bins of richness, as Mamon (1987) noticed that the values of $T_{1}$ and $T_{2}$ are underestimated for samples of groups with fewer than 30 galaxies. We checked this effect by building mock samples of groups, where the galaxy luminosities are drawn from a Schechter (1976) LF, with parameters taken from Blanton et al. (2003b): faint-end slope $\alpha=-1.05$ and $L / L^{*}>0.24$ (given our limiting absolute magnitude of $\mathcal{M}_{r}^{\text {Petro }} \leq$ -19.57 ). In each bin of richness, we built the groups with the same distribution of multiplicity. For each sample of groups, we computed the fraction of groups with high $\left(\Delta \mathcal{M}_{12} \geq 2\right)$ and low ( $\Delta \mathcal{M}_{12} \leq 0.3$ ) magnitude gaps, as well as the values of $T_{1}$ and $T_{2}$. We computed the fifth and $95^{\text {th }}$ percentiles of these values for 1000 random samples built in this manner, as shown in red arrows in Figs 5(c), (e), (g) and (i). This exercise led to typically 30-40 times fewer LGGs on
average (Fig. 5c), typically two to three times more SGGs (Fig. 5d), and much higher values of $T_{1}$ ( 95 percent of the mock samples have $T_{1}$ values that are $\sim 0.6$ units higher than measured in the SDSS in every bin of richness.) and $T_{2}$. Very similar results are found with $\alpha=-1.3$. This confirms that the fractions of LGGs and SGGs in the SDSS are inconsistent with a single Schechter form for the CLF.

The alternative is that the CLF has an additional component specific to the BGGs, as described in Section 3. We now go one step further and test whether these CLFs are consistent with the distribution of gaps observed in the SDSS. For each bin of $M_{\text {vir }}$ and $N_{\text {gals }}$ shown in Fig. 5, we sampled the CLFs (model 6 in Table 2) 1000 times, building a sample of groups with exactly the same characteristics (i.e. number of groups and distributions of $M_{\text {vir }}$ and $N_{\text {gals }}$ ) as that of the sample in the bin. The results obtained by sampling the CLFs are shown as the shaded areas in Fig. 5.

The 84th percentiles of the distributions of $\Delta \mathcal{M}_{12}$ in bins of $N_{\text {gals }}$ (Fig. 5a) and $\log M_{\text {vir }}$ (Fig. 5b) from the CLF sampling agree with the observations. However, the 16th percentiles are higher than the ones measured from our sample. This can be also seen in the fraction of LGGs and SGGs: while we find an agreement for the fraction of LGGs (Figs 5c and d), the sampling of the CLF underestimates the number of SGGs, and the discrepancy increases with decreasing $N_{\text {gals }}$ and halo mass (Figs 5c and d). Moreover, the CLF sampling results in $T_{1}$ and $T_{2}$ values that are, respectively, 0.02-0.1 (Figs 5g and h ) and $0.08-0.2$ (Figs 5 i and j ) lower than observed.

The discrepancies between statistics might be due to the fact that the $\Delta \mathcal{M}_{12}$ values of the sample in Fig. 5 are computed within $0.5 r_{\text {vir }}$, while we are sampling the CLFs that were fitted to all galaxies within $2 r_{\text {vir }}$ (see Section 3). However, the statistics of gap obtained with SBGGs defined within $2 r_{\text {vir }}$ are also discrepant with those predicted by these CLFs (see Section 4.4). In addition, if we consider the CLF of galaxies within $0.5 r_{\mathrm{vir}}$, we get similar results and are still not able to reproduce the gap statistics.

We also compared our results with the $T_{1}$ and $T_{2}$ values from studies by DG +12 , LS06, LOM10 and $\mathrm{S}+14$. Our $T_{1}$ values are much lower than those found by LS06 and LOM10. As indicated in Fig. 5(h), LS06 found $T_{1}=0.75 \pm 0.1$ for a sample of nearby rich SDSS clusters, while we obtain $T_{1} \sim 0.4$ for groups with halo masses similar to those of their sample. LOM10 also found much higher values of $T_{1}: 0.84 \pm 0.01$ and $0.70 \pm 0.1$ for low- and high-luminosity SDSS clusters, respectively. Interestingly, LOM10 obtain a higher $T_{1}$ for low-luminosity clusters compared to highluminosity clusters, in contrast to our results, which show that $T_{1}$


Figure 5. Statistics of the magnitude gap as a function of group richness (left) and halo mass (right): magnitude gap ( $\mathrm{a}, \mathrm{b}$ ); fractions of groups with $\Delta \mathcal{M}_{12}>2.0(\mathrm{c}, \mathrm{d})$ and $<0.3 \mathrm{mag}(\mathrm{e}, \mathrm{f}) ; T_{1}(\mathrm{~g}, \mathrm{~h})$ and $T_{2}(\mathrm{i}, \mathrm{j})$. In panels (a) and (b), the black squares and error bars represent the median, the 16th and the 84 th percentiles of the $\Delta \mathcal{M}_{12}$ distribution in bins of $N_{\text {gals }}$ and $\log M_{\text {vir }}$. The Kendall $(\tau)$ and Spearman $(\rho)$ rank correlation coefficients, as well as the $p$-values, are indicated in each panel. In panels $c-\mathrm{f}$, the error bars represent the binomial errors. In panels c and d , the errors were estimated by bootstrapping the sample 1000 times, and error bars indicate the 16th and 84th percentiles. The shaded grey areas represent the 16th and 84th percentiles of the results from sampling the CLF of model 6 in Table 2. The red symbols in panels $g$ and $i$ indicate the 5th (arrows pointing upward) or 95th (arrows pointing downward) percentiles of the values expected for the Schechter (1976) LF with $\alpha=-1.05$. We also show the $T_{1}$ and $T_{2}$ values from studies by DG+12, LS06, LOM10 and Shen et al. (2014, S+14). In the panels showing the statistics versus halo mass, the solid black lines indicate the equivalent $\pm 1 \sigma$ range of values estimated from the SAM of Henriques et al. (2015).
increases with increasing halo mass. These comparisons should be taken with caution because LS06 and LOM10 consider much wider bins of group mass that causes blurring between the locations of the central and satellite components of the CLFs. DG+12 found $T_{1}=0.51 \pm 0.06$ for a complete sample of compact groups drawn from the 2MASS catalogue, which is lower than the values obtained by the previous authors, but still higher than ours $\left(T_{1} \sim 0.4\right.$;

Fig. 5 g ). 80 per cent of their groups have richness $N_{\text {gals }}=4$, therefore these compact groups span a narrower range of halo masses than the cluster samples of LS06 and LOM10, yet a wider range than our own mass bins. This might explain why their $T_{1}$ and $T_{2}$ values are lower than other published studies, but still larger than ours.

On the other hand, the discrepancy between our results and those by $\mathrm{S}+14\left(T_{1} \sim 0.7-0.85\right.$; Fig. 5g) cannot be explained by the bin widths, since $S+14$ computed the Tremaine \& Richstone statistics in narrow bins of group richness. Their $\sigma\left(\mathcal{M}_{1}\right)$ decreases from $\sim 0.57$ to $\sim 0.4$ for groups with $N_{\text {gals }}=2$ to $\geq 21$. In contrast, we find that the dispersion increases from $\sigma\left(\mathcal{M}_{1}\right) \sim 0.30$ to 0.41 within the same richness interval, which is consistent with the values of $\sigma_{c}$ of Y08 and our CLF fits (Fig. 3). Our smaller values of $\sigma\left(\mathcal{M}_{1}\right)$ are at least partially due to the cleaning for spectroscopic incompleteness; if we do not require the BGG to be correctly identified, we obtain $\sigma\left(\mathcal{M}_{1}\right)=0.41$ for groups with $N_{\text {gals }}=2$. $\mathrm{S}+14$ also found smaller $\left\langle\Delta \mathcal{M}_{12}\right\rangle$ values than we do, varying from $\sim 0.8$ to $\sim 0.4$ for poor to rich groups. Instead, we find $\left\langle\Delta \mathcal{M}_{12}\right\rangle$ decreasing with richness from 1.0 to 0.9 for $R_{\max }=0.5 r_{\text {vir }}$ (Fig. 5a) and from 1.1 to 0.5 for $R_{\max }=2 r_{\text {vir }}$ (Fig. 7k). Since $\mathrm{S}+14$ selected the SBGGs among all group galaxies, their results should be compatible with ours for $R_{\max }=2 r_{\text {vir }}$. However, even in that case our $T_{1}$ values are much smaller than those obtained by $\mathrm{S}+14$ for groups with less than approximately five galaxies (Fig. 7q). Perhaps the discrepancy with $\mathrm{S}+14$ arises from their using model magnitudes, while we (and Y08) use Petrosian magnitudes.

Finally, as shown in Figs 5(i) and (j), the values of $T_{2}$ obtained by LS06 $\left(T_{2}=0.86 \pm 0.1\right)$, LOM10 $(0.77,0.79 \pm 0.01), \mathrm{DG}+12$ $(0.70 \pm 0.06)$ and $S+14(0.70-0.85)$ are in better agreement with our results $\left(T_{2} \sim 0.75-0.80\right)$.

### 4.2 Gap statistics from a semi-analytical model

We compared our results with predictions from the SAM of Henriques et al. (2015), run on the Millennium-II simulations (Boylan-Kolchin et al. 2009). We extracted the snapshot corresponding to $z=0$ from the Henriques2015a..MRIIscPlanck1 table in the Virgo-Millennium data base of the German Astrophysical Virtual Observatory $\left(\mathrm{GAVO}^{7}\right)$.

From the simulation box extracted from GAVO, we built a mock flux-limited, SDSS-like sample of groups and galaxies, following the steps of Duarte \& Mamon (2014). Since the simulation box is not large enough to produce the SDSS-like group catalogue, we replicated the simulation box along the three Cartesian coordinates, then placed an observer at some position and mapped the galaxies on the sky. The absolute magnitudes in the $r$ band (including internal dust extinction) were converted to apparent magnitudes, and the flux limit of the MGS of the SDSS, $m_{r}<17.77$, was applied. We added errors of 0.2 dex to the group halo masses and of 0.08 mag to the absolute magnitudes, as determined by Duarte \& Mamon (2015) for the galaxies in the SDSS/MGS. We then selected the galaxies and groups from the mock catalogue following the same selection criteria that we had applied to the observations and presented in Section 2.

We obtained the distribution of $\Delta \mathcal{M}_{12}$, fractions of LGGs and SGGs, as well as the $T_{1}$ and $T_{2}$ values as a function of halo mass (solid black lines in the right panels of Fig. 5). The median value of the $\Delta \mathcal{M}_{12}$ distribution from the SAM is $\sim 0.3 \mathrm{mag}$ lower (Fig. 5b)

[^4]and the SAM predicts fewer LGGs (Fig. 5d) than observed. On the other hand, the fractions of SGGs is compatible with the observations, at least for groups with $\log \left(M_{\text {vir }} / \mathrm{M}_{\odot}\right) \gtrsim 13.3$.

We observe much lower $T_{1}$ values than those predicted from the SAM (Fig. 5h). Although higher than those computed from observations, the $T_{1}$ values from the SAM are still lower than those obtained by LOM10 for groups with $\log \left(M_{\text {vir }} / \mathrm{M}_{\odot}\right) \lesssim 13.7$, but they are in agreement with LOM10 and LS06 for $\log \left(M_{\text {vir }} / \mathrm{M}_{\odot}\right)>$ 13.7. But, again, the comparison with LOM10 and LS06 is not a fair one since they consider much wider bins of group mass than we do. The discrepancies between the $T_{2}$ values from the SAM and SDSS sample are much smaller than that of $T_{1}$, with lower observed $T_{2}$ values than those from the SAM (Fig. 5j).

### 4.3 Magnitude gap and groups with two central galaxies

The fractions of SGGs from the CLF sampling are lower than observed, indicating that the SBGGs are brighter than predicted by the satellite CLF. Therefore, one may ask whether, in some SGGs, both the BGG and the SBGG come from the same distribution.

Indeed, when groups merge, the magnitude gap will be altered at the time of the group merger and again when the two brightest (and usually most massive) galaxies merge. Assume that group 1 is more massive than group 2, and that its BGG is more luminous and massive than that of group 2. Suppose also that the two groups merge before the time when the SBGG and BGG of group 1 merge together. This would suggest that, before the group merger, the BGG of group 2 is more massive, hence more luminous than the SBGG of group 1. Thus, when the groups merge, the magnitude gap will suddenly decrease to a lower value. At the same time, one can consider that the merged group has two centrals. Once the original BGG of defunct group 2 merges into the BGG of the merged group (that of the original group 1), after orbital decay by dynamical friction, the group will have a single central galaxy, and the gap will now correspond to the luminosity ratio of the galaxy created by the merger of the two BGGs with the SBGG of the original group 1. This gap will be larger than the original gap of group 1 before the group-group merger. This simple scenario is complicated by many features such as additional group-group and galaxy-galaxy mergers. But, to the first order, one expects that groups are more likely to have two centrals if they have small gaps.

We therefore propose a model where some groups have two galaxies whose luminosities follow the CLF of central galaxies (equation 5). We assume that the probability of having such a group is the cumulative distribution function of a Gaussian distribution of $\Delta \mathcal{M}_{12}$ :
$p_{2 \mathrm{c}}\left(\Delta \mathcal{M}_{12}\right)=\frac{1}{2}\left[1-\operatorname{erf}\left(\frac{\Delta \mathcal{M}_{12}-\mu}{\sqrt{2} \sigma}\right)\right]$.
We then repeat the exercise of sampling the CLFs allowing some groups to have two central galaxies. One may think that we would then need to recompute the CLF, but while the central and satellite CLFs will change, the total CLF will not, as suggested by the lack of sensitivity, seen in Fig. B1, of the gap statistics to the choice of free parameters in the analytical fits of the CLF.

We find that the observed gap statistics are now very well reproduced when we sample the CLF with equation (11) for $\mu=0.4$ and $\sigma=0.2$ (blue shaded regions in Fig. 6). Therefore, it seems that in groups with small gaps, the SBGG luminosities follow the same CLF as the BGG, i.e. they have two central galaxies. This can be a consequence of mergers of groups. Note that this functional form


Figure 6. Same as Fig. 5, but showing gap statistics when we allow some groups to have 'two central galaxies' according to equation (11) with $\mu=0.4$ and $\sigma=0.2$ (shaded blue areas).
does not have a physical motivation; it merely provides an adequate description of the observations.

### 4.4 Gap statistics versus $\boldsymbol{R}_{\text {max }}$

Since the SBGGs can be further than $0.5 r_{\text {vir }}$ from the BGG, the statistics of magnitude gap might depend on the maximum radius, $R_{\max }$, allowed us to obtain the SBGG. To investigate how different $R_{\text {max }}$ affects the statistics of the magnitude gap, we now repeat the analysis for the sample defined within 1 and $2 r_{\text {vir }}$ (see Section 2.3 and Table 1). The results are presented in Fig. 7.

As was the case for $R_{\max }=0.5 r_{\mathrm{vir}}$, we are not able to reproduce the observed gap statistics when we assume that all groups have only one central galaxy (grey shaded regions). However, when we allow some groups to have two central galaxies, as described in Section 4.3, the observed gap statistics are reproduced by those obtained by sampling the CLF, allowing one or two centrals (blue shaded regions).

In Fig. 8, we show a more complete picture of how the magnitudegap statistics vary with the maximum distance allowed to find the


Figure 7. Same as Fig. 6, but changing the maximum radius, $R_{\max }$, allowed for obtaining the SBGG to $R=1.0 r_{\text {vir }}$ (left) and $2.0 r_{\text {vir }}$ (right).

SBGG. We present the magnitude gap, the fractions of groups with $\Delta \mathcal{M}_{12}>2.0$ and $<0.3 \mathrm{mag}$, as well as the $T_{1}$ and $T_{2}$ values as a function of $R_{\max }$ in two bins of halo mass $\left[\log \left(M_{\text {vir }} / \mathrm{M}_{\odot}\right)<13.4\right.$ and $>13.4]$.
For low-mass groups, the gap statistics vary very slightly with $R_{\max }$. For $R_{\max } / r_{\text {vir }}=0.5$ to 2 , the median $\Delta \mathcal{M}_{12}$ value decreases 0.14 mag, the fraction of LGGs varies from 11.4 percent to 8.0 percent, SGGs from 18 percent to 20 percent, and $T_{1}\left(T_{2}\right)$ from 0.31 ( 0.78 ) to $0.35(0.82)$. On the other hand, for $\log \left(M_{\text {vir }} / \mathrm{M}_{\odot}\right)>$ 13.4, there is a strong variation from $R_{\max } / r_{\text {vir }}=0.5$ to 1.0 : we observe a decrease of 0.25 mag in the median $\Delta \mathcal{M}_{12}$ value, and the fractions of LGGs and SGGs vary from 7.6 percent to 2.4 percent and 16 percent to 26 percent, respectively. The value of $T_{1}\left(T_{2}\right)$ varies from $0.41(0.78)$ to $0.52(0.90)$. For $R_{\max }>$ $r_{\mathrm{vir}}$, all the quantities vary very little with $R_{\max }$. All these results are very well reproduced when we allow groups with two central galaxies.

The predictions from the SAM here are strikingly different from the observations, in particular for low-mass groups. However, the results based on the SAM may depend on the choice of the SAM.

Finally, our results remain inconsistent with the values expected for a single Schechter LF for all $R_{\max }$ values.

## 5 DISCUSSION

### 5.1 The bright end of the conditional luminosity functions

The differences that we find between our CLFs and those of Y08 can arise from differences in the sample definition. Y08 include galaxies up to $z=0.2$, while our sample contains only groups at $z<0.07$. In addition, their analysis was based on SDSS-DR4, while our sample was defined from DR7, and we use DR12 photometry (which is thought to be more accurate for massive ellipticals with shallow surface brightness profiles). Y08 used a flux-limited sample, while we use a doubly complete one, and they applied an evolution correction to the magnitudes, while we do not (although we estimate the correction from our sample, these estimates are not precise.). Finally, the methods used to obtain the best-fitting CLF are different: while Y08 appear to have applied $\chi^{2}$ minimization to the binned data, we use MLE with no binning.


Figure 8. Statistics of the magnitude gap versus the maximum radius, $R_{\max }$, allowed for obtaining the SBGG: magnitude gap (a, b); fractions of LGGs $\left(\Delta \mathcal{M}_{12}>2.0, \mathrm{c}, \mathrm{d}\right)$ and SGGs $\left(\Delta \mathcal{M}_{12}<0.3, \mathrm{e}, \mathrm{f}\right) ; T_{1}(\mathrm{~g}, \mathrm{~h})$ and $T_{2}(\mathrm{i}, \mathrm{j})$. The left and right panels show the results for groups with $\log \left(M_{\text {vir }} / \mathrm{M}_{\odot}\right)<$ 13.4 and $>13.4$, respectively. The shaded areas represent the 16 percent and 84 percent when we assume that all groups have only one central (shaded grey) and when we allow some groups to have two central galaxies according to equation (11) with $\mu=0.4$ and $\sigma=0.2$ (shaded blue areas). The dashed and solid black lines indicate the values and uncertainties estimated from the SAM by Henriques et al. (2015).

We obtain steeper $\alpha$ versus $\log M_{\text {vir }}$ relations than Y08 (Fig. 4). From $\log \left(M_{\text {vir }} / \mathrm{M}_{\odot}\right)=13$ to 15 , we find that $\alpha$ decreases (steepens) from $\sim-0.6$ to -1.8 (model 2), while $\alpha$ varies between $\sim-0.9$ and -1.6 for the Y08 model. This extra steepening of the faint-end slope in our fits may be a consequence of our cleaning the sample to ensure the correct identification of the SBGG (see Section 2.1). Indeed, if we do not clean the sample to ensure that SBGGs are correctly identified and redo our CLF fits (Fig. A1 and Table A1), we obtain a shallower relation between $\alpha$ and group halo mass (Fig. A2).

The characteristic luminosity of the satellite CLF (relative to the characteristic central luminosity) increases with $\log M_{\text {vir }}$, while the bright-end cut-off $\beta$ decreases with increasing halo mass (Fig. 4). Since high $\beta$ values can be compensated by $\operatorname{low} \log \left(L_{\mathrm{c}} / L_{\mathrm{s}}\right)$ ratios, one might ask whether there is degeneracy between these two pa-


Figure 9. Illustration of the anti-correlation between the bright-end shape of the satellite CLF and the characteristic luminosity (relative to the characteristic central luminosity) for different halo masses. The figure shows the best-fitting parameters of model 6 obtained by bootstrapping the group sample 1000 times.
rameters. In Fig. 9 we show that at a fixed and low halo mass, there is indeed an anti-correlation between these two parameters, but it is much smaller than the variation that we observe between low- and high-mass haloes [ $\left.\log \left(M_{\text {vir }} / \mathrm{M}_{\odot}\right)=13-15\right]$.

### 5.2 Magnitude-gap statistics versus group richness and halo mass

When comparing our statistics on the magnitude gaps with those of previous authors, one should be careful to note that two effects alter these statistics. First, samples (bins) of larger richness will naturally lead to smaller gaps. Secondly, the Tremaine-Richstone statistics $T_{1}$ and $T_{2}$, which measure departures from one-component LFs or CLFs, will be less efficient in wide ranges of group halo masses. Indeed, since the CLF depends on group mass with more luminous characteristic central and satellite luminosities at higher group mass, wider mass bins will wash out the separation between centrals and satellites, and will lead to higher values of $T_{1}$ and $T_{2}$.

Using SAMs of galaxy formation, Dariush et al. (2007), DíazGiménez, Muriel \& Mendes de Oliveira (2008) and Farhang et al. (2017) have shown that major progenitors of LGGs assemble half their mass earlier than do regular groups of similar final mass, as previously shown by D'Onghia et al. (2005) using hydrodynamical cosmological simulations. Since the major progenitors of presentday massive haloes assembled their mass later than those of lower mass haloes (van den Bosch 2002), one can conclude that early mass assembly is linked to both low final-mass haloes and to present-day LGGs, hence to a possible anti-correlation of magnitude gap with halo mass. Physically, given the flattening of the stellar mass versus halo mass relation at high halo masses (e.g. Yang et al. 2009), the dynamical friction times should be longer in high-mass clusters than in low-mass groups, so if the wide magnitude gaps of LGGs are caused by mergers, one indeed expects an anti-correlation of magnitude gap with group halo mass. Our results confirm this anticorrelation of magnitude gap with group mass (Fig. 5b): the negative
trend is significant ( $p=7-8 \times 10^{-4}$ ). The most massive LGG in our sample (defined within $0.5 r_{\text {vir }}$ ) has $\log \left(M_{\text {vir }} / \mathrm{M}_{\odot}\right)=14.2$, while the most massive systems among the SGGs have halo masses as high as $\log \left(M_{\text {vir }} / M_{\odot}\right)=15.1$.

Applying a SAM of galaxy formation (De Lucia \& Blaizot 2007) to the Millennium simulation (Springel et al. 2005), Díaz-Giménez et al. (2008) predict that only 5.5 percent of the systems with $\log \left(M_{\text {vir }} / \mathrm{M}_{\odot}\right)>13.7$ are LGGs, in agreement with the fraction that we find in the same mass range, which is $5.6 \pm 1.2$ percent ( 21 out of 377 groups). Tavasoli et al. (2011) also find an anti-correlation between the fraction of LGGs and group mass for both a sample of SDSS groups and from the outputs of a SAM (Bower et al. 2006) run on the Millennium simulation. In addition, our CLFs also lead to a decreasing $\Delta \mathcal{M}_{12}$ with increasing halo mass (see Section 4, in particular Fig. 5b).

On the other hand, this anti-correlation of gap with halo mass is in contradiction with the results by Hearin et al. (2013), who found that haloes with richness between 12 and 18 and $\Delta \mathcal{M}_{12} \leq 0.2 \mathrm{mag}$ are less massive than haloes of the same richness with $\Delta \mathcal{M}_{12} \geq$ 1.5 mag. They reached this conclusion by analysing a mock galaxy catalogue built by subhalo abundance matching of galaxy luminosities to the subhalo masses of the Bolshoi cosmological $N$-body simulations (Klypin, Trujillo-Gomez \& Primack 2011). If we repeat their selection by limiting our sample to groups that have 12 to 18 galaxies brighter than $\mathcal{M}_{r}^{\text {Petro }}=-19.5$, we find that the halo masses of groups with $\Delta \mathcal{M}_{12} \leq 0.2$ are compatible with those of groups with $\geq 1.5 \mathrm{mag}$, in disagreement with their results. This disagreement suggests that, while subhalo abundance matching reproduces the clustering of galaxies in a large range of masses and redshifts (Conroy, Wechsler \& Kravtsov 2006), it cannot match the tail end of the CLF. Part of this discrepancy may also lie in Hearin et al.'s use of a group FoF group finder to extract their groups that was nonoptimal in its adopted linking lengths (Duarte \& Mamon 2014) and non-optimal relative to the Yang et al. (2007) group finder (Duarte \& Mamon 2015).
Hearin et al. (2013) also analysed the FoF groups that Berlind et al. (2006) extracted from the SDSS. They find that, for fixed group velocity dispersion, groups with $\Delta \mathcal{M}_{12} \geq 1.5$ have fewer galaxies than groups with $\Delta \mathcal{M}_{12} \leq 0.2$. We confirm this result in Fig. 10(a) for $\Delta \mathcal{M}_{12}$ defined within $R_{\max }=2 r_{\text {vir }}$ (red versus blue points), but not for $R_{\max }=0.5 r_{\text {vir }}$ (orange versus green points). In addition, we show in Fig. 10(b) that the lower richness of LGGs is a consequence of their lower halo masses in a given $\sigma_{\mathrm{v}}$ bin. This is confirmed in Figs 10 (c) and (d), where we show that in bins of richness, the masses and velocity dispersions of LGGs are compatible with those of SGGs. Therefore, the conclusion by Hearin et al. that, for a fixed richness, LGGs are more massive than SGGs is based on their inversion of the richness versus velocity dispersion relation. In contrast, the differences between SGGs and LGGs in the direct relation of velocity dispersion versus richness are much weaker, and are even weaker in the Y08 mass versus richness relation. Moreover, in all cases any difference between SGGs and LGGs is weakened when we limit the samples to $0.5 r_{\text {vir }}$, to such an extent that the mass versus richness relations of SGGs and LGGs are identical (orange versus green points in Fig. 10d).

Motivated by the results of Hearin et al., More (2012) investigated analytically how the distribution of $\Delta \mathcal{M}_{12}$ varies with halo mass. He also concluded that, at fixed richness, SGGs tend to be more massive than LGGs. However, More assumes that the faint-end slope of the satellite CLFs is independent of halo mass (he adopted $\alpha=-1.17$ ). As shown in Fig. 4, $\alpha$ decreases with $\log M_{\text {vir }}$, and the faint end of the CLFs of massive haloes is steeper than what was assumed
by More. Although high-mass haloes have more satellite galaxies, for our (and Y08's) CLFs, the probability of these satellites having low luminosities is higher. Therefore, the assumption of a constant $\alpha$ results in a larger fraction of bright satellites than that obtained when $\alpha$ decreases with halo mass, leading to the results obtained by More.

### 5.3 Gap statistics from the conditional luminosity functions: variations with $\boldsymbol{R}_{\text {max }}$

It is clear from our results that the single Schechter-like LF is ruled out (Fig. 6), regardless of the $R_{\text {max }}$ adopted (Figs 7 and 8). In addition, we are able to reproduce the observed gap statistics with CLFs corresponding to the sum of a Schechter-like and a lognormal (Section 3) only if we allow some groups to have more than one central galaxy, with the probability of such a group being given by equation (11) with $\mu=0.4$ and $\sigma=0.2$. This result does not depend on the particular choice of the CLF model (see Appendix B).

Moreover, our results show that the statistics of the magnitude gap depend on the adopted $R_{\max }$ (Fig. 8). However, clusters display luminosity segregation at the bright end, beyond the central galaxy, in their inner regions (Adami, Biviano \& Mazure 1998), with higher giant to dwarf ratios in these cluster cores (Driver, Couch \& Phillipps 1998; Boué et al. 2008). It is thus surprising that the SBGGs of a large fraction ( $\sim 50$ percent) of our groups lie, in fact, outside $0.5 r_{\text {vir }}$.

### 5.3.1 Are gap statistics affected by inaccurate group identifications?

The poor identification of group members can lead to the fragmentation and merger of real-space groups. In the former case, the secondary fragments (which amount to $\sim 20$ percent of the groups obtained with the Yang et al. method) have random galaxy luminosities and are preferentially located at the outskirts of the true group. But this should not affect the gap as a function of the maximum projected radius. On the other hand, group merging by the group finder could lead to the observed results, since a different group is included within the real group and it is more likely to be on the outskirts of the group. The BGG of the merged group would be considered the SBGG of the extracted group, and it would affect the statistics of the bright end of the LF.

To investigate this issue, we compared the observed statistics of the gap with those predicted by a SAM viewed in projection. Since these groups are perfectly extracted, the effects of poor identification should not be present and the fractions of LGGs, as well as the statistics of $\Delta \mathcal{M}_{12}$, should be the same for all samples (i.e. regardless the $R_{\text {max }}$ used to define the gap). Yet, the comparison between Figs 6 and 7, which presents the results for the gap defined within $0.5,1.0$ and $2.0 r_{\text {vir }}$, shows that the statistics of the gap in the SAMs also depends on $R_{\max }$. This result is more evident in Fig. 8, where both the observations and the SAM show similar trends with $R_{\max }$ of LGG and SGG fractions, $T_{1}$, and $T_{2}$. However, the SAM predicts fewer LGGs and more SGGs than observed. In addition, the $T_{1}$ and $T_{2}$ values from the SAM are higher than those obtained from the SDSS groups.

### 5.3.2 Does the group luminosity function vary with R ?

Other possibility to explain the variation of gap statistics with $R_{\max }$ is that the CLFs vary with distance to the group centre. To


Figure 10. Relations between group velocity dispersion, mass and richness. SGGs with $\Delta \mathcal{M}_{12} \leq 0.2$ and LGGs with $\Delta \mathcal{M}_{12} \geq 1.5$ within $R_{\text {max }}=0.5 r_{\mathrm{vir}}$ $\left(2.0 r_{\mathrm{vir}}\right)$ are indicated as the green (blue) and orange (red) symbols. The error bars indicate uncertainties on the mean values from 1000 bootstraps.

Table 4. Comparison between the luminosity distributions of satellite galaxies at different distances from the group centre. $(1,3)$ Radial distance limits and $(2,4)$ number of galaxies in the two samples being compared; (5) KS and (6) AD p-values.

| Sample 1 |  | Sample 2 |  | $p$-values |  |
| :--- | :---: | :---: | :---: | :---: | :---: |
| $R / r_{\text {vir }}$ | $N_{\text {gals }}$ | $R / r_{\text {vir }}$ | $N_{\text {gals }}$ | KS | AD |
| $(1)$ | $(2)$ | $(3)$ | $(4)$ | $(5)$ | $(6)$ |
| $(0.0,0.5]$ | 4460 | $(0.5,1.0]$ | 3940 | 0.12 | 0.33 |
| $(0.5,1.0]$ | 3940 | $(1.0,1.5]$ | 2151 | $3 \times 10^{-8}$ | $2 \times 10^{-12}$ |
| $(1.0,1.5]$ | 2151 | $(1.5,2.0]$ | 544 | 0.02 | 0.05 |

investigate that, we compared the distribution of luminosities of satellite galaxies at $R \leq 0.5 r_{\text {vir }}, 0.5<R / r_{\text {vir }} \leq 1.0,1.0<R / r_{\text {vir }}$ $\leq 1.5$ and $1.5<R / r_{\text {vir }} \leq 2.0$. We applied the KS and AD tests, as summarized in Table 4.

We cannot rule out the possibility that the observed luminosities of satellites at $R \leq 0.5 r_{\text {vir }}$ and at $0.5<R / r_{\text {vir }} \leq 1.0$ come from the same distribution. On the other hand, the tests indicate that the distributions of satellite luminosities for samples at $0.5<R / r_{\text {vir }} \leq$ 1.0 and $1.0<R / r_{\text {vir }} \leq 1.5$ are very different from each other, with the outer regions lacking $L^{*}$ galaxies compared to the inner parts ( $p=2 \times 10^{-12}-3 \times 10^{-8}$ ). In addition, there is a marginal evidence that the distributions also change from $1.0<R \leq 1.5$ to $1.5<R$ $\leq 2.0(p=0.02-0.05)$. These results reproduce qualitatively the variation of LFs with position, with higher giant-to-dwarf ratios in the inner regions found by previous authors (Driver et al. 1998; Boué et al. 2008).

However, our analysis shows that the gap statistics for different $R_{\text {max }}$ can be reproduced by a single, global CLF, allowing for two centrals in some groups. Therefore, although variations of the LFs with $R / r_{\text {vir }}$ seem to exist, the variations of gap distributions with $R_{\max }$ are, in fact, a consequence of number statistics.

### 5.4 Statistics of the magnitude gap and groups with two central galaxies

Our results show that the observed statistics of the magnitude gap can be only reproduced with two-component CLFs if we allow SGGs to have two central galaxies (Fig. 8). This scenario is expected when two groups merge, as appears to be the case in the nearby rich Coma cluster [Abell 1656] (Biviano et al. 1996), which has $\Delta \mathcal{M}_{12}=0.21$.

Although mergers of groups with similar halo masses (therefore with BGGs with similar luminosities) are less frequent than 'minor' group mergers (Fakhouri, Ma \& Boylan-Kolchin 2010), the latter are less likely to fill the magnitude gap of the more massive group, i.e. the BGG of the smaller group is more likely to become less than 1 st or 2nd-ranked galaxy.

These mergers may also lead to groups where the brightest galaxy is not the central galaxy of the system. In a recent study, Lange et al. (2017) found that the CLF model also underpredicts the number of groups where the brightest galaxy is not the central (BNC), suggesting that the systems with two centrals might be closely related to BNC groups. An analysis of the SGG phase space would be desirable to address the relation between these two classes of objects, but this is beyond the scope of the present article.

Finally, we tested if the gap statistics can be reproduced if both the BGG and SBGG luminosities are sampled from the satellite CLFs, i.e. we allow the existence of groups with no centrals. We find that, with the same equation (11) with $\mu=0.5$ and $\sigma=0.1$, we are able to reproduce quite well the fractions of LGGs, SGGs and $T_{2}$ values. However, the distribution of $\mathcal{M}_{1}$ is wider than the observations, leading to $T_{1}$ values that are much higher than observed.

## 6 SUMMARY AND CONCLUSIONS

In this paper, we use samples of SDSS groups, extracted on the web site of X. Yang (following the algorithm of Yang et al. 2007),
to study the bright end of the galaxy CLFs. Our SDSS samples are cleaned of selection effects, and we make use of more accurate SDSS-DR12 photometry.

We found that the CLFs provided by Y08 fit well the distribution of luminosities if we allow some changes to their parameters: our best Bayesian evidence suggests that, in low-mass groups, the shape of the bright end of the satellite component is steeper, while the ratio of characteristic satellite to central luminosities is higher compared to high-mass groups.
The statistics of magnitude gaps provide a fine test of the accuracy of the bright end of the CLFs. We first notice that these statistics depend on the maximum radius where we select the second brightest group galaxies. We also find that sampling the CLF produces too few SGGs, regardless of our different analytical fits. This suggests that some groups have more than one central galaxy, as expected in groups that are ongoing mergers. Indeed, we find that preferentially allowing SGGs to have two centrals recovers very well all the statistics of the magnitude gap.
Finally, we test the hypothesis of Hearin et al. (2013) that the richness-mass relation is a function of magnitude gap. We conclude that this relation disappears when we fit mass or velocity dispersion directly to richness and when we limit our choice of the second brightest galaxies to maximum projected radii of 1 or $0.5 r_{\text {vir }}$.

## ACKNOWLEDGEMENTS

We thank the referee for comments that led to an improved version of the manuscript. MT acknowledges the financial support from CNPq (process 204870/2014-3) and thanks the hospitality from the Institut d'Astrophysique de Paris. This research has been supported in part by the Balzan foundation via IAP. MT thanks Johannes U. Lange for a very useful discussion. We acknowledge the use of SDSS data (http://www.sdss.org/collaboration/credits.html) and topcat Table/VOTable Processing Software (Taylor 2005, http://www.star.bris.ac.uk/mbt/topcat/).

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## APPENDIX A: CORRECT IDENTIFICATION OF THE SBGG

To select the group sample for our analysis of magnitude-gap statistics, we ensure that the SBGGs are correctly identified, as described in Section 2.1. To investigate whether these selection criteria bias our sample and, therefore, the determination of the CLFs, we fitted the CLFs using the sample without requiring the correct identification of the SBGGs. By applying criteria (i) to (iv) listed in Section 2 (assuming $R_{\max }=2.0 r_{\text {vir }}$ ), and after cleaning the sample for the spectroscopic incompleteness of the BGGs and for groups near bright stars or the edges of the survey (see Sections 2.1 and 2.2), we obtain a sample of 2350 groups with a total of 13997 satellites within $2.0 r_{\text {vir }}$.

In Fig. A1, we show the LF of central and satellites galaxies, computed as the best-fitting CLFs averaged over the halo masses of the groups in our sample. As in Section 3.2, we compared satellite CLF models with different degrees of freedom. The best-fitting linear relations with $\log M_{\text {vir }}$ of $\log \left(L_{\mathrm{c}} / L_{\mathrm{s}}\right), \alpha$ and $\beta$ are shown in


Figure A1. Same as Fig. 1 (with the same notation), but for the sample for which we do not require the correct identification of SBGGs.


Figure A2. Same as Fig. 4 (with the same notation), but for the CLFs shown in Fig. A1.

Table A1. Comparison between models. (1) Free parameters of the model; (2) $p$-values of a KS and (3) AD tests; (4) $\Delta \mathrm{BIC}$ and (5) $\Delta \mathrm{AIC}$ relative to the model with the lowest BIC and AIC (model 6) values. Since model 1 has no free parameters, the values of BIC and AIC correspond to $-2 \ln \mathcal{L}_{\mathrm{s}}$.

| Model | $p$-values |  |  | $\Delta$ BIC |
| :--- | :---: | :---: | :---: | :---: |
| KS-test | AD-test | $\Delta$ AIC |  |  |
| $(1)$ | $(2)$ | $(3)$ | $(4)$ | $(5)$ |
| 1. Yang et al. (2008) | 0.06 | 0.05 | 11 | 41 |
| 2. $\alpha$ | 0.23 | 0.18 | 20.9 | 36.0 |
| 3. $L_{\mathrm{S}}$ | 0.04 | 0.02 | 26.6 | 41.7 |
| 4. $\beta$ | 0.09 | 0.06 | 29.9 | 45.0 |
| 5. $L_{\mathrm{s}}, \alpha$ | 0.58 | 0.47 | 15.8 | 15.8 |
| 6. $L_{\mathrm{s}}, \beta$ | 0.21 | 0.12 | $\mathbf{0 . 0}$ | $\mathbf{0 . 0}$ |
| 7. $\alpha, \beta$ | 0.39 | 0.29 | 9.6 | 9.6 |
| 8. $L_{\mathrm{s}}, \alpha, \beta$ | 0.35 | 0.29 | 17.9 | 2.8 |

Fig. A2. Comparing these results with Fig. 4, we see that we obtain shallower relations of these parameters with halo mass.

In Table A1, we show $p$-values of KS and AD tests, as well as $\Delta \mathrm{BIC}$ and $\triangle \mathrm{AIC}$ relative to the model with the lowest BIC and AIC values. Differently from the results obtained with the clean sample
(Table 2), no model can be strongly rejected based on the KS and AD tests, with the model 1 (Y08) and 3 (free $\log L_{\mathrm{s}}$ ) being only marginally inconsistent with the data. On the other hand, both BIC and AIC indicate that model 6 (free $\log L_{\mathrm{s}}$ and $\beta$, as strongly preferred by BIC evidence for the cleaned sample) is strongly preferred over the other models.

## APPENDIX B: GAP STATISTICS FROM

## DIFFERENT CLF MODELS

In Fig. B1, we show the gap statistics from different CLF models described in Table 3. Each CLF model was sampled assuming that


Figure B1. Statistics of the magnitude gap for as a function of group richness and halo mass for $R_{\max } / r_{\mathrm{vir}}=0.5$ (left) and 1.0 (right). Predictions from sampling different CLF models are shown: the best-fitting models with two free parameters, i.e. when we fit the linear relation with log $M_{\mathrm{vir}}$ of only one of the parameters: $\alpha$ (red dashed), $\log L_{\mathrm{S}}$ (green dotted), or $\beta$ (blue dot-dashed lines); the models with 4 degrees of freedom (orange long-dashed, cyan solid and purple short-long dashed lines); and the model with six free parameters (black solid lines). The Y08 model is indicated as the black dashed lines.

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[^1]:    ${ }^{1}$ We used the catalogue petroB, which is available at http://gax. shao.ac.cn/data/Group.html.
    ${ }^{2}$ See appendix A in Paper I for the conversion from quantities relative to the mean density to those relative to the critical density.

[^2]:    ${ }^{3}$ http://hmf.icrar.org/
    ${ }^{4}$ The evolution correction to $z=0.1$ determined by Blanton et al. (2003b) $(E[z]=-1.62[z-0.1]$ in the $r$ band) leads to corrections ranging from $E(z=0.015) \sim-0.14$ to $E(z=0.07) \sim-0.05 \mathrm{mag}$ for our sample.
    ${ }^{5}$ Although the median redshift of our sample is 0.05 , we kept the median redshift of the SDSS MGS as the reference for the $k$-corrections. The difference between ${ }^{0.1} \mathcal{M}_{r}^{\text {Petro }}$ and ${ }^{0.05} \mathcal{M}_{r}^{\text {Petro }}$ is $\sim 0.1 \mathrm{mag}$, and has no effect on the results and conclusions of our study.

[^3]:    ${ }^{6}$ We used the file sdss_dr72safe0_res6d.pol, which can be downloaded from http://space.mit.edu/molly/mangle/download/data.html.

[^4]:    ${ }^{7}$ http://gavo.mpa-garching.mpg.de/portal/

