

Photochemical depletion of heavy CO isotopes in the Martian atmosphere

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Photochemistry on Mars reduces the abundance of heavy isotopes in atmospheric CO and their escape efficiency to space

3

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- 15

16 Abstract

17

18 The atmosphere of Mars is enriched in heavy isotopes with respect to Earth as a result of

- 19 the escape of the atmosphere to space over billions of years. Estimating the enrichment in
- 20 heavy isotopes due to atmospheric escape requires a rigorous understanding of all
- atmospheric processes that contribute to the evolution of isotopic ratios between the lower and upper atmosphere, where escape processes take place. Using a combination of vertical
- 22 and upper atmosphere, where escape processes take place. Using a combination of vertica 23 profiles from the Atmospheric Chemistry Suite (ExoMars Trace Gas Orbiter) with the
- 24 predictions of a photochemical model, we provide evidence for a new process of
- 25 photochemistry-induced fractionation that depletes the heavy isotopes of C and O in CO.
- 26 Accounting for this new source of fractionation, we find that only a fraction of $\approx 12\%$ of the
- 27 atmosphere needs to have been lost to space through CO photochemical escape to explain
- 28 the enrichment of ${}^{13}C/{}^{12}C$ in CO₂ measured by the Curiosity Rover.
- 29

30 **1. Main text**

31

32 Numerous lines of geomorphological and mineralogical evidence suggest that liquid water 33 was once abundant on Mars' surface (e.g., Baker (2001); Carr & Clow (1981); Head et al. 34 (1999); Murchie et al. (2009)), but it remains unclear what climatic conditions enabled this, or 35 what drove the transition in the climate to the dry, low pressure atmosphere we observe 36 today. Enrichment in the heavy isotopes of several species such as Ar, N and H suggests 37 that atmospheric escape has been an important mechanism shaping the climate and 38 composition of the atmosphere throughout history (e.g., McElroy et al. (1976); Owen (1982); 39 Owen et al. (1988)). Combining measured atmospheric isotope ratios with evolution models 40 allows the estimation of the abundance of species in the atmosphere of early Mars, which 41 demonstrates the value of a thorough understanding of the atmospheric isotope composition

- 42 (Hu et al., 2015; Hu & Thomas, 2022; Jakosky et al., 2017; Scheller et al., 2021).
- 43
- 44 Accurate estimations of the long-term evolution of the atmosphere from its isotope
- 45 composition rely on two important quantities: measurements of the past and present isotopic
- 46 ratios, and the net escape fractionation factor, which determines the efficiency of the heavy-
- 47 isotope enrichment as the atmosphere escapes to space (e.g., Krasnopolsky et al. (1998);
- 48 Yung et al. (1988)). The most accurate measurements of the isotopic composition of C and
- 49 O in the atmosphere of Mars were made by the Curiosity Rover, which showed an
- 50 enrichment in the heavy isotopes in CO₂ with respect to Earth-like standards (${}^{13}C/{}^{12}C = 1.046$
- \pm 0.004 VPDB and $^{18}\text{O}/^{16}\text{O}$ = 1.048 ± 0.005 VSMOW, see Methods for definition of

- standards), consistent with the hypothesis of substantial atmospheric loss (Mahaffy et al.,
 2013; Webster et al., 2013). On the other hand, while the fractionation factor of several
 escape processes of C and O has been estimated before (Fox & Hać, 1999; Fox & Hać,
 2010; Hu et al., 2015), the net escape fractionation factor, accounting for all processes from
 the lower atmosphere to the upper atmosphere as well, remains undetermined.
- 57

58 By applying this methodology to the C isotopes, the density of CO₂ in the atmosphere of early Mars has been estimated (Hu et al., 2015; Jakosky, 2019). The escape of C at present 59 60 is suggested to occur mostly in the form of hot C produced by a number of different 61 photochemical reactions, such as the direct photodissociation of CO₂ and CO or the 62 dissociative recombination of CO⁺ (Cui et al., 2019; Gröller et al., 2014; Lo et al., 2021; 63 McElroy, 1972). Hu et al. (2015) showed that the photodissociation of CO efficiently enriches 64 the atmosphere in ¹³C as it escapes to space, and whose contribution is an important factor for explaining the current enrichment in ¹³C in CO₂ measured by the Curiosity Rover. 65 66 However, whilst CO therefore seems to be an important species for understanding the isotopic evolution of C in the atmosphere of Mars, its present-day isotopic composition has 67 68 not yet been determined.

69

70 Carbon isotopic ratios can also reveal the nature of surface-atmosphere interaction

71 processes and mineral formation. Recently, carbon isotopic measurements by the Curiosity 72 Payor revealed anomalously large depletions in $\frac{13}{12}C$ ($\frac{13}{12}C$ = 0.86, 0.03 (JDPR) on

Rover revealed anomalously large depletions in ${}^{13}C$ (${}^{13}C/{}^{12}C \approx 0.86-0.93$ VPDB) on sedimentary organics potentially associated with a possible paleosurface (House et al.,

73 Sedimentary organics potentially associated with a possible paleosurace (nouse et al.,
 74 2022). One of the possible explanations suggested to produce ¹³C-depleted organic material

75 is the photochemical reduction of CO_2 into formaldehyde (CH₂O) with CO as an

⁷⁶ intermediate. In this production, the abundance of ¹³C in CO is suggested to be lower than ⁷⁷ that is CO as a result of the differential photoleuric of 1^{3} CO and 1^{3} CO. (Schmidt et al. 2012)

that in CO_2 as a result of the differential photolysis of ${}^{12}CO_2$ and ${}^{13}CO_2$ (Schmidt et al., 2013), transferring its isotopic signature into formaldehyde and other organics that accumulated on

79 the surface of Gale Crater (House et al., 2022). Direct measurements of the ${}^{13}C/{}^{12}C$ in CO

80 may therefore provide constraints on whether the photoreduction of CO_2 is the responsible 81 mechanism for the depletion of ¹³C on the surface organics.

82

83 The reactions that determine the abundance of CO in the atmosphere of Mars are the same 84 as the ones controlling the long-term stability of CO₂: the photolysis of CO₂ into CO, and the posterior recombination of CO back into CO_2 . The stability of CO_2 in Mars' atmosphere was 85 86 at first not well understood since the photodissociation of CO_2 molecules (CO_2 + hv \rightarrow CO + 87 O) is much faster than the recombination of the products (CO + O + M \rightarrow CO₂ + M). It was 88 later realised that the stability of CO₂ is instead controlled by catalytic reactions with oddhydrogen species (CO + OH \rightarrow CO₂ + H), which can convert CO into CO₂ at a much faster 89 90 rate (McElroy & Donahue, 1972; Parkinson & Hunten, 1972). This reaction network not only 91 determines the relative abundances of CO₂ and CO in the atmosphere of Mars, but also 92 controls the relative isotopic fractionation of C and O in these species.

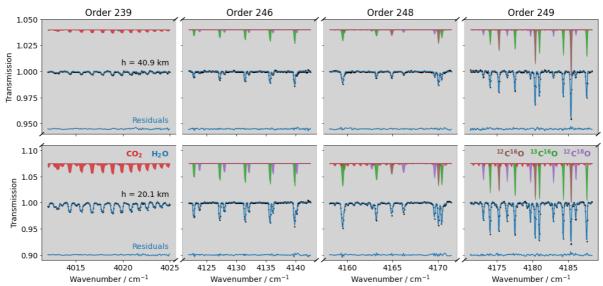
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94 Here, we investigate the isotopic composition of C and O in CO using solar occultation 95 measurements by the Atmospheric Chemistry Suite (ACS) onboard the ExoMars Trace Gas 96 Orbiter (TGO) (Korablev et al., 2018). ACS includes three spectrometers dedicated to 97 investigate the atmosphere using different spectral ranges: the near-infrared channel (NIR), 98 which operates between 0.7-1.7 µm; the mid-infrared channel (MIR) operating between 2.3-99 4.2 µm; and the thermal infrared channel (TIRVIM), which samples the atmosphere between 100 0.7-17 µm. In this study, we use observations from the MIR channel, which is dedicated to 101 solar occultation observations. For these observations, the instrument's boresight is pointed 102 towards the Sun as it rises or sets from behind the Mars disk. By making measurements 103 every 2.1 s, the instrument's line-of-sight crosses the atmosphere at different tangent 104 altitudes above the surface, allowing the reconstruction of the atmospheric vertical profiles. 105

106 ACS observations have been previously used to measure for the first time the vertical 107 distribution of CO and its seasonal variations in the atmosphere of Mars (Fedorova et al., 108 2022; Olsen et al., 2021). Here, we use the MIR data employed by Fedorova et al., but studying the relative abundances of ¹²C¹⁶O, ¹³C¹⁶O and ¹²C¹⁸O. In particular, we use the 109 ACS MIR measurements obtained using the secondary grating position 6, which includes 110 111 fourteen diffraction orders (236-249) encompassing a spectral range between 3950-4193 112 cm⁻¹. We perform the retrievals in four spectral windows in diffraction orders 239 (4012-4025 cm⁻¹), 246 (4122-4142 cm⁻¹), 248 (4158-4171 cm⁻¹) and 249 (4171-4188 cm⁻¹) (see Figure 113 114 1), to retrieve the pressure, temperature and volume mixing ratios of H_2O and the 115 aforementioned CO isotopes (see Methods). The retrieved dataset in this study comprises all ACS MIR secondary grating position 6 observations where the four presented diffraction 116 117 orders were measured, which includes 242 vertical profiles mainly measured in the second halves of Martian Year (MY) 34 (L_s = 166°-349°) and MY35 (L_s = 141°-366°). This large 118 119 range of observing conditions in terms of season, local time, location and dust-activity allows 120 a better characterisation of any potential variability in the isotopic ratios, as well as to derive 121 more reliable estimates of the measurement uncertainties.



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Figure 1: Example of ACS MIR spectra from secondary grating position 6 made during orbit 127 1896 (Latitude = 83.8° ; Longitude = 170.5° ; $L_s = 166^\circ$; Local time = 22.5 h). The panels show 128 the measured transmission spectra in four diffraction orders (black dots) and best fit (blue 129 lines) at 20 (bottom) and 40 km (top). The residuals between the measurement and best fit 130 in each diffraction order are shown in each panel, together with the contribution from each 131 gas to each spectrum, adding a constant offset in both to improve the visibility of the figure. 132

133

Figure 2 shows a summary of all ¹³C/¹²C and ¹⁸O/¹⁶O isotopic ratios in CO measured in this 134 dataset. We observe the ${}^{13}C/{}^{12}C$ ratio to be systematically depleted in the heavy isotopes 135 with respect to the Earth-like standard with a weighted averaged value of ${}^{13}C/{}^{12}C = 0.84 \pm$ 136 137 0.09 VPDB, where the uncertainties correspond to the standard deviation of the measurements. Comparing this value to that measured in CO₂ by the Curiosity Rover, it 138 corresponds to a depletion of ${}^{13}C/{}^{12}C$ in CO of approximately 20%. On the other hand, the 139 140 ¹⁸O/¹⁶O ratio is consistent at all altitudes with the Earth-like fractionation. The distribution of 141 measured values in this case yields a weighted average isotopic ratio of ${}^{18}O/{}^{16}O = 0.98 \pm$ 142 0.11 VSMOW, which compared to the isotopic ratio measured in CO₂ by the Curiosity Rover 143 corresponds to a depletion of the O isotopic ratio in CO of 6%.

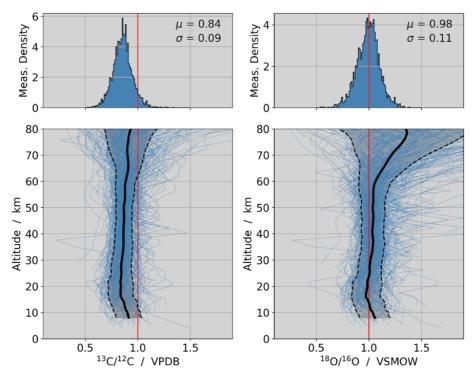




Figure 2: Vertical profiles of the C and O isotopic ratios in CO in the atmosphere of Mars. The top panels show the histogram of all measured ¹³C/¹²C and ¹⁸O/¹⁶O ratios in standard units with uncertainties lower than 0.4. The lower panels show the 242 measured vertical profiles of the isotopic ratios in this dataset (blue thin lines), as well as the weighted average (black line) and the standard deviation of the profiles (shaded areas). The red lines indicate the Earth-like standards given by Vienna Pee Dee Belemnite (VPDB) and Vienna Standard Mean Ocean Water (VSMOW).

154 155 In order to understand the physical and chemical processes that give rise to the depletion of 156 the C and O isotopic ratios measured in CO, we constructed a 1D photochemical model to 157 simulate the isotopic fractionation between CO and CO₂. In particular, we include isotopespecific rates for the two most important reactions relating these two species: the photolysis 158 of CO₂ into CO + O, and the recombination of CO + OH back into CO₂. The photolysis cross-159 sections for ¹²C¹⁶O₂, ¹³C¹⁶O₂ and ¹⁸O¹²C¹⁶O are taken from Schmidt et al. (2013), while the 160 161 fractionation during the recombination of CO into CO₂ is modelled using the data from the 162 laboratory measurements by Stevens et al. (1980) (see Methods). This photochemical 163 model extends from the surface up to an altitude of 200 km and is set to represent martian mean dayside conditions (i.e., solar zenith angle equal to 60° and Sun-Mars distance of 1.52 164 165 AU (Krasnopolsky, 2010)). The isotopic ratios in the model are initialised assuming Earth-166 like fractionation at all altitudes, and the model is then run for approximately 20 years, the 167 time by which the isotopic ratios in CO and CO_2 have converged to a steady-state solution 168 (see Extended Data Figure 2).

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171 Figure 3 shows the derived C isotopic abundances from the photochemical model scaled to

172 the isotopic composition of the atmosphere of Mars measured by the Curiosity Rover,

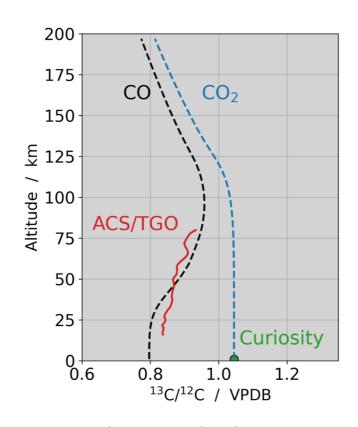
173 together with the measured values from the ACS observations. The C isotopic ratios in both

174 CO and CO₂ show a similar altitudinal trend above approximately 100 km, where they

- 175 continuously decrease with increasing altitude due to diffusive separation: above the
- 176 homopause, where molecular diffusion dominates, species are mixed according to their own

177 mass-dependent scale height. As a result of this, the isotopic ratios steadily decrease with 178 altitude, with the decrease in ${}^{13}C/{}^{12}C$ being approximately half of ${}^{18}O/{}^{16}O$ (Alday et al., 2021) 179 (see Extended Data Figure 2). While the isotopic ratios in CO₂ below the homopause altitude 180 are approximately constant, those in CO are depleted in the heavy isotopes due to photochemistry-induced fractionation, mainly due to the preferential photolysis of ¹²C¹⁶O₂ 181 over ¹³C¹⁶O₂ and ¹⁸O¹²C¹⁶O (see Extended Data Figures 3 and 4). In particular, the 182 183 photochemical model suggests that below 30 km the C isotopic ratio in CO is approximately 184 24% lower than in CO₂, increasing to about 9% at 100 km. The effect of the photochemistry 185 has a smaller impact in the O isotope ratios, with the depletion of the heavy isotopes near 186 the surface being approximately 9%. The lower value of the ¹³C/¹²C ratio in CO than that in 187 ¹⁸O/¹⁶O predicted by the photochemical model is consistent with the ACS measurements 188 presented in this study.

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190 191

191 Figure 3: Photochemistry-induced fractionation of the C isotopes in the atmosphere of Mars.

192 The ACS measurements (red) reveal a decrease of the ${}^{13}C/{}^{12}C$ isotope ratio in CO with

193 respect to that in CO_2 measured by the Curiosity Rover at the surface (green). The

194 predictions from the 1D photochemical model suggest that this decrease is produced by the

195 chemistry of the martian atmosphere, which preferentially favours the transfer of ¹³C from

196 CO (black dashed line) to CO_2 (blue dashed line). The isotope ratio in CO in the upper

197 atmosphere, responsible for a substantial fraction of the C escape to space, is substantially

198 lower than that of the bulk CO₂ atmosphere, suggesting a lower escape fractionation factor

199 than previously thought.

200

201 Both the ACS measurements and the predictions from the model suggest that the

202 photochemistry of the atmosphere of Mars produces a depletion of the heavy C and O

203 isotopes in CO, which represents the first direct detection of this effect in the atmosphere of

204 Mars. Isotopic fractionation also occurs during the photolysis of other species such as H₂O,

whose effect in the martian atmosphere has been extensively studied (Cheng et al., 1999;

- 206 Vladimir A. Krasnopolsky, 2002). The photolysis-induced fractionation in H₂O is responsible
- 207 for fractionating the D/H ratio in water, molecular hydrogen, and finally the escaping

products, giving rise to an efficient enrichment in D/H as the atmosphere escapes to space
 (e.g., Cangi et al. (2020); Krasnopolsky (2002)).

210

211 Since CO contributes a substantial fraction of the C escape from Mars, this new source of 212 fractionation has to be accounted for when estimating the net escape fractionation factor. 213 The escape of C from the atmosphere of Mars due to CO photodissociation and 214 photoionisation is most efficient at altitudes above 150 km (Cui et al., 2019; Lo et al., 2021). 215 At these altitudes, the isotopic composition of CO is affected by both the chemistry and 216 diffusive separation. Our model suggests that at 200 km, the ¹³C/¹²C isotopic ratio in CO is \approx 0.74 times lower than that of the bulk atmosphere, which combined with the fractionation 217 factor due to escape by CO photodissociation ($a \approx 0.6$, Hu et al. (2015)) yields a net escape 218 219 fractionation factor of ≈0.44 for this reaction. This lower fractionation factor indicates that the 220 C escape through CO photodissociation enriches the atmosphere in ¹³C more efficiently than 221 previously thought. In particular, considering a primordial isotopic ratio of ${}^{13}C/{}^{12}C = 0.975$ 222 VPDB (Wright et al., 1986), and a long-term enrichment of the isotopic ratio through 223 Rayleigh distillation, we estimate that only a loss of 12% of the atmosphere is required to 224 have escaped to space through this mechanism to explain the enrichment of the heavy 225 isotopes measured by the Curiosity Rover. While the total density of C on early Mars is also 226 subject to other processes such as carbonate formation or outgassing from the surface (Hu 227 et al., 2015), our results highlight the efficient role of this mechanism to enrich the 228 atmosphere in the heavy isotopes, suggesting that C escape may have been less important 229 than previously thought. 230

Since the atmospheric mass of ¹²C and ¹³C is conserved, the chemistry-induced fractionation 231 also implies a preferential transfer of ¹³C from CO to CO₂, which yields an enrichment in the 232 233 13 C/ 12 C ratio in CO₂. The near-surface isotopic ratio in CO₂ is found to be $\approx 3 \times 10^{-5}$ times 234 higher than that in the initial atmosphere, which corresponds to an enrichment of $\delta^{13}C \approx$ 235 0.03 ‰. This value is much smaller than the enrichment in the heavy isotopes observed by 236 the Curiosity Rover (δ^{13} C = 46 ± 4 ‰, Webster et al. (2013)), which suggests that 237 atmospheric escape remains as the preferred process contributing to the enrichment of ¹³C 238 in the atmosphere of Mars.

239

240 The depletion of ¹³C in CO observed with the ACS measurements and predicted by the photochemical model is in line with the measurements of the ¹³C/¹²C ratio in Martian 241 242 organics made by the Curiosity Rover (House et al., 2022). While other suggested formation 243 mechanisms such as the photolysis of CH₄ released from the subsurface or the deposition of 244 cosmic dust cannot be ruled out as responsible for the production of ¹³C-depleted organics 245 on Mars, our results confirm that the photolysis of CO₂ can produce strong depletions of the 246 ¹³C/¹²C ratio in CO, an isotopic signature that can be transferred to formaldehyde or other 247 organics that might accumulate on the surface of Mars.

248

249 The photochemical model also predicts a relative depletion of the ¹⁸O/¹⁶O isotopic ratio in 250 CO with respect to CO₂, although to a lesser extent than for the C isotopes. The implications 251 of this new source of fractionation to the long-term evolution of the O isotopes are 252 nevertheless more complicated due to the much more complex O chemistry in the 253 atmosphere of Mars. The dissociative recombination of O₂⁺, responsible for most of the O 254 escape from Mars (Jakosky et al., 2018), also efficiently enriches the atmosphere in ¹⁸O as 255 escape occurs (Fox & Hać, 2010). However, the overall escape fractionation factor will be subject to the ${}^{18}O/{}^{16}O$ ratio in O_2^+ and the source species of these ions (i.e., O, O⁺, CO₂, 256 CO₂⁺, (Fox & Hać, 2009)), which will likely be affected by the isotope photochemistry. Based 257 258 on the isotope-specific photolysis cross-sections included in this study, while the O isotopic 259 ratios of CO₂ in the ionosphere will mostly be affected by the diffusive separation above the 260 homopause, the isotopic composition of atomic O will be affected by the preferential photolysis of ¹²C¹⁶O₂ over ¹⁸O¹²C¹⁶O, similar to the observed depletion of ¹⁸O/¹⁶O in CO. 261

262 263 264 265 266	However, since atomic O is produced and lost in several other chemical reactions, a comprehensive analysis of the O isotopic composition in the upper atmosphere requires a model accounting for ¹⁸ O in other species, which is planned but beyond the scope of the work presented here.
267 268 269 270 271 272	The results from this study highlight the important role of photochemistry in fractionating the C and O isotopic composition of the atmosphere of Mars. Future measurements of the isotopic ratios in different species may not only provide key information about the photochemical cycles in the atmosphere of Mars, but also provide accurate estimations of the fractionation between the lower and upper atmospheres, which has important implications to our understanding of the long-term evolution of Mars' climate.
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305 2. Methods

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307 2.1 Retrieval scheme

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309 The inversion of the spectra is performed using the NEMESIS algorithm (Irwin et al., 2008), 310 applying the same two-step methodology as in Alday et al. (2021). First, the pressure and 311 temperature profiles are determined using the CO₂ absorption features in diffraction order 312 239 (see Figure 1). Secondly, the pressure and temperature are fixed, and the abundances of the trace gases (i.e., H₂O, ¹²C¹⁶O, ¹³C¹⁶O and ¹²C¹⁸O) are retrieved using all diffraction 313 314 orders. All gaseous absorption in this study is modelled using pre-computed line-by-line 315 look-up tables generated using the spectroscopic parameters from the 2020 edition of the 316 HITRAN database (Gordon et al., 2022). Apart from the retrieval of the atmospheric 317 parameters, the instrument line shape (ILS) is also fitted in each occultation using a double 318 Gaussian parameterisation (Alday et al., 2019).

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320 The pressure and temperature profiles are retrieved under the assumption of an atmosphere 321 in hydrostatic equilibrium and a known CO₂ volume mixing ratio. This approach has been 322 widely used in the past to determine the thermal structure of the atmosphere of Mars from 323 solar occultation observations (Forget et al., 2009; Montmessin et al., 2006; Quémerais et 324 al., 2006). The retrieval scheme from the NEMESIS algorithm works using the optimal 325 estimation framework, which aims to find an optimal solution that is consistent with the 326 observed spectra subject to a minimal departure from the *a priori*. Given that the absorption 327 features of CO₂ measured in ACS MIR secondary grating position 6 typically disappear at 328 lower tangent altitudes than the CO ones, we use temperature profiles from simultaneous 329 observations by ACS NIR (Fedorova et al., 2022) as the a priori in our retrievals, which 330 ensures an accurate description of the temperature field at all altitudes. The use of accurate 331 measurements of temperature is crucial to derive precise measurements of the CO isotopic 332 ratios, which are highly sensitive to this parameter (see Supplementary Material).

333

334 Once the pressure and temperature profiles are determined, these are fixed, and the 335 gaseous abundances of the CO isotopes and H_2O are retrieved. In this case, the *a priori* 336 information is taken from the OpenMARS reanalysis dataset (Holmes et al., 2020), which 337 assimilates the temperature profiles and dust opacities measured by the Mars Climate 338 Sounder onboard the Mars Reconnaissance Orbiter into the General Circulation Model. The 339 a priori isotopic composition is assumed to follow that of Earth-like conditions given by the 340 Vienna Pee Dee Belemnite $({}^{13}C/{}^{12}C = 1 \text{ VPDB} = 0.0112372)$ and the Vienna Standard Mean Ocean Water ($^{18}O/^{16}O = 1$ VSMOW = 0.0020052). Given that the sensitivity of the retrieval to 341 the CO isotopes can be different at different altitudes (i.e., the ¹²C¹⁸O absorption lines 342 typically disappear at lower tangent heights than the ¹²C¹⁶O), the influence of the *a priori* 343 344 guess in the retrieval of the isotopologues might also be different. Therefore, in order to 345 ensure no biases from the *a priori* are introduced into the derivations of the isotopic ratios, 346 the retrievals are run twice, the second time using the retrieved ¹²C¹⁶O abundances from the 347 previous iteration as the *a priori* and scaling the abundances of the minor isotopes with the 348 same Earth-like isotopic ratios.

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350 In each acquisition made by ACS MIR, the observed diffraction orders encompass about 20 detector rows, which correspond to the instantaneous field-of-view of the instrument, with 351 352 each row sampling slightly different tangent heights ($\Delta z \approx 150$ m per row). In addition, the 353 ILS varies from row to row, which can introduce systematic uncertainties in the retrievals. In 354 order to increase the confidence of the retrievals and better understand the uncertainties 355 arising from the ILS fitting, we apply the retrieval scheme to seven detector rows. The 356 retrieved profiles from each row are then combined using a weighted average. The 357 uncertainties are calculated considering both the error of the mean, and the standard 358 deviation of the profiles (Alday et al., 2021). We consider this method for calculating the

uncertainties to provide a more accurate representation of the true uncertainty of the retrieval in which not only the random uncertainties, but also the systematic ones are captured.

A detailed validation of the retrievals presented in this study is provided in the Supplementary Material.

2.2 Photochemical model

The 1D photochemical model aims to solve the continuity equation given by

 $\frac{\partial n_i}{\partial t} = P_i - L_i - \frac{\partial \phi_i}{\partial z},$

where n_i is the number density of a given species, P_i and L_i represent the production and loss of density due to the photochemical reactions, and the last term represents the diffusion between layers, where ϕ is the flux. The flux is calculated considering both Eddy and molecular diffusion and is given by

377
$$\phi_i = -(K+D_i)\frac{\partial n_i}{\partial z} - n_i \left[K\left(\frac{1}{H_0} + \frac{1}{T}\frac{\partial T}{\partial z}\right) + D_i \left(\frac{1}{H_i} + \frac{(1+\alpha)}{T}\frac{\partial T}{\partial z}\right) \right],$$

where K and D_i are the Eddy and molecular diffusion coefficients, H_0 is the mean scale height, H_i is the gas-dependent scale height, a is the thermal diffusion coefficient, and T is the temperature.

The diffusion term is calculated by finite differences, using an approach similar to Krasnopolsky (1993), which yields

$$\frac{\partial \phi_{i,j}}{\partial z} \sim -(k_j + d_{i,j}) \cdot n_{i,j+1} + (kl_j + dl_{i,j} + k_{j-1} + d_{i,j-1}) \cdot n_{i,j} - (kl_{j-1} + dl_{i,j-1}) \cdot n_{i,j-1}$$

$$387$$

where the subscripts *i* and *j* represent the gas and layer, respectively. The coefficients in the finite differences approximation are given by

391

$$k_j = \frac{K_{j+1/2}}{\Delta z^2},$$

392
 $d_{i,j} = \frac{D_{i,j+1/2}}{\Delta z^2},$

$$d_{i,j} = \frac{D_{i,j+1,j}}{\Delta z^2}$$

393
$$kl_{j} = k_{j} \cdot \left(1 - \frac{\Delta z}{H_{0_{j+\frac{1}{2}}}} - \frac{T_{j+1} - T_{j}}{T_{j+\frac{1}{2}}}\right),$$

394
$$dl_{i,j} = d_{i,j} \cdot \left(1 - \frac{\Delta z}{H_{i,j+\frac{1}{2}}} - (1+\alpha) \cdot \frac{T_{j+1} - T_j}{T_{j+\frac{1}{2}}}\right),$$

where Δz is the altitude of each layer.

The production and loss terms are calculated using the latest version of the photochemical scheme from the Mars Planetary Climate Model (PCM), which has been extensively used to

model the photochemical cycles in the atmosphere of Mars (Cariolle et al., 2017; Lefèvre et 401 402 al., 2021; Lefèvre et al., 2004). This scheme includes 14 chemical species and 39 chemical 403 reactions to model the photochemistry of H, C and O in the atmosphere of Mars. Here, we introduce four new species (13C16O2, 13C16O, 18O12C16O and 12C18O) and ten chemical 404 reactions to model the isotopic fractionation between CO₂ and CO. The reaction rates are 405 406 calculated assuming mean dayside conditions with a constant solar zenith angle of 60° and 407 a constant Sun-Mars distance of 1.52 AU. 408 409 The reactions included to model the ¹³C chemistry are: 410 $\label{eq:alpha} \begin{array}{ll} {}^{13}\text{CO}_2 + \text{hv} \rightarrow {}^{13}\text{CO} + \text{O} & (\lambda > 167 \text{ nm}) \\ {}^{13}\text{CO}_2 + \text{hv} \rightarrow {}^{13}\text{CO} + \text{O}({}^1\text{D}) & (\lambda < 167 \text{ nm}) \end{array}$ 411 412 413 414 $^{13}CO + O + M \rightarrow ^{13}CO_2 + M$ 415 $^{13}CO + OH \rightarrow ^{13}CO_2 + H$ 416 417 418 The reactions included to model the O fractionation between CO₂ and CO are: 419 ${}^{18}\text{OC}{}^{16}\text{O} + \text{hv} \rightarrow \text{C}{}^{18}\text{O} + {}^{16}\text{O}$ 420 (λ > 167 nm) $^{18}OC^{16}O + hv \rightarrow C^{16}O + {}^{16}O$ 421 $(\lambda > 167 \text{ nm})$ 422 $\label{eq:alpha} \begin{array}{ll} {}^{18}\text{OC}{}^{16}\text{O} + \text{hv} \rightarrow \text{C}{}^{18}\text{O} + {}^{16}\text{O}({}^{1}\text{D}) & (\lambda < 167 \text{ nm}) \\ {}^{18}\text{OC}{}^{16}\text{O} + \text{hv} \rightarrow \text{C}{}^{16}\text{O} + {}^{16}\text{O}({}^{1}\text{D}) & (\lambda < 167 \text{ nm}) \end{array}$ 423 424 425 $C^{18}O + {}^{16}O + M \rightarrow {}^{18}OC^{16}O + M$ 426 427 $C^{18}O + {}^{16}OH \rightarrow {}^{18}OC^{16}O + H$ 428 429 It must be noted that in the version for the ¹⁸O chemistry the total mass is not conserved, as 430 431 we do not include ¹⁸O or ¹⁸O(¹D) in the chemistry scheme. These species, together with the 432 definition of the ¹⁸O isotopologues in other O-bearing species are necessary to model the complete chemistry of ¹⁸O in the atmosphere of Mars. Here, we just aim to model the relative 433 fractionation of the ${}^{18}O/{}^{16}O$ isotopic ratio between CO₂ and CO, which is not expected to be 434 435 affected by this simplification, since CO molecules are solely produced by the photolysis of 436 CO₂, which is included in the scheme. 437 438 The continuity equation is solved using a second-order Rosenbrock algorithm (Hobbs et al., 439 2019; Verwer et al., 1999). The density at the next time step is calculated following 440
$$\begin{split} n^{t+1} &= n^t + 1.5\Delta t \cdot g_1 + 0.5\Delta t \cdot g_2, \\ & (I - \gamma \Delta t J)g_1 = f(n^t), \\ & (I - \gamma \Delta t J)g_2 = f(n^t - \Delta t g_1) - 2g_1, \end{split}$$
441 442 443 444 445 Where $\gamma = 1 + 1/\sqrt{2}$, Δt is the time step, f(n) is the right-hand side of the differential 446 equation (i.e., the chemistry and transport terms evaluated at the density n) and J is the 447 Jacobian matrix of the combined chemistry-transport system. Since the densities of a given 448 species at a given layer are only influenced by the densities of the other gases in the same laver (i.e., chemistry) and the densities of the same gas in the adjacent lavers (i.e., 449 450 diffusion), the Jacobian matrix takes the form of a block tridiagonal matrix, which is efficiently 451 inverted using the Thomas algorithm (Hu et al., 2012). 452

The choice of the time step is made based on the maximum local error, which is estimated by the difference between the first-order and second-order solutions, following

455 456

457

$$\varepsilon = n^{t+1} - n_1^{t+1} = (n^t + 1.5\Delta t \cdot g_1 + 0.5\Delta t \cdot g_2) - (n^t + \Delta t \cdot g_1)$$

458 The timestep is self-adjusted by $\Delta t^{t+1} \sim \Delta t^t / \sqrt{\varepsilon_{max}}$, which ensures that the convergence of 459 the system to a steady-state solution will capture the orders-of-magnitude difference in the 460 chemical and diffusion timescales while optimising the computational time.

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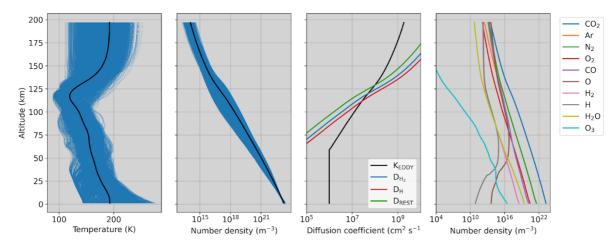
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475 476

462 The 1D photochemical model requires several inputs to represent the composition and mixing of the atmosphere of Mars. The initial temperature and density profiles are derived 463 464 using an average of vertical profiles generated from the Mars Climate Database (Forget et 465 al., 1999) at different times and locations. In particular, we generate vertical profiles of 466 temperature, number density and volume mixing ratios from 0 to 200 km in a grid of latitudes, local times and solar longitude L_s. The latitudes expand from -80° to 80° with a 467 step of 20°. The local time expands from 0 to 24 hours with a two-hour step. The L_s expands 468 469 the whole year, with a step of 10°. The calculated profiles are then averaged using the 470 atmospheric density as a weight (see Extended Data Figure 1). The abundances of CO₂ and 471 CO from the Mars Climate Database are then divided into the different isotopologues by 472 assuming a given isotopic ratio, and using

$$\begin{bmatrix} {}^{12}C^{16}O_2 \end{bmatrix} = \begin{bmatrix} CO_2 \end{bmatrix} / (1 + R)$$
$$\begin{bmatrix} {}^{13}C^{16}O_2 \end{bmatrix} = \begin{bmatrix} {}^{12}C^{16}O_2 \end{bmatrix} \times R$$

477 where the bracketed quantities represent the abundances of the different species, and *R* is 478 the assumed Earth-like isotopic ratio (i.e., ${}^{13}C/{}^{12}C = 1$ VPDB and ${}^{18}O/{}^{16}O = 1$ VSMOW). 479



480

481 Extended Data Figure 1: Overview of the main inputs for the photochemical model. The 482 initial profiles for the temperature, total number density and partial number density are

calculated by averaging profiles from the Mars Climate Database throughout different
seasons, locations and local time. The Eddy and molecular diffusion coefficients are
calculated from Krasnopolsky (1993) and using the coefficients from Hunten (1973),
respectively.

487

488 The diffusion coefficients assumed in our Mars 1D photochemical model are the same as 489 those used by Chaffin et al. (2017) and Cangi et al. (2020), which were taken from other

- 490 previous studies (e.g., Hunten (1973); Krasnopolsky (1993)). In particular, the Eddy diffusion
- 491 coefficient is given by
- 492

493 $K(z < 60 \ km) = 10^6 \ \text{cm}^2 \ \text{s}^{-1},$ 494 $K(z < 60 \ km) = 2 \cdot 10^{13} \cdot n(z)^{-1/2} \ \text{cm}^2 \ \text{s}^{-1},$

496 as shown in Extended Data Figure 1. The molecular diffusion coefficient is given by 497

$$D(z) = \frac{AT(z)^s}{n(z)}$$

499

503

498

495

500 The coefficients A and s for H₂ and H are given by Hunten, (1973), which are equal to $A_H =$ 501 8.4 x 10¹⁷, $s_H = 0.597$, $A_{H2} = 2.23 \times 10^{17}$, $s_{H2} = 0.75$. For the rest of the species, we adopt A =502 10¹⁷ and s = 0.75 (see Extended Data Figure 1).

The boundary conditions are similar to those used in Chaffin et al. (2017). We consider the top boundary conditions to be a zero flux for all species but for H₂, H and O. For H₂ and H, we fix their escaping velocities to be $v_{H2} = 3.4 \times 10^{1}$ cm s⁻¹ and $v_{H} = 3.1 \times 10^{3}$ cm s⁻¹, while for O we fix the escaping flux to be $\phi_{O} = 1.2 \times 10^{8}$ cm⁻² s⁻¹. For the lower boundary conditions, we set a zero flux for all species, but for H₂O, whose density is assumed to be

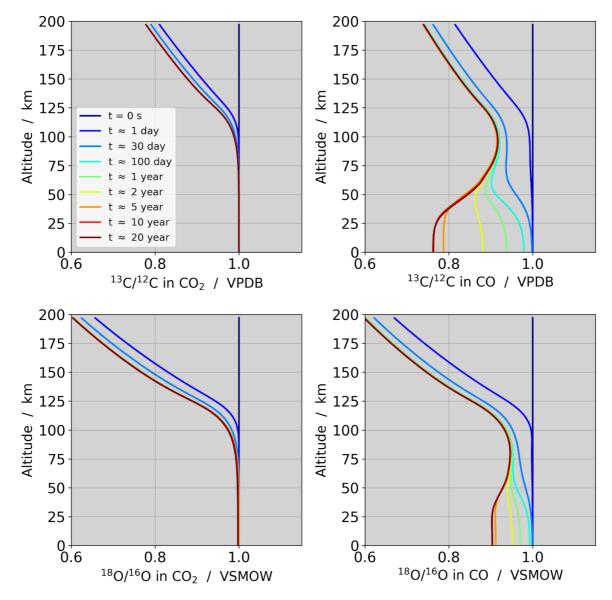
509 fixed below 50 km.

510

511 Using these inputs and the presented model setup, we run a simulation for 20 years, the

512 time by which the O and C isotopic ratios in both CO and CO₂ have converged to a steady-

- 513 state solution, as it is shown in Extended Data Figure 2.
- 514



515 516 Extended Data Figure 2: Evolution of the isotopic ratios in the 1D photochemical model. The 517 four panels show the carbon (top) and oxygen (bottom) isotopic ratios in CO₂ (left) and CO (right) as a function of time as the simulation from the photochemical model progresses. The 518 519 isotopic ratios in CO_2 converge rapidly, as they are mostly affected by diffusive separation, 520 which produces a decrease of the isotopic ratios above the homopause according to their 521 own mass-dependent scale heights. The isotopic ratios in CO are affected by both diffusive 522 separation and by the photochemistry-induced fractionation. The photochemistry of the atmosphere produces a depletion in the ¹³C/¹²C and ¹⁸O/¹⁶O ratios, whose effect is stronger 523 for the former of the two, consistent with the ACS observations.

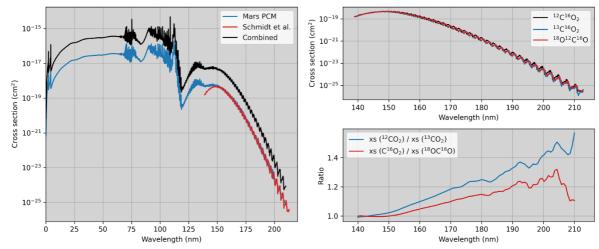
526 2.3 Fractionation during the photolysis of CO₂

527

528 The photolysis cross sections of CO₂ used in the chemistry scheme of the LMD Mars PCM 529 are a composite of different laboratory measurements in different spectral ranges (Huestis & 530 Berkowitz, 2010; Lewis & Carver, 1983; Parkinson et al., 2003; Stark et al., 2007; Yoshino et 531 al., 1996) and tabulated at temperatures of 195, 295 and 370 K (Lefèvre et al., 2021). In order to model the fractionation factor of this reaction, we use the photolysis cross sections 532 calculated by Schmidt et al. (2013) for ${}^{12}C^{16}O_2$, ${}^{13}C^{16}O_2$ and ${}^{18}O^{12}C^{16}O$ between 150-210 nm, 533

534 which are combined with the Mars PCM reference for <150nm at 195, 295 and 370 K. 535 Extended Data Figure 3 shows how the calculated cross sections for ${}^{12}C^{16}O_2$ compare with 536 those tabulated in the chemistry scheme, together with a comparison between the cross 537 sections calculated for the different isotopologues. These results suggest that the photolysis 538 cross sections of ${}^{12}C^{16}O_2$ are larger than those of the minor isotopes (i.e., ${}^{13}C^{16}O_2$ and 539 ${}^{18}O^{12}C^{16}O$) by a factor of 1.0-1.5 and 1.0-1.3, respectively. The ratio between the cross 540 sections of ${}^{12}C^{16}O_2$ and those of the minor isotopes also depends on the temperature, with 541 that at 195 K being on average 5% larger than that at 295 K.

542

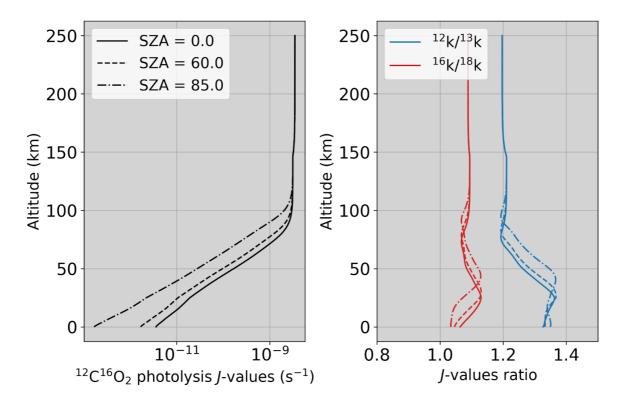


543

544 Extended Data Figure 3: Summary of the photolysis cross sections used in this study. The 545 left plot shows a comparison between the cross sections tabulated in the Mars PCM (blue 546 line) and those calculated by Schmidt et al. (2013) (red line) at 295 K. In addition, the combined cross sections used in this study are shown with a x10 offset. The top right plot 547 548 shows the cross sections calculated by Schmidt et al. (2013) for the different isotopologues. 549 The bottom right plot shows the ratio between the cross sections of the major and minor 550 isotopes, convolved with a Gaussian function with a FWHM of 2.5nm to smooth out high 551 frequency variations.

552

553 It is important to note that the cross sections for the minor isotopes are only available for a 554 limited spectral range (>150 nm), and these are lower than the cross sections elsewhere 555 (<150 nm), where we assume that they are equal for all isotopes. In order to evaluate where 556 the cross sections >150 nm are most effective in the photolysis of CO_2 molecules in the 557 atmosphere of Mars, we calculate the photolysis J-values as a function of altitude (see 558 Extended Data Figure 4). These results show that the *J*-values are approximately constant 559 above ≈100 km, where the atmosphere is optically thin. Below this altitude, the atmosphere 560 becomes optically thick for wavelengths lower than 150 nm, and the J-values decrease with 561 decreasing altitude as the atmosphere also becomes optically thick for wavelengths >150 562 nm. Therefore, while the cross sections <150 nm are important above 100 km, their effect 563 becomes negligible below 100 km, the region where the ACS observations can measure the 564 isotopic fractionation between CO₂ and CO.



566 567

Extended Data Figure 4: Photolysis J-values of ¹²C¹⁶O₂ and the isotope effect. The left-hand panel shows the calculated photolysis J-values of the $12^{12}C^{16}O_2$ at three different solar zenith 568 angles. The right-hand panel shows the ratio between the photolysis rates of ${}^{12}C^{16}O_2$ over 569 those of ${}^{13}C^{16}O_2$ and ${}^{18}O^{12}C^{16}O_2$. 570

572 The right-hand panel of Extended Data Figure 4 shows the ratio between the photolysis rates of ${}^{12}C^{16}O_2$ over those of ${}^{13}C^{16}O_2$ and ${}^{18}O^{12}C^{16}O_2$, which is higher than one at all 573 altitudes. These results suggest that the photolysis of ¹²C¹⁶O₂ is faster than that of the minor 574 isotopes, meaning that ¹²C¹⁶O will be more efficiently produced than ¹³C¹⁶O and ¹²C¹⁸O. The 575 net effect of this process is to produce a depletion of the ¹³C/¹²C and ¹⁸O/¹⁶O ratios in CO 576 with respect to those in CO₂, and the magnitude of this effect is larger for the C isotopes 577 578 than the O ones.

579

581

580 2.4 Fractionation during the recombination of CO + OH

582 The recombination of CO with the hydroxyl radical OH is an important reaction in the 583 chemistry of the atmosphere of Mars, which is responsible for the stability of CO₂ (McElroy & 584 Donahue, 1972; Parkinson & Hunten, 1972). In order to model the isotopic fractionation of 585 the C and O isotopes in this reaction, we use the data derived from laboratory 586 measurements by Stevens et al. (1980). These measurements determined the kinetic 587 isotope effect at a range of pressures from 0.2 to 1.35 atm (see Extended Data Figure 5). In 588 order to extrapolate the measurements to the low pressure of the martian atmosphere, we fit 589 this dependence using a second-order polynomial function. The fitted polynomial functions 590 for the C and O isotopes are respectively 591

592
$$k^{13}/_{\mu 12} = 1.00638 - 1.69300 \times 10^{-7} \cdot p + 4.69677 \times 10^{-13} \cdot p^2$$

593
$$k^{18}/_{k^{12}} = 1.01195 - 4.44340 \times 10^{-7} \cdot p + 1.93835 \times 10^{-13} \cdot p^2$$
,

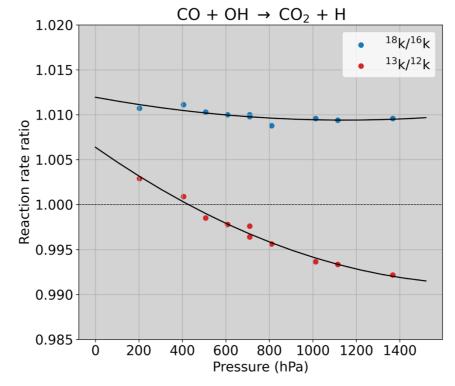
where the ratios ${}^{13}k/{}^{12}k$ and ${}^{18}k/{}^{16}k$ represent the ratio between the reaction rates of the minor isotopes (i.e., ¹³CO and C¹⁸O) with respect to that of the major one (i.e., ¹²C¹⁶O) and pis the pressure in Pascals.

At the low pressures of the martian atmosphere, this reaction is faster for the minor isotopes,

which yields a preferential transfer of ¹³C and ¹⁸O over ¹²C and ¹⁶O from CO to CO₂.

Therefore, the net effect of this reaction is to produce a depletion of the ${}^{13}C/{}^{12}C$ and ${}^{18}O/{}^{16}O$

- isotopic ratios in CO with respect to CO₂. However, given the small magnitude of this
- fractionation factor with respect to that produced by the photolysis of CO₂, this process is
- expected to have a much minor effect in the overall fractionation between CO₂ and CO.



606 Extended Data Figure 5: Isotopic fractionation during the recombination of CO and OH into CO₂ and H. The blue and red dots represent the ratio between the reaction rates of $C^{18}O/C^{16}O$ and $^{13}CO/^{12}CO$, measured by Stevens et al. (1988). The black lines are a

- second-order polynomial fit used to capture the pressure dependence of this fractionation effect.

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