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SOAR/Goodman Spectroscopic Assessment of Candidate Counterparts of the LIGO/Virgo Event GW190814

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








































































































































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SOAR/Goodman Spectroscopic Assessment of Candidate Counterparts of the LIGO/Virgo Event GW190814*

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* Based on observations obtained at the Southern Astrophysical Research (SOAR) telescope, which is a joint project of the Ministério da Ciência, Tecnologia, Inovações e Comunicações (MCTIC) do Brasil, the US National Science Foundation's National Optical-Infrared Astronomy Research Laboratory (NOIRLab), the University of North Carolina at Chapel Hill (UNC), and Michigan State University (MSU).

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Abstract

On 2019 August 14 at 21:10:39 UTC, the LIGO/Virgo Collaboration (LVC) detected a possible neutron star–black hole merger (NSBH), the first ever identified. An extensive search for an optical counterpart of this event, designated GW190814, was undertaken using the Dark Energy Camera on the 4 m Victor M. Blanco Telescope at the Cerro Tololo Inter-American Observatory. Target of Opportunity interrupts were issued on eight separate nights to observe 11 candidates using the 4.1 m Southern Astrophysical Research (SOAR) telescope’s Goodman High Throughput Spectrograph in order to assess whether any of these transients was likely to be an optical counterpart of the possible NSBH merger. Here, we describe the process of observing with SOAR, the analysis of our spectra, our spectroscopic typing methodology, and our resultant conclusion that none of the candidates corresponded to the gravitational wave merger event but were all instead other transients. Finally, we describe the lessons learned from this effort. Application of these lessons will be critical for a successful community spectroscopic follow-up program for LVC observing run 4 (O4) and beyond.

Unified Astronomy Thesaurus concepts: [Gravitational waves \(678\)](#); [Spectroscopy \(1558\)](#); [Neutron stars \(1108\)](#); [Black holes \(162\)](#)

1. Introduction

The 2017 discovery of the optical counterpart of a binary neutron star (BNS) merger—a kilonova (KN)—was one of the highlights of observational astrophysics of the early 21st century. This discovery, following on the 2015 discovery of the first ever detected gravitational wave (GW) event, GW150914 (Abbott 2016), was a significant leap forward for astrophysics. The detection of GW170817 in coincidence with a short gamma-ray burst by Fermi–GBM during the second observing run (O2) of the Advanced LIGO (The LIGO Scientific Collaboration et al. 2015) and Virgo (Acernese et al. 2015) network inaugurated the era of multi-messenger astronomy with GWs (Abbott et al. 2017a, 2017c). The optical counterpart was discovered 12 hr after the merger by several independent teams, including our own team, the Dark Energy Survey Gravitational Wave Search and Discovery Team (DESGW). DESGW utilizes the Dark Energy Camera (DECam) (Flaugher et al. 2015) on the Victor M. Blanco Telescope at the Cerro Tololo Inter-American Observatory (CTIO) in Chile (Soares-Santos et al. 2017). This discovery enabled panchromatic imaging and spectroscopy, which galvanized the astronomical community.

While this single event captured the focus of the entire astronomical community, the breadth and number of scientific analyses stemming from it are perhaps more astounding. Standard siren techniques enabled a direct measurement of the expansion rate of the universe today (Abbott et al. 2017b; Soares-Santos et al. 2019; Palmese et al. 2020) and in the future they will also be a useful probe of the growth of structure (Palmese & Kim 2021). Measuring elemental abundances in the merger ejecta using spectroscopic instruments led to an understanding of the origin of heavy elements synthesized during the merger (Chornock et al. 2017; Drout et al. 2017; Tanaka et al. 2018), and we note the unique wavelength coverage of the Very Large Telescope (VLT) X-Shooter in this task in particular (Pian et al. 2017; Smartt et al. 2017; Watson et al. 2019). X-ray and radio observations characterized the geometry of the explosion to be best described by a jet plus cocoon structure (Alexander et al. 2017; Hallinan et al. 2017; Margutti et al. 2017; Troja et al. 2017; Mooley et al. 2018; Ghirlanda et al. 2019). The gravitational waveforms tested and further bolstered the validity of the theory of general relativity, as verified by numerical relativity simulations (Shibata et al. 2017; Abbott et al. 2019), and several other studies explored the connection between BNS mergers and short gamma-ray bursts (e.g., Fermi-LAT Collaboration 2017; Fong et al. 2017; Savchenko et al. 2017; Xiao et al. 2017; Lyman et al. 2018; Ascenzi et al. 2020). These analyses, and many not listed, were enabled by the association of

the GW signal with its electromagnetic (EM) signal. Given that these events are such a rich source of astrophysical knowledge, finding counterparts to GW events related to compact object mergers remains a primary goal of the multimessenger-focused astronomical community.

On 2019 August 14 at 21:10:39 UTC, during its observing run 3 (O3), the LIGO/Virgo Collaboration (LVC) detected a binary merger initially designated as S190814bv and later given a final designation of GW190814. This was one of 56 event alerts from the LVC during O3 and was particularly interesting: GW190814 was at the time classified as a neutron star–black hole (NSBH) merger, the first high significance event of this kind ever observed (LVC 2019a, 2019b; Abbott et al. 2020). The LIGO/Virgo analysis found that this merger event occurred at a distance of 267 ± 52 Mpc. It had a 90% localization region of 23 deg^2 and a probability of being a NSBH merger of greater than 99%. Further, taking as an assumption that the GW170817 BNS KN (at a distance of 43 Mpc) had a typical luminosity for such an event and scaling by the inverse-square law, one could estimate that the optical counterpart to GW190814 could conceivably peak at a brightness of $i \sim 21$ (≈ 4 mag fainter than that of GW170817)—well within the range of DECam, as well as still within the range of medium-resolution spectrographs on 4 m class optical telescopes—simplifying the effort of following up any likely optical counterpart candidates. Thus, the DESGW team undertook an extensive search for a KN event that would form the optical counterpart to this potential NSBH merger event, making use of DECam observations within the high-probability region of the GW event. This search is described in detail in Morgan et al. (2020).

A number of other groups also searched for an EM counterpart to GW190814. Kilpatrick et al. (2021, many of whom are also members of the DESGW Collaboration) discuss searches for KN candidates using several 0.7–1 m class telescopes as well as Keck/MOSFIRE and also present spectroscopy of a number of candidates (including in their Figure 4 a copy of many of the spectra described here in the current paper). They also present limits on EM counterparts to GW190814 and consider scenarios in which an EM counterpart of an NSBH would be detected. The Australian Square Kilometre Array Pathfinder imaged 30 deg^2 at 2, 9, and 33 days after the event at a frequency of 944 MHz (Dobie et al. 2019). The Magellan Baade 6.5 m telescope was used to search on a selection of galaxies within the localization area out to limiting magnitude of $i = 22.2$ and found no counterparts (Gomez 2019). The MegaCam instrument on the Canada–

France–Hawaii Telescope (CFHT) was used to search much of the localization region. Although the CFHT team reached a depth of $i > 23.9$ at 8.7 days post-merger, no KN was found (Vieira et al. 2020). The GROWTH Collaboration used imaging from DECam along with other facilities for imaging and spectroscopy of possible KN candidates. Using simulations, they constrained possible ejecta mass from the merger to be $M_{\text{ejecta}} < 0.04 M_{\odot}$ at polar viewing angles (Andreoni et al. 2020). Watson et al. (2020) described limits on an EM counterpart to GW190814 using observations with optical imager DDOTI (at the Observatorio Astronómico Nacional in Mexico) and Swift/BAT observations. They showed that Swift/BAT should have detected an associated gamma-ray burst at the 98% level. Ackley (2020) described the ENGRAVE team search using the VLT as well as involvement with the ATLAS, GOTO, GRAWITA-VST, Pan-STARRS, and VINROUGE projects. Their observations covered the localization region to depths as faint as $r \approx 22$. Their limits suggest that it is likely the neutron star was not disrupted during the merger. DDOTI wide-field observations were also used along with the Lowell Discovery Telescope, the Reionization and Transients InfraRed, and spectroscopy from the Gran Telescopio Canarias to locate EM counterparts (Thakur et al. 2020). Their data suggest that there was no gamma-ray burst along the jet’s axis.

While searching for an optical counterpart to GW190814, the DESGW pipeline began with 33,596 events in the likelihood regions. Using the analysis pipeline we produced a final list of 11 candidates that passed our cuts and were bright enough for spectroscopy using a 4 m class telescope (Morgan et al. 2020; see also Section 4.2). For these candidates we proceeded to conduct spectroscopic typing at the Southern Astrophysical Research (SOAR) 4.1 m telescope⁹⁹ using the Goodman High Throughput Spectrograph (HTS; Clemens et al. 2004). (Spectroscopic typing is facilitated by the fact that, due to the fast ejecta velocities expected of KNe, 0.03–0.30c, their spectra are expected to be featureless or only have very broad, smooth spectral features, especially in the optical during the first few days after the merger event, which distinguishes their spectra from supernovae (SNe) and other optical transients; see, e.g., the KN models of Kasen et al. 2017.) The spectroscopic follow-up team submitted Target of Opportunity (ToO) observing requests to the SOAR telescope on eight separate nights in order to use the Goodman HTS on SOAR for spectroscopic typing of these 11 candidates.

After taking spectra for eight candidates (plus the host galaxies of three additional candidates which had faded beyond the straightforward capabilities of SOAR—i.e., $i \sim 21.5$), no optical counterpart was discovered for GW190814. Despite this null result, this paper serves several important functions. First, it serves as a companion to our other two papers (Morgan et al. 2020; Kilpatrick et al. 2021), providing a deep dive into the methodology and detailed results of a coordinated spectroscopic campaign of the first possible NSBH event ever detected, including the finding charts, light curves, and KN spectral fitting not covered in detail by the other two companion papers. Further, it describes and provides previously unpublished open-source tools that can be of use to similar future spectroscopic campaigns. Also, by comparing results from two separate SN spectrum fitters and a KN

spectrum fitter, this paper goes into some detail into the subtleties associated with spectroscopic classification of relatively faint SNe and KNe. Finally, although it does not change the conclusions of the companion papers, some of the final classifications of the candidate counterparts here are updates from what was seen in those papers.

In summary, we describe in this paper the DESGW collaboration’s spectroscopic follow-up campaign for the GW190814 gravitational merger event. We also describe our overall spectroscopic follow-up methods and strategy, how we employed them in this particular follow-up campaign, the lessons learned, and the prospects for the future. The paper is organized as follows: in Section 2 we describe the LIGO/Virgo observations of GW190814. In Section 3 we describe the DESGW search for candidate optical counterparts. In Section 4 we describe the selection and filtering of the candidates. In Section 5 we describe the SOAR observing strategy and the observations of counterpart candidates for GW190814. In Section 6 we discuss our results and address the population of objects we found. In Section 7 we summarize our conclusions. In addition, we provide in Section 8 a list of software packages used throughout our analysis.

In this paper we follow the cosmology given by Bennett et al. (2014), with flat Λ CDM cosmology with $\Omega_M = 0.286 \pm 0.008$ and $H_0 = 69.6 \pm 0.7 \text{ km s}^{-1} \text{ Mpc}^{-1}$.

2. LIGO/Virgo Observations

As noted above, on 2019 August 14 UTC, the LVC observed gravitational radiation at high statistical significance. The event, initially named S190814bv, occurred during a time that all three detectors (LIGO Hanford, LIGO Livingston, and Virgo) were operating normally, which enabled both a good angular localization of the source and more precise estimate of the source parameters. The false-alarm probability was calculated at $2.0 \times 10^{-33} \text{ Hz}$ —or once per 10^{15} Hubble times—suggesting a very high signal-to-noise ratio (S/N) event (LVC 2019b). Using the *bayestar* pipeline (Singer & Price 2016), the LVC team localized the source of the GW signal to a 38 (7) sq. deg. area at the 90% (50%) confidence level in the Southern Hemisphere on the night of the merger. The initial luminosity distance estimate was $276 \pm 56 \text{ Mpc}$ (LVC 2019a). Preliminary source classification via a machine-learning (ML)-based tool (Kapadia et al. 2020) identified the event as a “mass-gap” binary merger—i.e., a merger event in which at least one of the compact objects has a mass falling within the hypothetical mass gap between NSs and BHs (i.e., in the mass range 3–5 M_{\odot} ; Abbott et al. 2020; LVC 2020a). The small localization area and the potential of identifying an optical counterpart made this event interesting from the perspective of follow-up projects.

The following day, the LVC *LALInference* pipeline (LIGO Scientific Collaboration 2018) localized the source to 23(5) sq. deg. at the 90% (50%) confidence level, refined the classification to an NSBH merger, and estimated the luminosity distance of the event to be $267 \pm 52 \text{ Mpc}$ ($z = 0.059 \pm 0.011$ for a standard Λ CDM cosmology; Bennett et al. 2014; Wright 2006). S190814bv thus became the first possible NSBH system observed by a GW observatory and a prime target for follow-up by the EM astronomical community. However, the LVC parameter estimation indicated that the parameter *HasRemnant* was $< 1\%$. (*HasRemnant* is the probability that a nonzero mass was ejected during the collision and

⁹⁹ <https://noirlab.edu/science/programs/ctio/telescopes/soar-telescope>

remains outside the final remnant object; Foucart et al. 2018; LVC 2020b.) This suggested that there was a low probability that any ejecta was preserved outside the BH and thus that there was a small chance of there being an observable KN.

Well after searches for an EM counterpart were completed, the LVC published results from an updated offline analysis (Abbott et al. 2020), where the final luminosity distance was estimated to be 239_{-45}^{+41} Mpc (median and 90% credible interval), the 90% localization area was updated to 18.5 sq. deg., and the masses of the two merging objects were updated to $23.2 M_{\odot}$ (a BH) and $2.6 M_{\odot}$ (a mass-gap object—i.e., either an underweight BH or an excessively massive NS). It was also at this time that this GW event was renamed from its initial designation, S190814bv, to GW190814.

The nature of this GW190814 was recently debated and summarized by Abbott et al. (2020) and, since its discovery, only a couple more GW merger events with comparable properties have been identified (see The LIGO Scientific Collaboration et al. 2021 and the interactive plot at <https://ligo.northwestern.edu/media/mass-plot/index.html>). Particularly striking is the mass ratio of the GW190814 merger components—a value of 0.112—whereas the average mass ratio of more typical LIGO binary BH events is ~ 1 . As noted above, one of the components of the GS190814 merger was a $23.2 M_{\odot}$ BH, but the other was a $2.6 M_{\odot}$ “mass-gap” object. If this mass-gap object is an NS, this has ramifications for the NS equation of state, which is a determining factor in the maximum allowable mass of NSs (currently estimated to be $\lesssim 2.6 M_{\odot}$). Independent of whether the mass-gap object is an NS or a BH, if these types of mergers are more common than expected, there may be consequences for stellar population synthesis models, since these models tend to favor the merger of systems with components that are less asymmetric in mass, although stellar environment may also play a role: merger rates between NSs and BHs are low in globular clusters ($\sim 10^{-2}$ – 10^{-1} Gpc $^{-3}$ yr $^{-1}$; e.g., Ye et al. 2020), but likely higher in young stellar clusters ($< 10^{-1}$ Gpc $^{-3}$ yr $^{-1}$; Ziosi et al. 2014); thus, star clusters with young stellar populations might be the preferred location for mergers similar to GW190814. For the purposes of this paper, we will assume that GW190814 is a possible NSBH merger, as it was classified during the SOAR follow-up observing runs.

In the next section we describe the efforts of the DESGW Collaboration to identify transients that were possible KN candidates.

3. DECam Search Campaign

In searching for an optical counterpart to GW190814, the DESGW collaboration triggered ToO observations with the 570 megapixel DECam optical imager on the CTIO Blanco 4 m telescope. Together, the Blanco and DECam reach a 5σ limiting r -band magnitude of ~ 23.5 in a 90 s exposure in a 3 sq. deg. field of view (FoV) (Nielsen et al. 2019). The combination of deep imaging and a wide FoV make Blanco/DECam the ideal instrument for efficiently detecting optical transients localized to tens of square degrees.

Our follow-up efforts for GW190814 utilized the resources of the Dark Energy Survey (DES), which is a wide-field optical survey that covered a 5000 sq. deg. region of the southern sky from 2013 to 2019 using Blanco/DECam (Diehl et al. 2019). DES imaging of the DES footprint reaches a 10σ depth for point sources of $grizY = 25.2, 24.8, 24.0, 23.4, 21.7$ mag

(Mohr et al. 2012). The LVC 90% containment region for GW190814 is entirely within the DES footprint, enabling the use of high-quality DES images during difference imaging.

We performed DECam ToO follow-up observations of GW190814 for six nights following the LVC alert, namely nights 0, 1, 2, 3, 6, and 16. The early nights were chosen to look for rapidly evolving transients immediately following the merger. KNe from either BNS (Arcavi et al. 2017) or NSBH (Kawaguchi et al. 2016) events are expected to vary by about a magnitude over the course of a single night in the first days after the event. Observations 16 nights after the merger were used to exclude persisting SNe. Due to moon brightness, especially during the first nights of DECam follow-up, we opted to use the redder i and z bands to minimize the effect of sky brightness on our imaging depth.

The DECam images were processed by the DES Difference Imaging Pipeline (Herner et al. 2020), an updated version of the DES SN Program’s Pipeline described in Kessler et al. (2015), using coadded DES wide-field survey images (Abbott et al. 2018) as templates.

After image processing, candidate KNe were identified and then selected for spectroscopic follow-up. The selection process included eliminating moving objects (e.g., asteroids), known transients (e.g., variable stars and active galactic nuclei (AGNs)), and transients with colors and/or light curves characteristic of SNe. Visual inspection of the images was also important, especially in the first nights of DECam follow-up, when light curves for the candidates consisted of only one or two epochs. For GW190814 in particular, there were 33,596 candidates immediately after the image processing. KN candidates were found in DECam images after running them through the reduction pipeline. Objects were found by SExtractor (Bertin & Arnouts 1996). Objects that had good detections in SExtractor showed evidence of being transients by comparison to known object templates and passed visual inspection checks were considered. Other candidates were identified in alert notifications from the Gamma-ray Coordinates Network (GCN)¹⁰⁰ put out by other groups searching for kilonova KN candidates. A more rigorous process of object assessment was done later, described in more detail in Morgan et al. (2020) and summarized in Section 4.2. In the end, spectroscopic follow-up was performed using the SOAR Goodman HTS for 11 candidates (or their host galaxies).

In Table 1 we present candidates found and spectroscopically targeted by the DESGW team during DECam follow-up of GW190814. In this table we provide both the DESGW ID and the Transient Name Server (TNS) name, which we continue to use in this work. In the final two columns, we present the localization probability enclosed within the GW sky-map including each object location. For further details of the processing of the DECam data and the subsequent identification of possible candidates, refer to our companion paper (Morgan et al. 2020).

In Figure 1 we show both the initial and the final sky-localization maps issued by the LVC along with the locations of each of the 11 objects we observed. Note that in the smaller final probability regions, some of the objects we observed are outside the 90% probability area, but all are included within this area in the initial map.

¹⁰⁰ <https://gcn.gsfc.nasa.gov/>

Table 1
Candidates Found by the DESGW Team during the DECam Follow-up of GW190814 That Were Then Followed Up with SOAR ToO Observations

DESGW ID	TNS Name	R.A.(2000) (deg)	Decl.(2000) (deg)	GCN/ID	Mag at Discovery	Band	Prob Reg Initial	Prob Reg Final
624921	2019nqq	20.95506	−33.034762	25373/c	20.76	<i>i</i>	90%	o
624609	2019nqr	23.573539	−32.741781	25373/d	18.34	<i>i</i>	80%	90%
624690	2019noq	12.199493	−25.30652	25356 (Pan-STARRS)	19.93	<i>i</i>	30%	30%
624157	2019ntn ^a	23.722184	−31.380451	25393 (GROWTH)	20.8	<i>i</i>	90%	o
626761	2019npw	13.968327	−25.783283	25362/e	20.5	<i>i</i>	40%	60%
631360	2019num	13.881714	−22.968887	25393 (GROWTH)	21.3	<i>i</i>	90%	o
661833	2019ntr	15.007796	−26.714266	25393 (GROWTH)	21.2	<i>z</i>	80%	o
625839	2019omx	24.18436	−33.302678	25486/z	22.1	<i>z</i>	90%	o
626956	2019ntp	12.550247	−26.197878	25393 (GROWTH)	21.0	<i>i</i>	50%	60%
631484	2019nte	23.557358	−31.721700	25398/f	20.95	<i>i</i>	80%	o
635566	2019omw	12.234396	−23.170137	25486/y	22.8	<i>i</i>	50%	80%

Note. The DESGW ID is the internal identification number while the TNS name comes from the Transient Name Server (<https://wis-tns.weizmann.ac.il>). The coordinates are given here in degrees, along with the GCN announcing discovery of the transient. Magnitude at discovery is given in the band listed. The confidence probability enclosed within the GW sky-map including the object position is given both for the initial map issued by the LVC used during observing and for the final, smaller map. (The “o” means outside the 90% sky-localization probability region.)

^a AT2019mbq was accidentally targeted for SOAR spectroscopy instead of the intended target AT2019ntn, and this accident was not discovered until much later. This mistake has been traced to a copying error during the handoff of this target from the DECam processing and analysis team to the SOAR observing team. Candidate AT2019mbq is at R.A. = 10°835384, decl. = −25°883880, with a magnitude at discovery of $i = 18.75$. We note that AT2019mbq was not originally considered for spectroscopic follow-up since its host galaxy had a too high estimated photoz ($z_{\text{photo}} = 0.17 \pm 0.05$) and since there was evidence of a pre-merger detection for this candidate. As for AT2019ntn, although no spectrum was taken of it, the fact that it brightened in *z*-band about 4 days after the merger and the fact that it lay outside the 90% confidence contour of the LVC final map (Figure 1) make it unlikely that AT2019ntn was the optical counterpart.

4. SOAR Spectroscopic Candidate Selection

To achieve the maximum science, rapid spectroscopic follow-up of candidate KNe is a necessity: first, to discover the optical counterpart from among the list of potential candidates, and then, if discovered, to permit the longest possible timeline for optical monitoring of the evolution of the potential KN’s light curve and spectral energy distribution before it fades to obscurity. The constraints for our SOAR spectroscopic program, however, were two-fold: (1) to preserve each night’s main program as much as possible, as SOAR ToO interrupts are limited to 2.5 hr per night (including overheads); and (2) to achieve reasonable S/N ($\gtrsim 5$ –10) of a medium-resolution spectrum on SOAR within a reasonable amount of time. Due to these constraints, each observation is limited to objects with brightnesses of $i < 21$. (We pushed the limits for GW190814, relaxing this constraint to $i \lesssim 21.5$.) In Section 4.1 we present our baseline strategy for SOAR/Goodman spectroscopy in LVC O3. Then in Section 4.2 we describe our strategy for filtering transients found with DECam observing to find the candidates that should be followed up with spectroscopy.

4.1. SOAR Program Baseline Strategy for LVC O3

We designed our SOAR ToO program for rapid and robust identification and subsequent nightly follow-up of KN candidates to be coupled with the DECam wide-field search and discovery program (Soares-Santos et al. 2017; DES Collaboration et al. 2020; Herner et al. 2020; Morgan et al. 2020), which would be providing a selection of candidates for spectroscopy. This project was awarded time at the SOAR/Goodman HTS to observe GW optical candidates discovered during the entire year-long O3 run of the LIGO/Virgo campaign. Due to the transient nature of GW optical counterparts (KNe), SOAR spectroscopy must be carried out in ToO mode. We requested SOAR/Goodman HTS ToO time in

instant activation mode for a total of 10 hr or at least four ToO activations per semester. This way we took advantage of the fast survey confirmations from the DECam search and discovery program, which could be available within 1 hr, if the merger happened during the Chilean night. The LVC predicted that there would likely be roughly eight BNS mergers and one NSBH merger—the events most likely to yield an optical counterpart—over the course of the LVC O3 run (Abbott et al. 2017a; Chen et al. 2021). Thus we planned to use SOAR to follow up the two to three of these events likely visible from the Southern Hemisphere each observing semester.

The KN for the GW170817 BNS merger was exceptionally bright and easy to identify. It was expected that future events would on average be much farther away and thus likely to be much fainter and harder to distinguish from other transients (e.g., SNe Ia) in the larger volume encompassed by LVC O3 detection thresholds. We planned to use the SOAR Goodman HTS (1) to spectroscopically identify the optical counterpart to the GW event from among a small list of candidates provided by an initial DECam search and discovery program; (2) to obtain, once identified, a higher-S/N optical spectrum of the counterpart, suitable for detailed modeling; and (3) to obtain additional high-S/N spectra of the potential KN on successive nights until it was effectively too faint for useful follow-up on SOAR. We would employ an instrument setup almost identical to that of Nicholl et al. (2017), who were able to follow the GW170817 KN event at reasonable S/N using the Goodman HTS from day 1.5 to day 7.5 after the GW trigger. In that case the KN faded from magnitude $i \approx 18$ to 21 over 6 days; they used an integration time (IT) of 3×900 s with the 400 l mm^{-1} grating. Based on their Goodman spectra, we anticipated that we could achieve the S/N necessary to classify whether a given candidate was a true KN or just another transient using a single 900 s exposure for $i \leq 19$ candidates, a single 1200 s exposure for $i \approx 20$ candidates, and a single 1800 s exposure for $i \approx 21$

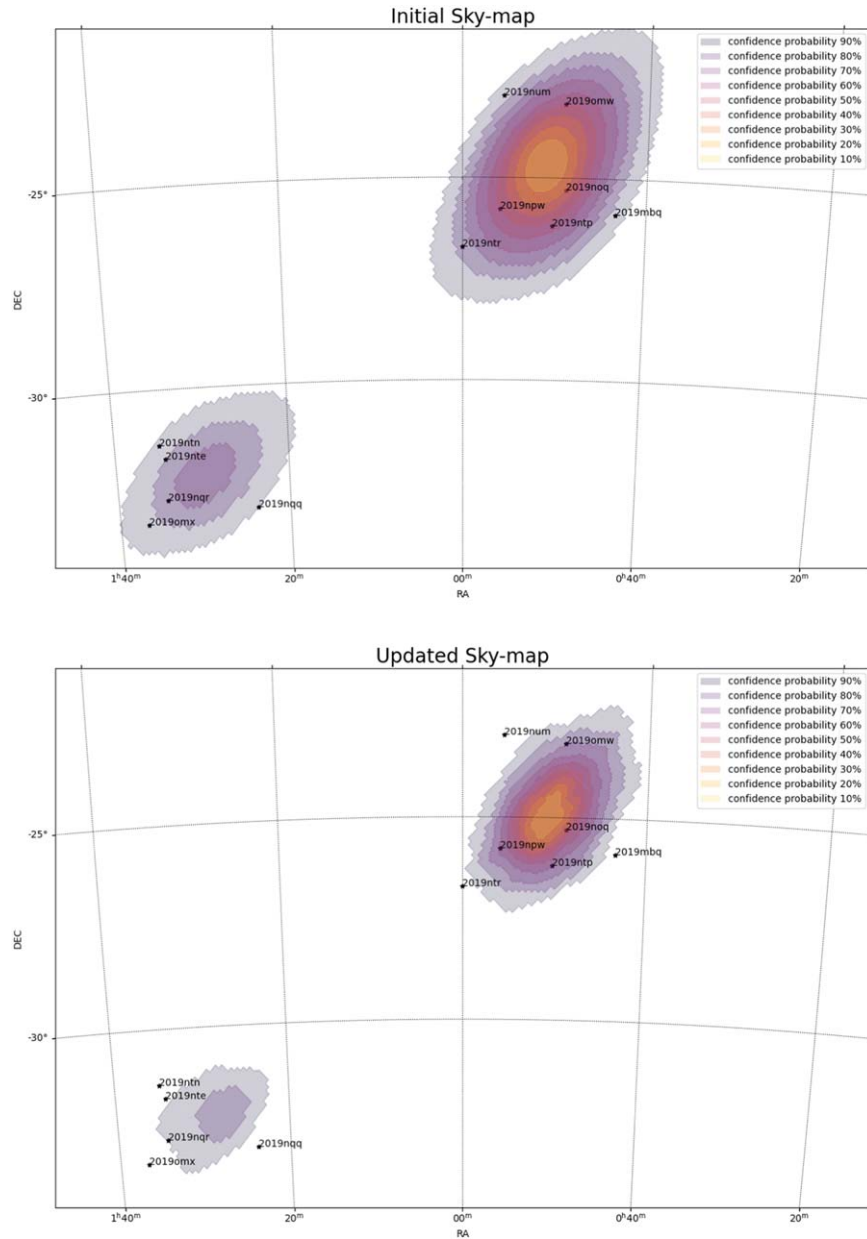


Figure 1. LIGO/Virgo Collaboration (LVC) sky-localization maps for GW190814; colors indicate confidence probability contours. The top figure is the initial sky-map, released shortly after event discovery on 2019 August 14. The bottom figure is the final sky-map, released after further analysis by the LVC collaboration. The locations of each of the 11 objects we describe in this paper are also given.

candidates. We would leave fainter candidates to programs on larger telescopes, like programs on VLT and Gemini-South.

We planned to follow up the list of candidates until we either finished the list (finding no KN) or identified the optical counterpart. For an identified KN, two additional exposures of the same integration time would allow us to build S/N suitable for model fitting. We planned to continue SOAR follow-up if a confirmed KN was brighter than $i = 20$ mag, requesting interrupts on all successive nights until it faded below that value. We ran 100,000 simulations of the SOAR search program. An average of 8.79 DECam candidates per LIGO event in the magnitude range $i = 16$ – 24 was assumed, where magnitudes were drawn randomly from the expected candidate distribution (see the `LC_SHAPE` row of Figure 2, where the numbers add up to 8.79). To estimate the time needed, we

included not only the expected exposure times, but also all relevant overheads (e.g., slewing, target acquisition, readout, standard star observations, etc.). To compensate for possibly worse sky transparencies (Nicholl et al. 2017 found clear skies), the science integration times were multiplied by a factor of 1.25. The simulations showed that, for a single GW event, 50% of the time a SOAR follow-up would be completed in 4.3 hr (two ToO interrupts), 95% of the time in 6.7 hr (three interrupts), and 100% of the time in 9.5 hr (four interrupts). Note that follow-up completion does not necessarily mean a guaranteed identification of the optical counterpart: it may just mean that the list of candidates bright enough to be observed by SOAR was exhausted without identifying the optical counterpart or even that the optical counterpart (if any) was too faint to be detected by the DECam imaging. Nonetheless, in our time

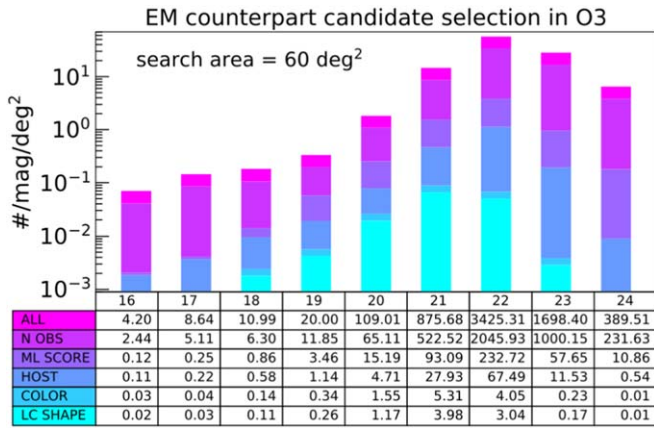


Figure 2. Baseline DECam search and discovery candidate selection for spectroscopic follow-up for LVC O3. The need for a robust classification pipeline to find kilonovae (KNe) in O3—as was uniquely done for GW170817 in Soares-Santos et al. (2017)—is shown here in the (i -band) magnitude distribution of all transient candidates expected to be found by a DECam search and discovery imaging sequence for a typical binary neutron star (BNS) gravitational wave (GW) trigger in LVC O3, assuming a typical search area of 60 sq. deg. (e.g., see Scolnic et al. 2018). The first row (“ALL”), which corresponds to the magenta histogram, is the distribution of candidates expected to be output from the DECam Difference Imaging Pipeline. In these simulations, we rejected moving objects and artifacts by requiring more than two observations (“N_OBS”) and machine-learning classification score >0.7 (“ML_SCORE”), rejected candidates with host galaxies at $z > 0.2$ (“HOST”), and performed a color cut using the fact that, unlike supernovae, the early evolution of a KN is blackbody-like (“COLOR”); as detection of a rising light curve would immediately pin-point the target, we applied a reduction of 25% assuming that, given DECam scheduling constraints, we would be able to get two epochs at <24 hr from merger for one in four events (“LC_SHAPE”). Thus, this last row (“LC_SHAPE”), which corresponds to the cyan histogram, is the expected distribution of candidates remaining after all the image-level culling procedures have been run. (Note: the numbers listed below the plot are the total per magnitude bin for the full 60 sq. deg. search area; the y -axis of the plot, however, is the number per magnitude bin *per square degree*. Also note: the results shown in the above plot and histogram are based on multiple simulations covering areas larger than 60 sq. deg.; scaling to a 60 sq. deg. localization area and averaging over the multiple simulations means that the numbers in these bins are not integers (e.g., why the number of candidates in the $i = 21$ bin in the “ALL” row is 875.68 and not, say, exactly 875).)

requests, we estimated approximately 10 hr per GW event to optimize our chances of spectroscopically identifying and monitoring a KN with SOAR during the LVC O3 run.

For spectroscopic classification, it was anticipated SOAR could go as faint as $i = 21$. In Figure 2 we visually represent the process for DECam search and candidate selection for spectroscopic follow-up. This figure shows the expected number of DECam candidates per magnitude per square degree in LVC O3, for a typical localization area of 60 sq. deg. The columns are arranged in order of magnitude, with magnitude getting dimmer to the right.

For continued monitoring of the evolution for the optical spectrum of an identified KN, it was thought that a higher S/N would be required, so additional monitoring was planned to be constrained to KNe brighter than $i = 20$. Candidates fainter than $i = 21$ and confirmed KNe fainter than $i = 20$ would be handed over for larger telescopes for spectroscopic follow-up. Via simple timing simulations, we estimated the amount of time to obtain SOAR spectra for typical KN candidates from a given LVC O3 event to take no more than ≈ 10 hr over the course of $\lesssim 5$ nights (recalling the maximum ToO “interrupt” time per night is 2.5 hr) The SOAR team

would meet with the DECam team once the DECam team had a set of candidates.

To elaborate, in Figure 3, panel A, we present a simplified flow chart for a simulated SOAR follow-up for the optical counterpart of a single LVC O3 event. N_{cand} is the total number of candidates from an imaging search and discovery program—i.e., the expected number of objects for which we would need to take spectroscopy from SOAR or, for fainter candidates, from other telescopes. If we run this flowchart over 100,000 realizations and compile the results, we get the histograms in panels B and C of Figure 3. Panel B shows the distribution—over 100,000 simulated realizations—of the total duration (in hours) of SOAR ToO interrupt time expected for a single LVC O3 event. Likewise, panel C shows the distribution over 100,000 simulated realizations of the total number of SOAR interrupts expected for a single LVC O3 event.

4.2. Candidate Filtering for GW190814

For GW190814, we selected targets for SOAR spectroscopy by reducing the DECam images in real time and monitoring the GCN for objects of interest detected by other follow-up teams. In both approaches, one important constraint is the brightness of the candidates. For accurate spectroscopic classification, we wanted a minimum S/N of 5–10 in the collected spectra. Therefore in typical observing conditions, with 45 min to 1 hr exposure times, objects fainter than 21.5 i -band mag are excluded. However, if the candidate’s host galaxy was brighter than the magnitude threshold, we targeted the host to obtain a precise redshift of the candidate.¹⁰¹

The candidate selection performed in real-time for the SOAR targets differs from the offline candidate selection presented in Morgan et al. (2020). One important difference is that all potential SOAR targets were selected before we began co-adding the DECam images within the same night and filter. The cuts applied to select spectroscopic targets were as follows.

1. *ALL*. Detected in DECam images by the DESGW Search and Discovery Pipeline.
2. *DETECTED 2x*. At least two detections by SExtractor with no errors and with an autoscan score of at least 0.7 separated by at least one hour (autoscan is an ML-based tool for differentiating between image artifacts and real objects; Goldstein & D’Andrea 2015).
3. *PHOTO z*. If a host-galaxy exists in the DES Catalog, the estimated photometric redshift and its error must be consistent with the LVC distance mean within three standard deviations.
4. *INSPECTION*. Pass visual inspection by the DESGW team.

Whether an object was first reported to the GCN by the DESGW team or by another follow-up team, it was still required to pass the same set of selection criteria prior to being targeted with SOAR. Technical details and motivations for these criteria are presented in Morgan et al. (2020). Remaining objects after the above selection criteria were sorted by their single-band average rate of change in flux to look for rapidly evolving transients. Finally, we triggered SOAR on objects passing the criteria and that had not already been ruled out by

¹⁰¹ We note that the host galaxy for each candidate was identified by matching the candidate’s coordinates with the DES Y3 galaxy catalog using both angular and galaxy photoz information. Details can be found in Section 3.3 of Morgan et al. (2020).

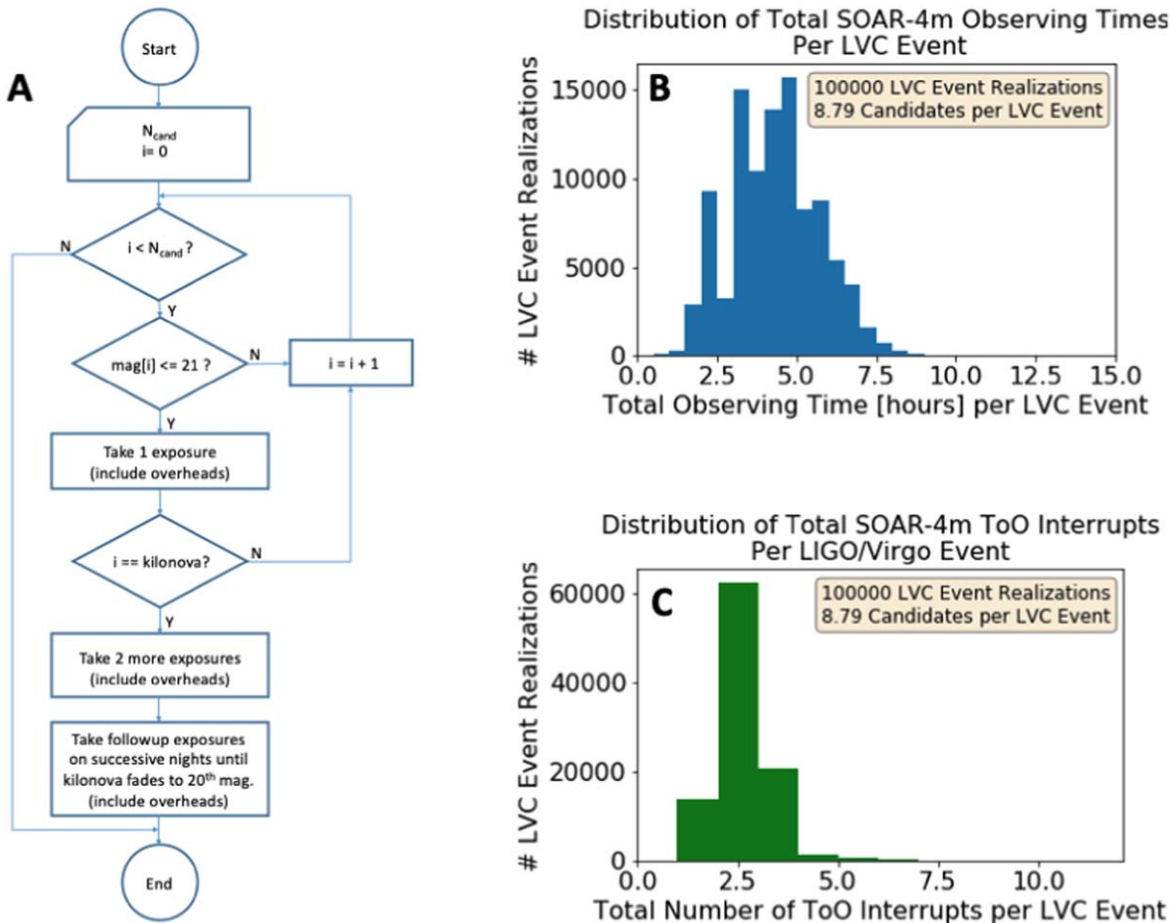


Figure 3. (A) Simplified flow-chart for a single realization of a simulated SOAR follow-up of a single GW event, where N_{cand} is the total number of candidates from an imaging search and discovery program. For the simulations here, N_{cand} is either 8 or 9, but averages overall to 8.79. The distribution of i -band magnitudes for the candidates is drawn from the “LC_SHAPE” row in Figure 2, and the overall average number of candidates (8.79) is just the sum of the entries in the “LC_SHAPE” row. (B) Results of the simulation (using 100,000 realizations): histogram of the total durations of SOAR Target of Opportunity (ToO) interrupt time (in hours) for a single LVC O3 event. (C) Results of the simulation (using 100,000 realizations): histogram of the total number of SOAR ToO interrupts for a single LVC O3 event. (Note that the number of interrupts does not scale exactly as the total duration of interrupt time, since the number of hours per interrupt will vary between the “search and discovery” phase and the follow-up phase of the observations for a given KN event.)

other teams in order of largest flux change to smallest flux change.¹⁰² The selection process for the specific case of GW190814 is illustrated in Figure 4.

In total, 11 objects were targeted with SOAR for either spectroscopic classification of the transient or to obtain a spectroscopic redshift of the host galaxy. These objects are cataloged in Table 2 and their times of photometric discovery and spectroscopic follow-up are shown visually in Figure 5. We note that the observed rate (11 candidates within 48 sq. deg.) well matches the anticipated rate (nine candidates within 60 sq. deg.), and are in fact identical within the Poisson errors.

In Figure 6 we show the expected incidence of each of several types of SN during a search for a KN. These data come from simulated full light curves using the SuperNova ANALysis software (SNANA; see Section 8). The models are the same as in the Photometric LSST Astronomical Time-

series classification challenge (Kessler et al. 2019). We start with ≈ 3300 SNe with a distribution of SN types at random points in their light curves—what one might net in a typical transient search by DECam covering several tens of square degrees—and then apply the selection (culling) steps detailed above, in the end yielding about a dozen SNe whose imaging and photometric properties mimic that of a KN closely enough that they would require follow-up spectroscopy (and/or a more robust photometry-based technique) to eliminate them as candidates in a KN search. This could be viewed as an estimate of the rough contamination rate by SNe in a real-time imaging search using similar candidate selection criteria. Finally, it is interesting to note that the distribution of SN types is very similar between the sample of 3346 SNe that were rejected by the above selection steps and the sample of a dozen SNe that successfully passed through all these steps. In other words, the selection steps do not seem to favor or disfavor any particular SN type.

5. SOAR Observations

In the following section (Section 5.1) we provide details of our ToO triggers and real-time (*not* final) classifications in

¹⁰² Those candidates ruled out by other teams included candidates observed on The Gran Telescopio Canarias (GTC; Castro-Tirado et al. 2019; Hu et al. 2019; Lopez-Cruz et al. 2019a, 2019b), The Southern African Large Telescope (Morgan et al. 2020), and The Giant Magellan Telescope (Morgan et al. 2020), and in general were too faint for SOAR ToO follow-up.

Table 2
Initially Reported Data for the 11 Candidates Described in This Paper

Candidate	Night	GCN	Classification Source	Classification	Redshift
AT2019nqq	Aug 16	25379	Astrodash	Type Ic-broad SN	0.3257
AT2019nqr	Aug 16	25379	Astrodash	Type Iib SN	0.0888
AT2019noq	Aug 20	25423	SNID	Type IIP SN	0.07
AT2019mbq	Aug 20	25423	SNID	Type Ia-CSM SN	0.10
AT2019npw	Aug 26	25484	Astrodash	Type Iib SN	0.163
AT2019num	Aug 26	25484	Astrodash	Type IIP SN	0.113
AT2019ntr	Aug 28	25540	Astrodash	Type II-L SN	0.2
AT2019omx	Aug 28	25540	H α emission line	host galaxy	0.275 ^a
AT2019ntp	Aug 31	25596	Astrodash	Type Ic-BL SN	0.3284
AT2019nte	Sep 13	25784	H α /[N II] emission lines	host galaxy	0.0704 ^a
AT2019omw	Oct 17	N/A	H α emission line	host galaxy	0.0467 ^a

Note. Data include candidate name as assigned by the TNS, night of observation, GCN in which spectral results were reported, source of initial classification and redshift, initial classification, and initial redshift. These are the values reported in the GCNs. (No GCN was submitted for AT2019omw.) These values were updated after full reduction and processing of data. Updated values are given in Table 3. (Astrodash and SNID are SN spectrum-fitting codes; see Section 6.2 and Section 8. Which fitting code was used in this initial classification for a given candidate depended heavily on which team member was available on that night to perform the classification, and the team member’s preference.)

^a Redshift of the host galaxy. Night = civil date of the start of the night of observation, the NOAO convention of designating an observing night. The asterisk to the right of several z values indicates that this is redshift for the host galaxy, as the transient was too dim to observe.

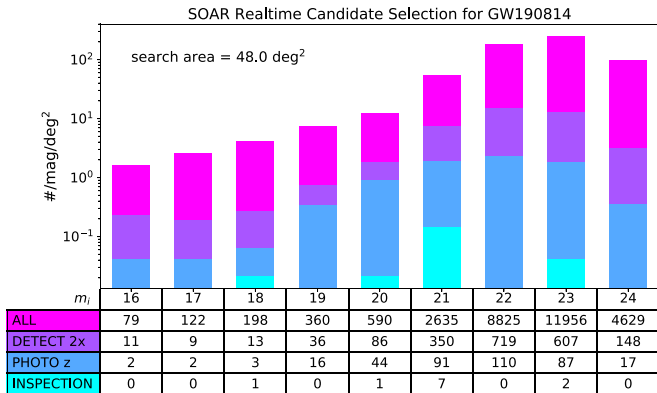


Figure 4. DECam search and discovery candidate selection for spectroscopic follow-up for GW190814. Whereas Figure 2 provided the typical distribution of DECam candidates expected for a typical LVC O3 BNS merger, here we show the corresponding i -band magnitude distribution of all transient candidates observed and visually inspected and identified within the observed area by DECam across the selection criteria of Section 4.2 specifically for the GW event GW190814. The final 11 candidates targeted with SOAR compose the cyan histogram and the “INSPECTION” row; four other candidates, which were in the $i = 21$ –22 range, were observed by other telescopes and are omitted from the cyan histogram and “INSPECTION” row. Note that at the time of SOAR follow-up on three of these transients, their magnitudes had faded below the SOAR detection limit, so we observed their host galaxies to measure their redshifts. (Note: the numbers listed below the plot are the total per magnitude bin for the full 48.0 sq. deg. search area; the y -axis of the plot, however, is the number per magnitude bin *per square degree*.)

search of the optical counterpart of GW190814. We explain how the methods described in Section 4 were executed when our SOAR 2019B ToO program was triggered to observe candidates for an optical counterpart of GW190814.

5.1. GW190814 Candidate Observations

Based on input from the DECam search and discovery program, we developed a list of candidates for spectroscopy as described in the previous section. For the objects possible to observe each night we developed nightly webpages with

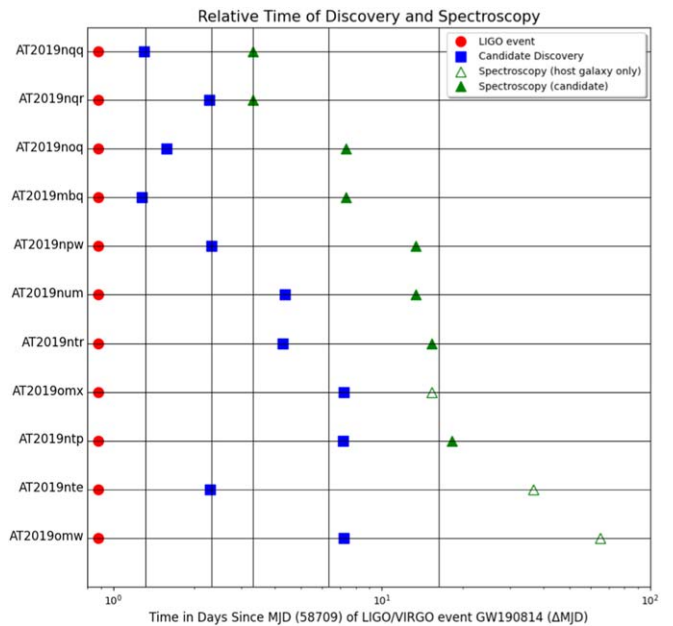


Figure 5. Observational timelines for each KN candidate. All dates are shown as number of days (Δ MJD) since 58,709.00, MJD corresponding to 2019 August 14, the day GW190814 was detected. The time of the neutron star–black hole merger event at MJD 58,709.88 is shown (using a red circle) on each. The date of transient discovery is shown as a blue square. The date of SOAR spectroscopy is shown as a green triangle for each KN candidate (open triangles indicate that spectroscopy was only done for the host galaxy). Vertical lines show the beginning times of DECam observations.

information on object airmasses, finding charts and other information that would be required once our ToO time began. On each night we issued a ToO interrupt, there were several possible KN candidates that could be observed. The selection of which ones were to be targeted for the night was based on observing conditions (e.g., low airmass) and brightest magnitude.

In order to complete data processing in real time, we employed a custom-made reduction pipeline that we developed, a Jupyter notebook we call the SOAR Goodman Quick

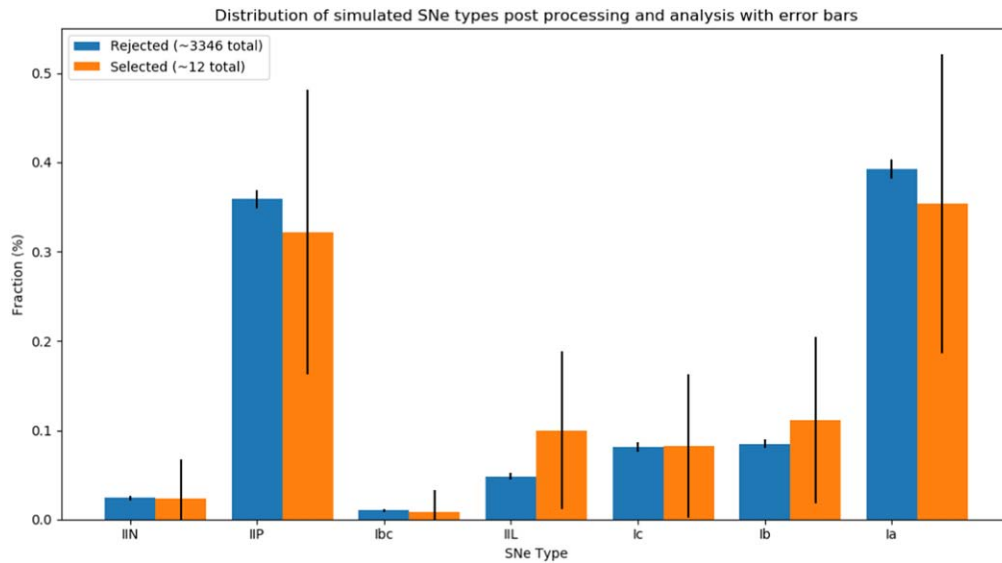


Figure 6. Predictions of the relative incidence of each of several types of SN within a spectroscopic follow-up KN candidate sample post-DECam processing and analysis. The predictions are based on simulated data using SuperNova ANALysis software light-curves and Photometric LSST Astronomical Time-series classification challenge models and run through the selection steps of Morgan et al. (2020). The blue histogram shows the relative distribution of SNe that were rejected by the selection steps; the orange histogram, the relative distribution of SNe that survived (i.e., were selected by) all the selection steps. Similar relative sizes of bars indicates no bias toward any particular SN type. The error largely comes from the Poisson counting statistics.

Table 3
Final Results for the Eight Transients and Three Host Galaxies for Which We Took Spectra.

Name/ID	S/N	AstroDash				SNID				Comments
		Type	r_{lap}	z	M_{abs}	Type	r_{lap}	z	M_{abs}	
AT2019nqc ^b	2.4	Ia-csm	0.14	0.071	-16.8	IIIn	5.3	0.070	-16.8	SNID preferred
AT2019nqr	32.6	Ia-csm	9.97	0.086	-19.6	Ia	4.36	0.101	-20.0	Seyfert 2 AGN @ $z = 0.083$
AT2019noq	7.7	IIIn	19.55	0.074	-17.7	IIP	13.11	0.072	-17.6	AstroDash preferred
AT2019mbq ^b	23.1	IIIn	15.96	0.102	-17.6	Ia	12.09	0.110	-17.8	AstroDash preferred
AT2019npw	6.4	IIP	4.76	0.148	-18.7	IIP	6.44	0.148	-18.7	SNID preferred
AT2019num ^b	7.5	IIL	7.95	0.123	-17.5	IIB	6.96	0.149	-18.0	AstroDash preferred
AT2019ntr ^b	1.8	Ic-broad	0.81	0.224	-19.0	Ia	4.01	0.861	-22.5	None preferred; unknown
AT2019omx ^{a, b}	2.3	host galaxy @ $z = 0.275$ ($M_{\text{abs}} = -18.7$)
AT2019ntp	11.8	Ia-pec	6.44	0.116	-17.7	Ia	12.22	0.114	-17.6	SNID preferred
AT2019nte ^{a, b}	5.8	host galaxy @ $z = 0.0704$ ($M_{\text{abs}} = -16.6$)
AT2019omw ^a	1.8	host galaxy @ $z = 0.0467$ ($M_{\text{abs}} = -13.8$)

Notes. Results include name from the TNS and the S/N of the spectrum calculated using the 6000–6100 Å region. Then we report the outputs from AstroDash and SNID, respectively, including SN type, r_{lap} values, redshift, and absolute magnitude (at DECam discovery; see Table 1). For spectra with $S/N < 5$ and for fits with $r_{\text{lap}} < 6.0$ (AstroDash) or $r_{\text{lap}} < 5.0$ (SNID), the classification may be unreliable.

^a Only the spectrum of the host galaxy was obtained; so it was not fit by either AstroDash or SNID.

^b This candidate lies outside the 90% confidence probability contours of the final LVC map for GW190814; see Figure 1.

Reduce (see Section 8), to obtain quick results immediately after the data are transferred from the SOAR telescope machines. The preliminary processing consists of a quick reduction of the spectra using an arc-lamp wavelength calibration frame and a calibration from a standard star taken at the start of ToO observing. This publicly available Jupyter notebook takes the 2D spectrum, extracts the 1D spectrum, and performs basic wavelength and spectrophotometric calibration with relatively simple and straightforward inputs. With a little practice, using the SOAR Goodman Quick Reduce is just as fast as using the IRAF implot task—but with the added

advantage of providing a quick calibrated spectrum. Generally, a “by eye” check of the calibrated spectrum indicates whether or not a candidate is a KN—usually due to the disqualifying presence of one or more relatively sharp emission lines or the spectral features of an SN—but, even so, each calibrated spectrum was also sent that same night to one of our SN-fitting experts, who would fit the spectrum to SN model spectra. The resulting spectra were intended to be analyzed with fast classification tools (see below) and the spectroscopic class and redshift of the transient to be published promptly to the community via a GCN circular. The list of objects for which

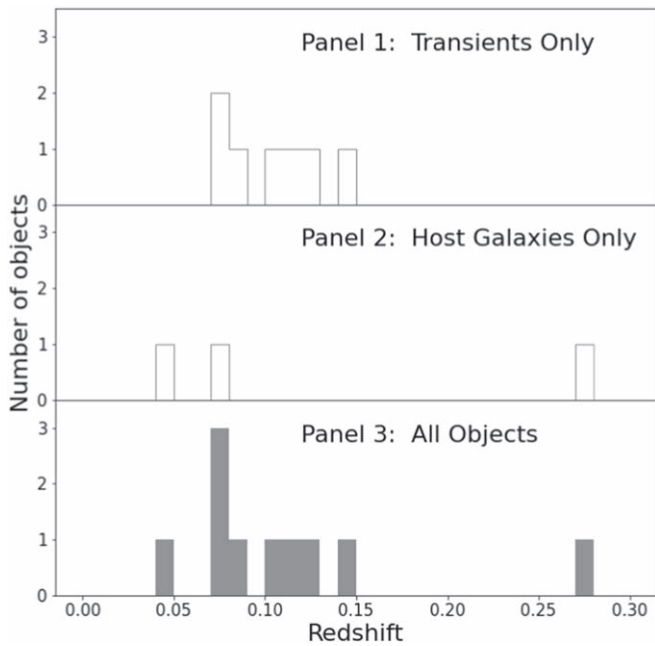


Figure 7. Histograms of the redshifts of the 11 candidates, using final preferred results from Table 3. The top panel is for the eight transient targets alone, the middle panel is for the three host galaxy targets alone, and the bottom panel is for all 11 SOAR targets combined (transients and host galaxies together).

spectra were taken, along with initial redshift and SN classifications and the GCNs the DESGW SOAR observing team issued, is given in Table 2.

To avoid fatigue, the DESGW SOAR spectroscopy task force was divided into four teams—one based in Brazil (PI M. Makler), one based in Chile (PI F. Olivares), one based at UC-Santa Cruz (PI C. Kilpatrick), and one based at Fermilab (PI D. Tucker)—each team signing up for multiple two-week shifts throughout the course of LVC O3. Our default plan was to use the Goodman HTS Blue camera, the 4001 mm^{-1} grating in its M1 configuration, and a slit width of $1''$, to yield a wavelength range of roughly $3000\text{--}7050\text{ \AA}$ at a resolution of $R \sim 930$ (e.g., see Nicholl et al. 2017) but, if the night’s main program that our ToO was interrupting was using a roughly similar configuration, we could also use that instead, minimizing issues with switch-overs from and to the main program.

5.1.1. Observations

We issued ToO interrupts on 2019 August 16, 20, 26, 28, and 31 (start dates, based on local time). On several other nights we attempted to conduct ToO observations, but found skies to be too cloudy to effectively observe and so we canceled the ToO interrupts. During the course of the 2019 August observations, the Fermilab and Chilean teams were on shift. In addition, spectra were taken for us by SOAR scientific staff during the SOAR engineering nights of September 13 (host galaxy for AT2019nte) and October 17 (host galaxy of AT2019omw). This information and the GCNs issued are summarized in Table 2.

In Figure 5 we graphically summarize our sequence of observations. In this figure we show a set of timelines indicating the dates of discovery and SOAR spectroscopy of each of the candidates we observed, using a log scale for the x -

axis. The first mark (red circle) on each timeline is the MJD of the GW190814 merger event. The second mark (blue square) is the date of discovery in DECam observations. The third mark (green triangle) indicates the date of SOAR spectroscopy. Vertical lines are also included that show the date of DECam observations, as described in Morgan et al. (2020). The marks denoting SOAR spectroscopy of AT2019nte, AT2019omw, and AT2019omx are unfilled, indicating that we did not take spectroscopy of the transient but of the host galaxy only. We report redshifts of these host galaxies in Table 2. The horizontal axis is given in ΔMJD , time in days since MJD 58709.

Even though none of these 11 candidates were determined to be the optical counterpart of GW190814, these results will permit important upper limits to be established in preparation for future searches for the optical counterparts of these types of mergers (see the next section).

6. Results and Discussion

In this section, we cover our final results from our SOAR observations of the GW190814 candidates. In Section 6.1 we describe the full reduction and analysis of spectra and present the spectra themselves. In Section 6.2 we present classifications of the SNe and consider our methods of analysis. In Section 6.3 we fit each spectrum with Kasen et al. (2017) KN models; as nearly all were found to be an SN, the KN models are generally not good fits. In Section 6.4, we discuss the three candidates for which we only obtained spectra for the host galaxy and the likelihood that either of these three candidates could be the optical counterpart for GW190814. Finally, in Section 6.5 we consider lessons learned in LVC O3 that can be applied as we prepare for LVC observing season O4.

6.1. Spectral Data from SOAR Telescope

For the final reduced spectra (shown in Figures 8–18)—unless otherwise noted¹⁰³—we employed the UCSC spectral pipeline (link to Github repository in Section 8). This pipeline consists of the standard steps for the processing of optical spectroscopic data: bias subtraction, flat-fielding, extraction of the 1D spectrum and flux, and wavelength calibration against a standard star, typically a Hamuy Tertiary Standard Star (Hamuy et al. 1992, 1994). These more careful reductions, performed later, are the same as those used in the recent GW190914 omnibus paper by Kilpatrick et al. (2021).

6.2. SN Classifications

Offline analysis of the spectra we obtained was performed using the public codes Super Nova IDentification (SNID; Blondin & Tonry 2007) and Deep Automated Supernova and Host classifier (DASH, a.k.a. AstroDash; Muthukrishna et al. 2019) (see Section 8). SNID is a template-fitting method based on the correlation techniques by Tonry & Davis (1979). AstroDash is a deep convolutional neural network used to train a matching algorithm. These analysis tools provide spectral matching, which allowed us to classify our spectra by means of a comparison against a spectral library of transients and other astrophysical sources. We chose these codes as SNID has been

¹⁰³ For the final reduced spectra for the host galaxies of AT2019nte and AT2019omw, we made use of standard IRAF reductions provided by the SOAR science staff.

AT2019noq

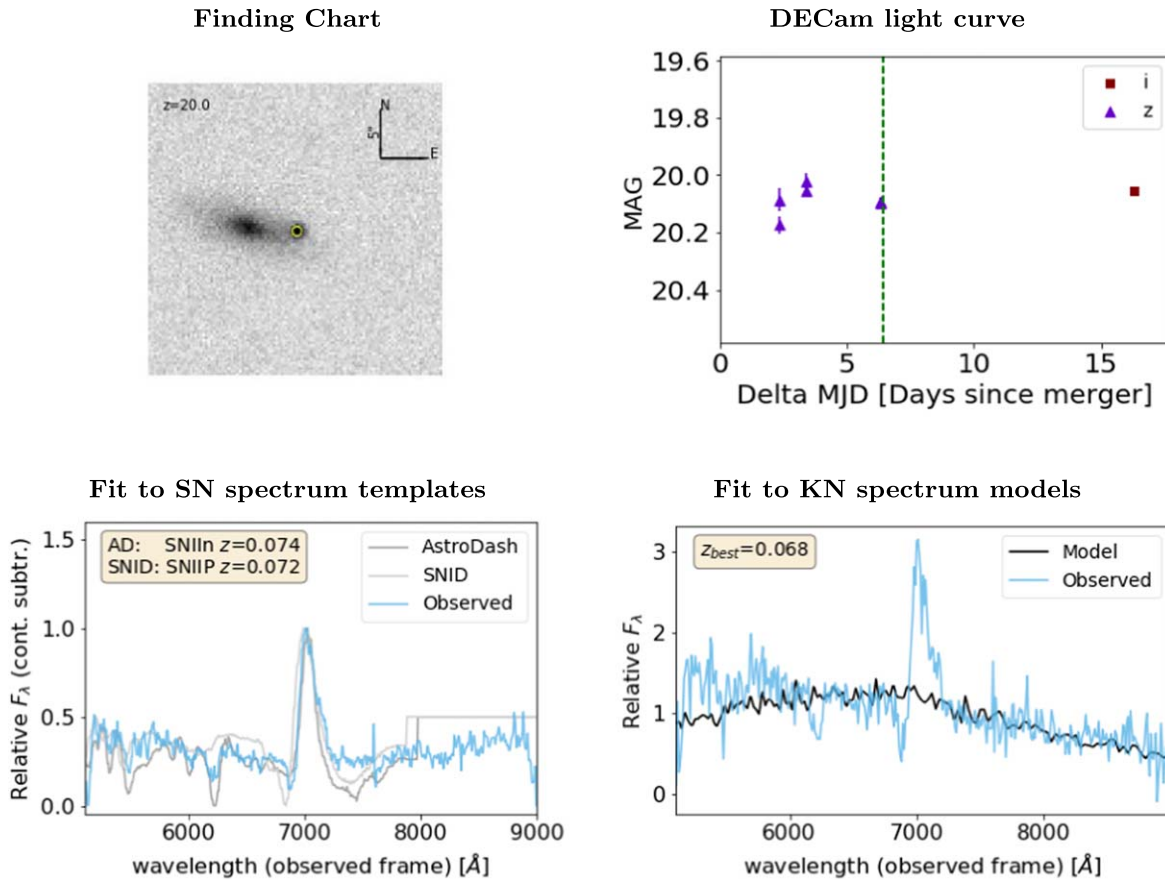


Figure 8. Top left: thumbnail-finding chart (using the DECAM imaging) for the AT2019noq KN candidate; the location of the candidate is marked by a small yellow circle. Top right: the candidate’s i - and z -band light curves from DECAM photometry; the vertical dashed green indicates when SOAR spectroscopy was obtained. Bottom left: observed and best-fit SN model spectrum for the candidate object. Light blue is the processed, calibrated, and continuum-subtracted observed spectrum, dark gray is the best-fit SN model from AstroDash, and light gray is the best-fit SN model from Super Nova IDentification (SNID). In the panel we provide the best-fit SN type and redshift from the two codes. Bottom right: observed and best-fit model KN spectra for the candidate objects. Light blue is the processed and calibrated observed spectrum; black is the best-fit Kasen et al. (2017) KN model. In the panel we provide the best-fit value of the redshift, z_{best} . Unlike in the AstroDash/SNID fits plot, the continuum has not been subtracted. Also, a slightly different smoothing technique is used for the SN fits and for the KN fits.

used extensively by the community and AstroDash makes use of a powerful deep learning technique. We discuss below the importance of using more than one SN typing package to check results.

For our AstroDash fits of the spectrum of each candidate, we applied an AstroDash smoothing length of 3 (unless otherwise stated), and we left the redshift a free parameter. We then visually inspected the 20 best SN template fits for that candidate, choosing the top two for further consideration. (The top two fits based on visual inspection also typically had among the highest r_{lap} values of the 20 best fits.¹⁰⁴) Unless there were other relevant considerations (e.g., the putative epoch in the light curve at which the spectrum was obtained), the SN template spectrum with the higher of the two r_{lap} values was chosen as the final best fit.

¹⁰⁴ r_{lap} is a measure of the quality of the fit that combines the correlation between the observed and the template spectrum with the amount of overlap in $\ln \lambda$ -space between the observed and the template spectrum. The higher the value of r_{lap} , the higher the quality of the fit. For the detailed definition, see Blondin & Tonry (2007).

For our SNID fits of the spectrum of each candidate, we applied the default SNID smoothing length of 1 pixel and, as with our AstroDash fits, we also fit for the redshift. We visually inspected the top five SN template fits for each candidate, but in the end chose the one with the highest r_{lap} as our SNID classification.

In Table 3 we present final measurements from AstroDash and from SNID for the eight transients of which we took spectra. (For completeness, we also include information on the three candidates for which we only obtained host galaxy spectra: AT2019omx, AT2019nte, and AT2019omw.) These results are based on the final reduced spectra. This table includes classification, the redshift, and a measure of the goodness of fit (r_{lap}) from these two SN spectrum-fitting codes. We kept redshifts as free parameters in the fitting; the photometric redshifts of the host galaxies were used during the selection process of candidate objects discussed in Section 3.

The distribution of the redshifts from the preferred fits in Table 3 is given in Figure 7; as expected, transients were found

AT2019mbq

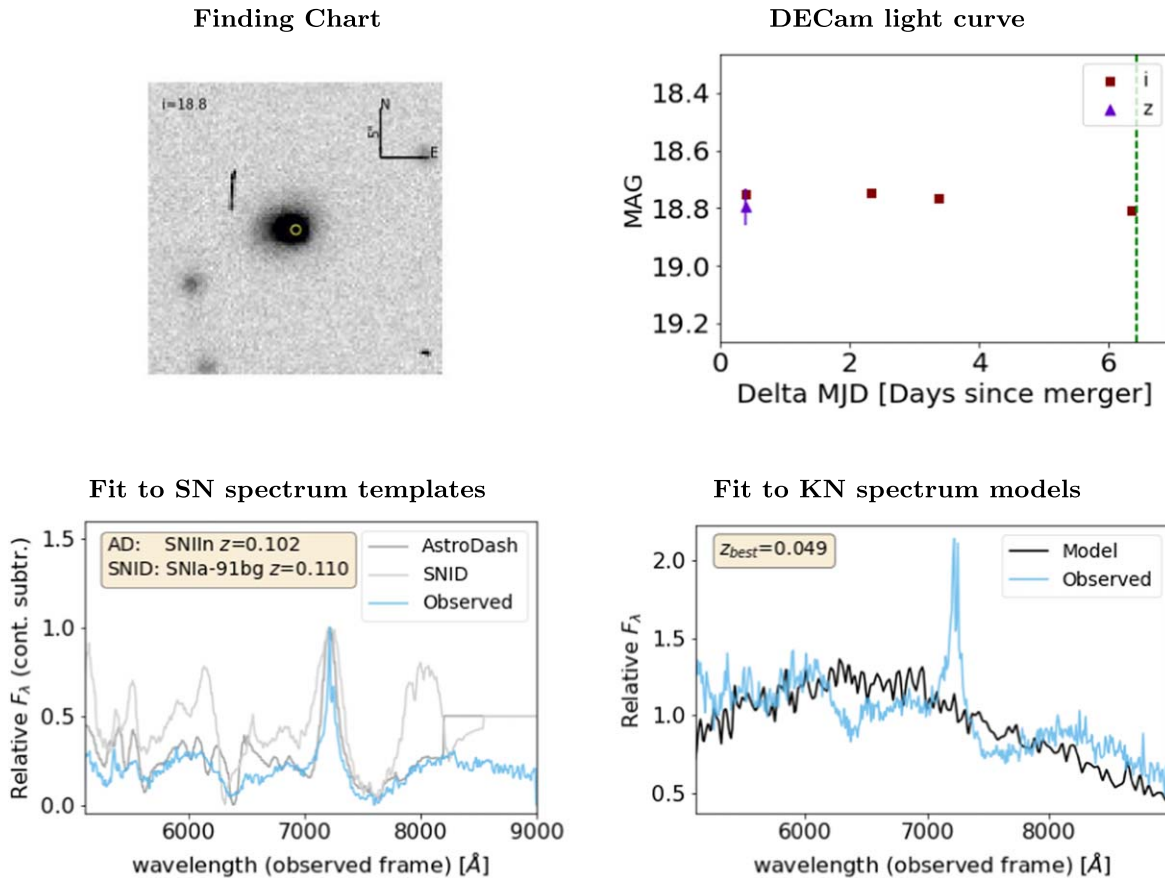


Figure 9. Same as Figure 8 but for the AT2019mbq KN candidate.

over a range of redshifts with a predominance of lower- z objects.

In Figures 8–18, we provide the following information for each candidate: a thumbnail-finding chart containing the host galaxy and marking the location of the transient; the DECcam-based i - and z -band light curves for the transients; and the final reduced observed spectrum. For the candidates for which we only obtained the host galaxy spectrum,¹⁰⁵ that is the sum of what we show in these figures. For candidates for which we took a spectrum of the transient candidate itself, we also include the best-fit SN templates from AstroDash and SNID and the best-fit KN model from Kasen et al. (2017) overlaid on the final reduced observed spectrum. As shown below, the interplay of these different types of data often helped in the final classification of a given candidate.

6.2.1. AT2019noq

For AstroDash, our two best fits were a $z=0.074$ SN IIn 42–46 days past maximum light ($r_{lap}=19.55$) and a $z=0.079$ SN IIP 2–6 days past maximum light

($r_{lap}=19.31$). The DECcam light curve was relatively flat over the period it was observed (Figure 8); so we chose the SN IIn classification as more likely. For SNID, our best fit was a $z=0.072$ SN IIP 9.8 days past maximum light ($r_{lap}=13.11$). Due to its higher r_{lap} value, the AstroDash fit is preferred; see Figure 8.

6.2.2. AT2019mbq

Recall that a spectrum of AT2019mbq was mistakenly observed by SOAR (the original target was AT2019ntn), and that there was evidence of a detection of AT2019mbq *before* the GW190814 merger event, making it highly unlikely that AT2019mbq is the optical counterpart.

For AstroDash, our two best fits were a $z=0.102$ SN IIn 46–50 days past maximum light ($r_{lap}=15.96$) and a $z=0.103$ SN IIn 42–46 days past maximum light ($r_{lap}=14.92$). The difference between the two classifications was small, and the DECcam light curve provided no strong motivation to choose one over the other (Figure 9); so we chose the template with the higher r_{lap} (a $z=0.102$ SN IIn 46–50 days past maximum light) as more likely. For SNID, our best fit was a $z=0.110$ SN Ia 45.9 days past maximum light ($r_{lap}=12.09$). Despite the SNID fit’s relatively high r_{lap} value, a visual inspection of both the AstroDash and the SNID spectral fits (Figure 9) leads us to prefer the AstroDash fit.

¹⁰⁵ Note that, within the 2.5 hr time constraint of a SOAR ToO interrupt, we were basically confined to observing targets that were $i \lesssim 21.5$; so, in some cases—especially for the later targets—we instead obtained spectra of the candidate’s host galaxy as a means of excluding the target by its redshift: i.e., if the redshift of the candidate’s host galaxy is substantially discrepant from the redshift expected for the luminosity distance of the GW event ($z_{GW} = 0.059 \pm 0.011$), we can exclude that candidate.

AT2019npw

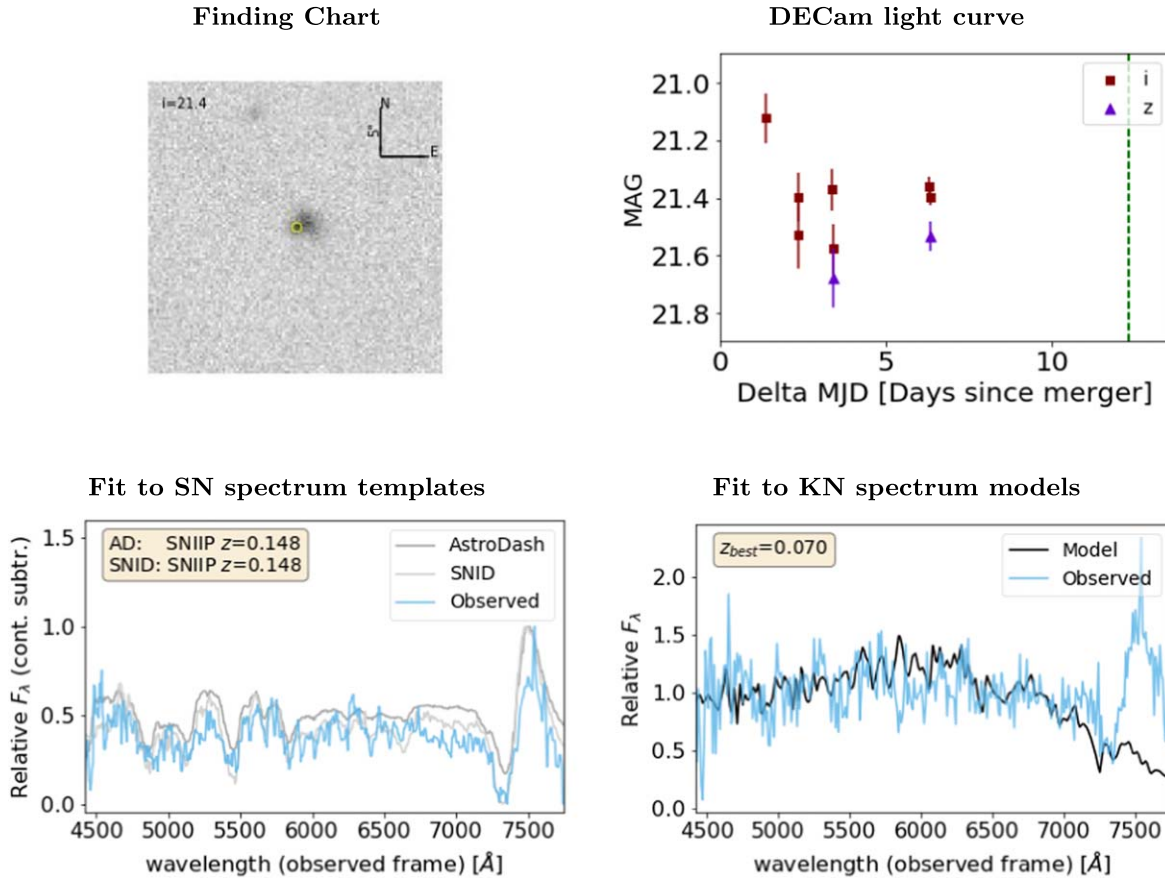


Figure 10. Same as Figure 8 but for the AT2019npw KN candidate.

6.2.3. AT2019npw

For AstroDash, our two best fits were a $z=0.148$ SN IIP 18–22 days past maximum light ($r_{\text{lap}}=4.76$) and a $z=0.147$ SN IIP 22–26 days past maximum light ($r_{\text{lap}}=4.72$). The difference between the two classifications was small, and the DECcam light curve provided no strong motivation to choose one over the other; so we chose the template with the higher r_{lap} (a $z=0.148$ SN IIP 18–22 days past maximum light) as more likely. The relatively low r_{lap} values ($r_{\text{lap}} < 6$), however, are of some concern. For SNID, our best fit was a $z=0.148$ SN IIP 44.3 days past maximum light ($r_{\text{lap}}=6.44$). Due to its higher r_{lap} value, the SNID fit is preferred, see Figure 10.

6.2.4. AT2019num

For AstroDash, our two best fits were a $z=0.123$ SN IIL 6–10 days past maximum light ($r_{\text{lap}}=7.95$) and a $z=0.239$ SN Ibn 22–26 days past maximum light ($r_{\text{lap}}=0.4$). Since the DECcam light curve for this candidate is rising noticeably 10–6 days before the SOAR spectrum was obtained (Figure 11), it appears that this candidate is a likely a young SN; that, combined with the substantial difference in r_{lap} values, led us to choose the $z=0.123$ SN IIL 6–10 days past maximum light template as the more likely classification. (We

note that, for AT2019num, we used a smoothing length of 6 instead of 3 for our AstroDash fits.) For SNID, our best fit was a $z=0.149$ SN IIb, 17.3 days *before* maximum light ($r_{\text{lap}}=6.96$). Due to its higher r_{lap} value (and the relative rarity of catching a SN so early before maximum light), the AstroDash fit is preferred; see Figure 11.

6.2.5. AT2019ntr

For AstroDash, our two best fits were a $z=0.224$ SN Ic-broad near maximum light (between 2 days before and 2 days after peak; $r_{\text{lap}}=0.81$) and a $z=0.264$ SN Ia-csm 6–10 days past maximum light ($r_{\text{lap}}=0.76$). The DECcam light curve seems to be slightly rising 11–8 days before the SOAR spectrum was taken (Figure 12), indicating a relatively young SN. Due to the low S/N of the spectrum (1.8) and the poor r_{lap} values for the fits, we are reluctant to assign a classification based on the AstroDash fits; that said, the $z=0.224$ SN Ic-broad template near maximum light appears to be marginally better.

For SNID, our best fit was a $z=0.861$ SN Ia 11.2 days *before* maximum light ($r_{\text{lap}}=4.01$). Given a discovery z -band magnitude of 21.2 (Table 1), a redshift of $z=0.861$ implies a z -band absolute magnitude of roughly $M_{\text{abs}}=-22.5$, or substantially more luminous than a typical SN Ia (Richardson et al. 2014). We therefore view the SNID fit as unreliable.

AT2019num

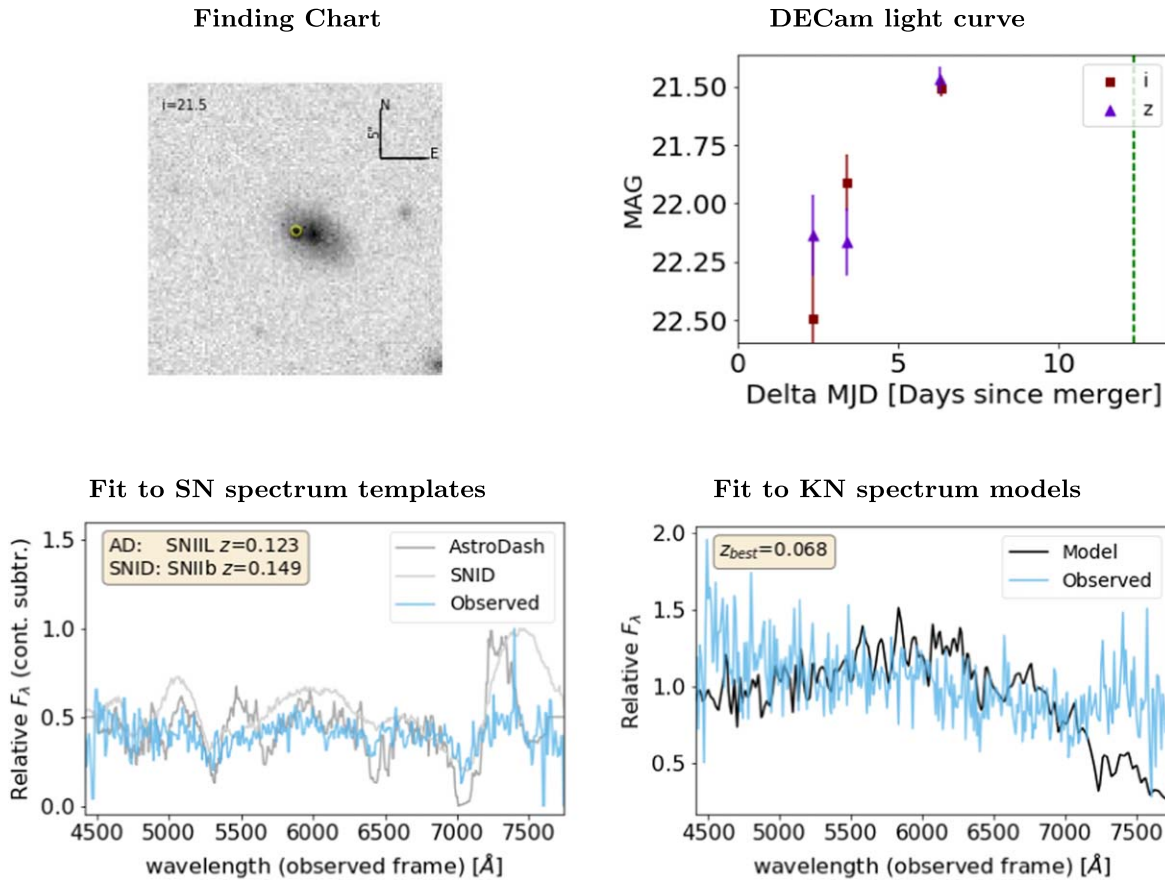


Figure 11. Same as Figure 8 but for the AT2019num KN candidate.

Due to the noisiness of this spectrum and the problems with both the AstroDash and the SNID fits, we prefer neither the AstroDash nor the SNID classifications. We therefore view AT2019ntr’s spectral classification as unknown; see Figure 12. In hindsight, AT2019ntr would have been a natural candidate for additional spectroscopy with a larger telescope.

6.2.6. AT2019ntp

For AstroDash, our two best fits were a $z = 0.116$ SN Ia-pec 34–38 days past maximum light ($r_{\text{lap}} = 6.44$) and a $z = 0.331$ SN Ic-broad 26–30 days past maximum light ($r_{\text{lap}} = 4.35$). The DECcam light curve provided no strong motivation to choose one over the other (Figure 13); so we chose the template with the higher r_{lap} (a $z = 0.116$ SN Ia-pec 34–38 days past maximum light) as more likely. For SNID, our best fit was a $z = 0.114$ SN Ia 45.8 days past maximum light ($r_{\text{lap}} = 12.22$). Due to its higher r_{lap} value, the SNID fit is preferred; see Figure 13.

6.2.7. AT2019nqr

For AstroDash, our two best fits were a $z = 0.086$ SN Ia-csm 46–50 days past maximum light ($r_{\text{lap}} = 9.97$) and a $z = 0.086$ SN II_n 46–50 days past maximum light ($r_{\text{lap}} = 7.85$). We chose the template with the higher r_{lap} value as the better fit,

despite that none of the SN templates did a reasonable job at fitting the narrow-but-strong emission lines at the observed wavelengths of 5371 and 5422 Å, and despite that the DECcam light curve indicated that the transient may have been near a maximum brightness when the spectrum was observed. For SNID, our best fit was a $z = 0.101$ SN Ia 5.7 days past maximum light ($r_{\text{lap}} = 4.36$). In the end, due to this candidate’s central location in a spiral galaxy and a spectrum that well fits that of a Seyfert 2 at $z = 0.083$, we classify AT2019nqr as a Seyfert 2 AGN; see Figure 14.

6.2.8. AT2019nqq

For AstroDash, our two best fits were a $z = 0.071$ SN II_n 14–10 days before maximum light ($r_{\text{lap}} = 0.57$) and a $z = 0.071$ SN Ia-csm 6–10 days past maximum light ($r_{\text{lap}} = 0.14$). The DECcam light curve appears to show a very slight fading over the short time it was monitored before the spectrum was taken (about 1 day before the SOAR spectrum was obtained; Figure 14); so we chose the second template (a $z = 0.071$ SN Ia-csm 6–10 days past maximum light) as more likely, even though it has a lower r_{lap} . We note that the observed spectrum contains a prominent H α emission line redshifted to 7028 Å and a less prominent [O III] 5007 emission line redshifted to 5362 Å, and an even less prominent H β emission line redshifted to 5205 Å. For SNID, our best fit

AT2019ntr

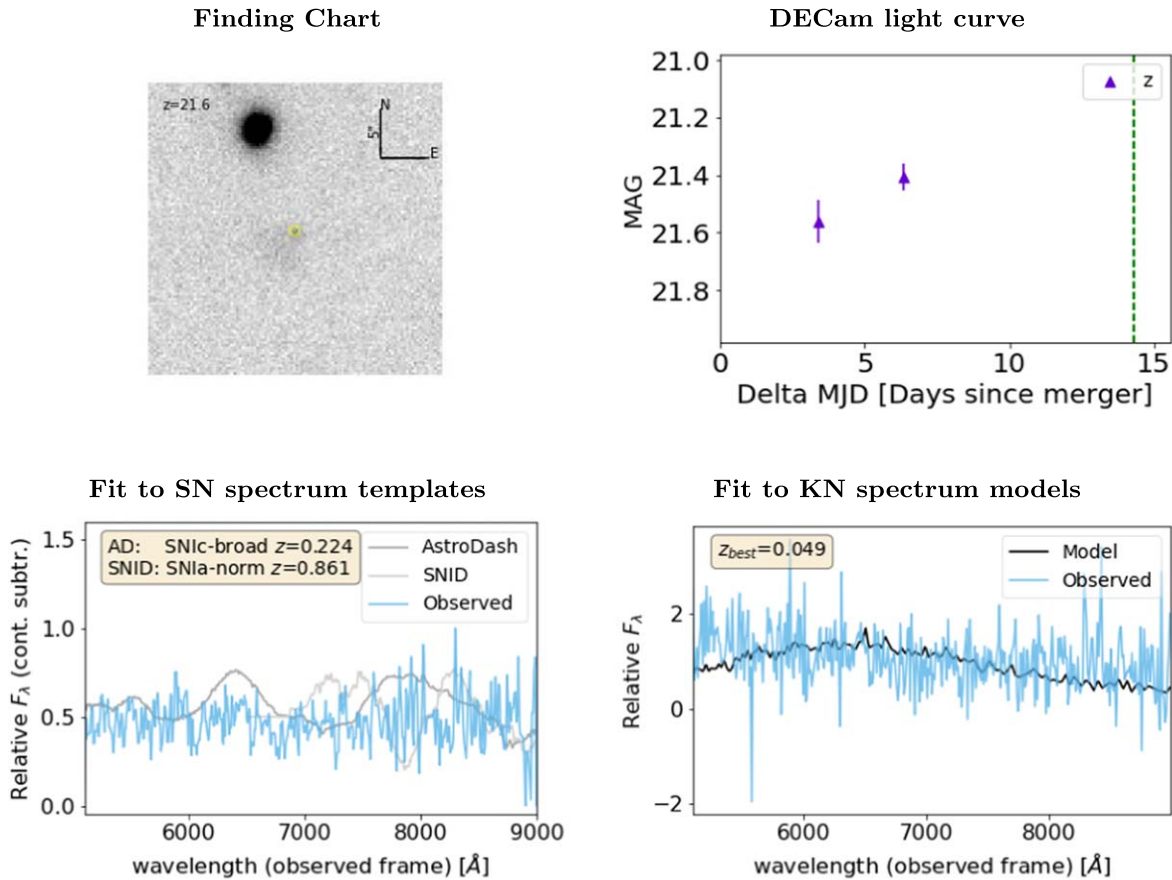


Figure 12. Same as Figure 8 but for the AT2019ntr KN candidate.

was a $z=0.070$ SN IIn 50.2 days past maximum light ($r_{\text{lap}}=5.3$). Due to its higher r_{lap} value, the SNID fit is preferred; see Figure. 15.

We note that AT2019nqq was one system for which we could compare results from another facility. It was also observed by the GTC 10.4 m (GCN25419), classified as a Type IIP SN at 4 days post maximum with $z_{\text{host}}=0.071$. Although the type classification differs from our result for this system (SN IIn), the redshift estimate is consistent with ours.

In closing, we found that some classifications from both AstroDash and SNID might be inconclusive. For one case, AT2019ntr, this is probably related to the low-S/N spectrum, in which the low value of r_{lap} from both SNID and AstroDash points toward a poor fit. It is also worth re-iterating that our methods of choosing the best fits differed for the two packages: for AstroDash, we depended more on a visual inspection of the 20 models with the highest r_{lap} values; for SNID, we basically chose the model with the highest r_{lap} value. This can lead to different classifications for the same object. In general, for a fit of a relatively high-S/N spectrum ($S/N \geq 5$) and a relatively high value for r_{lap} (≥ 6.0 for AstroDash; ≥ 5.0 for SNID), we view the classification (AstroDash or SNID) with the higher the value of r_{lap} as the preferred classification; in cases of a low-S/N spectrum ($S/N < 5$), we view neither AstroDash’s nor SNID’s classification as particularly reliable. These results enhance the

importance of using multiple methods to perform spectral classification.

6.3. Spectral Fitting with KN Models

KNe are expected to produce quasi-blackbody radiation. They are expected to have a rapidly changing lightcurve, a luminosity consistent with nuclear rapid neutron capture (r -process) heating, and a long-lived infrared emission. Analysis of the spectrum of AT2017gfo (the KN associated with GW170817) showed emission from both light r -process and heavy r -process components which led to a spectrum that appears as a superposition of two blackbodies at different temperatures. At early times the spectra are mostly featureless, while at later times there are distinct features in the infrared.

For our analysis, we used the set of synthetic KN spectra by Kasen et al. (2017) (see Section 8). This set of Kasen et al. models covers a regularly sampled grid in parameter space of ejecta mass ($M=0.001-0.1M_\odot$), ejecta velocity ($v_{\text{kin}}=0.03-0.40c$), and ejecta lanthanide mass fraction ($X_{\text{lan}}=10^{-9}-10^{-1}$). At each of these grid points in ($M, v_{\text{kin}}, X_{\text{lan}}$)-space is a time series of synthetic spectra spaced in units of 0.1 day from ≈ 2 days pre-merger out to ≈ 25 days post-merger. Each of these synthetic spectra covers a rest-frame wavelength range from the ultraviolet ($\approx 150 \text{\AA}$) through the infrared ($\approx 10 \mu\text{m}$).

AT2019ntp

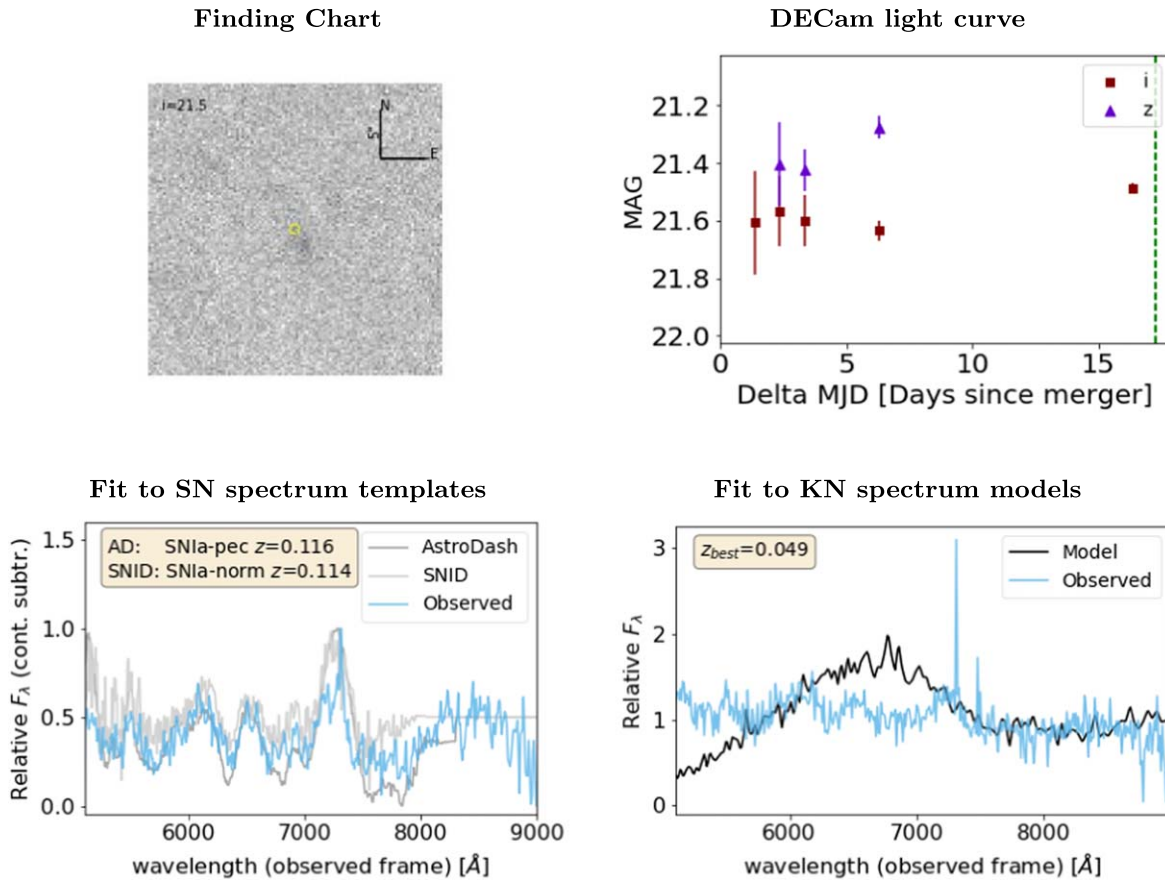


Figure 13. Same as Figure 8 but for the AT2019ntp KN candidate. (Due to the additional smoothing in the SN-fitting plot, the strong narrow emission line seen in the KN-fitting plot is mostly washed out.)

We took the processed and calibrated observed spectrum for each of our KN candidates and performed a least-squares fit to the Kasen et al. (2017) grid of synthetic spectra for the appropriate time post-merger when the candidate’s spectrum was observed. In this fit, the redshifts of the synthetic spectra were also allowed to float within a 1σ range centered on the estimated redshift of the LVC source ($z = 0.059 \pm 0.011$), yielding a best-fit spectrum at a best-fit redshift.

In Figures 8–15 we show the results of these fits for our sample of observed KN candidate spectra. With the possible exception of AT2019ntp, none of these candidates has an observed spectrum that is a particularly good fit to the Kasen et al. (2017) models—mostly due to the appearance of one or more strong emission features in the observed spectrum—which is consistent with our conclusion that none of these objects is a KN, but rather each is an SN from one of several types. What of AT2019ntp? For this object the best-fit redshift ($z_b = 0.049$) is on the low end, but still within the 1σ errors from the redshift based on the original LVC O3 distance estimate ($z = 0.059 \pm 0.011$). Furthermore, this is one of the cases where the AstroDash and SNID fits are both poor (low `rlap`) and inconsistent with each other (see Table 3). So, is AT2019ntp the optical counterpart to GW190814? Unfortunately, we cannot provide a definite conclusion based on the SOAR data alone. As it turns out, though, it is unlikely that

AT2019ntp is the KN we were seeking: first, its sky coordinates lie outside the final LVC 90% confidence contour for GW190814 (see Figure 1); second, and more importantly, in their analysis of the DECcam data for these candidates, Morgan et al. (2020) applied a light-curve-based ML classifier—a combination of Sako et al. (2011)’s `PSNID` fitting code and a random forest classifier—to the photometric time series data for AT2019ntp, and this yielded a 96% probability that AT2019ntp is an SN.

Finally, it might be asked whether it would not be more efficient to add the Kasen templates into AstroDash/SNID so one could directly compare the likelihood that an object is a classical SN versus a KN. One of the first things AstroDash/SNID does is to fit the continuum of the spectrum and remove it. KN spectra—especially early on in their light curves—are continuum dominated, with few prominent emission/absorption features. Thus, there would be little left to fit in the case of the KNe models. Maybe a version of AstroDash/SNID that did *not* subtract off the continuum during the fit would work, but that would be a future project.

6.4. Spectra of Host Galaxies

Finally, there were three candidates which were too faint for us to target effectively with SOAR (AT2019nte, AT2019omw,

AT2019nqr

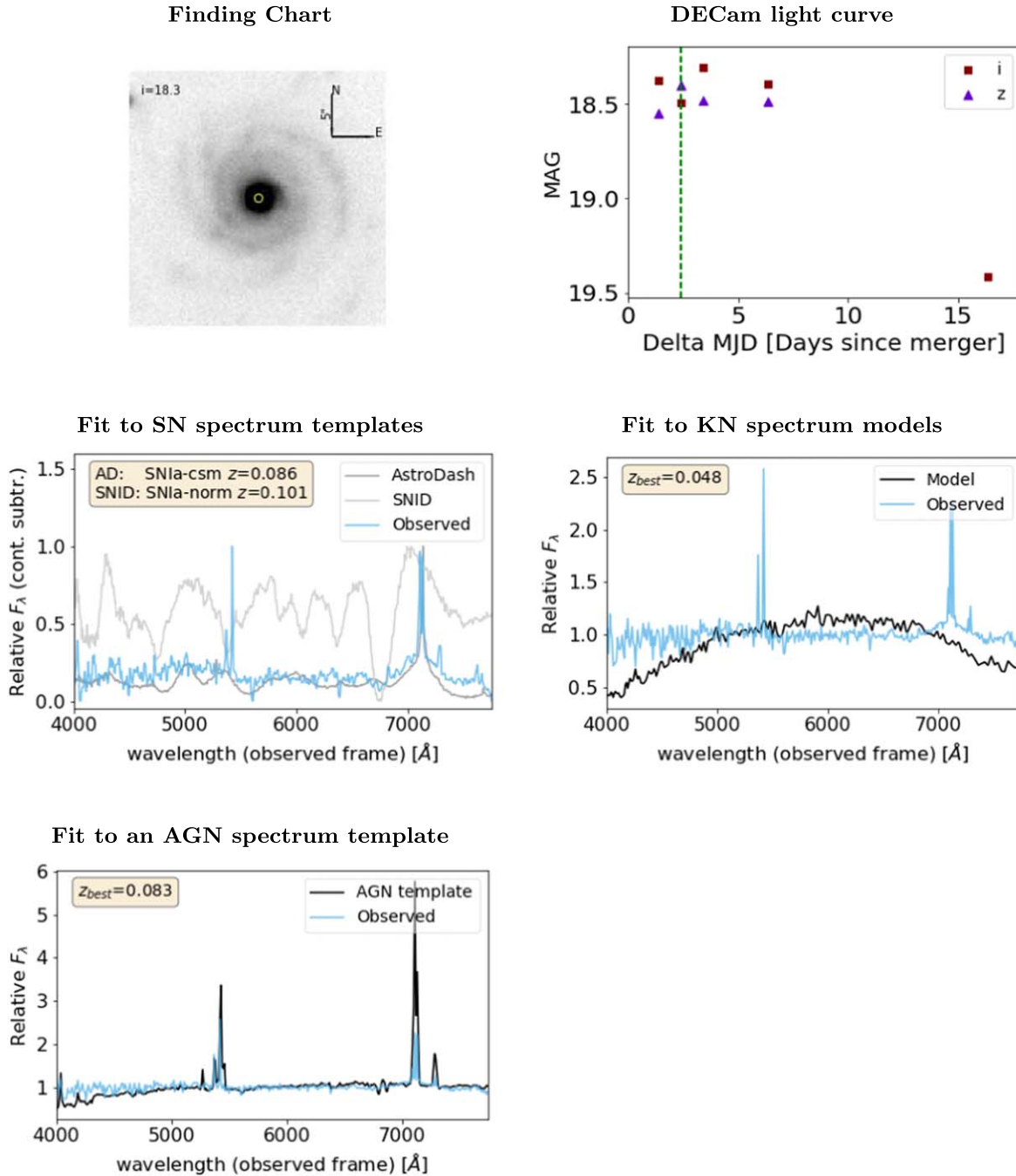


Figure 14. Same as Figure 8 but for the AT2019nqr KN candidate. We also show the best fit to active galactic nucleus (AGN) template spectra, which is that of a Seyfert 2.

AT2019omx). We instead targeted the host galaxy, with the idea that, if the host galaxy’s redshift was significantly discordant with that of the distance estimated from the GW signal, that would rule out that candidate as a possible counterpart to GW190814. We found that only one (AT2019omx) had a truly discordant redshift ($z = 0.275$); see Figure 16. The host galaxies of the other two candidates, AT2019nte ($z = 0.070$; Figure 17) and AT2019omw

($z = 0.047$; Figure 18) have redshifts that are consistent with the redshift corresponding to the GW distance at about the 1σ level. As it turns out, in the end both AT2019nte and AT2019omw failed the DESGW Search and Discovery *offline* imaging pipeline criteria for a good candidate: AT2019nte because it did not meet a sufficiently high detection threshold in the DECam imaging, and AT2019omw because it did not survive the offline visual inspection of candidates (Morgan

AT2019nqq

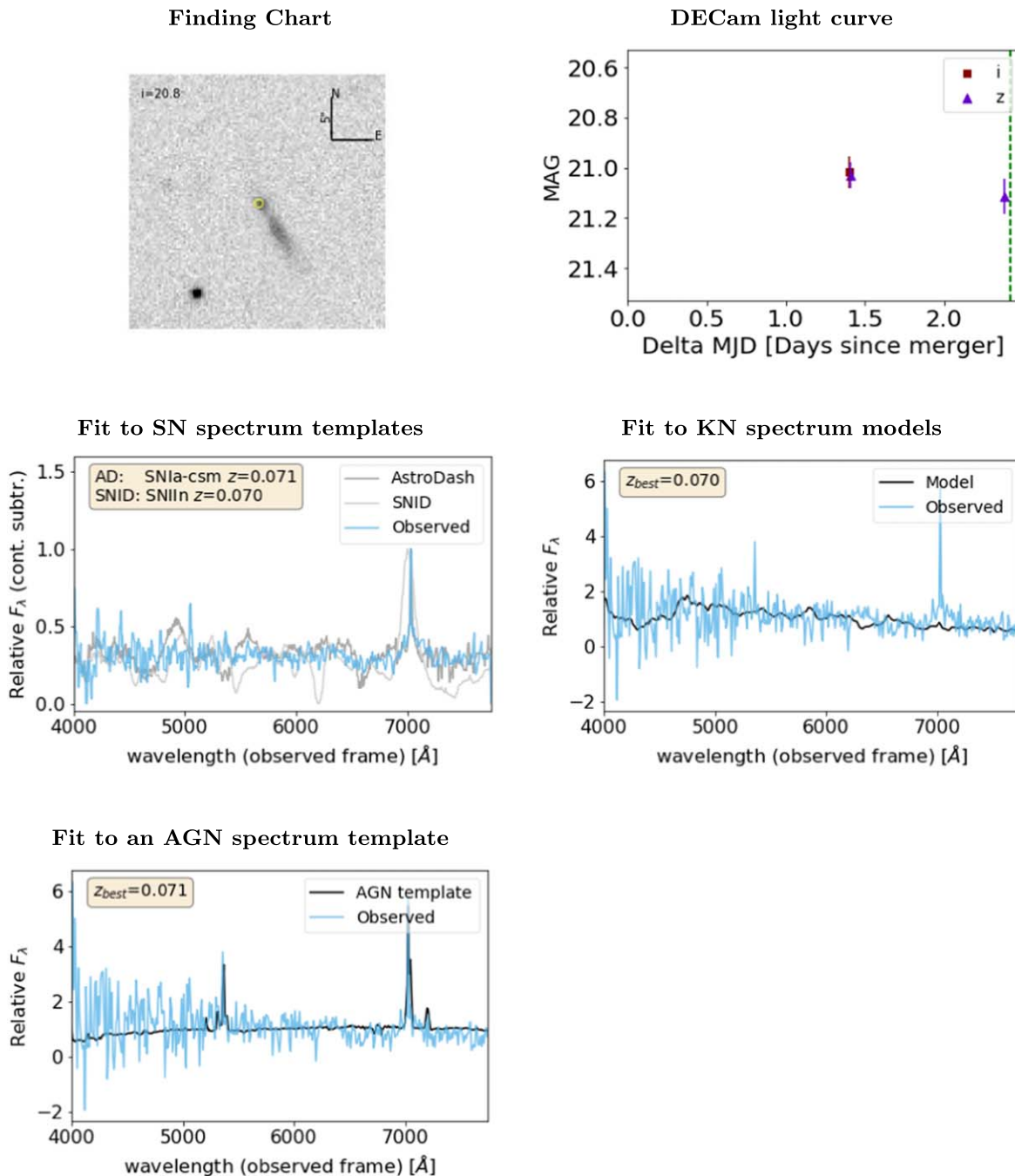


Figure 15. Same as Figure 8 but for the AT2019nqq KN candidate. We also show the best fit to AGN template spectra, which is that of a Seyfert 2.

et al. 2020). Thus, we consider all three of these candidates as being ruled out.

6.5. Lessons Learned from DESGW Spectroscopy in O3

One of the final results we would like to discuss are those of “lessons learned” during the concerted effort by the DESGW imaging and spectroscopic follow-up teams during the follow-

up of GW190814 candidates, particularly as the spectroscopic follow-up of this LVC event may be viewed as a template for future spectroscopic follow-ups in LVC O4 and beyond since, as the LVC becomes increasingly more sensitive, the optical counterparts of future LVC events will likely be relatively distant and faint, unlike the very nearby and bright BNS KN GW170817.

AT2019omx

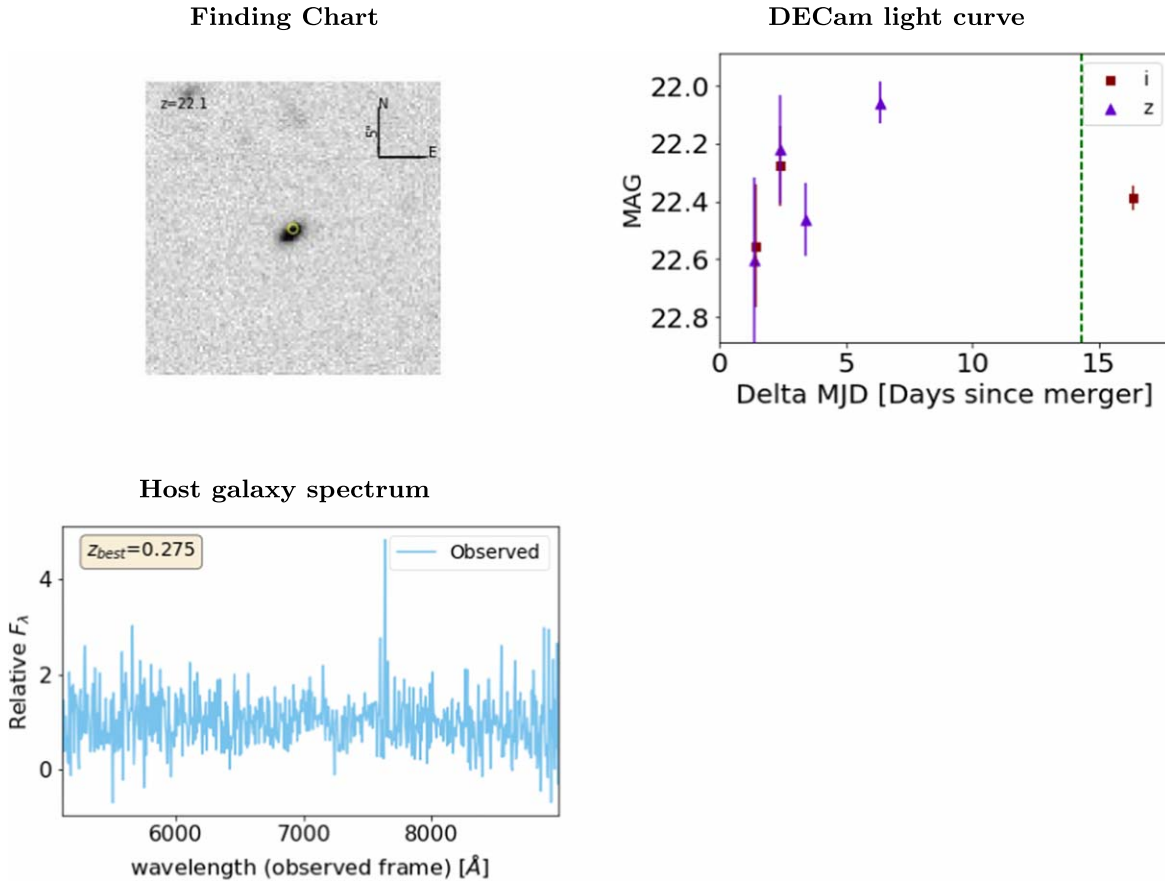


Figure 16. Top left and top right: same as Figure 8 but for the AT2019omx KN candidate. Bottom: spectrum of the host galaxy.

First, we found that our SOAR spectroscopic follow-up effort benefited from being a loose confederation of semi-independent teams that could operate the telescope remotely: a team based at Fermilab, one at University of California-Santa Cruz, one in Chile, and one in Brazil. Each of these teams signed up to be “on-call” for two-week blocks throughout LVC O3. The team “on-call” when an LVC O3 alert went out would have the responsibility for preparing and carrying out any SOAR spectroscopic follow-up during their watch. That said, the “on-call” team could request help from the other teams, and the other teams were welcome to follow along during the night of a follow-up observation. In the case of GW190814, the Fermilab team was the on-call team for most of the time of the spectroscopic follow-up, but other teams also provided help during Fermilab’s time block (in particular, the Chilean team took over for a couple of nights when the Fermilab team was unable to observe). This relatively loose structure of our spectroscopic follow-up effort seemed to work well, especially over the full course of LVC O3.

Second, especially as SOAR is primarily run as a remote observing facility, it is vital to have good communications with the SOAR scientific and technical staff. We were able to easily communicate with the SOAR staff and on several occasions they provided invaluable help to us in obtaining spectra of dimmer objects that required a longer process for target acquisition. Further, long after the optical signature of any expected KN should have faded, the SOAR staff

obtained the spectra of the host galaxies of two remaining candidates (AT2019nte and AT2019omw) during engineering time, in order to check if these candidates had redshifts that fell within the distance estimates measured by LVC for the GW event.

Third, it became clear early on that it is very difficult to obtain sufficiently high-S/N spectra with SOAR for candidate KNe fainter than about $i \approx 21$ in the allotted time for a SOAR ToO interrupt. For spectroscopic follow-up in LVC O4, candidates fainter than $i \approx 21$ should either be pursued by 6–10 m class telescopes, or have their host galaxies targeted as a means to qualify them or to rule them out.

Finally, we stress the importance of being able to reduce and analyze the data at the telescope for quick classification of the candidate as a KN or not. If there are obvious features in the spectrum indicating that a given candidate is not a KN (e.g., sharp emission or absorption lines or features typical of an SN spectrum), one can quickly move on to the next target in the candidate list; if, however, the spectrum indicates that the candidate is indeed the KN, the rest of the astronomical community can be quickly alerted. At the telescope during the observations for this paper, we typically made use of our SOAR Quick Reduce Pipeline or IRAF routines to process and calibrate the spectra on the fly, and classified the spectra by eye or by running them through the AstroDash and/or the SNID SN typing software that same night. A later, more refined reduction and analysis were performed offline, as described in

AT2019nte

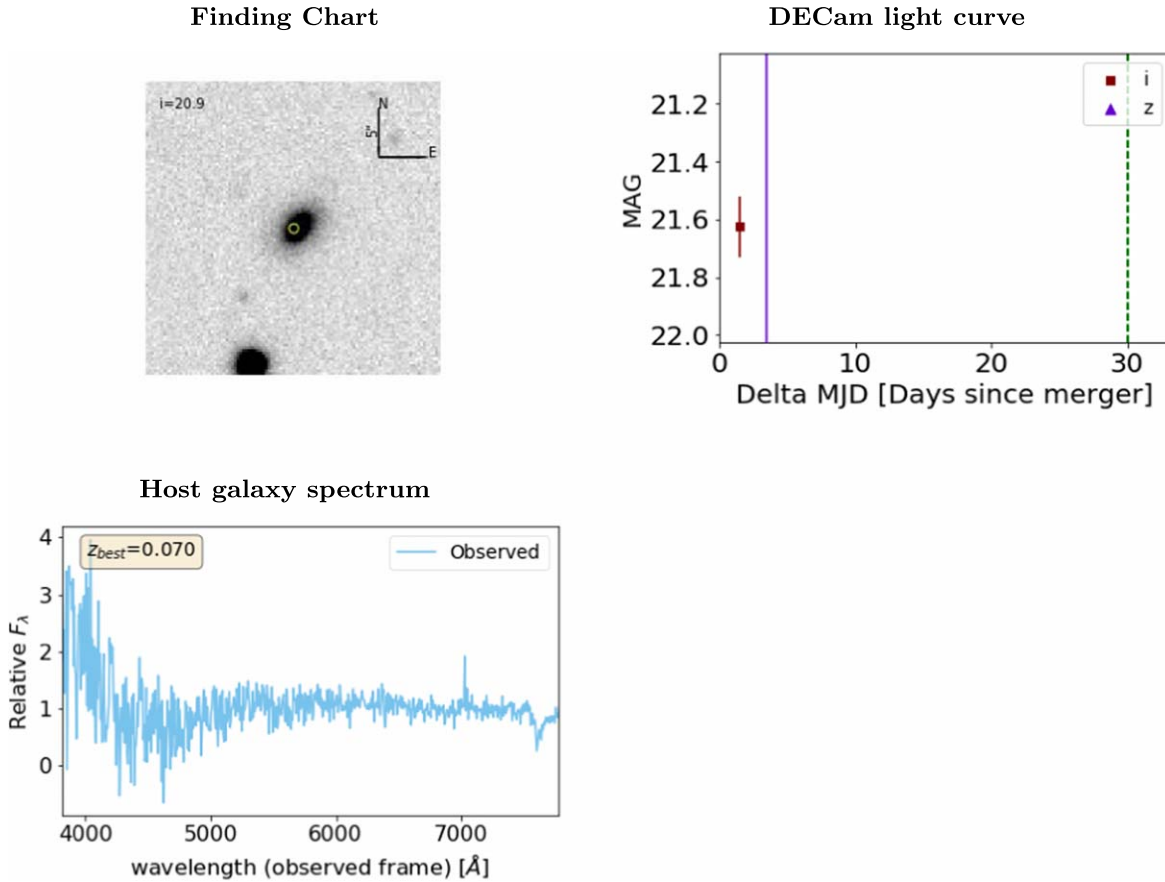


Figure 17. Same as Figure 16 but for the AT2019nte KN candidate. (The vertical purple line in the light-curve plot is just a very large error bar for the z -band observation.)

Sections 6.1 and 6.2. We note that, however, whereas some of the classifications changed between the real-time and off-line analysis, none of the resulting spectra—with the possible exception of the very low-S/N AT2019nte spectrum—was ever seriously considered to be that of a KN: i.e., the quick reductions are sufficient for the purpose. One weakness during our O3 observations of GW190814 candidates was the lack of an analog of our Quick Reduce pipeline to fit a candidate’s spectrum to a grid of KN model spectra on the fly at the telescope. Since then, we have developed an initial version of our publicly available DESGW KN spectrum fitter (DLT_DESGW_KNfit; see Section 8), which can be run at the telescope with the output of our SOAR Quick Reduce pipeline and should be useful for spectroscopic follow-up in LVC O4.

7. Conclusions

In the era of multi-messenger astronomy, we have demonstrated that we can perform a deep, one-of-its-kind spectroscopic follow-up campaign for possible NSBH events. We have reported on the SOAR/Goodman spectroscopy of 11 KN candidates associated with the LIGO/Virgo event GW190814. For eight of these we have reported the redshift and spectroscopic typing of the transient itself, and for the other three we have reported the redshift of the host galaxy. We concluded

that none of these candidates was the optical counterpart associated with the compact object binary merger. This SOAR/Goodman spectroscopy was done through SOAR ToO observations on a series of nights following the LVC discovery of gravitational waves from GW190814. These targeted observations were performed after KN candidate identification and culling by the DESGW collaboration following observations using DECcam on the Blanco telescope, and they have allowed us to place interesting constraints on the properties of the binary (Morgan et al. 2020) and to use this event as a dark standard siren (that is, as a constraint on H_0 using GWs) (Palmese et al. 2020).

We have also described the DESGW spectroscopic pipeline, part of the DESGW KN search process and candidate assessment, and our process and timeline for creating a spectroscopic follow-up candidate list. In addition, we have presented our QuickReduce software (for quick-look spectroscopic reduction) and the UCSC Reduction Pipeline software (for offline spectroscopic reduction). Furthermore, we have shown our use of AstroDash, SNID, and a least-squares KN model-fitting software for the process of candidate spectrum classification. Finally, we have demonstrated the effectiveness of our program and these tools within DESGW and are prepared for more extensive searches for KNe in LVC O4.

AT2019omw

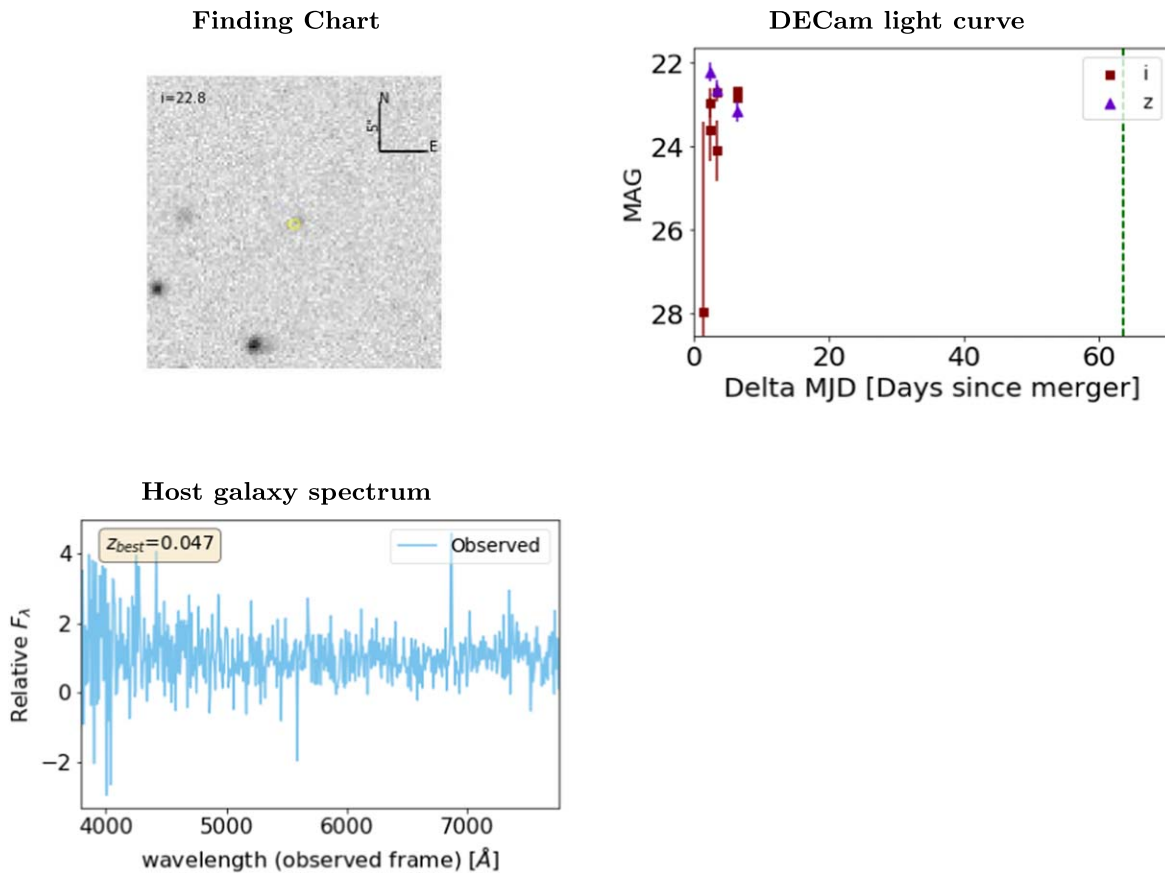


Figure 18. Same as Figure 16 but for the AT2019omw KN candidate.

8. Software

We present here links to the software packages mentioned in the text.

1. Quick Reduce Pipeline, used for reduction and analysis of spectra immediately after observing. https://github.com/DouglasLeeTucker/SOAR_Goodman_QuickReduce/blob/master/notebooks/SOAR_Goodman_QR_Notebook.ipynb
2. UCSC spectral pipeline, used for data reduction and analysis: https://github.com/msiebert1/UCSC_spectral_pipeline
3. AstroDash supernova typing software: <https://github.com/daniel-muthukrishna/astrodash>
4. Image Reduction and Analysis Facility (IRAF). IRAF had been distributed by the National Optical Astronomy Observatory, which was operated by the Association of Universities for Research in Astronomy (AURA) under a cooperative agreement with the National Science Foundation. The software is currently maintained and distributed by the IRAF Community: <https://iraf-community.github.io/>
5. SNID supernova typing software: <https://people.lam.fr/blondin.stephane/software/snid/>
6. Kasen KN models: https://github.com/dnkasen/Kasen_Kilonova_Models_2017
7. DESGW KN spectrum fitting software: https://github.com/cdebom/DLT_DESGW_KNfit
8. SNANA SuperNova ANALysis software <https://snana.uchicago.edu/>
9. matplotlib (Hunter 2007),
10. numpy (Van Der Walt et al. 2011),
11. scipy (Jones et al. 2001),
12. astropy (Astropy Collaboration et al. 2013),
13. TOPCAT (Taylor 2005).

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



































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

















Based on observations obtained at the Southern Astrophysical Research (SOAR) telescope, which is a joint project of the Ministério da Ciência, Tecnologia e Inovação (MCTI) da República Federativa do Brasil, the U.S. National Optical Astronomy Observatory (NOAO), the University of North Carolina at Chapel Hill (UNC), and Michigan State University (MSU).

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References

- Abbott, B. P., Abbott, R., Abbott, T. D., et al. 2017a, *PhRvL*, **119**, 161101
Abbott, B. P., Abbott, R., Abbott, T. D., et al. 2017b, *Natur*, **551**, 85
Abbott, B. P., Abbott, R., Abbott, T. D., et al. 2017c, *ApJL*, **848**, L12
Abbott, B. P., Abbott, R., Abbott, T. D., et al. 2019, *PhRvD*, **100**, 104036
Abbott, B. P., Abbott, R., Abbott, T. D., et al. 2020, *LRR*, **23**, 3
Abbott, B. P. 2016, *PhRvL*, **116**, 061102
Abbott, R., Abbott, T. D., Abraham, S., et al. 2020, *ApJL*, **896**, L44
Abbott, T. M. C., Abdalla, F. B., Allam, S., et al. 2018, *ApJS*, **239**, 18
Acerese, F., Agathos, M., Agatsuma, K., et al. 2015, *CQGra*, **32**, 024001
Ackley, K. 2020, *A&A*, **643**, A113
Alexander, K. D., Berger, E., Fong, W., et al. 2017, *ApJL*, **848**, L21
Andreoni, I., Goldstein, D. A., Kasliwal, M. M., et al. 2020, *ApJ*, **890**, 131
Arcavi, I., Hosseinzadeh, G., Howell, D. A., et al. 2017, *Natur*, **551**, 64
Ascenzi, S., Oganessian, G., Salafia, O. S., et al. 2020, *A&A*, **641**, 61
Astropy Collaboration, Robitaille, T. P., Tollerud, E. J., et al. 2013, *A&A*, **558**, A33
Bennett, C. L., Larson, D., Weiland, J. L., & Hinshaw, G. 2014, *ApJ*, **794**, 135
Bertin, E., & Arnouts, S. 1996, *A&AS*, **117**, 393
Blondin, S., & Tonry, J. L. 2007, *ApJ*, **666**, 1024
Castro-Tirado, A. J., Valeev, A. F., Hu, Y. D., et al. 2019, *GCN Circ.*, 25543, 1
Chen, H.-Y., Holz, D. E., Miller, J., et al. 2021, *CQGra*, **38**, 055010
Chornock, R., Berger, E., Kasen, D., et al. 2017, *ApJL*, **848**, L19
Clemens, J. C., Crain, J. A., & Anderson, R. 2004, *Proc. SPIE*, **5492**, 331
DES Collaboration, Garcia, A., Morgan, R., et al. 2020, arXiv:2007.00050
Diehl, H. T., Yanny, B., Tucker, D. L., Paz-Chinchón, F., & Neilsen, E. 2019, The Dark Energy Survey and Operations: Year 6 – The Finale FERMILAB-TM-2720-AE, SLAC National Accelerator Lab
Dobie, D., Stewart, A., Murphy, T., et al. 2019, *ApJL*, **887**, L13
Drout, M. R., Piro, A. L., Shappee, B. J., et al. 2017, *Sci*, **358**, 1570
Fermi-LAT Collaboration 2017, arXiv:1710.05450
Flaugher, B., Diehl, H. T., Honscheid, K., et al. 2015, *AJ*, **150**, 150
Fong, W., Berger, E., Blanchard, P., et al. 2017, *ApJL*, **848**, L23
Foucart, F., Hinderer, T., & Nissanka, S. 2018, *PhRvD*, **98**, 081501
Ghirlanda, G., Salafia, O. S., Paragi, Z., et al. 2019, *Sci*, **363**, 968
Goldstein, D. A., & D'Andrea, C. B. 2015, *AJ*, **150**, 82
Gomez, S., Hosseinzadeh, G., & Cowperthwaite, P. S. 2019, *ApJ*, **884**, L55
Hallinan, G., Corsi, A., Mooley, K. P., et al. 2017, *Sci*, **358**, 1579
Hamuy, M., Suntzeff, N. B., Heathcote, S. R., et al. 1994, *PASP*, **106**, 566
Hamuy, M., Walker, A. R., Suntzeff, N. B., et al. 1992, *PASP*, **104**, 533
Herner, K., Annis, J., Brout, D., et al. 2020, *A&C*, **33**, 100425
Hu, Y. D., Castro-Tirado, A. J., Valeev, A. F., et al. 2019, *GCN Circ.*, 25588, 1
Hunter, J. D. 2007, *CSE*, **9**, 90
Jones, E., Oliphant, T., Peterson, P., et al. 2001, SciPy: Open source scientific tools for Python, <http://www.scipy.org/>
Kapadia, S. J., Caudill, S., Creighton, J. D. E., et al. 2020, *CQGra*, **37**, 045007
Kasen, D., Metzger, B., Barnes, J., Quataert, E., & Ramirez-Ruiz, E. 2017, *Natur*, **551**, 80
Kawaguchi, K., Kytoku, K., Shibata, M., & Tanaka, M. 2016, *ApJ*, **825**, 52
Kessler, R., Marriner, J., Childress, M., et al. 2015, *AJ*, **150**, 172
Kessler, R., Narayan, G., Avelino, A., et al. 2019, *PASP*, **131**, 094501
Kilpatrick, C. D., Coulter, D. A., Arcavi, I., et al. 2021, *ApJ*, **923**, 258
LIGO Scientific Collaboration 2018, LIGO Algorithm Library—LALSuite, free software (GPL), doi:10.7935/GT1W-FZ16
Lopez-Cruz, O., Castro-Tirado, A. J., Macri, L., et al. 2019a, *GCN Circ.*, 25419, 1
Lopez-Cruz, O., Castro-Tirado, A. J., Macri, L., et al. 2019b, *GCN Circ.*, 25571, 1
LVC 2019a, *GCN Circ.*, 25324, 1
LVC 2019b, *GCN Circ.*, 25333, 1
LVC 2020a, Public Alerts User Guide, <https://emfollow.docs.ligo.org/userguide/content.html#inference>
LVC 2020b, Public Alerts User Guide, <https://emfollow.docs.ligo.org/userguide/analysis/inference.html#diskmass>
Lyman, J., Lamb, G., Levan, A., et al. 2018, *NatAs*, **2**, 751
Margutti, R., Berger, E., Fong, W., et al. 2017, *ApJL*, **848**, L20
Mohr, J. J., Armstrong, R., Bertin, E., et al. 2012, *Proc. SPIE*, **8451**, 84510D
Mooley, K. P., Deller, A. T., Gottlieb, O., et al. 2018, *Natur*, **561**, 355
Morgan, R., Soares-Santos, M., Annis, J., et al. 2020, *ApJ*, **901**, 83
Muthukrishna, D., Parkinson, D., & Tucker, B. E. 2019, *ApJ*, **885**, 85
Neilsen, E. J., Annis, J. T., Diehl, H. T., et al. 2019, arXiv:1912.06254
Nicholl, M., Berger, E., Kasen, D., et al. 2017, *ApJL*, **848**, L18
Palmese, A., deVicente, J., Pereira, M. E. S., et al. 2020, *ApJL*, **900**, L33
Palmese, A., & Kim, A. G. 2021, *PhRvD*, **103**, 103507
Pian, E., D'Avanzo, P., Benetti, S., et al. 2017, *Natur*, **551**, 67
Richardson, D., Jenkins, Robert, L., Wright, I., & Maddox, L. J. 2014, *AJ*, **147**, 118
Sako, M., Bassett, B., Connolly, B., et al. 2011, *ApJ*, **738**, 162
Savchenko, V., Ferrigno, C., Kuulkers, E., et al. 2017, *ApJL*, **848**, L15
Scolnic, D., Kessler, R., Brout, D., et al. 2018, *ApJL*, **852**, L3
Shibata, M., Fujibayashi, S., Hotokezaka, K., et al. 2017, *PhRvD*, **96**, 123012
Singer, L. P., & Price, L. R. 2016, *PhRvD*, **93**, 024013
Smartt, S. J., Chen, T. W., Jerkstrand, A., et al. 2017, *Natur*, **551**, 75
Soares-Santos, M., Holz, D. E., Annis, J., et al. 2017, *ApJL*, **848**, L16
Soares-Santos, M., Palmese, A., Hartley, W., et al. 2019, *ApJL*, **876**, L7
Tanaka, M., Kato, D., Gaigalas, G., et al. 2018, *ApJ*, **852**, 109
Taylor, M. B. 2005, in ASP Conf. Ser. 347, Astronomical Data Analysis Software and Systems XIV, ed. P. Shopbell, M. Britton, & R. Ebert (San Francisco, CA: ASP), 29
Thakur, A. L., Dichiera, S., Troja, E., et al. 2020, *MNRAS*, **499**, 3868
The LIGO Scientific Collaboration, Aasi, J., Abbott, B. P., et al. 2015, *CQGra*, **32**, 074001

- The LIGO Scientific Collaboration, The Virgo Collaboration, The KAGRA Collaboration, et al. 2021, arXiv:2111.03606
- Tonry, J., & Davis, M. 1979, *AJ*, 84, 1511
- Troja, E., Piro, L., van Eerten, H., et al. 2017, *Natur*, 551, 71
- Van Der Walt, S., Colbert, S. C., & Varoquaux, G. 2011, *CSE*, 13, 22
- Vieira, N., Ruan, J. J., Haggard, D., et al. 2020, *ApJ*, 895, 96
- Watson, A. M., Butler, N. R., Lee, W. H., et al. 2020, *MNRAS*, 492, 5916
- Watson, D., Hansen, C. J., Selsing, J., et al. 2019, *Natur*, 574, 497
- Wright, E. L. 2006, *PASP*, 118, 1711
- Xiao, D., Liu, L.-D., Dai, Z.-G., & Wu, X.-F. 2017, *ApJL*, 850, L41
- Ye, C. S., Fong, W.-f., Kremer, K., et al. 2020, *ApJL*, 888, L10
- Ziosi, B. M., Mapelli, M., Branchesi, M., & Tormen, G. 2014, *MNRAS*, 441, 3703