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# The absence of an effect of nickel on iron isotope fractionation during core formation

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## Abstract

The Fe isotopic compositions of mantles of differentiated inner solar system bodies are similar to, or heavier than those of chondritic meteorites. Core–mantle differentiation is a potential contributor to planetary isotopic fractionation. However, previous metal–silicate experiments provide only equivocal evidence for such fractionation, and have been used to argue that the Ni content of core-forming metal influences the extent of Fe isotopic fractionation. Here, we complement existing data with twenty-two novel metal–silicate

25 equilibrium experiments with varying Ni content to better quantify the effect of Ni on the  
26 vector and magnitude of Fe isotopic fractionation during core formation. We find no  
27 statistically resolvable effect of the Ni content in the metallic phase on the metal–silicate Fe  
28 isotopic fractionation factor over a wide range of Ni concentrations (0 to 70 wt.% in the  
29 metal). In particular, the Fe isotopic composition of alloys from two experiments performed  
30 with 70 wt.% of Ni ( $\delta^{56}\text{Fe}_{\text{metal}}=0.27\pm 0.04\text{‰}$  and  $0.32\pm 0.03\text{‰}$ , 2 standard deviations  $\sigma$ ) are  
31 identical to the bulk experimental starting material ( $\delta^{56}\text{Fe}_{\text{bulk}}=0.27\pm 0.10\text{‰}$ ,  $2\sigma$ ). Our data  
32 across all experiments yield an average isotopic fractionation factor  $\Delta^{56}\text{Fe}_{\text{met-sil}}=0.05\pm 0.22\text{‰}$   
33 ( $2\sigma$ ) at 1873 K and 1–2 GPa, suggesting that little to no isotopic fractionation of Fe is  
34 expected to occur during core formation at low pressures. As such, our data does not support  
35 core formation as the main mechanism causing the observed variability in Fe isotope ratios  
36 between the silicate Earth, Moon, Vesta and other differentiated asteroids. A combination of  
37 multiple accretion-related processes—including condensation from the solar nebula, volatile-  
38 depleting events such as giant impacts, and the disproportionation of ferrous iron to ferric iron  
39 and iron metal in larger bodies—as well as deep mantle and recycling processes could explain  
40 the heavier-than-chondritic signatures in the silicate Earth and Moon. Furthermore, our results  
41 support the ideality of mixing among Fe–Ni alloys, as previously demonstrated for physical  
42 properties but less conclusively evidenced for chemical properties.

## 43 **Introduction**

44 Experimental determination of isotopic fractionation of siderophile elements between  
45 metal and silicate has been used to investigate core formation on planetary bodies (*e.g.* [Georg  
46 et al., 2007](#); [Bourdon et al., 2018](#)). Isotopic fractionation recorded in laboratory experiments  
47 can be compared to natural isotope signatures, providing a powerful tool for interpretation of  
48 isotopic variations among planetary materials and refining our knowledge of their history in

49 the early solar system. In principle, these experiments involve equilibration of molten silicate  
50 and molten metal at high pressure and high temperature until elemental and isotopic  
51 equilibrium between the two phases is reached. The silicate glass and metal grains from each  
52 quenched experiment are physically separated and their isotope compositions are determined.  
53 Such studies on various siderophile elements can be used to trace the conditions of planetary  
54 differentiation—its temperature (Hin et al., 2014), its pressure (Labidi et al., 2016), redox  
55 conditions (Dalou et al., 2019), subsequent rehomogenisation mechanisms (Kempl et al.,  
56 2013), and the composition of the core (Savage et al., 2015; Mahan et al., 2017; Xia et al.,  
57 2019). Applied to volatile elements, they open the possibility to trace the timing and origin of  
58 planetary volatile delivery (Kubik et al., 2021b). Such results allow the isotopic  
59 characterisation of the building blocks that eventually accreted to form planetary bodies (Hin  
60 et al., 2014), permitting assessment of whether core formation is the main planetary process  
61 controlling the abundances and isotopic compositions of certain elements in the accessible  
62 mantle of the body. Alternatively, other possible mechanisms such as magmatic  
63 differentiation (Bonnand et al., 2016) or volatilisation processes (Xia et al., 2019) may be  
64 required. Numerous key questions may thus be addressed through experimental evidence.  
65 However, experimentally-determined isotopic fractionation factors from different studies are  
66 often in disaccord with one another, namely, Zn (Bridgestock et al., 2014; Mahan et al., 2017;  
67 Xia et al., 2019), Si (Shahar et al., 2009,2011; Hin et al., 2014; Kempl et al., 2013), Ni (Lazar  
68 et al., 2012; Guignard et al., 2020), Cu (Savage et al., 2015; Xia et al., 2019). These  
69 discrepancies are partially attributable to differences in the experimental approach. In this  
70 respect, replication of such experiments combined with independent, analytical determination  
71 of the isotopic consequences of core formation (*e.g.* Georg et al., 2007; Moynier et al., 2011)  
72 can only be beneficial to the general understanding of planetary differentiation.

73 Planetary materials show significant variability in their Fe isotopic compositions.  
74 Ordinary, enstatite and carbonaceous chondrites have an average  $\delta^{56}\text{Fe}$  (per mil variation of  
75 the  $^{56}\text{Fe}/^{54}\text{Fe}$  ratio relative to IRMM-014) of  $0.00\pm 0.01\text{‰}$  ( $2\sigma$ ) (Craddock and Dauphas, 2011;  
76 Schoenberg and von Blanckenburg, 2006; Wang et al., 2013, 2014). The terrestrial mantle is  
77 estimated to have a  $\delta^{56}\text{Fe}=0.03\pm 0.02\text{‰}$  (Craddock et al., 2013; Sossi et al., 2016a; Weyer and  
78 Ionov, 2007) while mid-ocean ridge basalts (MORB) are isotopically heavier with average  
79  $\delta^{56}\text{Fe}=0.10\pm 0.01\text{‰}$  (Teng et al., 2013; Craddock et al., 2013; Craddock and Dauphas, 2011;  
80 Nebel et al., 2018; Weyer and Ionov, 2007). Even after accounting for the effect of olivine  
81 crystallisation and correcting their Fe isotopic composition back to that of a primary melt in  
82 equilibrium with mantle peridotite, MORBs define an average  $\delta^{56}\text{Fe}=0.07\pm 0.01\text{‰}$  (Sossi et  
83 al., 2016a; Nebel et al. 2018), heavier than chondritic meteorites. By contrast, primary melts  
84 of the martian and vestan mantles, determined respectively from the SNC (shergottite-  
85 nakhite-chassignite) and HED (howardite-eucrite-diogenite) groups of meteorites, have near-  
86 chondritic isotopic signatures (Poitrasson et al., 2004; Schoenberg and von Blackenburg,  
87 2006; Wang et al., 2012; Sossi et al., 2016b) whereas angrites are markedly heavier  
88 ( $\delta^{56}\text{Fe}=0.12\pm 0.04\text{‰}$ ,  $n=7$ , Wang et al., 2012). Lunar basalts yield a range of isotopic  
89 compositions from near-chondritic to  $0.3\text{‰}$  (Weyer et al., 2005; Poitrasson et al., 2004; Liu et  
90 al., 2010; Wang et al., 2015; Sossi and Moynier, 2017; Poitrasson et al., 2019), but ferroan  
91 anorthosites and Mg suite rocks from the lunar highlands define an estimated bulk lunar Fe  
92 isotope composition similar to that of the Earth's mantle ( $\delta^{56}\text{Fe}=0.05\pm 0.03\text{‰}$ , Sossi and  
93 Moynier, 2017; Poitrasson et al., 2019).

94 Several processes have been proposed to account for this Fe isotopic variability  
95 including (1) volatile loss by vaporisation during the Moon-forming giant impact (Poitrasson  
96 et al., 2004; Poitrasson 2007; Liu et al., 2010), (2) accretion from the solar nebula associated  
97 with various degrees of nebular or post-nebular Fe loss in planetary building blocks (Sossi et

98 [al., 2016b](#)), (3) disproportionation of ferrous iron ( $\text{Fe}^{2+}$ ) into  $\text{Fe}^{3+}$  and metallic iron by  
99 bridgmanite crystallisation during accretion in bigger planetary bodies such as the Earth  
100 ([Williams et al., 2012](#); [Frost and McCammon, 2008](#)), (4) core–mantle equilibration related  
101 isotopic fractionation at either high ([Polyakov, 2009](#)) or low pressures ([Elardo and Shahar,](#)  
102 [2017](#)) and (5) deep mantle recycling processes ([Smith et al., 2021](#); [Soderman et al., 2021](#)).  
103 [Leshner et al. \(2020\)](#) argued that the upper part of the Earth’s core could be enriched in heavier  
104 Fe isotopes by thermo-diffusion over time. If this isotopically heavier Fe is, as suggested in  
105 the study, transported to the upper mantle by ascending plumes, it could explain the heavy Fe  
106 isotopic signature of the BSE (bulk silicate Earth). However, core–mantle equilibration  
107 remains the largest chemical mass transfer process that occurred on Earth, suggesting that it  
108 may have a significant contribution in defining the observed isotopic signatures of its present-  
109 day core and mantle. It is therefore paramount to understand the isotopic behaviour of Fe  
110 during core–mantle equilibration.

111         However, the degree and direction of Fe isotope fractionation between silicate and  
112 core-forming alloy is influenced by a variety of parameters. Early experimental work  
113 addressing this question did not find any resolvable Fe isotopic fractionation between S-free  
114 metallic and silicate phases over a range of pressures, from 1 to 60 GPa ([Poitrasson et al.,](#)  
115 [2009](#); [Hin et al., 2012](#); [Shahar et al., 2015](#); [Liu et al., 2017](#)), leading to the conclusion that, at  
116 the high temperatures required for core formation on terrestrial bodies ( $\geq 2000$  K, *e.g.*,  
117 [Righter and Drake, 1996](#)), equilibrium isotope fractionation between metal and  
118 silicate—which is proportional to  $1/T^2$ —should be negligible compared to the observed  
119 variations. [Liu et al. \(2017\)](#) reported evidence up to 60 GPa of the insensitivity of Fe force  
120 constant to pressure in both metal and silicate materials, which supports the relevance of low-  
121 pressure studies using the piston cylinder and isotopic measurements on the synthesised  
122 phases to address Fe isotopic fractionation during core formation. Although some of the

123 metallic phases investigated contained a small fraction of Ni (~5–9 wt.%) as is relevant to  
124 planetary cores (5 wt.% in the Earth’s core; Palme and O’Neill, 2014), the effect of Ni on iron  
125 isotope fractionation was not explicitly investigated prior to the study of [Elardo and Shahar](#)  
126 [\(2017\)](#). They reported piston cylinder experiments conducted at 2123 K and 1 GPa in which  
127 the Ni content of the metal varied between 1 and 26 at.%, and observed a systematic increase  
128 of the Fe isotopic fractionation factor with Ni content, which defines the following relation:  
129  $\Delta^{57}\text{Fe}_{\text{met-sil}}=0.011(\text{Ni, at.}\%)+0.04\text{‰}$ . Thus, adding Ni to the metal phase enhances the Fe  
130 isotopic fractionation between metal and silicate, leading to an enrichment of the silicate in  
131 the lighter Fe isotopes relative to the metal. This result differs from that obtained with nuclear  
132 resonant inelastic X-ray scattering spectroscopy (NRIXS) results by [Liu et al. \(2017\)](#)  
133 suggesting that the force constant of iron bonds in Fe alloys is unaffected by the addition of  
134 Ni (up to 8 wt.%). Because the sense of this isotopic fractionation is opposite to that observed  
135 in the mantles of differentiated bodies, [Elardo and Shahar, \(2017\)](#) invoked significant isotopic  
136 fractionation during partial melting to restore the light subchondritic mantle isotopic  
137 composition inferred from their experiments to the superchondritic values observed in basaltic  
138 achondrites (*e.g.* [Sossi et al. 2016a](#)). Importantly, such an isotopic effect implies that Fe–Ni  
139 alloys may not mix ideally, where ideal mixing is expected to induce no change in physical  
140 and chemical properties with varying Ni contents in the alloy. Therefore, a dependence of Fe  
141 isotopic fractionation on the Ni content is at odds with the constancy of density,  
142 compressibility, and wave velocities of Fe–Ni alloys with increasing Ni content (*e.g.* [Lin et](#)  
143 [al., 2003](#); [Kantor et al., 2007](#); [Martorell et al., 2013](#); [Badro et al., 2014](#)). In other words, the  
144 properties of Fe–Ni alloys provide no evidence that the bonding environment of Fe in Fe–Ni  
145 alloys is modified by varying the proportion of Ni that would otherwise be necessary to  
146 generate an isotopic fractionation. Moreover, Fe–Ni ideality has been assumed in numerous  
147 experimental studies on core formation processes in which the effect of Ni was neglected. [Xia](#)

148 [et al. \(2019\)](#) reported Zn and Cu metal–silicate isotope fractionation experiments with  
149 significant sensitivity to the Ni content of the metal. However, the positive Ni-induced  
150 isotopic fractionations were measured for Fe ([Elardo and Shahar, 2017](#)), Zn and Cu ([Xia et](#)  
151 [al., 2019](#)), were based on fewer than 10 experiments for each element. More experimental  
152 results, involving a broader range of Ni contents, are needed to establish whether there is a  
153 resolvable effect of the Ni content of an Fe–Ni alloy on various chemical and isotopic  
154 parameters. Therefore, assessing the Fe isotopic fractionation between metal and silicate has  
155 profound implications for our understanding of core formation processes and for the physical  
156 and chemical properties of core-forming alloys.

157 To address this question, twenty-two novel metal–silicate equilibrium experiments  
158 were performed in order to characterise the effect of Ni on the Fe isotope fractionation. In  
159 particular, a wider range of Ni content in the metal (from 0 to 70 *wt.%*) was covered  
160 compared to previous studies (0 to 25 *wt.%*, [Elardo and Shahar, 2017](#)).

## 161 **Method**

### 162 **Experimental method**

163 Mixtures composed of 70 *wt.%* of a crushed mid-ocean ridge basalt (MORB) ([Kubik](#)  
164 [et al., 2021a](#)) with varying proportions of pure Fe and Ni were prepared by weighing each  
165 component and grinding them in an agate mortar until complete homogenisation was  
166 achieved. All experiments were performed in a 150-ton end-loaded piston cylinder apparatus  
167 at the Institut de Physique du Globe de Paris (IPGP, Université de Paris), in ½” talc pyrex  
168 assemblies using graphite furnaces. Pressure calibration of the apparatus was performed using  
169 the alumina concentration in orthopyroxene in equilibrium with pyrope in a MgO–Al<sub>2</sub>O<sub>3</sub>–  
170 SiO<sub>2</sub> system ([Perkins and Newton, 1981](#)). The friction coefficients reported by [McDade et al.](#)  
171 ([2002](#)) were compared and found to be identical to the coefficients measured during the



172 calibration of the apparatus, allowing similar corrections to be applied for friction loss for  
173 assembly materials. During the experiments, the temperature was measured using a type D  
174 ( $W_{97}Re_3/W_{75}Re_{25}$ ) thermocouple contained in a 4-hole alumina sleeve just above the capsule  
175 lid. Uncertainties on pressure and temperature are estimated to be around 0.1 GPa and 50 K  
176 (Siebert et al. 2011). All experiments were performed at 1 to 2 GPa and at 1873 K in order to  
177 induce complete fusion and promote equilibrium between metal and silicate in the system.  
178 Four capsule types—boron nitride, crushable MgO, single-crystal MgO and graphite—were  
179 tested in order to assess their effect on the Fe isotopic fractionation factor.

### 180 **Sample collection and purification**

181 The metal and silicate phases were separated mechanically after each experiment.  
182 Clean pieces of each phase were selected carefully under magnifying glass. Metal pieces were  
183 directly dissolved in 6 N HCl. Silicate samples were sorted in order to select clean pieces  
184 devoid of metal beads or capsule material. These pieces were roughly crushed and a magnet  
185 was used to remove any stray pieces of metal (Kubik et al., 2021b). The silicate was finely  
186 crushed in an agate mortar and digested in a 1:3 solution of HF:HNO<sub>3</sub> in closed Teflon bombs  
187 on a hotplate at 100°C. Silicate samples were then dried and further digested in aqua regia on  
188 a hot plate at 150°C to remove any residual fluoride complex. They were dried again and  
189 dissolved in 6 N HCl for purification. The purification of Fe in the samples was performed  
190 following the method described in Sossi et al. (2015), by anion exchange chromatography  
191 using AG1-X8 (200-400 mesh) and 0.4 × 7 cm custom-made Teflon columns. Columns were  
192 filled with 1 mL AG1-X8 resin and cleaned with 5 mL 3 N HNO<sub>3</sub>, 5 mL H<sub>2</sub>O and conditioned  
193 with 5 mL 6 N HCl. The samples were subsequently loaded in 0.5 mL of 6 N HCl. The matrix  
194 was eluted with 11 mL of 6 N HCl. The Fe cuts were eluted with 3 mL of 0.5 N HCl, and  
195 dried to be redissolved in 0.5 N HNO<sub>3</sub>. They were diluted to 4 ppm Fe solutions and spiked

196 with 8 ppm of Ni, following the method described in [Sossi et al. \(2015\)](#). This protocol yielded  
197 procedural blanks inferior to 10 ng (n=2) ([Kubik et al., 2021c](#)).

## 198 **MC-ICPMS measurements**

199 Iron isotopic measurements were performed on a Neptune Plus MC-ICPMS at IPGP  
200 (Université de Paris), used in medium resolution mode. The cups were configured as  
201 described in [Sossi et al. \(2015\)](#) in order to collect the intensity on the masses 53, 54, 56, 57,  
202 60 and 61. The mass 53 was used to correct for any  $^{54}\text{Cr}$  interference on  $^{54}\text{Fe}$ , even though  
203 chromatographic separation was shown to be very efficient, meaning the correction was  
204 insignificant ( $^{53}\text{Cr}$  signals were  $\sim 5 \times 10^{-5}$  V, indistinguishable from 2%  $\text{HNO}_3$  blanks). The  
205  $^{61}\text{Ni}/^{60}\text{Ni}$  ratios measured in all samples were used to correct for mass bias effects.  
206 Measurements of IRMM-014 performed in between samples were used to calculate the delta  
207 values reported in this study. International reference materials were measured along with the  
208 experimental samples with Fe isotopic compositions reproducing the previously published  
209 values (*e.g.* for BHVO-2:  $\delta^{56}\text{Fe} = 0.10 \pm 0.04\%$ ,  $2\sigma$ ).

## 210 **Evidence of metal–silicate equilibrium**

211 Two methods are typically used to assess metal–silicate isotopic equilibrium in high-  
212 temperature experiments: time series and the “three-isotope method”. Elemental equilibrium  
213 can be reached very quickly at superliquidus conditions in metal–silicate experiments (5  
214 seconds at 2300 K, [Thibault and Walter, 1995](#)). As diffusion rates for elemental exchange  
215 cannot necessarily be extrapolated to those for isotopic exchange rate, several experimental  
216 studies report that isotopic exchange can be both faster than chemical diffusion ([Leshner et al.,](#)  
217 [1990](#); [Van der Laan, 1994](#)) and slower (*e.g.* [Guignard et al., 2020](#); [Shahar et al., 2008](#)). In this  
218 context, numerous studies report relatively short equilibration time experiments (*e.g.* [Hin et](#)  
219 [al., 2013](#); [Kubik et al., 2021b](#); [Poitrasson et al., 2009](#); [Shahar et al., 2011](#)). The attainment of  
220 isotopic equilibrium after a few minutes can be explained by the superliquidus temperatures

221 reached, at which the entire sample is molten. The molten state promotes equilibration  
222 because (1) diffusion rates are significantly higher in liquids state than in a solid phase, (2)  
223 convection further stirs the melt inside the capsule, owing to the small thermal gradient  
224 endemic to high pressure experiments, (3) homogenisation is additionally favoured by using a  
225 finely crushed powder as a starting material. Time series experiments are an empirical way to  
226 establish the time needed to achieve equilibrium. In this study specifically, (a) thorough SEM  
227 and EDX characterisation of each phase provided proof that the entirety of the sample  
228 material was fully molten and resulted in (b) chemically homogeneous phases (see Figures  
229 S1, S2, S3 and supplementary information) as well as (c) textural equilibrium evidenced by  
230 one large spherical metal blob (Néri et al., 2021). (d) The Fe isotopic fractionation factors  
231 follow a systematic trend that can be statistically determined to good precision due to an  
232 extensive data set comprising 22 experiments. (e) Repeated time series experiments show no  
233 statistically resolvable evolution of the Fe isotopic fractionation factor between runs of 3- and  
234 30-minute duration (see supplementary information and Figure S4). Experiments with  
235 durations of 30 minutes or more exhibiting greater variability, with both positive and negative  
236 isotopic fractionation factors observed. This enhanced variability in longer experiments is  
237 attributed to progressive reaction of the silicate melt with the capsule material over time.  
238 Moreover, the isotopic results reported in this study are in agreement with several other  
239 studies performed in a range of laboratory facilities with different methodologies and  
240 apparatuses (Poitrasson et al., 2009; Hin et al., 2012; Liu et al., 2017; Elardo and Shaha,  
241 2017, Figure 1a). The results from our study are also in agreement with numerous pieces of  
242 evidence of Fe–Ni alloy ideality (Lin et al., 2003; Kantor et al., 2007; Martorell et al., 2013;  
243 Badro et al., 2014; Huang and Badro, 2018). With respect to demonstration of isotopic  
244 equilibrium, the three-isotope method was not included in the protocol of this study. Although  
245 Cao and Bao (2017) and Bourdon et al. (2018) reported artifacts that could induce non-linear

246 trajectories, leading the authors to delineate limitations of the method, other studies have  
247 argued for its robustness (Shahar et al., 2008; Shahar and Young, 2020), and in particular its  
248 usefulness in verifying that the system remained closed with respect to mass exchange  
249 throughout an experiment. In this study, the use of single-crystal MgO capsules was  
250 employed to ensure minimal chemical and isotopic interaction with the capsule (see  
251 supplementary information). Moreover, the “lever rule method”, consisting in the  
252 recalculation of the bulk isotope composition of the experiments based on measured  $\delta^{56}\text{Fe}$  of  
253 each phase by mass balance yields a good reproducibility (on average  $\delta^{56}\text{Fe}_{\text{final}}$   
254  $\text{bulk}=0.40\pm 0.08\%$ ,  $1\sigma$ ). The three-isotope method is predicated on experiments that preserve a  
255 record of the approach to equilibrium between two phases, and thus represent a form of time-  
256 series in which equilibration times are necessarily slow. Because equilibrium is achieved  
257 within a few minutes in our experiments, the three-isotope method is unfeasible from an  
258 experimental standpoint. Namely, the timescale of equilibration of the piston cylinder  
259 assembly to ensure stable conditions (*e.g.*, Matjuschkin et al., 2015) is of the same order of  
260 that required for the three-isotope method to produce dispersion in isotope composition. Each  
261 bulk starting material used in this study was also measured for Fe isotopes and are found to  
262 reproduce the calculated bulk value within error for 12 of them which further supports closed  
263 system and achievement of equilibrium.

## 264 **Iron isotopic behaviour in metal–silicate equilibrium experiments**

265 Twenty-two metal–silicate equilibrium experiments were performed at 1 and 2 GPa  
266 and 1873 K. Run durations were varied between 3 and 60 minutes in order to assess the  
267 timescales required to approach isotopic equilibrium. Petrographic observations indicate that  
268 the starting material was fully molten in all experiments (see Figure S1). The metal coalesced  
269 into a discrete sphere in most experiments, and subsequent to quenching, was separated

270 mechanically from the silicate phase. All experimental parameters and isotopic measurements  
271 are available in Table 1.

272 Iron isotope fractionation results for all experiments from this study and literature are  
273 presented in Figure 1a. It should be noted that experiments from other studies were carried out  
274 at different temperatures and were therefore recalculated at 1873 K for comparison with the  
275 dataset reported herein by using the relationship  $\Delta^{56}Fe_{metal-silicate} = a \cdot \frac{1}{T^2}$  (e.g. Urey,  
276 1947). Our results show no correlation between the Ni content of the metal and the Fe  
277 isotopic fractionation factor between metal and silicate (Figure 1b). There is also no  
278 resolvable difference in isotopic fractionation factor between the experiments using the  
279 different types of capsules (boron nitride, crushable MgO, single-crystal MgO and graphite),  
280 although from petrographic observations, use of the single-crystal MgO capsules minimises  
281 the extent of reaction between capsule and sample by containing the interactions within the  
282 first 30  $\mu\text{m}$  of the capsule. Isotope measurements on multiple parts of experiment E375 are  
283 presented in Figure 2. This short-duration experiment exhibits a glassy silicate in the center  
284 and partially crystallised silicate towards the capsule boundary where skeletal olivine crystals  
285 grew as a result of MgO enrichment and the relatively low quenching rates achievable in the  
286 piston cylinder (e.g. Green et al., 1975). Measurements on these separated parts show no  
287 difference in Fe isotopic composition between glass and crystalline silicates  
288 ( $\delta^{56}\text{Fe}=0.30\pm 0.08\text{‰}$  and  $0.29\pm 0.11\text{‰}$  respectively,  $2\sigma$ ) which means that MgO diffusion  
289 towards the center of the sample does not affect the Fe isotope composition of the silicate.  
290 Therefore, in this particular case, MgO capsules are well-suited for isotope fractionation  
291 experiments with respect to MgO diffusion. In this experiment, the isotopic composition of  
292 the metal is identical to that of the silicate. The first 30  $\mu\text{m}$  of the capsule in contact with the  
293 melt containing less than 4 wt.% of Fe, however, exhibits a clear enrichment in light Fe  
294 isotopes, even after a few minutes ( $\delta^{56}\text{Fe}=-0.72\pm 0.02\text{‰}$ ,  $2\sigma$ ). However, we estimate that this

295 Fe-bearing rim in the capsule contains less than 1.5% of the total Fe budget introduced in the  
296 capsule. In this context, light Fe contained in the rim, which represents less than 3% of the  
297 inner capsule volume, creates a shift of the bulk sample Fe isotopic composition of 5.5%  
298 relative at most, which is within the analytical uncertainty and therefore not resolvable by our  
299 measurements. At such superliquidus temperatures, chemical reactions between the capsule  
300 and the sample progress with increasing run duration. Therefore, in the long term, measured  
301 isotopic fractionations could be driven by interactions between the capsule and the sample,  
302 even if no correlation is observed between the isotopic fractionation and the time duration in  
303 the reported data set (see Figure S4). However, in this experiment, even with a limited  
304 reservoir of Fe on the boundaries of the sample that results in light isotopic composition  
305 created by Fe loss, the final bulk isotopic composition of the experiment calculated by mass  
306 balance of the silicate and metal compositions ( $0.29 \pm 0.06\text{‰}$ ,  $2\sigma$ ) is identical to the isotopic  
307 composition of the starting material ( $0.30 \pm 0.08\text{‰}$ ,  $2\sigma$ ). This suggests that the limited Fe loss  
308 resulted in insignificant and unresolvable isotopic fractionation of the experimental sample.  
309 The detection of Fe only in the first 30  $\mu\text{m}$  of the capsule in all single-crystal MgO  
310 encapsulated experiments suggest that such capsules permit near-closed system conditions in  
311 the experimental sample.

312 Iron isotopic fractionation factors from this study yield an average of  $\Delta^{56}\text{Fe}_{\text{met-sil}} = 0.05 \pm 0.22\text{‰}$  ( $2\sigma$ ). This means that these results are not significantly different from  $0.0\text{‰}$ ,  
313 and no dependence on the Ni content of the metal is detected over a wide range of Ni  
314 concentrations. The scatter of the data set, although noticeable, is of the same order of  
315 magnitude or lower than previous studies using the same method (see Supplementary  
316 Information and Table S1).

318 **Figure 1: Results of Fe isotopic fractionation in metal–silicate equilibrium experiments.** Error bars correspond to 2  
319 standard deviations ( $2\sigma$ ) propagated from the errors on metal and silicate Fe isotopic measurements. **a.**  $\Delta^{56}\text{Fe}_{\text{metal-silicate}}$  of all  
320 experiments of this study and from literature recalculated at 1873 K (Poitrasson et al., 2009; Hin et al., 2012; Elardo and  
321 Shahar, 2017) as a function of the Ni content of the metallic phase. The horizontal dashed line with orange error represents  
322 the average fractionation factor of all experiments presented on the graph, *i.e.*  $0.04 \pm 0.09\text{‰}$  ( $1\sigma$ ). The dotted line represents

323 the trend of Ni content dependency on Fe isotopic fractionation factor determined by Elardo and Shahar (2017). **b.**  $\Delta^{56}\text{Fe}_{\text{metal-silicate}}$   
324 for all experiments of this study. Squares are experiments carried out in graphite capsules, circles correspond to the ones  
325 in crushable MgO capsules, triangles for experiments using boron nitride capsules and diamonds for experiments carried out  
326 in single-crystal MgO capsules. Experiments are color coded according to their time durations: white for 1-5 minutes, pale  
327 green for 15 minutes, light green for 30 minutes, and dark green for 60 minutes. The horizontal dashed line with grey error  
328 represents the average of all experiments presented on graph a., *i.e.*  $0.04 \pm 0.09\%$  ( $1\sigma$ ).

329 **Figure 2: Results of Fe isotopic compositions on a single experiment.** E375 was performed at 2 GPa and 1873 K in a  
330 single-crystal MgO capsule with a time duration of 5 minutes. Its metal phase contains 5 wt.% of Ni. The displayed errors  
331 correspond to  $2\sigma$  on the measurements. Iron isotopic measurements were performed on two silicate samples—a central  
332 glassy part and an outlying crystallised part—as well as the metal phase and the capsule. The two silicate parts and the metal  
333 present identical isotopic composition to the starting material whereas the inner part of the capsule (inner 30  $\mu\text{m}$ ) is enriched  
334 in light Fe isotopes. The diffusion of MgO from the capsule has no effect on the Fe isotopic composition of the silicate phase  
335 and the Fe loss as detected by a light isotopic composition in the capsule is not sufficient to induce a resolvable fractionation  
336 of the sample towards heavier Fe isotopic compositions.

337 The Fe isotopic compositions of the bulk starting materials were measured to assess  
338 potential Fe loss during the experiments (Table 1). These measured Fe isotope compositions  
339 of the starting material (initial bulk) were compared to bulk values calculated by mass balance  
340 from the Fe isotopic composition of the metal and silicate from the run products (final bulk,  
341 Figure 3). Although there is no clear correlation between the fractionation factor and the  
342 agreement between the initial and final bulk values (see also Figure S5), both estimations of  
343 the bulk experimental Fe isotope composition coincide within error for 12 out of the 22  
344 experiments of the reported data set. This indicates that experiments largely approached  
345 closed system behaviour, resulting in limited loss of Fe through the capsule. Thorough  
346 characterisation of the Fe content of experimental capsules after the experiments showing  
347 moderate Fe diffusion into the capsules also corroborates closed system behaviour, a  
348 necessary condition to properly investigate isotope fractionation processes experimentally.  
349 Iron loss quantified by the difference between the measured and calculated bulks shows that  
350 single-crystal MgO and graphite capsules present reduced Fe loss with respect to other  
351 capsule types (Supplementary Figure S5).

352 **Figure 3: Comparison of initial and final Fe isotope compositions of bulk experiments.** The thick blue line and  
353 associated error represents the average Fe isotope composition measured from the powder mixtures used as starting materials  
354 in the experiments and  $2\sigma$  on the measurements. Because the MORB contains c.a. 6 wt.% Fe with a different Fe isotopic  
355 composition from the Fe dopant, the starting material presents an intermediate Fe isotopic composition between that of the  
356 two Fe-bearing components, namely the MORB (light blue line) and the Fe dopant (dark blue line). The symbols correspond  
357 to calculated Fe isotope compositions of final bulk experiments from mass balance using metal and silicate Fe isotopic  
358 measurements for each experiment (equation given on the figure). The error bar of each symbols corresponds to  $2\sigma$   
359 propagated from uncertainties on the Fe isotopic and Fe elemental measurements. The symbol shapes correspond to different  
360 capsule types. Symbols presenting a horizontal-line pattern correspond to experiments producing a metal-silicate  
361 fractionation factor superior to 0.06‰. Final bulk value symbols are color-coded according to the run duration: white for 1-5  
362 min, pale green for 15 min, light green for 30 min and dark green for 60 min.

363 Fig. 4 displays the Fe isotope composition of single experimental phases from this  
364 study and [Elardo and Shahar \(2017\)](#) as a function of metal Ni content. The Ni-induced Fe  
365 isotope fractionation proposed by [Elardo and Shahar \(2017\)](#) implies that such a fractionation  
366 is driven by the composition of the metallic phase as Ni is almost exclusively hosted in the  
367 metallic phase. Therefore, a clear variation of the metal Fe isotope composition should be  
368 observed as a function of the metal Ni content. However, no systematic change in the metal  
369 Fe isotope composition is reported in the data set from this study (Figure 4a). In fact, the  
370 composition of both experiments performed with 70 wt.% of Ni ( $\delta^{56}\text{Fe}_{\text{metal}}=0.27\pm 0.04\text{‰}$  and  
371  $0.32\pm 0.03\text{‰}$ ,  $2\sigma$ ) are strictly identical to the measured bulk experimental starting material  
372 ( $0.27\pm 0.10\text{‰}$ ,  $2\sigma$ ). It should be noted that these two experiments also underwent negligible  
373 Fe loss (Figure 3). This constitutes strong evidence for a very limited effect of Ni on the Fe  
374 isotopic fractionation at the experimental conditions of this study. The silicate Fe isotopic  
375 compositions present a more scattered distribution (Figure 4b), likely due to moderate  
376 interactions with the capsule, and equally show no correlation with the metal Ni content.  
377 Experiments contained in crushable MgO and BN have on average heavier Fe isotopic  
378 signatures in both the metal and the silicate, compared to single-crystal MgO and graphite  
379 capsules. This is likely due to higher Fe loss permitted by faster diffusion in these materials.  
380 However, we show that Fe loss is limited in all of our experiments (Figure 3) and importantly  
381 does not correlate with the Fe isotopic fractionation results (Figure S5). Therefore, the limited  
382 Fe loss occurring during the experiments presented in this study is either i) not significant  
383 enough and/or ii) occurs over a timescale slower than for metal–silicate equilibrium. As a  
384 result, the degree of Fe loss appears not to affect the Fe isotopic fractionation factor. Iron  
385 isotope compositions of single experimental phases from [Elardo and Shahar \(2017\)](#) show a  
386 clear correlation with Ni content (Figures 4c and d). However, this is due to the fact that the  
387 silicate fraction of the starting mixtures was spiked with  $^{54}\text{Fe}$ , resulting in starting materials





E378	2	0	30	SC MgO	4	0.37	0.03	0.50	0.02	4	0.28	0.07	0.37	0.10	0.10	0.07	0.15	0.10
E378 replicate										4	0.25	0.04	0.34	0.11				
E379	2	10	30	SC MgO	4	0.38	0.05	0.55	0.08	4	0.15	0.04	0.21	0.11	0.23	0.06	0.35	0.13
E380	2	20	30	SC MgO	4	0.34	0.09	0.50	0.16	4	0.51	0.09	0.73	0.16	-0.17	0.12	-0.23	0.23
E381	2	30	30	graphite	4	0.30	0.05	0.41	0.01	4	0.26	0.05	0.38	0.04	0.04	0.07	0.03	0.04
E382	2	30	60	SC MgO	4	0.30	0.09	0.44	0.12	4	0.54	0.05	0.81	0.01	-0.24	0.11	-0.37	0.12
Material					n	$\delta^{56}\text{Fe}$	$2\sigma$	$\delta^{57}\text{Fe}$	$2\sigma$									
Fe dopant					5	0.47	0.08	0.69	0.06									
MORB					5	0.09	0.05	0.11	0.08									
Fe-Ni <sub>0</sub>					4	0.33	0.10	0.47	0.24									
Fe-Ni <sub>10</sub>					4	0.28	0.04	0.45	0.14									
Fe-Ni <sub>20</sub>					4	0.26	0.03	0.38	0.17									
Fe-Ni <sub>30</sub>					4	0.30	0.08	0.45	0.23									
Fe-Ni <sub>70</sub>					4	0.27	0.10	0.42	0.20									

405

## 406 Iron isotopic fractionation during core formation

407 Our experimental results include novel metal–silicate equilibrium Fe isotope  
408 fractionation data over a wide range of Ni content. There is presently no evidence that  
409 pressures above 2 GPa could generate an Fe isotopic fractionation between metal and silicate  
410 (Liu et al., 2017; Poitrasson et al., 2009). This allows us to interpret our experimental data in a  
411 context of core formation on large bodies such as the Earth, but also on smaller bodies such as  
412 Vesta without the need for pressure extrapolation. The large number of experiments in this  
413 study provides enhanced statistical accuracy and shows that Ni proportion in the metal has a  
414 limited effect on the Fe isotope fractionation between metal and silicate. While our  
415 conclusions may differ, the results presented here yield a similar range of Fe isotopic  
416 fractionation factors (Figure 1a) to those reported in the study of Elardo and Shahar (2017)  
417 (average and  $1\sigma$  of  $0.05\pm 0.11$  and  $0.04\pm 0.10$  respectively) where the observed trend in Elardo  
418 and Shahar (2017) can be attributed to a relatively smaller number of experiments. In fact, the  
419 data presented herein, combined with that of Elardo and Shahar (2017) yields an average  
420 fractionation factor of  $\Delta^{56}\text{Fe}_{\text{met-sil}}=0.04\pm 0.11\%$  ( $1\sigma$ ), indicating no isotopic fractionation

421 within uncertainty. Our results are also in agreement with previous experimental studies on Fe  
422 isotopic fractionation (Poitrasson et al., 2009; Hin et al., 2012; Liu et al., 2017) demonstrating  
423 the absence of Fe isotope fractionation at metal–silicate equilibrium as well as the absence of  
424 a resolvable fractionation with addition of limited amounts of Ni (Poitrasson et al., 2009; Liu  
425 et al., 2017). Significant S contents in the alloy may induce a positive metal–silicate isotope  
426 fractionation with light isotope enrichment in the silicate (Shahar et al., 2014; Liu et al., 2017;  
427 Pinilla et al., 2021). A similar trend of isotopic fractionation has been observed for other light  
428 elements such as Si and C (Elardo et al., 2019; Shahar and Young, 2020). Nevertheless, this  
429 vector of Fe isotopic fractionation is opposite to that observed in Earth’s mantle and in the  
430 mantles of other differentiated bodies relative to chondrites. Moreover, iron meteorites  
431 present Fe isotopic signatures heavier than chondrites by 0.1–0.2‰ ( $\delta^{57}\text{Fe}$ ) (e.g. Williams et  
432 al., 2006), such that heavy signatures in planetary mantles are unlikely to be established solely  
433 by differentiation of chondritic material. Instead, the heavy signatures observed in iron  
434 meteorites have been explained by core crystallisation with significant Fe isotopic  
435 fractionation measured experimentally between solid and liquid iron alloys (Ni et al., 2020;  
436 Young 2021). As such, to a first order there is no direct evidence that core–mantle segregation  
437 causes significant isotopic fractionation of Fe, and hence is unlikely to control the isotopic  
438 compositions of planetary mantles. Indeed, equilibrium isotope fractionation depends on  $1/T^2$   
439 so that enhanced fractionation is expected from low temperature core formation. However,  
440 chondritic Fe isotopic values are observed in small bodies that underwent lower temperature  
441 core formation such as Mars and Vesta (~1900 K, Righter and Drake, 1996) compared to the  
442 non-chondritic Earth whose core equilibrated with the mantle at 2500–3500 K (Wood et al.,  
443 2006). Therefore, core formation alone cannot explain the diversity of Fe isotope  
444 compositions in the inner solar system. Other processes including various accretion and deep

445 mantle mechanisms need to be further investigated as potential causes for the Fe isotopic  
446 variability between inner solar system planetary bodies.

447 An increasing body of stable isotopic measurements on lunar rocks suggests that the  
448 Earth and the Moon share a common Fe isotope signature (Wang et al., 2015; Sossi and  
449 Moynier, 2017; Poitrasson et al., 2019), and isotopic differences are only observed among  
450 volatile elements that were depleted from the Moon. This result is in agreement with the  
451 widely-held view that volatilisation during Moon formation and evolution resulted in isotopic  
452 fractionation of multiple isotopic systems in the lunar mantle on a much larger scale than in  
453 Earth's mantle (e.g. Gargano et al., 2020; Kato et al., 2015; Pringle and Moynier, 2017; Sharp  
454 et al., 2010; Sossi et al., 2018a,b; Wang and Jacobsen, 2016; Wang et al., 2019; Wimpenny et  
455 al., 2022). The observation of isotopically heavy Fe isotope signatures in ureilites and angrites  
456 suggests that collisions between accreting bodies in the inner solar system may be common  
457 and that the observed signatures in inner planetary bodies are the result of contrasted  
458 accretional histories possibly involving multiple complex volatile-depleting events (Sossi et  
459 al., 2016b). The corollary states that Mars and Vesta suffered very limited syn- and post-  
460 accretion Fe loss as inferred from their chondritic signatures. The crystallisation of  
461 bridgmanite in large planetary bodies such as the Earth may enrich the mantle in heavy Fe  
462 isotopes (Williams et al., 2012) by disproportionation of  $\text{Fe}^{2+}$  into  $\text{Fe}^{3+}$  and metallic iron.  
463 Because metallic iron is isotopically lighter than ferric iron-bearing bridgmanite, its  
464 sequestration into Earth's core could produce a heavy Fe signature in the mantle without  
465 invoking volatilisation processes. Note however that such effects have not been confirmed so  
466 far by *in situ* NRIXS measurements and modelling of lower mantle crystallisation (Yang et  
467 al., 2019). Furthermore, this process would not explain the heavy isotope compositions of  
468 achondritic meteorites, notably the angrites, whose parent bodies were too small to have  
469 stabilised bridgmanite or  $\text{Fe}^{3+}$ -rich silicate melt (Armstrong et al., 2019) in equilibrium with

470 iron metal. Therefore, taken together, the aforementioned mechanisms provide possible  
471 means to explain the diverse Fe isotope signatures observed amongst inner solar system  
472 bodies.

473 [Leshner et al. \(2020\)](#) argued that the core–mantle boundary could be enriched in the  
474 heavier Fe isotopes and that pockets of isotopically heavy core-derived material could be  
475 entrained in plumes towards the upper mantle and crust. High temperature diffusion  
476 experiments over a thermal gradient of 250°C could provide a mechanism for the enrichment  
477 of the outer core and lowermost mantle in heavy Fe isotopes. This process is predicted to be  
478 slower than the mantle convection and recycling according to the study. Therefore, the  
479 proposed mechanism to transport heavier Fe core diapirs in magmatic plumes towards the  
480 upper mantle could be tested by measurements on OIBs spanning a wide age range. Although  
481 a subsequent study on OIBs does not support the entrainment of core-derived liquids as the  
482 source of observed heavy OIB signatures ([Soderman et al., 2021](#)), the authors present a  
483 variety of potential processes explaining Fe isotopic variations in the mantle (*e.g.* subduction  
484 processes, hydrothermal alteration, contribution from sediments, peridotite and pyroxenite  
485 melting). Three main processes contributing to the establishment of heavy Fe isotope  
486 signatures in magmatic rocks are identified: (1) processes enriching eclogite in heavy Fe  
487 isotopes such as metamorphism, (2) lithospheric processing and (3) melting processes in  
488 conditions preventing the dilution of heavy Fe isotope-enriched melts with lighter isotope-  
489 enriched materials ([Soderman et al., 2021](#)). Diamond metallic inclusions with isotopically  
490 heavy signatures ([Smith et al., 2021](#)) that could not be explained by iron disproportionation  
491 were also proposed to have originated from deep recycling processes. Despite the propensity  
492 for these processes to generate isotopic variability locally, they are incapable of shifting the  
493 composition of the entire mantle over geological time. This conclusion is founded on the  
494 observation of the constancy in the Fe isotope compositions of komatiites through time

495 (Hibbert et al. 2012; Dauphas et al. 2010), and their concordance with those of peridotites that  
496 reflect the modern-day mantle (e.g. Weyer and Ionov, 2007).

## 497 **On ideality in Fe–Ni alloys**

498 The amount of Ni in the core is estimated to be around 5 wt.% based on geochemical  
499 and cosmochemical models (Palme and O'Neill, 2014). Numerous studies show that the  
500 physical properties such as density, compressibility, and wave velocities of Fe–Ni alloys vary  
501 only marginally as a function of their Ni content (e.g., Lin et al. 2003; Kantor et al. 2007;  
502 Asker et al., 2009; Martorell et al. 2013; Wakamatsu et al., 2018). Importantly, the constant  
503 density of such alloys does not uniquely constrain the amount of Ni in the core but implies  
504 that the entirety of the observed density deficit detected by seismological techniques (Birch,  
505 1952) must be due to the presence of ~10% of light elements in the core (e.g. O, Si; Badro et  
506 al., 2015). Although the insensitivity of the physical properties of Fe–Ni alloys to the Ni  
507 content has been verified from multiple experimental studies (Lin et al., 2003; Kantor et al.,  
508 2007) and *ab initio* calculations (Martorell et al., 2013), this does not necessarily hold for  
509 their chemical properties.

510 Presently, studies investigating experimentally the elemental and isotopic behaviour of  
511 elements between metal and silicate that have considered the effect of the amount of Ni in the  
512 metal are very scarce, largely because it has been widely assumed to be negligible for the  
513 conditions relevant to core formation. Through the transpiration method and integration of the  
514 Gibbs-Duhem relation, early studies illustrated that the dissolution of Ni had little effect on  
515 the activity of Fe in binary alloys (Oriani, 1953; Zellars et al. 1959). As such, it is frequently  
516 assumed that the alloying of Fe with small quantities of Ni has a negligible effect on the  
517 behaviour of the metal. More recent Knudsen effusion mass spectrometric investigations of  
518 the activity–composition relations of liquid Fe–Ni alloys (Fraser and Rammensee, 1982) show  
519 that the activity coefficient of Fe,  $\gamma_{\text{Fe}}$ , varies between 0.9 at  $X_{\text{Ni}}=0.5$  to 1 at  $X_{\text{Ni}}=0.05$  (with X

520 the mole fraction), with  $\gamma_{\text{Fe}}$  remaining essentially  $\sim 1$  up to  $X_{\text{Ni}} = 0.4$  (Swartzendruber et al.  
521 1991), exceeding the range for the metallic cores of rocky bodies. This therefore highlights  
522 that for all likely core compositions, Fe dissolution in Fe–Ni alloys is ideal at ambient  
523 pressures. Translation of Fe–Ni ideality to isotopic fractionation can be examined from a  
524 crystal chemical standpoint. The activity coefficient is a manifestation of the energetics of  
525 mixing between two compounds. That two phases (here, Ni (l) and Fe (l)) mix ideally implies  
526 little to no energy penalty with respect to a mechanical mixture of these two phases. How  
527 readily Ni substitutes into Fe depends upon the degree to which Ni distorts its structural and  
528 electronic arrangement. Such a distortion can be described by the lattice strain model, and be  
529 used to predict activity coefficients, where  $\ln\gamma(\text{Fe}) = \Delta G_{\text{Fe-Ni}}^{\text{strain}}/RT$ . Here,  $R$  is the gas  
530 constant,  $T$  the absolute temperature and  $\Delta G_{\text{Fe-Ni}}^{\text{strain}}$  the Gibbs free energy change of reaction  
531 associated with the strain induced on the structure by the substitution (e.g. Wood et al. 2019).  
532 Importantly, the quantity  $\Delta G_{\text{Fe-Ni}}^{\text{strain}}$  relates to the difference in bonding environment  
533 (electronic charge & radius; cf. Blundy and Wood, 1994) between Fe and Ni, where a value of  
534 0 leads to  $\gamma = 1$  (i.e., ideal behaviour). These same properties induce equilibrium stable  
535 isotope fractionation (e.g., Dauphas et al. 2014; Sossi and O'Neill, 2017). Waseda and Ohtani  
536 (1974) measured the nearest-neighbour distances in pure liquid Fe and liquid Ni metals,  
537 obtaining values very similar to one another at 1500 °C, 2.58 Å and 2.53 Å, respectively. That  
538  $\gamma_{\text{Fe}} \sim 1$  implies little distortion of the electronic and physical state of liquid Fe metal by the  
539 substitution of up to  $\sim 40$  mol % of Ni, meaning there is no driver for isotopic fractionation  
540 associated with changing Ni content.

541 At higher pressures (up to 94 GPa), Huang and Badro (2018) demonstrated that Ni  
542 (from 3.5 to 48.7 wt.% Ni) has no effect on the metal–silicate partitioning of Ni, Cr and V.  
543 Our study suggests that the Ni fraction in an Fe–Ni alloy has no effect on Fe isotopic  
544 behaviour in a metal–silicate equilibrium, also pointing towards chemical Fe–Ni ideality. Xia

545 [et al. \(2019\)](#) report contrasting observations, in which the magnitude of Zn and Cu isotope  
546 fractionation between metal and silicate was shown to be sensitive to the Ni content of the  
547 alloy, which could be due to specific interactions between Ni and these elements, owing to  
548 their disparate electronic arrangement and mean bond lengths. The 22 experiments reported  
549 herein not only enhance statistical accuracy but were also designed to explore a much broader  
550 range of Ni concentrations in the metal, from 0 to 70 wt.% (~25 wt.% maximum [in Xia et al.,](#)  
551 [2020](#) and [Elardo and Shahar, 2017](#)). The assumption of chemical Fe–Ni ideality has extended  
552 implications regarding our understanding and knowledge of core formation related processes.

## 553 **Conclusion**

554 We performed 22 metal–silicate equilibrium experiments with the objective of  
555 establishing the effect of Ni content of Fe–Ni alloys on the Fe isotopic fractionation factor  
556 between metal and silicate. All experiments were carried out at 1–2 GPa and 1873 K, at which  
557 a variety of capsule materials were tested. An experimental time series, as well as  
558 petrographic analyses of the samples were performed in order to assess the conditions at  
559 which chemical equilibrium was achieved. No significant effect of Ni concentration on Fe  
560 isotope fractionation was detected over a wide range of Ni contents (0–70 wt. % in the metal  
561 fraction), supporting the common assumption of physical and chemical ideality of Fe–Ni  
562 alloys. Because the observed fractionation factor between metal and silicate is  
563 indistinguishable from zero, our data does not support core formation as the main mechanism  
564 causing the diversity in Fe isotope signatures observed amongst inner solar system planetary  
565 bodies. A combination of multiple accretion-related processes—including condensation from  
566 the solar nebula, volatile-depleting events such as giant impacts, and disproportionation of  
567 ferrous iron at high pressure in larger bodies—could together contribute to the observation of  
568 heavier-than-chondritic signatures in the Earth and the Moon’s mantles.



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## 580 **Appendix A. Supplementary material**

581 Research data and supplementary information to this article regarding isotopic equilibrium,  
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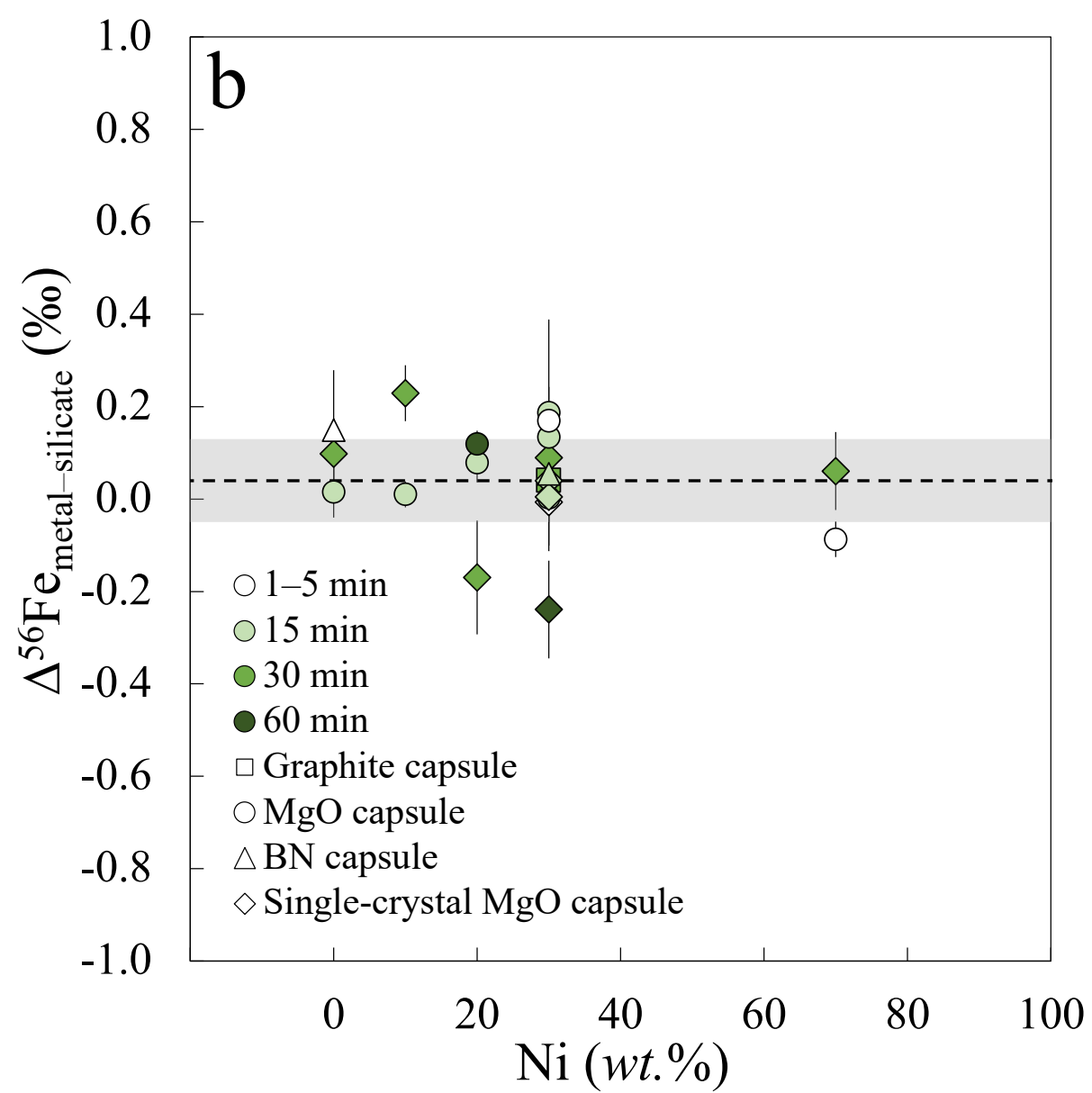
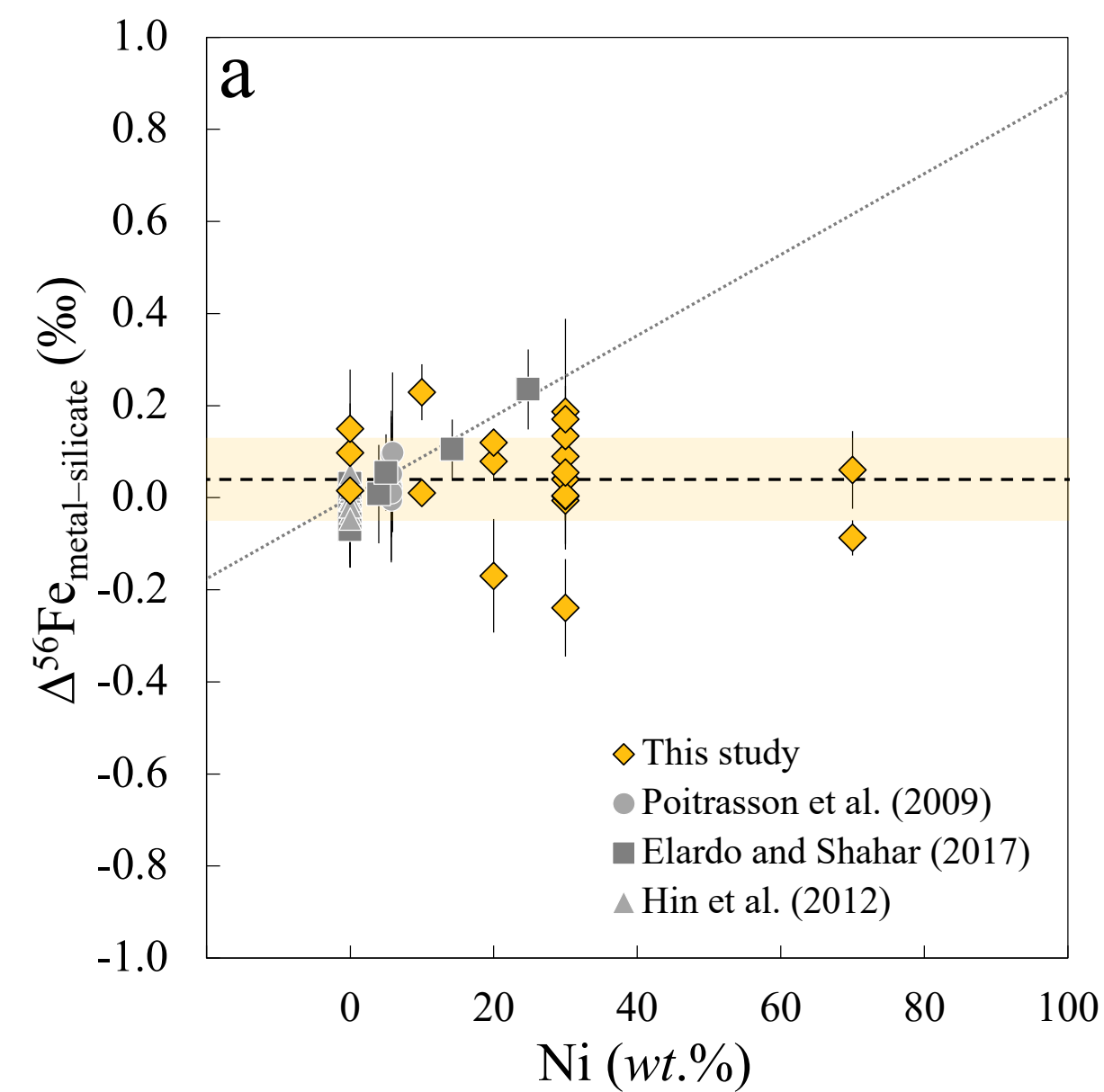
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Single-crystal MgO capsule

