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The formation of early-type galaxies through monolithic collapse of gas clouds in Milgromian gravity

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ABSTRACT

Studies of stellar populations in early-type galaxies (ETGs) show that the more massive galaxies form earlier and have a shorter star formation history. In this study, we investigate the initial conditions of ETG formation. The study begins with the collapse of non-rotating post-Big-Bang gas clouds in Milgromian (MOND) gravitation. These produce ETGs with star-forming time-scales (SFT) comparable to those observed in the real Universe. Comparing these collapse models with observations, we set constraints on the initial size and density of the post-Big-Bang gas clouds in order to form ETGs. The effective-radius–mass relation of the model galaxies falls short of the observed relation. Possible mechanisms for later radius expansion are discussed. Using hydrodynamic MOND simulations, this work thus for the first time shows that the SFTs observed for ETGs may be a natural occurrence in the MOND paradigm. We show that different feedback algorithms change the evolution of the galaxies only to a very minor degree in MOND. The first stars have, however, formed more rapidly in the real Universe than possible just from the here studied gravitational collapse mechanism. Dark-matter-based cosmological structure formation simulations disagree with the observed SFTs at more than 5σ confidence.

Key words: galaxies: formation – galaxies: evolution – galaxies: elliptical – galaxies: star formation – galaxies: stellar content.

1 INTRODUCTION

Present-day early-type galaxies (ETGs) host old stars without much gas and with a very low star formation activity. Studying the stellar properties of these galaxies is very important for understanding the cosmic history of the Universe. Glazebrook et al. (2017) notes that extreme star formation events in the early Universe are not rare events and that they play a significant role in the early mass assembly of these galaxies. They report the discovery of a quiescent galaxy with stellar mass $\approx 1.7 \times 10^{11} M_{\odot}$ at $z = 3.17$. Martín-Navarro et al. (2018), in their stellar population analysis of ETGs, report that the bulk of the stars form in a short time-scale with star formation lasting longer in the central regions, and they suggest that an early monolithic formation is highly likely. The observational evidence is reviewed in Yan, Jeřábková & Kroupa (2021) and Kroupa et al. (2020), who show how the formation of supermassive black holes is a natural outcome of the formation of ETGs within the integrated galaxy-wide initial mass function (IGIMF; Kroupa & Weidner 2003; Weidner & Kroupa 2006; Jeřábková et al. 2018; Yan et al. 2021) theory.

The observations thus point towards putative progenitors of ETGs undergoing very rapid early formation (Cowie et al. 1996; Nelan et al.

2005; Thomas et al. 2005; Recchi, Calura & Kroupa 2009; McDermid et al. 2015; Liu et al. 2016). Studies based on elemental abundances of the ETGs indeed constrain them to have formed rapidly on a time-scale of ≈ 1 Gyr. Also, studies from stellar population models with observed line indices show that ETGs must have formed the bulk of their stars in a short time-scale. Cowie et al. (1996) suggested the name ‘downsizing’ to describe that the less massive galaxies have more extended SFHs compared to the massive ones.

The duration of the formation, i.e. the star-forming time-scale (SFT) for ETGs, is expressed in Thomas et al. (2005), Recchi et al. (2009), and McDermid et al. (2015) in terms of the downsizing time which is a function of the mass of the galaxy. These downsizing time-scales correspond to approximately 0.34 Gyr for a galaxy with a present-day baryonic mass of $10^{12} M_{\odot}$, suggesting that it should have formed under a monolithic cloud collapse scenario rather than hierarchical merging that requires longer time-spans (Ricciardelli et al. 2010; Wuyts et al. 2010). On average, more massive galaxies have been found to have formed the bulk of their stars on a shorter time-scale and have completed their star formation activity at a higher redshift (Daddi et al. 2005; Wiklind et al. 2008; Castro-Rodríguez & López-Corredoira 2012) which points to a strong dependence of the SFT in a galaxy on its present-day stellar-mass content.

Massive ETGs are thought to be important tracers of the cosmic history of the stellar mass assembly and of galaxy evolution. A theoretical model that describes ETGs should explain how and why they have these observed SFTs. The standard model of cosmology

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(SMoC) has been successful in forming ETGs (see Naab & Ostriker 2017 for a detailed review on the theory and numerical simulations in the SMoC) but has not been able to reproduce ETGs with SFTs (see Fig. B1) similarly short as the observed ETGs (Thomas et al. 2005; Recchi et al. 2009; McDermid et al. 2015). With the existence of dark matter (DM) remaining a hypothesis, an alternate theory of cosmology should be considered (Kroupa, Pawlowski & Milgrom 2012; Kroupa 2015).

By combining observational constraints on the dynamics in galaxies with constraints from the Solar system, Milgrom (1983a) corrected the theory of gravitation at low acceleration, which could be a consequence of the quantum vacuum (Milgrom 1999; Pazy 2013; Smolin 2017; Verlinde 2017). In analogy with Newtonian dynamics, a non-relativistic Milgromian gravity theory (MOND) can be constructed by setting up a Lagrangian which, upon extremization of the action, yields a generalized Milgromian Poisson equation. Two such Lagrangians have been proposed, called AQUAL (Bekenstein & Milgrom 1984) and QUMOND (Milgrom 2010). The latter, which we adopt hereafter, yields the following generalized Poisson equation,

$$\Delta\Phi(\vec{x}) = 4\pi G\rho_b(\vec{x}) + \vec{\nabla} \cdot [\tilde{v}(|\vec{\nabla}\phi|/a_0)\vec{\nabla}\phi(\vec{x})], \quad (1)$$

or,

$$\Delta\Phi(\vec{x}) = 4\pi G[\rho_b(\vec{x}) + \rho_{ph}(\vec{x})], \quad (2)$$

where $\rho_b(\vec{x})$ is the baryonic density, $\phi(\vec{x})$ is the Newtonian potential that fulfils the standard Poisson equation, $\Delta\phi(\vec{x}) = 4\pi G\rho_b(\vec{x})$ and Milgrom's constant $a_0 \approx 1.2 \times 10^{-10} \text{ ms}^{-2} \approx 3.8 \text{ pc Myr}^{-2}$. The phantom dark matter (PDM) density, $\rho_{ph}(\vec{x})$, is not a real density distribution but a mathematical function that arises out of the non-linearity of the Poisson equation. $\Phi(\vec{x})$ is the total gravitational potential from which the accelerations follow, $\vec{a} = -\vec{\nabla}\Phi$, and $\tilde{v}(y)$ is a transition function characterizing the theory (see Milgrom 2008, 2010, 2014; Famaey & McGaugh 2012; and Banik & Zhao 2022 for detailed reviews on the theory). $\tilde{v}(y)$ has the limits,

$$\tilde{v}(y) \rightarrow 0 \text{ for } y \gg 1 \text{ and } \tilde{v}(y) \rightarrow y^{-1/2} \text{ for } y \ll 1. \quad (3)$$

The above formulation deals only with linear differential equations and is shown to emerge as a natural modification of a Palatini-type formulation of Newtonian gravity. It is a member of a larger class of bipotential theories (quasi-linear formulation of MOND; QUMOND Milgrom 2010). Bekenstein & Milgrom (1984) also noted a correspondence with some theories of quark confinement using a different form of the function \tilde{v} .

MOND predicted galaxy scaling relations obeyed by galaxies (Milgrom 1983b) such as the Baryonic Tully Fisher Relation (BTFR; McGaugh et al. 2000; McGaugh 2005, 2012; Liu et al. 2016) and the Radial Acceleration Relation (RAR; Sanders 1990; McGaugh 2004; Lelli et al. 2017).

Any realistic theory of galaxy formation has to address the question of when and how the stars form in them. There exists a large body of literature on the formation of ETGs in the SMoC (Naab & Ostriker 2017 for a review). Here, we concentrate on the question whether ETGs can form in Milgromian gravitation from collapsing post-Big-Bang gas clouds and how these models compare to the observed ETGs. The aim of this work is to use simulations of collapsing and non-rotating post-Big-Bang gas clouds to study the duration of their SFTs and the resulting size–mass relation. Galaxy-scale fluctuations in a MOND Universe should grow by the application of MOND to the peculiar accelerations allowing for large density contrasts and for large galaxies to form as early as $z = 10\text{--}25$ from the collapse of almost isolated gas clouds. Nipoti, Londrillo & Ciotti (2007) and Sanders (2008) have therefore studied disipationless collapse

in a MOND context. Here, we model for the first time the collapse of massive non-rotating gas clouds including the physics of star formation. The results are found to be very close to the observations. Previous work has shown the important bulk properties of disc galaxies to follow naturally from the collapse of initially rotating post-Big-Bang gas clouds (Wittenburg, Kroupa & Famaey 2020).

The layout of this paper is as follows. In Section 2, the numerical hydrodynamical code used is discussed briefly and in Section 3, the details of the simulations are described. The results from the work are presented in Section 4. Section 5 contains the discussion and Section 6 has conclusions along with potential future work.

2 PHANTOM OF RAMSES

The Phantom of RAMSES (POR) code (Lüghausen, Famaey & Kroupa 2015; Nagesh et al. 2021) is a customized version of RAMSES (Teyssier 2002), applying the adaptive mesh refinement (AMR) method to solve the Milgromian Poisson equation (equation 1). It uses a multigrid and a conjugate gradient solver to solve the generalized Poisson equation.¹

There are several transition functions used in the literature to accommodate the non-linearity of the Poisson equation in MOND. POR uses

$$\tilde{v}(y) = -\frac{1}{2} + \left(\frac{1}{4} + \frac{1}{y}\right)^{\frac{1}{2}}, \quad (4)$$

where $y = |\vec{\nabla}\phi|/a_0$ (see Lüghausen et al. 2015).

The POR working scheme involves first solving the standard Poisson equation to compute the Newtonian potential, and then the PDM density is calculated using a discrete scheme (see Lüghausen et al. 2015). Then, the Poisson equation is solved for a second time with the Newtonian density and PDM density to compute the total gravitational potential.

As for now, POR has been successfully applied in the simulations of Antennae-like galaxies (Renaud, Famaey & Kroupa 2016), simulations of the Sagittarius satellite galaxy (Thomas et al. 2017), simulations of the Local Group producing the planes of satellites (Bílek et al. 2018, 2021b; Banik et al. 2022), simulations of streams from globular clusters (Thomas et al. 2018), simulations of the formation of exponential disc galaxies (Wittenburg et al. 2020), the global stability of M33 (Banik et al. 2020), the evolution of globular-cluster systems of ultra-diffuse galaxies due to dynamical friction (Bílek et al. 2021a), and polar-ring galaxies have also been successfully modelled with a pre-POR code (Lüghausen et al. 2013). A detailed user-guide can be found in Nagesh et al. (2021).

3 MODELS

We use the POR numerical code introduced in Section 2 to set-up post-Big-Bang gas clouds in an isolated environment in-order to allow a first assessment of how ETGs would have formed in MOND. Here, we start with initial conditions comparable to those in Wittenburg et al. (2020), but with no initial rotation.

Stellar particles are created in POR if the gas mass-volume-density exceeds a user-defined threshold value. The code checks at every time-step whether any cell exceeds the density threshold, and then a stellar particle is created. Details on the conditions for the star formation threshold value and various other code parameters are

¹Candlish, Smith & Fellhauer (2015) developed a similar code independently called RAYMOND.

described in Lüghausen et al. (2015), Wittenburg et al. (2020), and Nagesh et al. (2021).

In total 42 different model galaxies were calculated for this study to understand how different initial masses, M_{initial} , and initial radii, R_{initial} , of the gas cloud can affect the SFT (Table 1). Only the simple cooling/heating feedback algorithm is used in the simulation of all the model galaxies except models e34c and e39c, where a more complex feedback algorithm is used (see Section 4.3). The computations are run for 10 Gyr. The starting temperature of the gas is 10^4K and the size of the simulation box is 1000 kpc. The minimum and maximum refinement levels for the models² are 7 and 12, respectively, which sets a limit to the minimum spatial resolution of 7.81 kpc and maximum spatial resolution of 0.24 kpc, respectively, for each simulation (Teyssier 2002; Lüghausen et al. 2015).

In all computations, the star formation rate (SFR) is calculated by separating all stellar particles in bins of $\delta t = 10\text{ Myr}$ according to their age, summing up the stellar mass in every time bin and dividing this by the length of δt . The SFR increases sharply at the time of the collapse of the sphere until the maximum is reached and then it decreases approximately exponentially (Fig. 1). The SFH of a galaxy is constructed by the consecutive SFRs of each δt time-step. The SFH can be represented by a Lorentz function (dashed grey line in Fig. 1),

$$\text{SFR}/M_{\odot}\text{yr}^{-1} = \frac{\text{SFR}_{\text{peak}}/M_{\odot}\text{yr}^{-1}}{\pi} + \left(\frac{0.5 \cdot \Delta\tau_{\text{m}}/\text{Gyr}}{(t/\text{Gyr} - t_{\text{peak}}/\text{Gyr})^2 + (0.5 \cdot \Delta\tau_{\text{m}}/\text{Gyr})^2} \right), \quad (5)$$

where t_{peak} is the time in the simulation where the peak of the SFH is observed and $\Delta\tau_{\text{m}}$ is the SFT for the model. The full-width at half maximum (FWHM) of the SFH gives the duration for which the bulk of the star formation takes place, so we use the FWHM to define the SFT (i.e. $\Delta\tau_{\text{m}}$) for the models in this work.

The collapse for models with the same M_{initial} and varying R_{initial} shows that $\Delta\tau_{\text{m}}$ increases with increasing R_{initial} . The sum of all formed stellar particles at the end of each simulation is the final stellar mass,³ M_{final} , of each model, independently of the chosen feedback algorithm. The models here do not lose mass through outflows.

We also calculate the time taken for the first stellar particle to form, t_{start} , and compare it to the MOND free-fall time,

$$t_{\text{ff,mond}} = \left(\frac{\pi}{2} \right)^{0.5} \cdot \frac{R_{\text{initial}}}{(G \cdot M_{\text{initial}} \cdot a_0)^{0.25}}, \quad (6)$$

(equation 27 in Zonoozi et al. 2021) in Fig. C1. The free-fall time can be used as an approximation for the time the gas needs to collapse into stars. The models have t_{start} comparable to the MOND free-fall time (Fig. C1).

4 RESULTS

We compare the SFTs, ages, and effective radii of our simulated galaxies with the observed galaxy properties in this section. The average ages estimated by Thomas et al. (2005) from absorption line indices of 124 ETGs, t_{thomas} for both high-density and low-density

environments (quantities in brackets are the values derived for the high-density environment) can be written as

$$t_{\text{thomas}} = 1.33(2.67) \cdot (M_{\text{final}}/M_{\odot})^{0.07(0.05)}, \quad (7)$$

(equation 3 in Thomas et al. 2005) and the relation deduced for the SFTs is

$$\Delta\tau_{\text{thomas}}/\text{Gyr} = 3.67 - 0.37 \cdot \log_{10}(M_{\text{final}}/M_{\odot}) \quad (8)$$

(equation 5 in Thomas et al. 2005), being the same for high- and low-density environments.

4.1 Downsizing SFH

The SFTs are found to follow a similar downsizing behaviour as documented in Thomas et al. (2005, 2010), where, for a given R_{initial} , the galaxies with larger masses (higher density) form at an earlier epoch compared to galaxies with lower masses (lower density). Fig. 2 shows that the model galaxies with high initial cloud density, ρ_{initial} , form earlier and quicker (shorter $\Delta\tau_{\text{m}}$) than model galaxies with low ρ_{initial} . We use a power law function to obtain the scaling relation in Fig. 2,

$$\Delta\tau_{\text{m}}/\text{Gyr} = 3.03 \cdot (\rho_{\text{initial}}/M_{\odot}\text{kpc}^{-3})^{-0.30}. \quad (9)$$

A downsizing SFH can be well explained using Fig. 3 where the SFT is plotted against M_{final} for each model galaxy along with the observations from Thomas et al. (2005), Recchi et al. (2009), and McDermid et al. (2015).

We use a power-law function to fit the models and obtain the scaling relation for the models in Fig. 3,

$$\Delta\tau_{\text{m}}/\text{Gyr} = A \cdot (M_{\text{final}}/M_{\odot})^B, \quad (10)$$

where A and B are fit parameters for each R_{initial} given in Table A1. From Fig. 3, it can be seen that models with $R_{\text{initial}} = 500\text{ kpc}$ lead to similar $\Delta\tau_{\text{m}}$ values as the observational results from Thomas et al. (2005) and models with $R_{\text{initial}} = 200\text{ kpc}$ and $R_{\text{initial}} = 300\text{ kpc}$ are similar to the observational data from Recchi et al. (2009) but steeper. The models with $R_{\text{initial}} = 50\text{ kpc}$ and $R_{\text{initial}} = 100\text{ kpc}$ do not agree with any of the observed relations. We note that our models do not agree with the new downsizing relation obtained by Yan et al. (2021) and this is discussed in Section 5.

Fig. 4 shows that only models with $R_{\text{initial}} = 500\text{ kpc}$ ($\Delta\tau_{\text{m}}/\Delta\tau_{\text{thomas}} \approx 1$) are comparable to the results from Thomas et al. (2005). From here, we only consider models with $R_{\text{initial}} = 500\text{ kpc}$ for discussions. Using the basic assumption that the standard age of the Universe is 13.8 Gyr and the age dating deduced by Thomas et al. (2005) for ETGs is valid, we translate the SFH of our models to a time related to the start of the Big Bang. Since our models have SFTs comparable to Thomas et al. (2005), we assume that the peaks of the SFHs of our models should coincide with the average ages computed by Thomas et al. (2005). This allows us to place our models to a time relative to the Big Bang (equation 7). We shift the SFHs of our model galaxies to a time corresponding to the average ages as deduced by Thomas et al. (2005). To achieve this, we first transform the observed average formation times (equation 7) to time since the Big Bang,

$$t_{\text{av,th}}/\text{Gyr} = 13.8 - t_{\text{thomas}}/\text{Gyr}, \quad (11)$$

and the time the first stellar particles would have formed according to the models presented here becomes

$$t_{\text{first}}/\text{Gyr} = t_{\text{av,th}}/\text{Gyr} - (t_{\text{peak}}/\text{Gyr} - t_{\text{start}}/\text{Gyr}), \quad (12)$$

where t_{start} is the time when the first stellar particle is formed in the simulation and t_{first} is the time when the first star would have

²Models e34c and e39c have maximum refinement levels of 14 and 13, respectively, which sets the maximum spatial resolution to 0.06 and 0.12 kpc.

³We do not list the M_{final} since the final mass of the model galaxy is equal to the initial mass of the gas cloud.

Table 1. The initial conditions and results of all the models.

Model Name	M_{initial} ($10^9 M_{\odot}$)	R_{initial} (kpc)	$\Delta\tau_m$ (Gyr)	t_{start} (Gyr)	t_{peak} (Gyr)	r_{eff} (kpc)	SFR_{peak} ($M_{\odot} \text{ yr}^{-1}$)
e1	0.6	50	0.62	1.18	2.15	0.43	6.71E + 00
e2	1	50	0.53	1.02	1.86	0.47	1.34E + 01
e3	6.4	50	0.17	0.51	0.82	0.39	2.66E + 02
e4	10	50	0.13	0.44	0.71	0.6	5.20E + 02
e5	30	50	0.07	0.29	0.49	0.84	2.06E + 03
e6	50	50	0.05	0.23	0.44	0.73	4.40E + 03
e7	64	50	0.05	0.21	0.41	0.67	7.22E + 03
e8	70	50	0.04	0.2	0.39	0.77	7.94E + 03
e9	100	50	0.04	0.17	0.35	0.83	1.45E + 04
e10	0.6	100	0.91	2.46	3.5	0.4	4.88E + 00
e11	1	100	0.78	2.11	3.95	0.53	1.03E + 01
e12	5	100	0.37	1.19	1.73	0.47	9.67E + 01
e13	10	100	0.26	0.95	1.41	0.47	2.77E + 02
e14	20	100	0.19	0.73	1.21	1.41	6.05E + 02
e15	30	100	0.16	0.63	1.07	0.56	1.12E + 03
e16	50	100	0.13	0.52	0.92	0.68	2.52E + 03
e17	64	100	0.11	0.48	0.87	0.75	3.97E + 03
e18	70	100	0.11	0.46	0.85	0.77	4.40E + 03
e19	100	100	0.09	0.4	0.76	0.99	7.78E + 03
e20	0.6	200	1.62	5.12	6.4	0.4	3.15E + 00
e21	1	200	1.38	4.37	7.57	0.56	6.38E + 00
e22	10	200	0.53	1.95	2.98	0.55	1.27E + 02
e23	30	200	0.34	1.31	2.18	0.68	6.25E + 02
e24	50	200	0.28	1.11	1.90	1.21	1.32E + 03
e25	70	200	0.25	0.99	1.76	0.93	2.18E + 03
e26	100	200	0.22	0.87	1.59	0.93	3.62E + 03
e27	0.6	300	1.65	7.91	9.29	0.31	2.54E + 00
e28	5	300	0.79	3.81	5.09	0.46	4.51E + 01
e29	10	300	0.65	2.99	4.48	1.03	9.41E + 01
e30	30	300	0.48	2.02	3.31	0.76	4.48E + 02
e31	50	300	0.42	1.68	2.9	0.73	8.96E + 02
e32	70	300	0.39	1.49	2.66	0.83	1.38E + 03
e33	100	300	0.35	1.32	2.42	1.36	2.34E + 03
e34	6.4	500	1.41	5.97	8.44	0.41	3.29E + 01
e34c	6.4	500	1.43	5.96	8.41	0.39	3.25E + 01
e35	10	500	1.18	5.09	7.45	0.47	6.41E + 01
e36	30	500	0.75	3.43	5.58	0.6	2.85E + 02
e37	50	500	0.62	2.85	4.88	0.84	5.95E + 02
e38	64	500	0.55	2.58	4.48	0.95	8.60E + 02
e39	70	500	0.54	2.48	4.4	1.04	9.39E + 02
e39c	70	500	0.53	2.47	4.44	1.09	9.48E + 02
e40	1000	500	0.24	1.03	1.9	1.7	1.48E + 04

Notes. M_{initial} is the initial mass of the post-Big-Bang gas cloud, R_{initial} is the initial radius of the post-Big-Bang gas cloud, $\Delta\tau_m$ is the resulting SFT, t_{peak} is the time after start of the simulation when the first stellar particle is formed, t_{peak} is the resulting time when the peak of the SFH is observed in the simulation, r_{eff} is the resulting projected effective radius, and SFR_{peak} is the resulting SFR at the peak of the SFH. All the models here have a maximum resolution of 0.24 kpc except for models e34c and e39c, which have a maximum resolution of 0.06 and 0.12 kpc, respectively. Models e34c and e39c are computations with complex feedback (see Section 4.3).

formed in the ETG relative to the Big Bang that is consistent with the Thomas et al. (2005) age–mass relation. Fig. 5 shows the average age and when the first star forms in the models for both low- and high-density environments. Fig. 6 is a representation of the SFH of the model galaxies that follow the downsizing timescales as deduced by Thomas et al. (2005). That is, a massive ($\approx 10^{12} M_{\odot}$) post-Big-Bang gas cloud would have started to form an ETG around 1.5 Gyr after the Big Bang in a high-density environment and 3.6 Gyr after the Big Bang in a low-density environment. The observed quasar activity only ≈ 0.70 Gyr after the Big Bang suggests that there may have been significantly more overdense regions (Kroupa et al. 2020). The transparent SFHs in Fig. 6 are on the time axis

where $t = 0$ is the start of the simulation. Thus, the real Universe appears to have accelerated the start of the collapse even in the low-density regions. Rather than the real Universe having accelerated the onset of star formation, another interpretation of Fig. 6 is possible though. Fig. 6 shows that the low-mass models agree with the observed galaxies in terms of the time when the first stars formed in low-density environments, while the massive models agree with the observed timings in high-density regions. This suggests that low-density regions form preferentially low-mass galaxies, while high-density regions form preferentially massive galaxies. This will be testable with cosmological structure formation simulations in MOND.

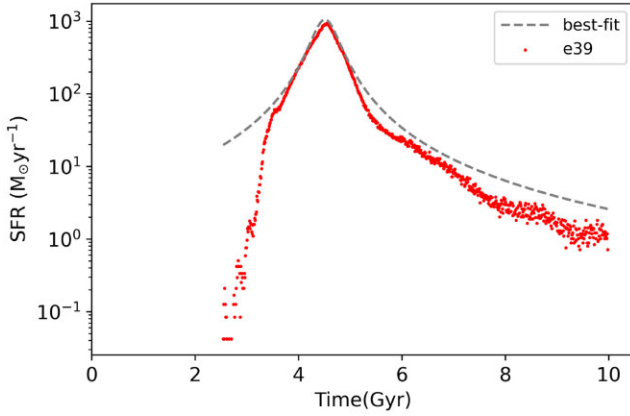


Figure 1. SFH of model e39 (Table 1). The red dots are the SFR at each δt and the dashed grey line is the calculated best fit (equation 5). The x -axis is the time in the computational domain where the computation and thus the collapse starts at 0 and the calculation ends at 10 Gyr. The fit is done only to obtain the FWHM of the SFH and the discrepancy at large t is not important for the purpose of this study.

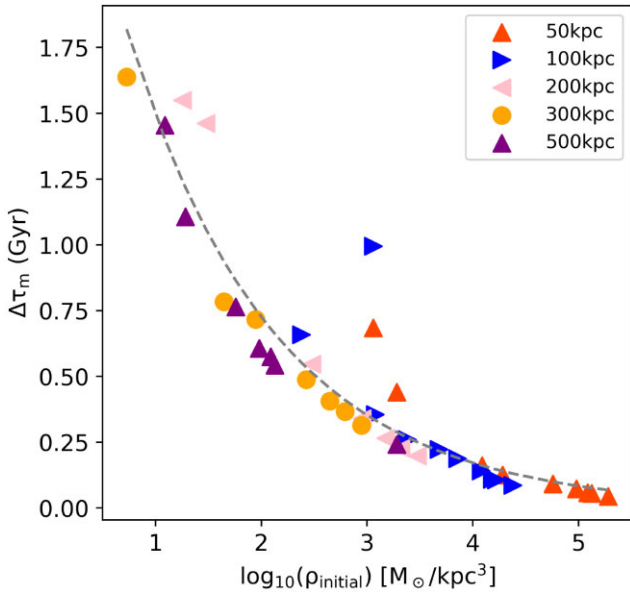


Figure 2. SFT–density relation for all models listed in Table 1. The coloured markers are model galaxies and the dashed grey line is the calculated best fit (equation 9).

4.2 Effective radii of the model galaxies

The effective radius is defined as the projected radius within which half of the final stellar mass of the model galaxy is found, assuming that the mass density follows the luminosity density, this can be compared to the projected half-light radius from observations. The effective-radius is calculated by projecting the model galaxy into the XY plane (Eappen et al. in preparation will further discuss the effective-radius in other projections as well). The shape of the model galaxy is found to be disc-like similar to Wittenburg et al. (2020). The collapse occurs along the Z-direction such that the mid-plane of the formed galaxy is the XY plane. The effective-radius, r_{eff} , is plotted against M_{final} of the model galaxies and compared to the real galaxies from Dabringhausen & Fellhauer (2016) in Fig. 7.

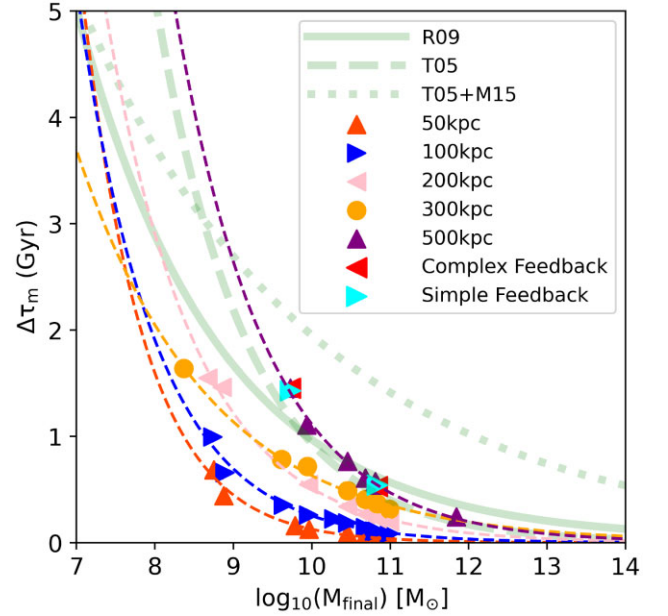


Figure 3. SFT–mass relation. The coloured markers are model galaxies listed in Table 1 with different R_{initial} as indicated by the legend. The observationally constrained relations are from Thomas et al. (2005) (dashed green line), Recchi et al. (2009) (solid green line), and Yan et al. (2021) (which is a combination of equation 3 of Thomas et al. 2005 and equation 3 of McDermid et al. 2015, dotted green line). The red left-pointing triangles are the models with complex feedback (see Section 4.3).

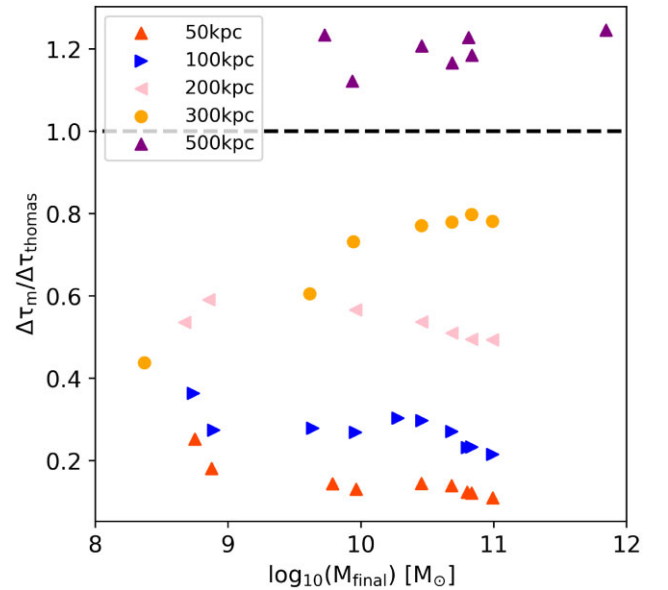


Figure 4. A plot showing the ratio of SFTs of the model galaxies and observed galaxies ($\Delta\tau_m / \Delta\tau_{\text{thomas}}$) versus final stellar mass of the galaxy. The coloured markers are model galaxies as listed in Table 1 with different R_{initial} as indicated by the legend. The dashed black line is where $\Delta\tau_m = \Delta\tau_{\text{thomas}}$.

The MOND radius, r_M , at which the Newtonian radial acceleration equals a_0 (Milgrom 2014) is

$$r_M = \left(\frac{G \cdot M_{\text{final}}}{a_0} \right)^{0.5}, \quad (13)$$

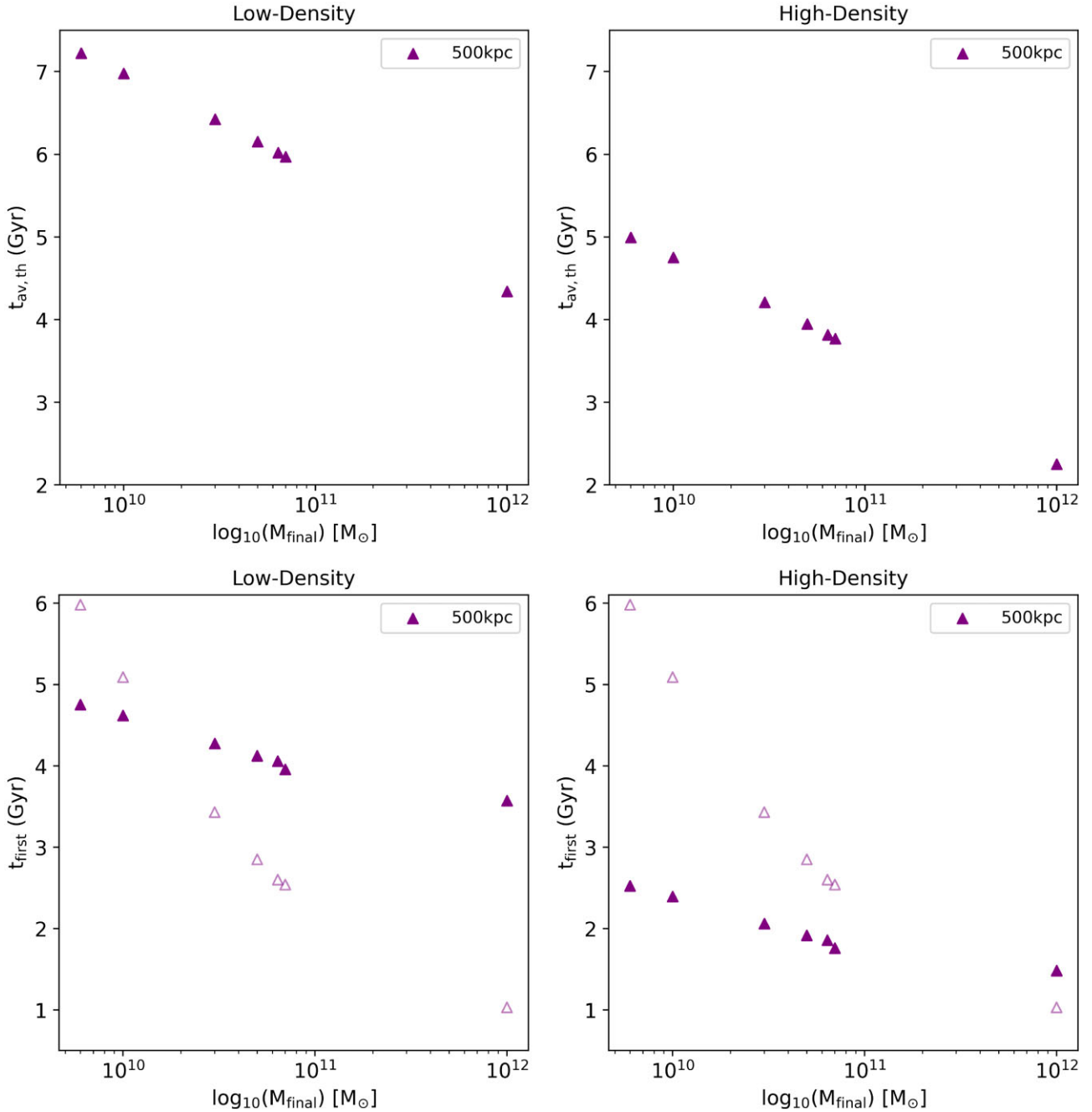


Figure 5. $t_{av,th}$ and t_{first} versus M_{final} for models with $R_{initial} = 500$ kpc in the low-density (*left*) and high-density (*right*) environment (equations 11 and 12). The times are computed since the nominal Big Bang. The unfilled upward pointing purple triangles in the second panel are relative to the start of the computation at $t = 0$.

and is plotted as the solid red line in Fig. 7. Sanders (2008) notes that the effective radius of the dissipationless models in his work are comparable to r_M ; that is, galaxy mass ($\approx 10^{11} M_{\odot}$) objects naturally collapse to a radius of about 10 kpc. The hydrodynamical dissipative simulations of the models in this work form model galaxies similar in shape, size and kinematical properties to the compact relic galaxies like NGC 1277 (Trujillo et al. 2014) and this will be discussed in a companion paper (Eappen et al. 2022b, in preparation).

The best-fitting relation for the Dabringhausen & Fellhauer (2016) galaxies in Fig. 7 is given by

$$r_D/\text{pc} = 0.46 \cdot (M_*/M_{\odot})^{0.35}, \quad (14)$$

where r_D is the observed effective-radius calculated by Dabringhausen & Fellhauer (2016) and M_* is the total stellar mass of the galaxy. The best fit for the model galaxies in Fig. 7 is calculated using a power-law function,

$$r_{\text{eff}}/\text{pc} = 0.12 \cdot (M_{\text{final}}/M_{\odot})^{0.36}. \quad (15)$$

It can be seen in Fig. 7 that the effective-radius–mass relation (equation 15) of the model galaxies has a power-law index comparable to the effective-radius–mass relation (equation 14) of the observed ETGs from Dabringhausen & Fellhauer (2016) but that it falls below the observed relation by a factor of 4.

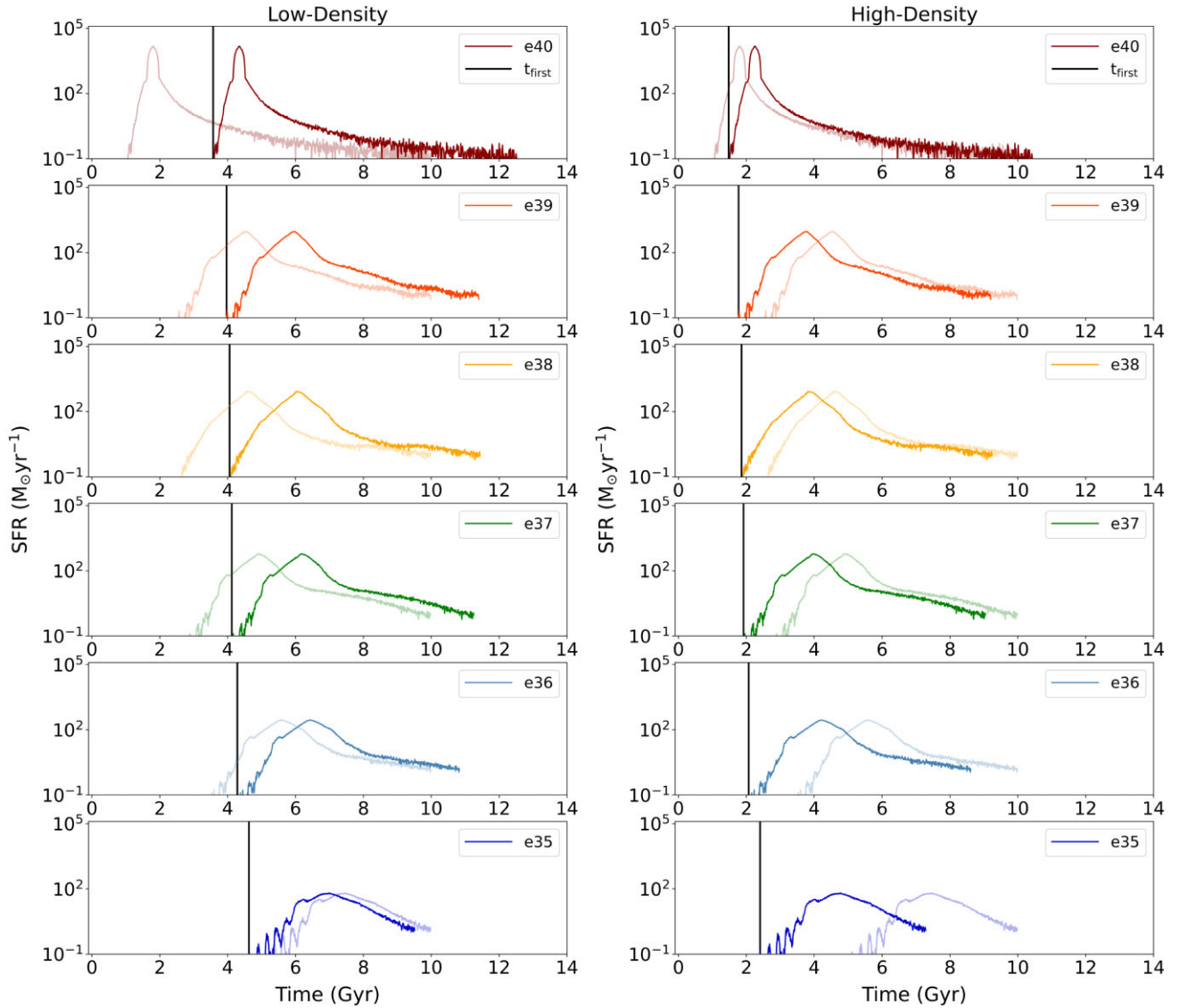


Figure 6. A plot showing the SFH of model e40, e39, e38, e37, e36, and e35 (Table 1) in the low-density (*left*) and high-density (*right*) environment. The black line is the time when the first star is formed in the models. Each model has been shifted in time according to the Thomas et al. (2005) age dating (equations 11 and 12). The nearly transparent SFHs are relative to the start of the computation at $t = 0$.

It is possible that the monolithic collapse that forms individual compact ETGs occur in the inner regions of local density enhancements that form neighbouring gas clouds which collapse monolithically, forming binary or even more ETGs (similar to collapse of stars in a very dense configuration, e.g. Joseph & Wright 1985; Wright et al. 1988; Stolte et al. 2014). If these are on radial orbits they would merge, leaving a merger remnant which is likely to have a larger effective-radius (for further empirical evidence that ETGs formed first at their innermost regions, see section 2.2 in Kroupa et al. 2020). Trujillo et al. (2007) study an effective size evolution of compact galaxies where they find that the massive ETGs were formed in the early Universe, and have subsequently evolved passively until today. They find an effective-size evolutionary mechanism would be able to evolve their compact galaxies to the observed local relation with just two major mergers. To investigate this, we study the merger scenario where two model galaxies listed in Table 1 are placed 100 kpc apart in a box of 1 Mpc. The models only consist of stellar particles. The spin angular momentum of each model galaxies is retained as they

are placed in the new box and they merge under MOND gravity to form a new merged model galaxy in 1.5 Gyr. Table 2 lists the initial conditions and results of the merged models. Five different merger studies were done placing the models at different axes and at different orientation. M1, M2, M3, and M4 are merged models formed from the merger of the same model galaxies and M5 is formed by the merger of model galaxies of different masses (e39 and e36) that merge at different angles with respect to the XY plane (where parent models e36 and e39 are placed tilted at an angle of 45 deg with respect to the XY plane). The effective-radius–mass relation of these merged models is plotted in Fig. 7. It can be seen that the merged models lie around the same dashed black line as the real ETGs in Fig. 7. A detailed study of the morphology and other structural properties of these merged galaxies will be discussed in a follow-up paper (Eappen et al. in preparation).

Another possible mechanism for the model galaxies to reach large radii is through stellar evolution mass-loss, which is directly linked to the galaxy-wide initial mass function (IMF) assumed. It has been

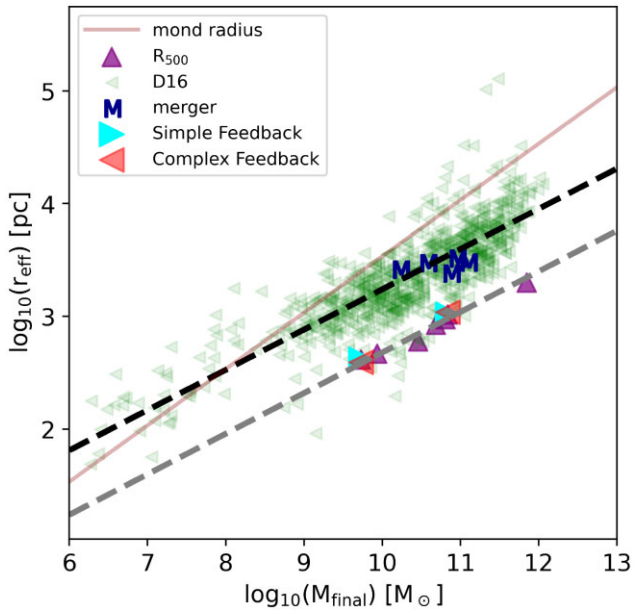


Figure 7. The effective-radius–mass relation. The green triangles are the real ETGs from Dabringhausen & Fellhauer (2016), purple triangles are the model galaxies e34 to e40 with $R_{\text{initial}} = 500$ kpc listed in Table 1 and the blue M’s are the merged models listed in Table 2. The black dashed line is the best fit for the real galaxies (equation 14), and the grey dashed line is the best fit for the model galaxies (equation 15). The solid red line is the MOND radius (equation 13) and the red left-pointing triangles are the models with complex feedback (see Section 4.3).

Table 2. Initial conditions of the merger models.

Model Name	Parent models	r_{eff} (kpc)	M_{final} (M_{\odot})
M1	e35 + e35	2.6	0.17E11
M2	e36 + e36	2.9	0.38E11
M3	e37 + e37	3.2	0.9E11
M4	e39 + e39	3.0	1.25E11
M5	Re36 + Re39	2.4	0.75E11

Notes. Column 1: name; column 2: the models from Table 1 that merge; column 3 (r_{eff}): the effective-radius of the merged model; and column 4 (M_{final}): the mass of the merged model. Models Re36 and Re39 are models e36 and e39 but rotated (see the text for details).

shown that the stellar remnant population is comparable to the mass of living stars in the IGIMF theory, but using a canonical invariant, IMF would decrease the remnant mass to be much lower than the stellar mass (Yan et al. 2019). Yan et al. (2021) finds that the ratio of dynamical mass and stellar mass, M_{dyn}/M_{\star} , is higher for more massive galaxies, because they have a more top-heavy galaxy-wide IMF and more stellar remnants. Given the high SFRs ($> 10^2 M_{\odot} \text{ yr}^{-1}$) the model galaxies reach (Table 1), it is expected that the IGIMF is top heavy (Yan, Jerabkova & Kroupa 2017; Jeřábková et al. 2018; Yan et al. 2021), leading to significant mass loss from the evolving stars. This mass lost would be in the form of gas and if it is heated sufficiently enough it would expand and lead to the expansion of the galaxy (Suzuki et al. 2022).

4.3 Effects of complex feedback mechanism

The above results (except for models e34c and e39c) have been obtained using the simple feedback algorithm (simple cooling/heating and no additional more complex baryonic physics, such as supernovae). Models e34c and e39c were computed with additional feedback process to investigate how complex baryonic physics affects the results obtained.

The simple star formation algorithm only considers the heating/cooling modules in POR. Cooling and heating of the gas are computed by using tables that are included in the code, which describe the cooling and heating rates due to several physical processes. POR uses Courty & Alimi (2004) look-up tables for the different cooling/heating processes. The complex feedback algorithm includes more initial inputs such as supernovae, radiative transfer, sink particles, and a higher refinement level which increases the resolution of the simulation. The supernova (and in fact all feedback scheme) scheme is implemented as in RAMSES where the kinetic part of the supernova energy is injected as a spherical blast wave with the size of galactic superbubbles of radius 150 pc. The radiative transfer option allows to compute the radiative transfer between stellar particles with the addition of photon fluxes which uses a first-order Godunov solver (Rosdahl et al. 2013). A sink particle scheme is used to stabilize the simulation by artificially stopping the collapsing gas cloud if a certain density is exceeded and the gas is condensed into a point mass (sink particle, see Teyssier 2002; Lüghausen et al. 2015; Wittenburg et al. 2020; Nagesh et al. 2021 for a detailed description on how RAMSES and POR handles feedback).

The model e39c has the same R_{initial} and M_{initial} as model e39 but is computed with complex baryonic feedback of supernovae, sink particles, and radiative transfer with a higher resolution (maximum refinement level set to 13 instead of 12). Similarly, model e34c has the same initial conditions as model e34 but is computed with a higher resolution (maximum refinement level set to 14 instead of 12). Fig. 8 shows the SFH of the models with simple and complex feedback. Table 3 shows the effect of feedback on the models which have the same R_{initial} and M_{initial} . The percentage change (PC) is calculated using equation (D1) in Appendix D and the results obtained (Table 3) show that the percentage change of the results from the complex feedback is within 5 per cent of the results obtained from simple feedback. The complex feedback is significantly more CPU intensive but is found to affect the results negligibly (see Fig. 3, 7, and 8 and Table 3), as already shown in the disc–galaxy formation simulations by Wittenburg et al. (2020). This demonstrates again that differences in the feedback processes play a minor role in the evolution of galaxies and that the physics of galaxy evolution are pre-dominantly defined by the Milgromian law of gravitation (Kroupa 2015).

5 DISCUSSION

The monolithic post-Big-Bang collapse of non-rotating gas clouds in MOND results in a SFT–mass relation (equation 10) which is comparable to the observed relation for ETGs (Thomas et al. 2005, 2010). Collapsing post-Big-Bang gas clouds in Milgromian gravitation thus behave as the observed ETGs in terms of the collapse timing (i.e. more massive clouds collapse faster and form their stellar particles more quickly as shown in Fig. 2). The grid of models with different R_{initial} and M_{initial} allows us to constrain $R_{\text{initial}} \approx 500$ kpc as best fitting the observed SFT–mass relation. This suggests that typically a region spanning almost a Mpc across collapsed to form ETGs if sufficiently massive. We note that field elliptical galaxies are slightly younger (consistent with the results found by Thomas

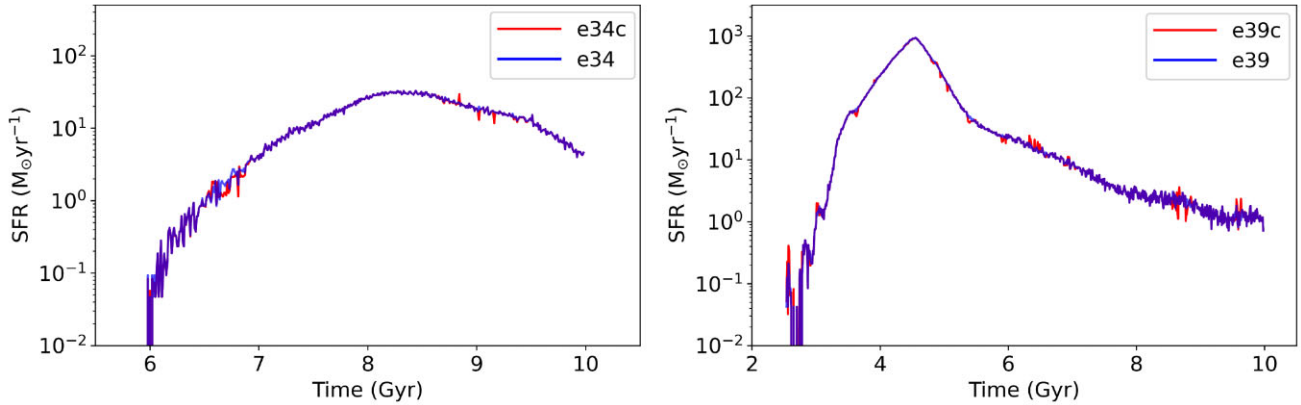


Figure 8. A plot comparing the SFH of models with simple and complex feedback. Models e34c and e39c are models with complex feedback.

Table 3. Effects of feedback.

	Simple feedback	Complex feedback	
	e34	e34c	PC
r_{eff}	0.41	0.39	-4.87%
$\Delta\tau_m$	1.41	1.43	+1.41%
	e39	e39c	PC
r_{eff}	1.04	1.09	+4.80%
$\Delta\tau_m$	0.54	0.53	-1.85%

Note. PC shows the percentage change (equation D1) in the results obtained from models with same R_{initial} and M_{initial} but with simple (e34 & e39) and complex (e34c & e39c) feedback.

et al. 2005) but otherwise do not differ significantly from elliptical galaxies in clusters (Saracco et al. 2017). This suggests that also in the field the bulk stellar population of elliptical galaxies formed through monolithic collapse on the downsizing time-scale.

How do these results compare to the expectations from the standard hierarchical structure formation theory as quantified through the Illustris TNG (Pillepich et al. 2018a; Nelson et al. 2018, 2019a) and Eagle (Schaye et al. 2015; McAlpine et al. 2016) projects? It is noted that Fontanot & Monaco (2010) write: ‘The so-called downsizing trend in galaxy formation (e.g. more massive galaxies forming on a shorter timescale and at higher redshift with respect to lower mass counterparts) is not fully recovered.’ On the other hand, based on their analysis of the Millenium Simulation, De Lucia et al. (2006) claim ‘These findings are consistent with recent observational results that suggest “down-sizing” or “anti-hierarchical” behaviour for the star formation history of the elliptical galaxy population, despite the fact that our model includes all the standard elements of hierarchical galaxy formation and is implemented on the standard, Λ CDM cosmogony.’ Given these discrepant results based on the same theory and observational data, we evaluate the formation time-scales of ETGs that form in the Illustris TNG and EAGLE projects (see Figs B1 and B2), noting that both of these rest on very different simulation codes (AMR versus SPH, respectively) and also very different baryonic physics algorithms. The details are provided in Appendix B and the result is, in summary, that the formation time-scales of ETGs formed in both of these simulation projects do not concur with the observed population.

The MOND radius (equation 13) proposed by theory is larger but follows the same trend as the observed r_{eff} . This similarity of the MOND radius and the effective-radius was first noted by Sanders (2008), where he concludes that a dissipationless initially expanding gas cloud collapse would naturally form galaxies with effective-radii

that are comparable to their MOND radii. The models in our study do not consider an expanding gas cloud collapse and fall short of this expected effective-radius–mass relation (see Fig. 7), but this is possibly also due to a non-variant galaxy-wide IMF which makes the post-Big-Bang clouds to collapse too deeply. If, on the other hand, the galaxy-wide IMF varies such that it becomes top-heavy at a high SFR (e.g. Jeřábková et al. 2018; Yan et al. 2021) then the collapse is likely less deep and r_{eff} will end up larger. We do not address the chemical evolution of the models here as the chemical enrichment has not yet been incorporated into POR and would also require a systematically varying galaxy-wide IMF (the current simulations effectively assume the canonical IMF for the formation of stellar particles). The SFHs of the models presented here are however comparable, by virtue of them matching the formation timescales, with the constraints from abundance and stellar population observations as analysed by Yan et al. (2021). It is to be expected that the bulk chemical enrichment of the models computed here would follow the closed box models of Yan et al. (2021). The problem of spatially-resolved chemical enrichment of the forming model ETGs will be studied in the future in connection with the systematically varying galaxy-wide IMF.

The model galaxies in this work do not agree with the SFT–mass relation found in Yan et al. (2021) (dotted green curve in Fig. 3). This is possibly due to the non-variant IMF that is assumed in the simulation as mentioned above. The other method through which the model galaxies could reach the Yan et al. (2021) SFT–mass relation is if the gas cloud collapse calculations are considered in an expanding cosmological volume which will be investigated in the future.

In Fig. 7, we show that mergers of the model galaxies form models that lie on the observationally deduced effective-radius–mass relation. These kinds of mergers could happen in the initially densest parts of the Universe, for e.g. in central regions of galaxy clusters, where the galaxies would have initially formed through monolithic collapse of post-Big-Bang gas clouds and then slightly later merge with another galaxy of comparable mass in a time-scale somewhat larger than the formation timescale of these galaxies to reach the observed effective-radius–mass relation. The morphology and structural properties of the model galaxies and the mergers using different projections will be discussed in a follow-up paper (Eappen et al. in preparation) where we find that the shape of these merged models is triaxial and very similar to the observed ETGs. Thus, the formation time-scales and other observed chemical properties of ETGs could have been embedded in them before they undergo one or two major mergers such that only the structural properties and the morphology of these ETGs changed through the mergers. Trujillo et al. (2007) report a similar result in their study of size evolution

of massive galaxies at different redshifts. They find that at $z \approx 1.5$ massive spheroid like objects were a factor of 4 smaller than we see today and suggest that just two major mergers of equal mass would evolve the size of the galaxy to what we observe today. Similarly, Pipino & Matteucci (2008) in their study of ETGs through chemical evolution models find that a series of dry mergers of galaxies is not the way to recover downsizing and suggest 1–3 major-dry mergers to be in agreement with the observations. Our results show that it is possible to evolve the size and mass of the galaxy without changing the intrinsic properties of the galaxy such as the SFT without a large number of mergers.

6 CONCLUSION

In this work, we discuss the formation of ETGs in comparison with observational results. The isolated monolithic collapse of post-Big-Bang clouds of different initial radii at fixed mass is computed in Milgromian gravitation and is found to produce model galaxies with SFTs similar to the observed ETGs.

The main results are as follows:

(i) This work for the first time shows that the SFTs observed for ETGs is a natural occurrence in the MOND paradigm if ETGs form from monolithically collapsing non-rotating post-Big-Bang gas clouds. We were able to reproduce the SFTs as noted by observations (Fig. 3). The models with $R_{\text{initial}} = 500$ kpc are similar to the relation of Thomas et al. (2005, 2010) and the models with $R_{\text{initial}} = 200$ kpc fit fairly closely to the relation of Recchi et al. (2009). This work thereby gives a rough estimate of the initial conditions of the post-Big-Bang clouds in order to produce ETGs. The results indicate that the observed Universe started to form ETGs earlier than the here studied collapse models (assuming $t = 0$ is the Big-Bang), although the SFTs of the models match those observed for gas clouds with initial radii in the range of 200–500 kpc. Higher resolution simulations and different cosmological expansion models will be needed to address this issue. Lambda cold dark matter (Λ CDM) models cannot account for the observed time-scales (Figs B1 and B2).

(ii) The effective-radius–mass relation of the model galaxies falls short of the expected value but the models with only one major merger (Fig. 7) are found to be consistent with observations (Dabringhausen & Fellhauer 2016).

(iii) We find that different feedback processes affect the results negligibly (as already shown by Wittenburg et al. 2020 in their formation of disc-galaxies), suggesting that the physics of galaxy formation is pre-dominantly defined by the Milgromian law of gravitation (Kroupa 2015).

The future aspects of this project will analyse the structural properties and shapes of both the model galaxy and merged galaxies, using different projections to calculate the sizes (Eappen et al. 2022c, in preparation). Then the next step would be to investigate how a possible inclusion of a systematically varying galaxy-wide IMF (Yan et al. 2017) would lead to the expansion of the models also taking into account the cosmological expansion of space.

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ANR-19-CE31-0017), and from the European Research Council (ERC) under the European Union’s Horizon 2020 Framework programme (grant agreement number 834148). We thank the DAAD Eastern European grant at Bonn University for supporting the research visits. The IllustrisTNG project used in this work have been run on the HazelHen Cray XC40system at the High Performance Computing Center Stuttgart as part of project GCS-ILLU of the Gauss centres for Supercomputing (GCS).

DATA AVAILABILITY

All data used here have been generated as described using the publicly available POR code and the cited literature.

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Table A1. Fit parameters for equation (10).

R_{initial} (kpc)	A	B
50	40747.40	-0.55
100	6102.39	-0.43
200	2989.50	-0.37
300	227.87	-0.25
500	6253.16	-0.37

APPENDIX A: FIT PARAMETERS

The fit parameters for equation (10) are provided in Table A1.

APPENDIX B: SFTS FOR Λ CDM SIMULATIONS

Using the TNG50-1 and TNG100-1 simulation of the Illustris The Next Generation (TNG) project (Nelson et al. 2018, 2019a, b; Pillepich et al. 2018a, b, 2019; Weinberger et al. 2017; Marinacci et al. 2018; Naiman et al. 2018; Springel et al. 2018) and the Ref-L100N1504 simulation (hereafter EAGLE100) of the EAGLE project (Crain et al. 2015; Schaye et al. 2015; McAlpine et al. 2016; The EAGLE team 2017), the star formation time-scales, SFTs, of the formed model galaxies are investigated within the Λ CDM framework. These projects are sets of different self-consistent cosmological galaxy formation simulations. The TNG project is consistent with the Planck-2015 (Planck Collaboration XIII 2016) results: $H_0 = 67.74 \text{ km s}^{-1} \text{ Mpc}^{-1}$, $\Omega_{\text{b},0} = 0.0486$, $\Omega_{\text{m},0} = 0.3089$, $\Omega_{\Lambda,0} = 0.6911$, $\sigma_8 = 0.8159$, and $n_s = 0.9667$. Here, the TNG50-1 has a box side of 51.7 co-moving Mpc (cMpc) an initial baryonic mass resolution of $m_{\text{b}} = 8.5 \times 10^4 M_{\odot}$, and a dark matter particle mass of $m_{\text{dm}} = 4.5 \times 10^5 M_{\odot}$. The TNG100-1 simulation has a box side of 110.7 co-moving Mpc (cMpc), $m_{\text{b}} = 1.4 \times 10^6 M_{\odot}$, and $m_{\text{dm}} = 7.5 \times 10^6 M_{\odot}$. An overview of the physical and numerical parameters is given in table 1 of Nelson et al. (2019). The EAGLE project assumes the Planck-2013 (Planck Collaboration I 2014) measurements with the cosmological parameters being $H_0 = 67.77 \text{ km s}^{-1} \text{ Mpc}^{-1}$, $\Omega_{\text{b},0} = 0.04825$, $\Omega_{\text{m},0} = 0.307$, $\Omega_{\Lambda,0} = 0.693$, $\sigma_8 = 0.8288$, and $n_s = 0.9611$ (see also table 1 of Schaye et al. 2015). The EAGLE100-1 run has a box size of 100 cMpc, an initial baryonic particle mass of $m_{\text{b}} = 1.81 \times 10^6 M_{\odot}$, and a dark matter particle mass of $m_{\text{dm}} = 9.70 \times 10^6 M_{\odot}$ (table 2 of Schaye et al. 2015).

We select present-day ($z = 0$) subhaloes with a stellar mass $M_{\star} > 10^9 M_{\odot}$ and a SFR of $\text{SFR} < 10^{-3} M_{\odot} \text{ yr}^{-1}$. In addition, we remove subhaloes with a subhaloflag value of 0 in the TNG simulation, ensuring that only galaxies or satellites of a cosmological origin are considered. This gives a final sample of 418 (45 central galaxies), 4384 (803 central galaxies), and 2916 (373 central galaxies) galaxies in the TNG50-1, TNG100-1, and EAGLE100 run, respectively.

The SFT in the cosmological Λ CDM simulations is calculated by using three different methods. In the first approach, we extract the formation time and the masses of the stellar particles at $z = 0$ of the above selected galaxies. The SFT, $\Delta\tau$, for the TNG100-1 and EAGLE100 data is calculated using the same method as described in Section 3, where the FWHM of the SFH gives the duration for which the bulk of the star formation takes place. The SFR is calculated by separating all stellar particles in bins of $\delta t = 10 \text{ Myr}$ according to their age, summing up the stellar mass in every time bin and dividing this by the length of δt . Their SFT–mass relation is shown in the top panels of Fig. B1.

In the second method, we extract again the stellar particles at $z = 0$ of the selected galaxies but using the mass of the star particle at the

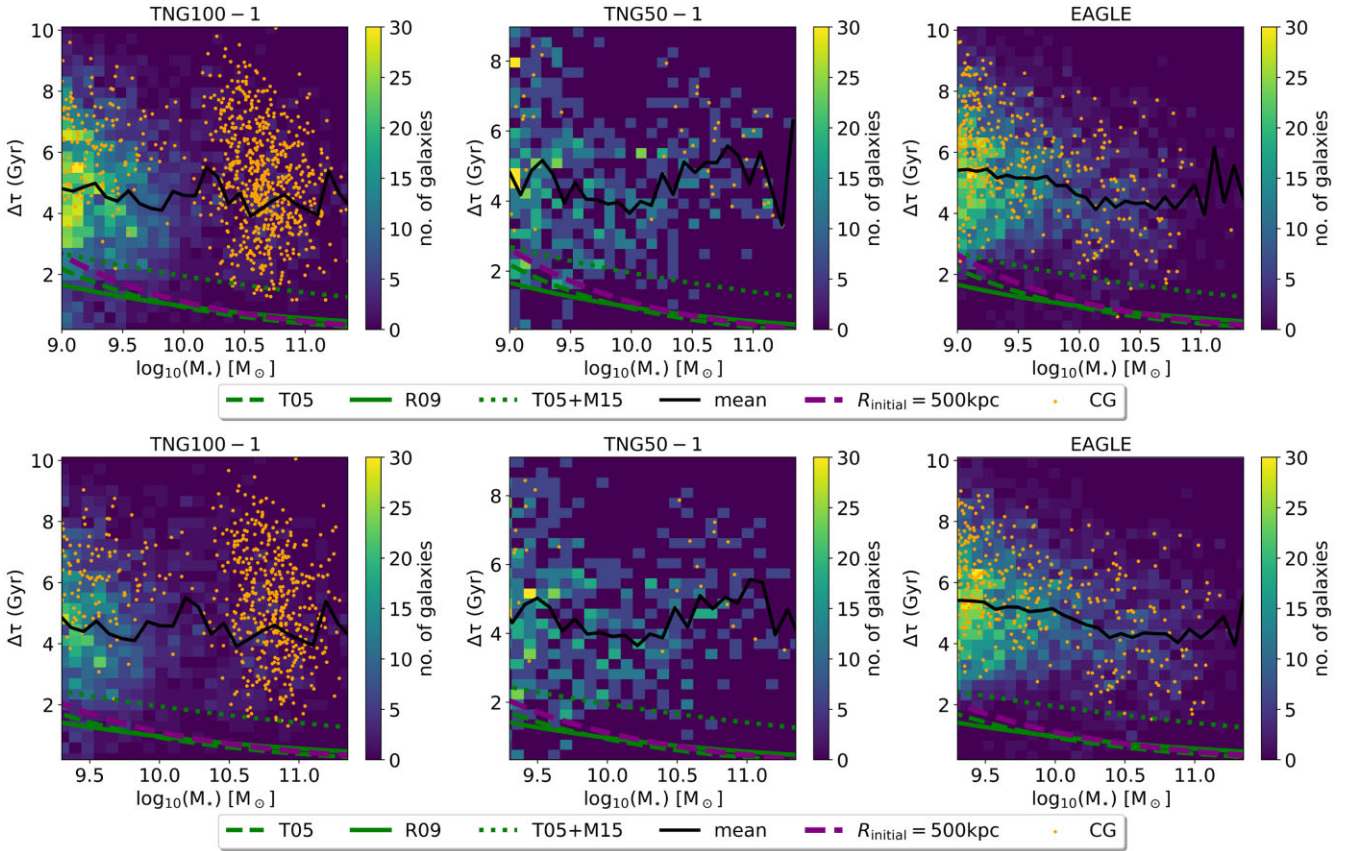


Figure B1. SFT–mass relation of model galaxies formed in the TNG100-1 (*left*), TNG50-1 (*centre*), and EAGLE100 (*right*) (SMoC) simulation by using the method 1 (top panels) and method 2 (bottom panels) compared with observationally constrained relations for ETGs. The orange coloured dots labelled as CG are the central galaxies in each simulation. There are a total of 30 bins on the x- and y-axis and the binsize is 0.08 dex.

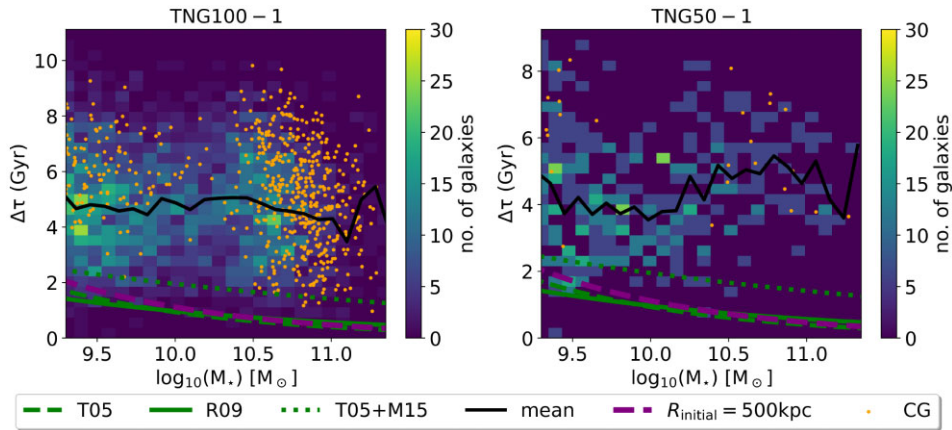


Figure B2. Same as Fig. B1 but using method 3 as described in Appendix B.

time of its formation in order to calculate the SFT (bottom panels of the Fig. B1). We use the same method as described above to calculate the SFRs and obtain the SFT–mass relation.

The mean of $\Delta\tau$ for all three simulation runs lies around 5 Gyr and a trend in M_* is absent for both methods (except for EAGLE) but is in complete disagreement to observed ETGs (Thomas et al. 2005; Recchi et al. 2009; McDermid et al. 2015). The Wilcoxon Signed-Rank test (Bhattacharyya & Johnson 1977) shows the SMoC models to disagree with the observed SFTs at more than 5σ confidence

because the latter lie consistently and significantly below the former.

In the third method, we trace the selected galaxies back in cosmic time by using the merger tree catalogues (Rodríguez-Gomez et al. 2015) and extract their SFR from the Subfind Subhalo catalogues at different snapshots. For this, we only select subhaloes with a dark matter mass of $M_{\text{dm}} > 0$ in order to guarantee that these galaxies are included in the merger trees catalogues. This gives 417 and 4296 subhaloes in the TNG50-1 and TNG100-1 run, respectively. Their SFT–mass relation is shown in Fig. B2.

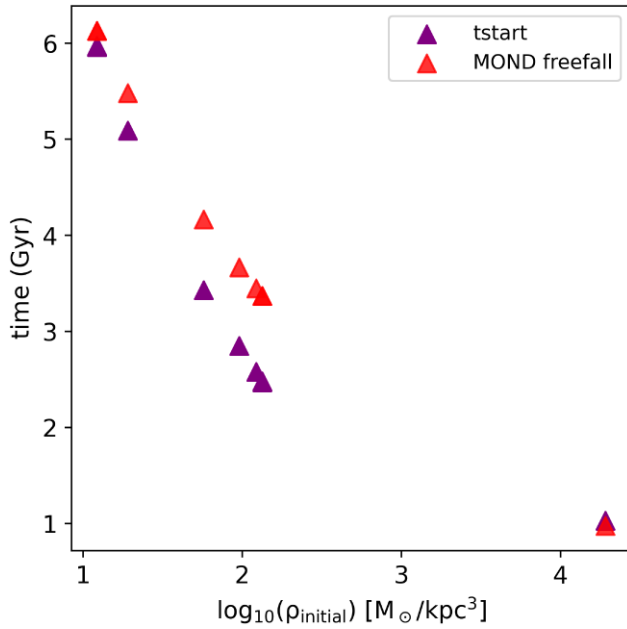


Figure C1. Collapse time–cloud density relation for all models with $R_{\text{initial}} = 500$ kpc. The purple triangles are the time when the first stellar particle is formed in the simulation, t_{start} , and the red triangles are the MOND free-fall time (equation 6).

The third method cannot be applied for the EAGLE run because only 29 snapshots outputs (see table C.1 of McAlpine et al. 2016) are publicly available and a fit cannot be made to get the SFT.

Given the test performed here which relies on quantifying the present-day ($z = 0$) galaxies in the cosmological simulations, the results unambiguously show that the models do not have the

observed downsizing. This is in contradiction to other Λ CDM models (see section 2.3 of Wright et al. 2019) where they use their own quantification of measuring quenching times of model galaxies. All 3 methods result in a 5σ tension (Figs B1 and B2).

APPENDIX C: COLLAPSE TIME

The MOND free-fall time is the time taken for the gas cloud to collapse under its own gravitational attraction to form stars. This is calculated using equation (6) (see Zonoozi et al. 2021). We compare the time taken for the first stellar particle to form in the simulation, t_{start} , with this theoretical value (Fig. C1). The t_{start} , and the MOND free-fall time are comparable in our simulation demonstrating consistency between the numerical and analytical results.

APPENDIX D: PERCENTAGE CHANGE

The percentage change (PC) quantifies the change from an initial value to the new value and expresses the change as an increase or decrease,

$$\text{PC} = \frac{(X_2 - X_1)}{|X_1|} \times 100. \quad (\text{D1})$$

Here, X_1 is the initial value and X_2 is the new value. If the PC obtained is a positive value then it is quantified as an increase from the initial value and if the obtained PC is negative then it is quantified as a decrease from the initial value. In Table 3, the results obtained for the simple feedback (SF) models are used as the initial values and the results obtained for the complex feedback (CF) are used as the new values.

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