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# Geophysical Research Letters®

## RESEARCH LETTER

10.1029/2022GL098111

### Key Points:

- First identification of  $\text{H}_2^+$  in the Jovian magnetosphere from 13 to 18  $R_J$ , with a density ratio of  $\text{H}_2^+/\text{H}^+ = 8 \pm 4\%$
- $\text{H}_2^+$  in Jupiter's magnetosphere is predominantly produced by Europa's neutral toroidal cloud
- Europa's total  $\text{H}_2$  neutral loss rate is  $1.2 \pm 0.7 \text{ kg s}^{-1}$

### Supporting Information:

Supporting Information may be found in the online version of this article.

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## Water-Group Pickup Ions From Europa-Genic Neutrals Orbiting Jupiter

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**Abstract** Water-group gas continuously escapes from Jupiter's icy moons to form co-orbiting populations of particles or neutral toroidal clouds. These clouds provide insights into their source moons as they reveal loss processes and compositions of their parent bodies, alter local plasma composition, and act as sources and sinks for magnetospheric particles. We report the first observations of  $\text{H}_2^+$  pickup ions in Jupiter's magnetosphere from 13 to 18 Jovian radii and find a density ratio of  $\text{H}_2^+/\text{H}^+ = 8 \pm 4\%$ , confirming the presence of a neutral  $\text{H}_2$  toroidal cloud. Pickup ion densities monotonically decrease radially beyond 13  $R_J$  consistent with an advecting Europa-genic toroidal cloud source. From these observations, we derive a total  $\text{H}_2$  neutral loss rate from Europa of  $1.2 \pm 0.7 \text{ kg s}^{-1}$ . This provides the most direct estimate of Europa's  $\text{H}_2$  neutral loss rate to date and underscores the importance of both ion composition and neutral toroidal clouds in understanding satellite-magnetosphere interactions.

**Plain Language Summary** Jupiter's moons Europa, Ganymede, and Callisto all have icy surfaces which interact with their local environments. From this interaction, water-group atoms and molecules are released from the icy surfaces and orbit Jupiter as a collection of material in “neutral toroidal clouds.” The material in these toroidal clouds interact with the local charged particle environment, where neutrals in the toroidal clouds can become charged and incorporated into Jupiter's charged particle environment. Here, we highlight observations of Jupiter's charged particle environment and present the first detections of  $\text{H}_2^+$  in this environment. These  $\text{H}_2^+$  ions are shown to be originally produced from  $\text{H}_2$  lost from Europa and the abundance of detected ions allows us to determine Europa is losing  $1.2 \pm 0.7 \text{ kg s}^{-1}$  of neutral  $\text{H}_2$ . This provides the most direct estimate of Europa's  $\text{H}_2$  neutral loss rate to date and underscores the importance of both ion composition and neutral toroidal clouds in understanding how satellites interact with their local charged particle environments.

## 1. Introduction

Each of the outer three Galilean satellites Europa, Ganymede, and Callisto have surficial ice and atmospheric water group species ( $\text{H}$ ,  $\text{H}_2$ ,  $\text{O}$ ,  $\text{OH}$ ,  $\text{H}_2\text{O}$ , and  $\text{O}_2$ ). These species are abundant to varying degrees due to differing surface compositions and atmospheric source/loss processes (e.g., Ip, 1996; Johnson, 1990, 2004; Marconi, 2007; McGrath et al., 2004; Roth, 2021; Saur et al., 1998; Smyth & Marconi, 2006), where atmospheric losses could sustain neutral toroidal clouds. Of these three satellites, only Europa's neutral toroidal cloud has been inferred from observations (Lagg et al., 2003; Mauk et al., 2003) and is assumed to be produced by sublimation and sputtering from its icy surface material (Smyth & Marconi, 2006). Depending on the lifetimes of orbiting neutrals, these clouds may or may not completely populate all longitudes throughout their parent body's orbits. Hence, we refer to them as “neutral toroidal clouds” in lieu of the historically common “neutral tori” terminology.

As observed by the Galileo spacecraft, energetic neutral atoms originating just outside Europa's orbit, born via charge exchange between a cold neutral toroidal cloud and energetic charged particle environment, suggest the densities in Europa's neutral gas toroidal cloud are comparable to Io's neutral gas toroidal cloud (Mauk

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et al., 2003, 2020). Unlike near Io's orbit, however, where the dense Io plasma population quickly ionizes neutrals, the less dense plasma environment at Europa and beyond is more conducive to longer neutral lifetimes (Bagenal & Dols, 2020). Separately, observations of depletions in energetic protons that would primarily reside at low magnetic latitudes (pitch angles near  $90^\circ$ ) were proposed to be due to charge exchange with the very same neutral Europa toroidal cloud (Lagg et al., 2003). Neither observation was able to directly determine the composition of the neutral cloud, however, ultraviolet observations ruled out a dominantly oxygen-rich cloud (Hansen et al., 2005). Europa's neutral toroidal cloud was subsequently modeled using multiple atmospheric constraints, and a cloud predominantly composed was proposed to explain both indirect cloud observations (Smith et al., 2019; Smyth & Marconi, 2006).  $H_2$  is understood to have the highest escape rate because of its enhanced surface abundance from radiolysis, large-scale height, light mass, and non-condensability as a diatomic species from Europa (Smyth & Marconi, 2006), Ganymede (Marconi, 2007), and Callisto (Mogan et al., 2021).

While remote observations did not detect an oxygen-rich toroidal cloud and subsequent modeling efforts have suggested a dominantly  $H_2$ -rich Europa cloud,  $H_2$  had yet to be definitively measured and direct constraints on its abundances or loss rates were not possible given the prior observations. Here, we present direct observations of  $H_2^+$  pickup ions sourced from a neutral  $H_2$  toroidal cloud in the Europa-Ganymede region. We use these observations to constrain the total abundances and loss rate of neutral  $H_2$  from Europa. In Section 2, we discuss how and where  $H_2^+$  pickup ions are observed from the Juno spacecraft. In Section 3, we use these observations to determine numerical densities of both  $H^+$  and  $H_2^+$  and compare these to a plasma advection model to estimate loss rates from Europa in Section 4. We conclude in Section 5 with a discussion on implications for loss rates from the other Galilean satellites and future work that could build on these results.

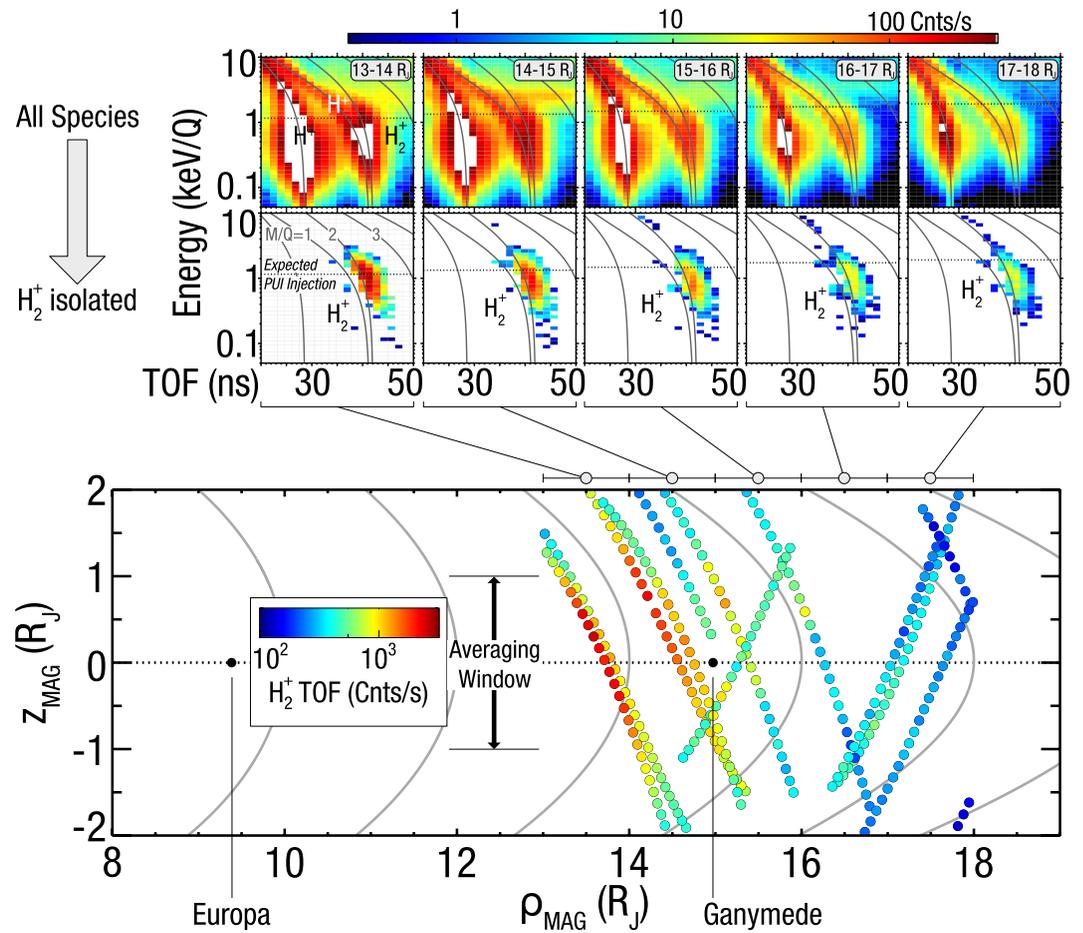
## 2. Observations of $H_2^+$

The Juno mission (Bolton et al., 2017) carries the Jovian Auroral Distributions Experiment (JADE; McComas et al., 2017), a suite of plasma instruments that includes an ion time-of-flight (TOF) mass spectrometer. JADE's ion sensor measures the flux of ions with energy-per-charge of 10 eV/Q to 46 keV/Q and a mass-per-charge of  $\sim 1$ –64 AMU/Q. Juno's eccentric, nearly polar orbit precesses such that its furthest distance (apojove) is sequentially more southward after each close approach (perijove). Therefore, the location where Juno intersects Jupiter's equatorial plane moves inward each orbit, enabling ion composition observations by JADE's TOF mass-spectrometry of the equatorial plasma environments at varying radial distances (Kim et al., 2020a, 2020b).

By late-2021 up to and including its 37th perijove, Juno transited the equatorial plasma environments at the orbital distances of Europa, Ganymede, and Callisto. A JADE signal consistent with  $H_2^+$  is present near the equatorial region inside  $\sim 20 R_J$  ( $1 R_J \equiv 71,492$  km). Isolating  $H_2^+$  in the JADE data set requires removing multiple background and foreground sources that overlap with the  $H_2^+$  signal in the combined TOF versus energy space. Notably, the  $H_2^+$  feature overlaps with  $H^+$  in this space, where  $H^+$  is more abundant, hence correctly subtracting the proton foregrounds is critical. We apply three background subtractions, and additionally remove the signature of  $H^+$  foreground to isolate the count rates due to  $H_2^+$ . To remove the  $H^+$  signal, we use an empirical response from a unique event where Juno was connected to Io's Main Alfvén Wing (Szalay et al., 2020). During this period, large proton abundances were observed and energized throughout the entire observable range across multiple Juno instruments (Clark et al., 2020; Sulaiman et al., 2020) enabling a reliable determination of the in-flight response to protons alone. The subtraction schemes are described in detail in Supporting Information S1.

We manually identified regions where  $H_2^+$  could be reasonably separated, producing a conservative subset of JADE data with high-fidelity  $H_2^+$  data. In the region past  $10 R_J$  near the magnetic equator, an unambiguous  $H_2^+$  signature is observed from  $\sim 13$  to  $18 R_J$ . Outside these locations,  $H_2^+$  signatures are still observed, but cannot be fully separated from foregrounds and backgrounds with the current techniques.

Figure 1 shows these  $H_2^+$  TOF observations, where each of the top TOF spectra are the superposed average count rate spectra within  $\pm 1 R_J$  from the magnetic equator for each  $1 R_J$  radial bin shown in the bottom panel. Times



**Figure 1.** Count rates of  $\text{H}_2^+$  from 13 to 18  $R_J$  (top) Ion energy per charge as a function of time-of-flight (TOF) for all species with nominal background subtraction (upper row) and with  $\text{H}_2^+$  isolated (lower row) for all observations within an averaging window of  $\pm 1 R_J$  from the magnetic equator.  $\text{H}_2^+$  is unambiguously observed when foregrounds and backgrounds are removed, in the vicinity of the second gray TOF trace for a mass per charge of 2. The horizontal line indicates the expected detected energy for a freshly created pickup ion at the spacecraft's position (bottom) Locations of each observation in cylindrical coordinates with the  $+z$  axis are aligned with Jupiter's magnetic dipole.  $\text{H}_2^+$  is observed with the strongest magnitude in the inner-most detected locations.

used in these combined TOF spectrograms are given in Table S1 in Supporting Information S1. Without background and foreground  $\text{H}^+$  subtraction, the counts along the  $M/Q = 2$  track are dominated by  $\text{H}^+$  and to varying degrees, a long TOF tail from heavy ions that manifests as a horizontal bar around 2–4 keV notable in the second TOF spectrogram (14–15  $R_J$ ). As indicated in Figure 1, which shows the expected energy and TOF tracks for different species, protons are detected along two separate tracks due to the internal workings of the instrument. See Text S1 in Supporting Information S1 for additional description on how this is accounted for in the  $\text{H}_2^+$  identification scheme.

The  $\text{H}_2^+$  count rates occur near and mostly below the expected local corotation energy (Text S2 in Supporting Information S1), consistent with a population of pickup ions corotating with Jupiter's magnetosphere sourced from a neutral gas moving slowly with respect to local corotation speed. Pickup ion fluxes are expected to peak in the detected energy range of 100s eV to a few keV when Juno is within 10–30  $R_J$ . Therefore, the instrument's energy range of 10 eV/Q–46 keV/Q is sufficient to capture any pickup ions generated in this region. The bottom panel of Figure 1 shows where these observations occurred along with magnetic field lines every 2  $R_J$  in equatorial distance using the JRM09 + CAN models (Connerney et al., 1981, 2018). As shown here, the largest  $\text{H}_2^+$  count rates are observed near the magnetic equator and closest to Jupiter over this distance range.

### 3. Numerical Densities

While the TOF data set does not have any information on directionality, JADE observed the full-sky over each  $\sim 30$ s spin and we can calculate a partial numerical density from the count rates as a function of energy over the JADE energy bandpass. After all foregrounds and backgrounds are subtracted, we sum count rates over all TOFs corresponding to AMU/Q between 1.5 and 2.5 to determine a total count rate  $R_{obs}$  as a function of energy. JADE instantaneously observes an angular range of  $270^\circ$  extending from the anti-sunward spin axis, such that for each spacecraft rotation it records counts from a total angular extent of  $6\pi$  sr, double-counting half the sky. We must reduce  $R_{obs}$  by an appropriate factor to determine the “true” average count rate  $R = \eta R_{obs}$  corresponding to the  $4\pi$  sr full sky (Text S3 in Supporting Information S1). Due to the instrument mounting and orbit geometry, for each  $H_2^+$  observation by JADE, plasma corotating with Jupiter is predominantly observed on the hemisphere where JADE double-counts incident populations, which also gives improved counting statistics. Table S1 in Supporting Information S1 gives values of  $\eta$ ; for nearly all data analyzed here,  $\eta \approx 0.5$ .

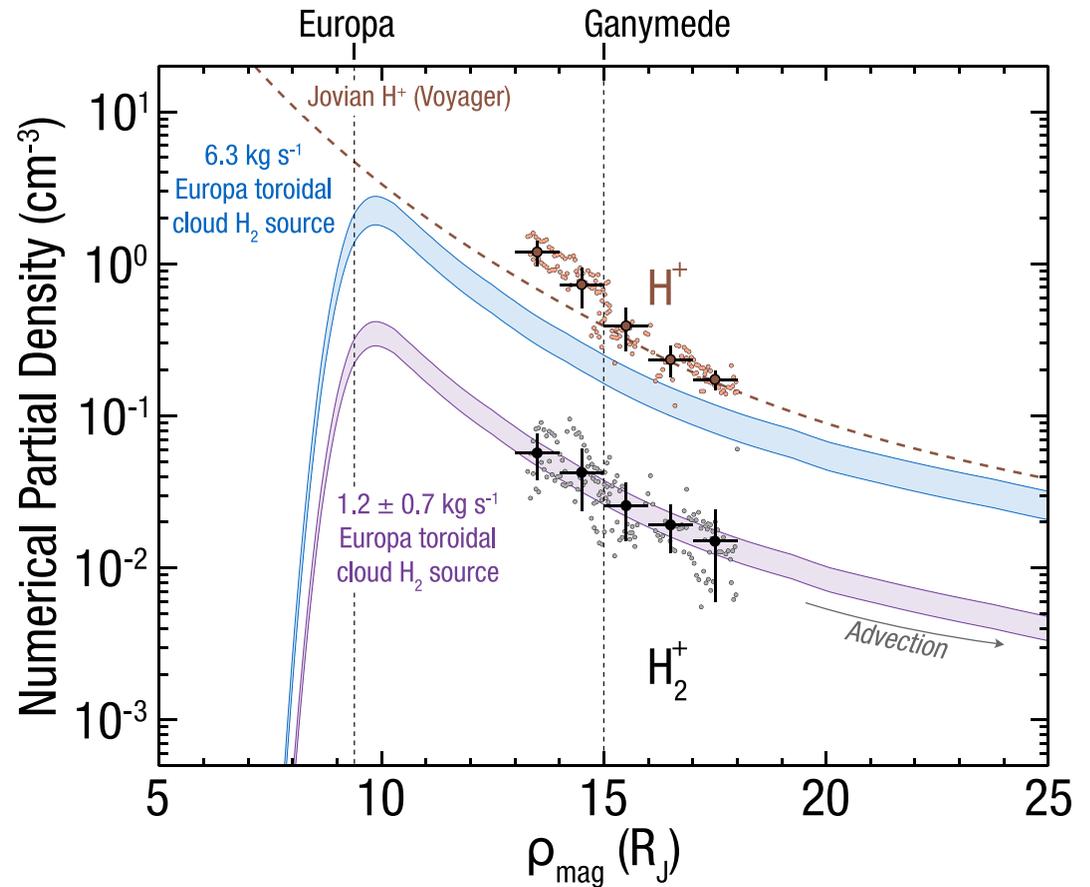
We convert count rate  $R$  into phase space density  $f$  via  $f(v) = R / (G_{eff}^v(v)v^4)$  where  $G_{eff}^v = G_{eff}^E/2$  is the energy dependent geometric factor (Kim et al., 2020a), with a factor of two between the energy geometric factor (e.g., McComas et al., 2021) and velocity geometric factor, and  $v$  is the measured energy per charge converted to speed for a molecule weighing two proton masses. In turn, the number density derived from a one-dimensional phase space density is  $n_{num} = 4\pi \int_0^\infty f(v)v^2 dv$ . For JADE data with count rates in discrete energy bins, the numerical partial number density is given by  $n_{num} = \frac{4\pi}{3} \sum_i (v_{i,max}^3 - v_{i,min}^3) f(v_i)$ , where  $i$  is the energy bin and each energy bin spans from  $v_{i,min}$  to  $v_{i,max}$  in velocity space and  $v_{i,max} = v_{i+1,min}$ . We calculate numerical densities below 10 keV/Q to avoid additional backgrounds above this energy. This energy range is broad enough to well-capture the  $H_2^+$  features analyzed in this study. Finally, we must scale the numerical densities by a correction factor  $n = \epsilon n_{num}$  to account for a minority portion of  $H_2^+$  counts in TOF  $\times$  E space that are over-subtracted in the vicinity of the  $H^+$  fork. See Text S1 in Supporting Information S1 and Kim et al. (2020a) for additional description on the proton “fork.” Text S4 in Supporting Information S1 describes this scaling factor, where we use a range of  $\epsilon = 1.0$ , corresponding to no over-subtraction, to  $\epsilon = 1.6$ , a conservative upper bound on the correction factor.

Figure 2 shows these numerical densities for  $H^+$  and  $H_2^+$  along with average values every 1  $R_J$  for  $\epsilon = 1.2$ . We include the  $H^+$  numerical densities to compare the method to model expectations, as there are no published estimates for the densities of  $H_2^+$  pickup ions in the Jovian magnetosphere. The dark orange dashed line in Figure 2 shows the average expected Jovian proton density from an empirical model based on a reanalysis of Voyager data (Bodisch et al., 2017). Numerical densities for  $H^+$  are commensurate with previous JADE TOF-derived densities in this region (Kim et al., 2020), and are also consistent with empirical model results for protons, lending confidence that the same numerical method applied on  $H_2^+$  yields accurate estimates of the partial densities in these locations.

Using these methods, we find the  $H_2^+$  ion densities monotonically decrease from 13 to 18  $R_J$ , where  $H_2^+/H^+ = 8 \pm 4\%$  for  $\epsilon = 1.0$ –1.6. Such a low density ratio is consistent with the lack of detection by Voyager or Galileo plasma observations, which did not have the sensitivity or TOF mass-spectroscopy capabilities that JADE has. We do not see any appreciable enhancement in the vicinity of Ganymede's orbit, although we did not include data taken very near to Ganymede itself. In the next section, we compare these densities with expectations from pickup ions generated by Europa's neutral toroidal cloud.

### 4. Comparison With a Europa Toroidal Cloud Source

For an inclusive range of plasma and solar ultraviolet conditions near Europa's orbit, electron impact ionization of  $H_2$  neutrals,  $H_2 + e^- \rightarrow H_2^+ + 2e^-$ , is expected to be the dominant loss process for neutral  $H_2$  (Smith et al., 2019).  $H_2^+$  pickup ions generated by this process will be radially transported outwards. The integrated transport time from 6 to 10  $R_J$  is 10–40 days, while the transport time from 10  $R_J$  to the outer magnetosphere ( $\sim 50 R_J$ ) is on the order of a few days with radial outflow speeds approaching 10s of  $\text{km s}^{-1}$  (Bagenal & Delamere, 2011). Additionally, a reanalysis of Galileo plasma data from PLS showed the radial gradient of flux tube content beyond 10  $R_J$  is significantly reduced, precluding significant diffusion (Bagenal et al., 2016). Hence, we follow the assertion



**Figure 2.** Partial density of  $\text{H}_2^+$  pickup ions (gray) and  $\text{H}^+$  (orange), with average values every  $1 R_J$  for  $\epsilon = 1.2$ . The blue curves show expected  $\text{H}_2^+$  pickup ion densities for a neutral Europa loss  $\text{H}_2$  rate of  $6.3 \text{ kg s}^{-1}$  (Smyth & Marconi, 2006). The purple curves show the best fit solution to the observed  $\text{H}_2^+$  partial densities. These observations are consistent with a Europa neutral  $\text{H}_2$  loss rate of  $1.2 \pm 0.7 \text{ kg s}^{-1}$ . The orange dashed line shows an empirical model for proton densities based on Voyager observations (Bodisch et al., 2017).

that while slow diffusive transport is valid in the inner magnetosphere ( $<10 R_J$ ), transport in the middle magnetosphere is better modeled with advection (Ng et al., 2018).

At these low magnetic latitudes, we assume  $\text{H}_2^+$  pickup ions are transported radially outward via cylindrically symmetric advection. From conservation of mass, we solve the 1D time-independent continuity equation in cylindrical coordinates  $\frac{1}{\rho} \frac{\partial(n\rho v_\rho)}{\partial \rho} = P$ , where  $\rho$  is cylindrical radial distance,  $n$  is the number density of  $\text{H}_2^+$ ,  $v_\rho$  is the radial transport speed, and  $P$  is the production rate of pickup ions. This is solved using explicit finite differencing via  $F_{i+1} = F_i + P_i \rho_i \Delta \rho$ , where  $F = n\rho v_\rho$ , and a small value of  $\Delta \rho = 0.025 R_J$  is chosen such that the results are insensitive to the exact value of radial step size. We use a radial transport speed profile (Bagenal & Delamere, 2011) corresponding to a total Jovian plasma source of  $0.7 \text{ ton/s}$ , which represents the typical magnetospheric configuration (Delamere et al., 2004; Nerney & Bagenal, 2020). For the pickup ion input source, we approximate the modeled  $\text{H}_2^+$  pickup ion input given in Figure 10 of Smyth and Marconi (2006) as  $P_E(R) = P_{0E} e^{-\frac{(R-R_E)^2}{2\sigma_E^2}}$  where  $P_{0E}$  is the production rate of  $\text{H}_2^+$  and  $R_E = 9.4 R_J$ , and  $\sigma_E = 0.4 R_J$ . We assume transport occurs within a vertical extent of  $2 R_J$ .

To link the production rate of  $\text{H}_2^+$  pickup ions to neutral  $\text{H}_2$  loss from Europa, we use existing estimates that 50%–77% of neutral  $\text{H}_2$  in the toroidal cloud is lost to  $\text{H}_2^+$  pickup ions (Smith et al., 2019). The purple curves in Figure 2 show the best fit advection solutions to the binned  $\text{H}_2^+$  data. Blue curves show expectations for the initial estimates of  $\text{H}_2^+$  production, discussed in the following section. As shown here, the data are consistent with an advecting Europa-genic toroidal cloud source of  $\text{H}_2^+$  with a monotonically decreasing radial density profile

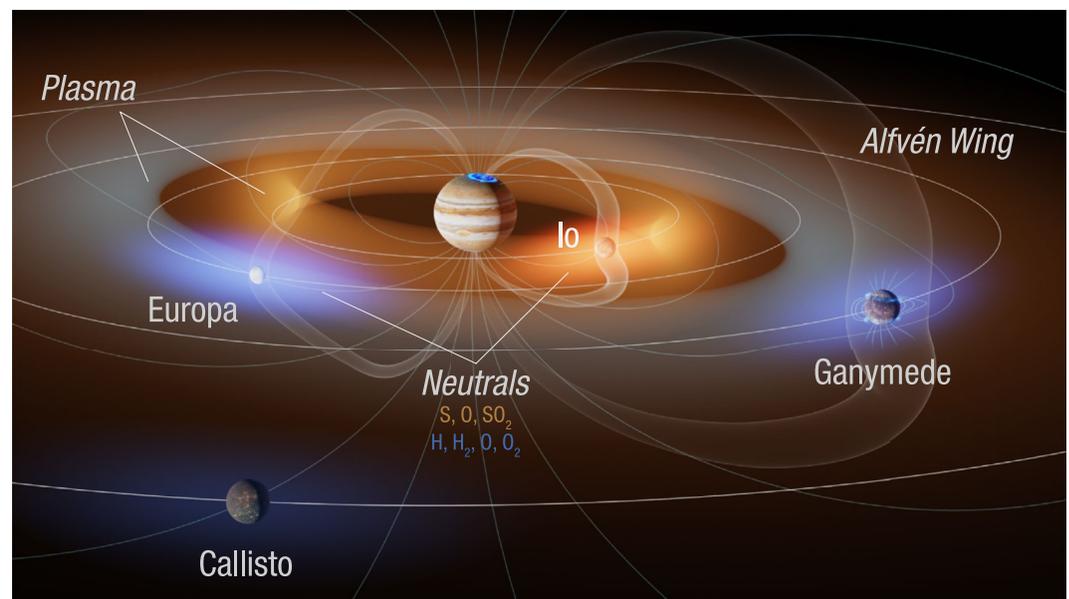
**Table 1**  
Estimates for the Rate of Neutral H<sub>2</sub> Mass Loss (kg s<sup>-1</sup>) and Number Loss (s<sup>-1</sup>) From Europa

H <sub>2</sub> Loss (kg s <sup>-1</sup> )	H <sub>2</sub> Loss (s <sup>-1</sup> )	Method	Reference
6.3	1.9 × 10 <sup>27</sup>	Sputtering and atmosphere model	Smyth and Marconi (2006)
1.5	4.6 × 10 <sup>26</sup>	Sputtering model	Plainaki et al. (2012)
0.7	2 × 10 <sup>26</sup>	Sputtering model	Cassidy et al. (2013)
1.9	5.7 × 10 <sup>26</sup>	Plasma-atmosphere charge exchange model	Dols et al. (2016)
1.8	5.5 × 10 <sup>26</sup>	Sputtering and atmosphere model	Vorburger and Wurz (2018)
1.2 ± 0.7	3.6 ± 2 × 10 <sup>26</sup>	Pickup ion observations and advection model	This work

past ~13 R<sub>J</sub>. These fits are consistent with a total H<sub>2</sub><sup>+</sup> pickup ion production rate of 0.7 ± 0.3 kg s<sup>-1</sup> for the full range of ε = 1.0–1.6. Hence, dividing this range by 0.5–0.77, a neutral H<sub>2</sub> production rate within a 1-sigma range of 1.2 ± 0.7 kg s<sup>-1</sup> from the Europa toroidal cloud, lost from Europa's surface, is derived from the H<sub>2</sub><sup>+</sup> pickup densities observed. Table 1 gives these estimates along with those in the literature to date.

## 5. Discussion and Conclusions

Using plasma observations from the JADE instrument onboard Juno, we detected a persistent population of H<sub>2</sub><sup>+</sup> pickup ions in Jupiter's magnetosphere near the magnetic equator from 13 to 18 R<sub>J</sub>. Using 1D numerical moments, we determine densities of H<sub>2</sub><sup>+</sup> to be in the range of 5 × 10<sup>-3</sup> to 1 × 10<sup>-1</sup> cm<sup>-3</sup> throughout this region. The densities monotonically decrease as a function of radial distance and do not show any clear enhancement near Ganymede's orbit. The overall radial profile and total densities are consistent with pickup ions generated by a Europa-genic neutral toroidal cloud with a loss rate of H<sub>2</sub> from Europa of 1.2 ± 0.7 kg s<sup>-1</sup>. Hence, these observations confirm the existence and composition of Europa's neutral H<sub>2</sub> toroidal cloud, which is expected to be the dominant component (Smith et al., 2019). Figure 3 shows a summary of the Jovian magnetosphere highlighting the existence of H<sub>2</sub> neutrals and H<sub>2</sub><sup>+</sup> pickup ions associated with Europa.



**Figure 3.** Overview of Jupiter's magnetosphere in the vicinity of the Galilean satellites. H<sub>2</sub><sup>+</sup> pickup ions (blue) originate from Europa's neutral toroidal cloud (brighter blue near Europa). Io and Europa contribute plasma pickup ions of different compositions to Jupiter's magnetosphere. Alfvén wings connected to the moons due to their interaction with corotating plasma are also shown in gray.

The total loss rate of neutral H<sub>2</sub> from Europa estimated here is a factor of 3–13 lower than the original estimated value of 6.3 kg s<sup>-1</sup> (Smyth & Marconi, 2006), where the blue curves in Figure 2 show corresponding H<sub>2</sub><sup>+</sup> densities for this estimate. This comparison also shows that were the source rates from Europa to be this large, H<sub>2</sub><sup>+</sup>/H<sup>+</sup> ≳ 30% such that H<sub>2</sub><sup>+</sup> pickup densities would nearly rival those of protons in the magnetosphere. This is certainly not the case as previous measurements would have identified such a population (Voyager, Galileo, Juno, and New Horizons), and we find the density ratio to be H<sub>2</sub><sup>+</sup>/H<sup>+</sup> = 8 ± 4%.

While we derive our estimates from direct observations of the expected primary loss mechanism for Europa's neutral toroidal cloud (pickup ions), the values calculated in this work of 0.5–1.9 kg s<sup>-1</sup> are commensurate with more recent estimates of 0.7–1.9 kg s<sup>-1</sup> (Cassidy et al., 2013; Dols et al., 2016; Plainaki et al., 2012; Vorbürger & Wurz, 2018) based on constraints of the loss rate from the surface due to sputtering inputs (Table 1). Both the sputtering estimates and our pickup ion calculations here require a synthesis of data and modeling to derive estimates; however, given that the two very different methods of estimating Europa's neutral H<sub>2</sub> loss rate yield similar results, we can be reasonably confident that the loss rate is on the order of ~1 kg s<sup>-1</sup>.

We refrain from estimating neutral H<sub>2</sub> densities. Linking pickup ion densities to neutral cloud densities requires an additional set of model-dependent assumptions above those already used in this analysis. Namely, the absolute lifetime of the H<sub>2</sub> neutrals needs to be estimated to constrain neutral densities and doing so would incorporate additional model-dependent uncertainty. Constraining neutral toroidal cloud densities using accurate lifetime estimates requires more complex modeling and is suggested as a future line of study.

Protons are also produced as pickup ions from Europa's neutral toroidal cloud, where the production rate of H<sup>+</sup> is expected to be ~15 times lower than that for H<sub>2</sub><sup>+</sup> (Smyth & Marconi, 2006). Therefore, we anticipate an input rate of protons to the Jovian magnetosphere from Europa to be less than 0.1 kg s<sup>-1</sup>. This value is at least an order of magnitude less input than the 2.5–13 kg s<sup>-1</sup> required to sustain magnetospheric proton abundances inside 30 R<sub>J</sub> (Bodisch et al., 2017), further pointing to Jupiter as the likely source for magnetospheric protons, where Juno observations have shown proton outflow contributes at least 1–5 kg s<sup>-1</sup> to the magnetosphere (Szalay et al., 2021).

The largest source of uncertainty in our analysis lies in the use of a scaling factor to correct for over-subtraction of H<sub>2</sub><sup>+</sup> counts (Text S4 in Supporting Information S1). Our conservative range of scaling factors leads to a large error bar in H<sub>2</sub> mass loss from Europa that could be narrowed with improved methods to isolate H<sub>2</sub><sup>+</sup> count rates in JADE measurements. We also do not estimate densities in the vicinity of Europa's orbit in this study. Inside 13 R<sub>J</sub>, particularly near Europa's orbit, two issues make measuring H<sub>2</sub><sup>+</sup> challenging. First, JADE experiences particularly large fluxes of penetrating radiation, adding to the background that must be removed. Second, and more importantly, locally produced pickup ions are expected to be observed around 0.5 keV and below, where the overlap between the H<sup>+</sup> “fork” in TOF and H<sub>2</sub><sup>+</sup> in the JADE TOF data is particularly pronounced and ε may be significantly larger near Europa using the current method. Due to these reasons, we do not include observations in this region; however, future efforts involving forward-modeling and/or fitting may help reveal the H<sub>2</sub><sup>+</sup> content in this region.

Ganymede is also a source of neutrals (Huang & Siscoe, 1987) and like Europa, is expected to predominantly shed H<sub>2</sub> to generate a population of co-orbiting neutrals (Marconi, 2007). However, such a neutral source would be more dispersed in a larger total volume as the moon orbits at 15 R<sub>J</sub> compared to Europa's 9.4 R<sub>J</sub> distance. Additionally, the radial transport speed is approximately an order of magnitude higher at 15 R<sub>J</sub> compared to 9.4 R<sub>J</sub> (Bagenal & Delamere, 2011), leading to shorter pickup ion lifetimes and densities. Hence even with equal neutral loss source rates from the two bodies, lower densities of neutrals and significantly lower peak densities of subsequent pickup ions would be expected near each of their orbits. While the monotonically decreasing densities as a function of radial distance are entirely consistent with a single source from Europa, we can very roughly constrain the contribution from Ganymede. Assuming Ganymede's neutral toroidal cloud has the same pickup

ion profile  $P_G(R) = P_{0G} e^{-\frac{(R-R_G)^2}{2\sigma_G^2}}$ , where  $R_G = 15.0 R_J$  and  $\sigma_R = \sigma_E = 0.4 R_J$ , we find source rates above ~50% of Europa's would present a detectable departure from the approximately power-law-like radial profile observed for H<sub>2</sub><sup>+</sup>, and therefore do not expect Ganymede's toroidal cloud H<sub>2</sub><sup>+</sup> pickup ion production rate to be larger than 50% of Europa's. If the conversion of H<sub>2</sub> neutrals to H<sub>2</sub><sup>+</sup> pickup ions near Ganymede's orbit is similar to Europa's orbit, H<sub>2</sub> neutral production at Ganymede would also be lower than ~50% of Europa's. This effect is even more dramatic for Callisto at 26.3 R<sub>J</sub>, which also likely loses neutrals (Mogan et al., 2021) to an even more dispersed toroidal

cloud. In the vicinity of Callisto's orbit, the current foreground and background removal techniques were not able to unambiguously identify a clear  $H_2^+$  signal, however, future methods may help constrain the  $H_2^+$  densities in this region and in the region near Europa.

These plasma measurements could directly tie into observations by the high-energy particle detector onboard Juno, JEDI (Mauk et al., 2017), which can search for pitch-angle depletions (e.g., Kollman et al., 2016; Lagg et al., 2003; Nénon & André, 2019) and cross-compare with inferred neutral densities from these pickup ion observations. As the observed  $H_2^+$  pickup ions likely originate from neutrals produced via radiolysis due to energetic particle bombardment and subsequent thermal desorption (e.g., Johnson et al., 2004), JEDI measurements of this charged particle environment could also be cross-compared to these pickup observations and loss rate estimates. Constraints such as these on satellite water species loss processes and composition provide key information on the chemistry and interaction with their local charged particle environments, where present and future spacecraft will continue to shed light on these important processes over a broad range of energies (Futaana et al., 2015; Grasset et al., 2013; Paranicas et al., 2021; Phillips & Pappalardo, 2014).

Thus, the JADE observations presented here for the first time directly measure ions originating from a neutral  $H_2$  toroidal cloud at Jupiter, prove the cloud provides an additional plasma source in Jupiter's magnetosphere, and provide the most direct constraints on Europa's loss of neutral  $H_2$  via observations of the neutral toroidal cloud's primary loss process: pickup ions. Future analyses of JADE  $H_2^+$  pickup ion data may further constrain loss rates and satellite-magnetospheric interactions at Europa, Ganymede, and Callisto. These results underscore the importance of both ion composition and neutral toroidal clouds in understanding the evolution of planetary bodies, not just in the Jovian system, but at all outer planetary systems (e.g., Richardson et al., 1986; Smith & Richardson, 2021), and at exoplanetary systems.

## Data Availability Statement

The JNO-J/SW-JAD-3-CALIBRATED-V1.0 data presented in this manuscript and Supporting Information S1, doi:10.1007/s11214-013-9990-9, can be obtained from the Planetary Data System (PDS) at <https://pds-ppi.igpp.ucla.edu/mission/JUNO/JNO/JAD>.

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