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Bar properties as a function of wavelength: a local baseline with S⁴G for high-redshift studies

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ABSTRACT

The redshift evolution of bars is an important signpost of the dynamic maturity of disc galaxies. To characterize the intrinsic evolution safe from band-shifting effects, it is necessary to gauge how bar properties vary locally as a function of wavelength. We investigate bar properties in 16 nearby galaxies from the *Spitzer* Survey of Stellar Structure in Galaxies (S⁴G) at ultraviolet, optical, and mid-infrared wavebands. Based on the ellipticity and position angle profiles from fitting elliptical isophotes to the two-dimensional light distribution, we find that both bar length and ellipticity – the latter often used as a proxy for bar strength – increase at bluer wavebands. Bars are 9 per cent longer in the *B* band than at 3.6 μ m. Their ellipticity increases typically by 8 per cent in the *B* band, with a significant fraction (>40 per cent) displaying an increase up to 35 per cent. We attribute the increase in bar length to the presence of star-forming knots at the end of bars: these regions are brighter in bluer bands, stretching the bar signature further out. The increase in bar ellipticity could be driven by the apparent bulge size: the bulge is less prominent at bluer bands, allowing for thinner ellipses within the bar region. Alternatively, it could be due to younger stellar populations associated with the bar. The resulting effect is that bars appear longer and thinner at bluer wavebands. This indicates that band-shifting effects are significant and need to be corrected for high-redshift studies to reliably gauge any intrinsic evolution of the bar properties with redshift.

Key words: methods: data analysis – techniques: photometric – galaxies: spiral – galaxies: structure – infrared: galaxies.

1 INTRODUCTION

Stellar bars are present in roughly two-thirds of all nearby galaxies (de Vaucouleurs 1963; Eskridge et al. 2000; Knapen, Shlosman & Peletier 2000; Whyte et al. 2002; Marinova & Jooe 2007; Menéndez-Delmestre et al. 2007; Sheth et al. 2008; Erwin 2018). Their formation, evolution, and potential destruction (the latter still under debate; e.g. Bournaud & Combes 2002; Bournaud et al. 2005; Athanassoula, Machado & Rodionov 2013) can dramatically affect their host galaxies. This is due to the fact that a bar introduces a non-

axisymmetric component to the gravitational potential of the host galaxy disc, inducing large-scale streaming motions in the gas that may have a major impact in the chemical, dynamic, and structural evolution of the host galaxy itself (Athanassoula 1992; Martin & Roy 1994; Zaritsky, Kennicutt & Huchra 1994; Sheth et al. 2003, 2008; Grand, Kawata & Cropper 2015). Simulations have shown that a galaxy disc will succumb to the bar instability (i.e. form a bar) within a Hubble time unless it is dynamically hot or its potential is dominated by dark matter (Athanassoula 2003). Work by Sheth et al. (2012) provided strong observational evidence of this by showing that bars are preferentially found in massive, rotation-dominated galaxies, indicating that mass and dynamic coldness of a disc are necessary – albeit not sufficient – conditions for bar

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formation. Indeed, Romeo, Agertz & Renaud et al. (2023) provided further observational evidence that we are still missing a proper bar instability criterion: bars form in disc galaxies of all morphological types, from lenticulars to blue compact dwarfs, and classical bar instability criteria such as the Efstathiou, Lake & Negroponte (1982) and the Mo, Mao & White (1998) criteria fail to predict this in about 55 per cent of the cases. For cosmological studies, the redshift evolution of bars is thus an important signpost on the growth and dynamic maturity of disc galaxies (Sheth et al. 2008; Kim et al. 2021).

Early claims of a sudden reduction of the bar fraction beyond $z \sim 0.8$ (Abraham et al. 1999) suscitated an immediate concern for the impact that band-shifting effects – the progressive shift of the photometric band to bluer rest-frame wavebands – could have on the reported bar fraction (Sheth et al. 2003). Could we be losing track of existing bars simply by observing them at bluer bands? The measure of the local bar fraction had been extended to longer wavelengths with the advent of large near-infrared (near-IR) surveys, including the Ohio State University Bright Spiral Galaxy Survey (Eskridge et al. 2002) and the Two Micron All Sky Survey (Skrutskie et al. 2006) Large Galaxy Atlas (Jarrett et al. 2003). Although bars are structures primarily comprising the older stellar population (Gadotti & de Souza 2006; Sánchez-Blázquez et al. 2011; de Sá-Freitas et al. 2022), better traced at longer wavelengths, it has been well established that – although individual weak bars in the optical become more conspicuous in the near-IR – the overall bar fraction remains roughly the same from the *B* band to the near-IR (Eskridge et al. 2002; Menéndez-Delmestre et al. 2007; Barway, Wadadekar & Kembhavi 2011; Buta et al. 2015). Hence, when observations are limited to the optical-through-IR wavebands, band-shifting effects do not have a significant impact on the measured bar fraction.

With the advent of the Advanced Camera for Surveys (ACS) on the *Hubble Space Telescope* in 2002, the availability of high-resolution optical imaging motivated the detailed exploration of stellar structure of galaxies at high redshift, including a characterization of the bar fraction out to $z \sim 0.8$ –1 (e.g. Abraham et al. 1999; Sheth et al. 2003; Elmegreen, Elmegreen & Hirst 2004; Jogee et al. 2004; Sheth et al. 2008; Melvin et al. 2014). Even with detailed correction for magnitude survey limits, stellar mass, inclination, and galaxy-type selection, Sheth et al. (2008) showed that reports on the bar fraction evolution beyond $z \sim 0.85$ based on optical data (e.g. Jogee et al. 2004; Elmegreen, Elmegreen & Hirst 2004) fall prey to band-shifting effects; beyond $z \sim 0.85$, even the red ACS *I*-band traces ultraviolet (UV) light, at which point bars are difficult to recognize (Sheth et al. 2003). More recent studies exploiting Wide-Field Camera 3 (WFC3) near-IR data pushed the redshift limit of the bar fraction characterization beyond $z \sim 1$ based on rest-frame optical images (Simmons et al. 2014). With the advent of the *JWST*, a handful of barred galaxies have been detected beyond $z = 2$ (Guo et al. 2023; Costantin et al. 2023). Furthermore, Le Conte et al. (2023) find that the fraction of barred galaxies beyond $z \sim 1$ as measured with *JWST* at rest-frame near-IR is ~ 3 –4 times that found earlier in studies based on *HST* WFC3 data. These results suggest the importance of bar studies at rest frame near- and mid-IR bands (see also Méndez-Abreu, Costantin & Kruk 2023; Bland-Hawthorn et al. 2023).

Considering local bar properties is of critical importance when assessing the completeness of high-redshift bar fraction studies, as weak bars are more difficult to pick out (Sheth et al. 2003) and small bars become undetectable as a consequence of resolution limits (see fig. 6 in Menéndez-Delmestre et al. 2007). Beyond an interest in characterizing the completeness of high-redshift studies, studying the distribution of bar properties allows us to test the accuracy of bar formation models (Athanasoula, Machado & Rodionov 2013). The-

oretical models predict that when the bar evolves, it becomes longer and stronger (e.g. Athanasoula 2003; Villa-Vargas, Shlosman & Heller 2010; Athanasoula, Machado & Rodionov 2013). Hydrodynamical simulations have started to look at the redshift evolution of bar properties showing this very behaviour and – more importantly – have started to establish predictions of how bar properties change over cosmological time (e.g. Okamoto, Isoe & Habe 2015). The evolution of bar fraction and other bar properties have been studied using recent cosmological hydrodynamical simulations of galaxy formation such as EAGLE (Algorry et al. 2017), Illustris TNG100 (Rosas-Guevara et al. 2020; Zhao et al. 2020; Zhou et al. 2020), Illustris TNG50 (Izquierdo-Villalba et al. 2022; Rosas-Guevara et al. 2022; Zana et al. 2022), NewHorizon (Reddish et al. 2022), and Auriga (Fragkoudi et al. 2021). In such simulations, bars typically start developing after redshifts $z \sim 1$ –2, although massive discs are already present earlier than that. This redshift of formation depends on mass, size, and other selection criteria. Theoretical predictions can thus be made for the evolution of bar properties, but simulation studies generally rely on stellar mass rather than light to measure structural properties of the bar.

A number of optical studies have looked at the distribution of bar properties in optically selected local bars (Kormendy 1979; Elmegreen & Elmegreen 1985; Martin 1995; Erwin 2005; Aguerri, Méndez-Abreu & Corsini 2009; Barazza et al. 2009; Marinova et al. 2009; Gadotti 2011; Li et al. 2011; Hoyle et al. 2011), while others have extended the analysis to the near- and mid-IR based on statistically large samples (e.g. Menéndez-Delmestre et al. 2007; Erwin 2018, 2019). A few studies have started to venture out to intermediate redshifts to characterize bar properties based on optical data (Abraham et al. 1999; Barazza et al. 2009; Kim et al. 2021), but the role of band-shifting effects has not yet been well established in the interpretation of their results. With ample optical through near-IR space-based data with sufficient spatial resolution now available (e.g. ACS/WFPC2/WFC3) and within the context of new high-resolution MIR imaging (e.g. *JWST*), we are at a stage that to reliably characterize the redshift evolution of bar properties, a local baseline is necessary to establish how bar properties vary with wavelength. Such a baseline is crucial for comparison, as it allows high-redshift studies to determine if observed changes in bar properties with redshift are simply band-shifting effects or genuine signs of secular evolution. Accurately understanding how bars evolve with redshift will provide valuable information about the astrophysical processes that lead to bar formation and an observational confirmation of what simulations point as necessary conditions for a galaxy to succumb to the bar instability (e.g. Athanasoula, Machado & Rodionov 2013). The interactions between bars and galactic discs will, in turn, give us a better understanding of how galaxies evolve over cosmological time-scales.

In this paper, we present a detailed study of bar properties – bar *length* and bar *strength*, as given by the bar’s ellipticity – as a function of wavelength for a sample of 16 large nearby barred galaxies, spanning the wavelength range from the UV to the mid-infrared (MIR). The MIR waveband mitigates dust-obscuration effects and provides the optimal window to probe the details of stellar structure in galaxies (see fig. 23 in Driver et al. 2016). By extending our investigation to the UV light we project the applicability of our results to ongoing and future studies of bar properties beyond $z \sim 0.8$ with the ample high-resolution optical data provided by current space telescopes. With this work we further envision preparing for the MIR inspection of bar properties at sub-kpc scales out to cosmic noon, in preparation for new and future space-based observing facilities (e.g. *JWST* and *Euclid*).

We describe our sample selection and observations in Section 2, summarize our analysis in Section 3, and state our results in Section 4. In Section 5, we discuss our results and give our main conclusions in Section 6. We assume a Λ CDM cosmology, with $H_0 = 71 \text{ km s}^{-1} \text{ Mpc}^{-1}$, $\Omega_M = 0.27$, and $\Omega_\Lambda = 0.73$.

2 SAMPLE SELECTION AND OBSERVATIONS

The rest-frame MIR light in normal galaxies is dominated by the old stellar population; by mitigating dust-obscuration effects, it provides the optimal window to probe the details of stellar structure in galaxies, in particular bars and their properties. We thus base our sample selection on the *Spitzer* Survey of Stellar Structure in Galaxies (S⁴G; Sheth et al. 2010), a volume- ($D < 40 \text{ Mpc}$), magnitude- ($m_b < 15.5$), and size-limited ($D_{25} > 1 \text{ arcmin}$) survey of 2331 galaxies in the nearby Universe with the 3.6 and 4.5 μm bands of the *Spitzer Space Telescope*'s Infrared Array Camera (IRAC; Fazio et al. 2004). Our sample comprises spiral galaxies with clear bar signatures from the optical through the MIR: classified as either barred (SB) or intermediately barred (SAB) in the optical by de Vaucouleurs (1963) and confirmed as bars in the near-IR (Menéndez-Delmestre et al. 2007) and the MIR (Buta et al. 2010). We required the availability of high-quality galaxy images in the UV, optical, and MIR to ensure adequate comparison. Taking into consideration that the galaxy data base from the Spitzer Infrared Galaxy Survey (SINGS; Kennicutt et al. 2003) includes excellent optical coverage in the B and R bands, we selected S⁴G galaxies that were also part of the SINGS sample. We exclude barred spirals with high disc inclinations, $i > 70^\circ$. These selection criteria resulted in a sample of 16 nearby barred spirals, listed in Table 1. We show the distribution in RC3 bar type (SAB/SB), Hubble T -type, inclination and distance in Fig. 1, taken from the Hyperleda Extragalactic data base (Makarov et al. 2014).¹

For our MIR analysis, we use *Spitzer* IRAC 3.6 μm images. These were first reduced by the Spitzer Science Center S18.5 pipeline and further processed by the S⁴G pipeline (see Muñoz-Mateos et al. 2015 for details). The resulting images, with a pixel scale of 0.75 arcsec and spatial resolution of $\sim 1.5 \text{ arcsec}$, reach a 1σ brightness limit of $27 \text{ mag arcsec}^{-2}$ and are part of the S⁴G data products publicly available at <http://irsa.ipac.caltech.edu/data/SPITZER/S4G>.

We used B - and R -band images from the SINGS fifth data release.² Originating from observations made with the 2.1 m telescope at Kitt Peak National Observatory (KPNO) and the 1.5 m telescope at Cerro Tololo Inter-American Observatory (CTIO) down to a uniform depth of $\sim 25 \text{ mag arcsec}^{-2}$, these images present pixel scales of 0.434 and 0.304 arcsec, respectively, with spatial resolutions dictated by the seeing at the time of observations, in the range of $\sim 1\text{--}2 \text{ arcsec}$.

In order to cover the UV emission for the galaxies in our sample, we rely on GALEX images in the far- and near-UV bands ($\lambda_{FUV,NUV} = 1516\text{\AA}, 2267\text{\AA}$, respectively) from the Nearby Galaxies Survey (NGS; Gil de Paz et al. 2004; Bianchi et al. 2003a, b) down to a surface brightness of $\mu_{FUV,NUV} = 27.25 \text{ mag arcsec}^{-2}$, $27.35 \text{ mag arcsec}^{-2}$, respectively (Thilker et al. 2007). With a spatial resolution of 4.0–4.5 arcsec (FUV) and 5.0–5.5 arcsec (NUV), these images establish our coarsest spatial resolutions, at a pixel scale of 1.5 arcsec (Morrissey et al. 2007).

To accommodate the large range in angular resolutions presented by this combined data set, we divide our analysis into two parts: a *higher resolution optical-through-MIR study*, where

we compare B - and R -band optical images with the MIR S⁴G-processed *Spitzer*/IRAC images; and a *low-resolution UV-through-MIR study*, where we compare the MIR images with the coarser resolution GALEX images. To ensure an adequate comparison of bar parameters across different bands, we resampled and smoothed all images onto a common pixel scale and resolution within each of our two studies. To set a common pixel scale we used the IDL routine *congrid* adapted for python.³ For the *higher resolution optical-through-MIR study* we resampled the B - and R -band optical images onto a pixel grid of $0.75 \text{ arcsec pixel}^{-1}$ set by the S⁴G-processed IRAC observations. For the *low-resolution UV-through-MIR study* we adopted the coarser 1.5 arcsec pixel size characteristic of the GALEX images. We smoothed higher-resolution images down to the limit set by the lowest-resolution images within each study. For this we relied on the IRAF⁴ GAUSS task, which convolves images with a Gaussian function parametrized by a user-fed *sigma* value in pixel units. In this manner we set a common resolution of 6 arcsec for the images within the *low-resolution UV-through-MIR study* and a resolution of 1.5 arcsec for the images within the *higher resolution optical-through-MIR study*.

3 DATA ANALYSIS

To trace the spatial extent of the bar and the degree of the influence that it has on the dynamics of the host galaxy, two bar properties are of fundamental importance: the bar length and the bar strength. Many approaches to measuring the bar length have been proposed in the past (e.g. Athanassoula & Misiriotis 2002). The *strength* of the non-axisymmetric influence of the bar's gravity on the otherwise axisymmetric gravitational potential of the galactic disc can be quantified in a variety of ways, including the bar-interbar contrast (e.g. Elmegreen & Elmegreen 1985; Regan & Elmegreen 1997), the Q_b parameter that measures the maximum gravitational bar torque relative to the galactic disc (e.g. Combes & Sanders 1981; Buta & Block 2001; Laurikainen & Salo 2002; Speltincox, Laurikainen & Salo 2008) and a variety of approaches that use a two-dimensional (2D) fast Fourier transform method on galaxy images (see Garcia-Gómez et al. 2017 and references therein). An alternative, simple method is based on measuring the maximum ellipticity of the bar isophotes (e.g. Martin 1995; Aguerri 1999; Shlosman, Peletier & Knapen 2000; Knapen, Shlosman & Peletier 2000; Laine et al. 2002; Erwin 2004; Menéndez-Delmestre et al. 2007; Marinova & Jogee 2007; Gadotti 2011). High ellipticity indicates a *thin*, well-defined bar in the equatorial plane of the galaxy, imposing sharp non-axisymmetry, while a low ellipticity is associated with a *broad* bar and weaker non-axisymmetrical deviation of the disc potential. Considering that bar ellipticity has been shown to correlate well with Q_b measures (e.g. Laurikainen & Salo 2002), bar ellipticity represents the simplest, yet accurate approach to characterize bar strength. Its simplicity is crucial for application to large galaxy samples, in particular when dealing with distant barred galaxies, where surface-brightness dimming prevent us from implementing other more elaborate methods (e.g. Sheth et al. 2008). In this work we rely on the bar's ellipticity as a measure of bar strength.

³CONGRID is a publicly available routine distributed as part of the open-source software database at <http://wiki.scipy.org/Cookbook/Rebinning>.

⁴IRAF is distributed by the National Optical Astronomy Observatory, which is operated by the Association of Universities for Research in Astronomy (AURA) under a cooperative agreement with the National Science Foundation.

¹<http://leda.univ-lyon1.fr/>

²<https://irsa.ipac.caltech.edu/data/SPITZER/SINGS/>

Table 1. Bar properties.

Galaxy Band	ϵ_{\max} $\delta\epsilon = 0.01$	a_{\max} (arcsec)	S ⁴ G [RC3] classification ^a $a_{\Delta\epsilon}$ (arcsec)	$a_{\Delta PA}$ (arcsec)	ϵ_{\max} $\delta\epsilon = 0.01$	Inclination ^b (deg) $a_{\epsilon\max}$ (arcsec)	Distance ^b (Mpc) $a_{\Delta\epsilon}$ (arcsec)	Visual bar length ^c (arcsec) $a_{\Delta PA}$ (arcsec)
Optical-to-MIR high-resolution study					UV-through-MIR low-resolution study			
NGC0337			SAB(s)cd: pec [SB(s)dl]			50.6	22.49	19.4
3.6 μm	0.66	26 \pm 2	35 \pm 3	31 \pm 2	0.45	21 \pm 6	27 \pm 10	33 \pm 6
R	0.59	26 \pm 2	33 \pm 2	33 \pm 2	0.46	27 \pm 6	33 \pm 11	33 \pm 11
B	0.70	31 \pm 2	35 \pm 2	35 \pm 2	0.57	27 \pm 6	33 \pm 8	33 \pm 6
NUV	–	–	–	–	0.72	27 \pm 6	39 \pm 14	39 \pm 14
FUV	–	–	–	–	0.75	27 \pm 6	39 \pm 6	39 \pm 6
NGC1097			(R)SB(rs,lr)ab pec [SB(s)b]			37.0	15.21	94.0
3.6 μm	0.65	94 \pm 4	115 \pm 2	117 \pm 2	0.62	94 \pm 6	112 \pm 11	118 \pm 6
R	0.63	99 \pm 9	99 \pm 2	119 \pm 2	0.61	92 \pm 9	118 \pm 14	118 \pm 6
B	0.69	112 \pm 2	128 \pm 2	142 \pm 2	0.66	112 \pm 6	124 \pm 6	124 \pm 6
NUV	–	–	–	–	0.79	120 \pm 6	141 \pm 6	141 \pm 6
FUV	–	–	–	–	0.84	123 \pm 6	141 \pm 6	141 \pm 6
NGC1291			(R)SAB(l,rb)0 ⁺ [(R)SB(s)0/a]			29.4	8.47	97.2
3.6 μm	0.40	94 \pm 6	99 \pm 2	124 \pm 3	0.39	88 \pm 6	106 \pm 19	124 \pm 15
R	0.39	94 \pm 6	108 \pm 13	94 \pm 2	0.38	88 \pm 6	106 \pm 18	124 \pm 11
B	0.39	90 \pm 4	108 \pm 15	92 \pm 3	0.37	88 \pm 6	106 \pm 18	100 \pm 9
NUV	–	–	–	–	0.31	90 \pm 12	120 \pm 6	126 \pm 6
FUV	–	–	–	–	0.29	100 \pm 6	118 \pm 6	124 \pm 6
NGC1512			(RL)SB(r,lr)a [SB(r)a]			68.3	9.16	71.4
3.6 μm	0.66	74 \pm 6	90 \pm 2	90 \pm 2	0.62	76 \pm 6	88 \pm 6	106 \pm 9
R	0.65	73 \pm 2	92 \pm 2	92 \pm 2	0.62	82 \pm 6	88 \pm 6	100 \pm 6
B	0.72	90 \pm 2	90 \pm 2	90 \pm 2	0.67	82 \pm 6	94 \pm 6	106 \pm 6
NUV	–	–	–	–	–	–	–	–
FUV	–	–	–	–	–	–	–	–
NGC1566			(R ¹)SAB(s)b [SAB(s)bc]			47.9	17.46	35.8
3.6 μm	0.49	32 \pm 2	36 \pm 2	43 \pm 3	0.34	34 \pm 6	34 \pm 6	40 \pm 6
R	0.42	32 \pm 2	43 \pm 2	43 \pm 2	0.30	34 \pm 6	34 \pm 6	40 \pm 6
B	0.57	34 \pm 2	36 \pm 2	47 \pm 4	0.30	34 \pm 6	34 \pm 6	40 \pm 6
NUV	–	–	–	–	–	–	–	–
FUV	–	–	–	–	–	–	–	–
NGC3049			SB(s)ab: [SB(rs)ab]			58.1	21.28	36.2
3.6 μm	0.77	38 \pm 3	52 \pm 2	65 \pm 2	0.61	40 \pm 6	52 \pm 6	76 \pm 6
R	0.80	38 \pm 6	54 \pm 2	68 \pm 2	0.64	44 \pm 6	50 \pm 6	74 \pm 6
B	0.82	43 \pm 8	58 \pm 2	70 \pm 2	0.65	46 \pm 6	52 \pm 6	70 \pm 6
NUV	–	–	–	–	0.59	50 \pm 6	50 \pm 9	50 \pm 9
FUV	–	–	–	–	0.65	52 \pm 6	58 \pm 6	88 \pm 6
NGC3198			SAB(rs)bc [SB(rs)c]			70.0	12.13	28.0
3.6 μm	0.77	68 \pm 3	81 \pm 2	83 \pm 2	0.71	70 \pm 6	88 \pm 6	82 \pm 7
R	0.73	76 \pm 9	86 \pm 2	86 \pm 2	0.69	76 \pm 6	88 \pm 9	82 \pm 17
B	0.73	74 \pm 10	88 \pm 2	88 \pm 2	0.70	76 \pm 6	88 \pm 8	82 \pm 6
NUV	–	–	–	–	–	–	–	–
FUV	–	–	–	–	–	–	–	–
NGC3351			(R ¹)SB(r,lr)a [SB(r)b]			41.5	8.79	51.4
3.6 μm	0.46	60 \pm 6	68 \pm 2	68 \pm 2	0.42	58 \pm 6	64 \pm 6	64 \pm 6
R	0.42	54 \pm 3	68 \pm 2	68 \pm 2	0.38	58 \pm 6	64 \pm 6	64 \pm 6
B	0.52	68 \pm 2	76 \pm 2	76 \pm 2	0.45	64 \pm 6	70 \pm 6	70 \pm 6
NUV	–	–	–	–	–	–	–	–
FUV	–	–	–	–	–	–	–	–
NGC3627			SB(s)b pec [SAB(s)b]			57.3	11.27	66.4
3.6 μm	0.75	61 \pm 3	81 \pm 2	81 \pm 2	0.69	64 \pm 6	76 \pm 6	82 \pm 6
R	0.66	63 \pm 2	79 \pm 2	72 \pm 2	0.60	64 \pm 6	76 \pm 9	88 \pm 9
B	0.69	63 \pm 2	72 \pm 2	65 \pm 2	0.62	64 \pm 6	82 \pm 6	82 \pm 6
NUV	–	–	–	–	–	–	–	–
FUV	–	–	–	–	–	–	–	–
NGC4321			SAB(rs,lr)bc [SAB(s)bc]			30.0	23.99	59.2
3.6 μm	0.59	59 \pm 2	70 \pm 2	70 \pm 2	0.55	63 \pm 6	63 \pm 6	75 \pm 6
R	0.53	66 \pm 4	66 \pm 2	75 \pm 2	0.49	60 \pm 6	72 \pm 6	72 \pm 6
B	0.66	76 \pm 2	79 \pm 2	79 \pm 2	0.62	78 \pm 6	78 \pm 6	78 \pm 6
NUV	–	–	–	–	0.78	105 \pm 6	105 \pm 6	105 \pm 6
FUV	–	–	–	–	–	–	–	–
NGC4559			SB(s)cd [SAB(rs)cd]			64.8	13.93	12.6
3.6 μm	0.58	14 \pm 2	19 \pm 3	21 \pm 2	–	–	–	–
R	0.59	28 \pm 2	32 \pm 2	32 \pm 2	–	–	–	–
B	0.79	30 \pm 2	37 \pm 2	37 \pm 2	0.51	24 \pm 6	30 \pm 6	30 \pm 6
NUV	–	–	–	–	0.48	20 \pm 6	26 \pm 6	26 \pm 6
FUV	–	–	–	–	0.61	20 \pm 6	26 \pm 6	26 \pm 6
NGC4579			(R)SB(rs)a [SAB(rs)b]			39.0	22.91	40.7
3.6 μm	0.47	43 \pm 2	52 \pm 3	52 \pm 2	0.40	40 \pm 6	46 \pm 6	52 \pm 6
R	0.48	45 \pm 2	50 \pm 2	52 \pm 2	0.40	44 \pm 6	46 \pm 8	52 \pm 8
B	0.51	45 \pm 2	52 \pm 2	52 \pm 2	0.42	46 \pm 6	46 \pm 6	52 \pm 6
NUV	–	–	–	–	–	–	–	–
FUV	–	–	–	–	–	–	–	–
NGC4625			(R)SAB(rs)m [SAB(rs)m pec]			46.1	11.75	7.7
3.6 μm	0.31	9 \pm 2	11 \pm 2	27 \pm 2	–	–	–	–
R	0.30	7 \pm 2	9 \pm 3	27 \pm 2	–	–	–	–
B	0.41	7 \pm 2	9 \pm 3	34 \pm 2	–	–	–	–
NUV	–	–	–	–	–	–	–	–
FUV	–	–	–	–	–	–	–	–
NGC4725			SAB(r,lr)ab [SAB(r)ab pec]			54.4	19.50	123.3
3.6 μm	0.68	137 \pm 8	151 \pm 2	148 \pm 2	0.67	130 \pm 6	148 \pm 10	148 \pm 10
R	0.68	137 \pm 7	144 \pm 7	148 \pm 2	0.66	130 \pm 6	148 \pm 12	142 \pm 12
B	0.73	137 \pm 2	151 \pm 2	162 \pm 2	0.71	136 \pm 6	160 \pm 10	154 \pm 10
NUV	–	–	–	–	–	–	–	–
FUV	–	–	–	–	–	–	–	–
NGC5713			(R)SB(rs)ab: pec [SAB(rs)bc pec]			48.2	28.05	17.2
3.6 μm	0.59	20 \pm 2	22 \pm 2	27 \pm 2	0.29	20 \pm 6	22 \pm 6	28 \pm 6
R	0.62	20 \pm 2	22 \pm 2	29 \pm 12	0.31	16 \pm 6	22 \pm 6	34 \pm 6
B	0.74	20 \pm 2	20 \pm 2	27 \pm 6	0.33	16 \pm 6	22 \pm 8	40 \pm 8
NUV	–	–	–	–	–	–	–	–
FUV	–	–	–	–	–	–	–	–
NGC7552			(R)SB(rs,lr)a [(R)SB(s)ab]			23.6	20.23	52.2
3.6 μm	0.63	46 \pm 8	58 \pm 2	70 \pm 5	0.56	51 \pm 6	57 \pm 14	63 \pm 14
R	0.63	45 \pm 4	54 \pm 5	68 \pm 5	0.57	45 \pm 6	57 \pm 17	69 \pm 17
B	0.68	40 \pm 2	54 \pm 2	70 \pm 8	0.61	45 \pm 6	51 \pm 13	69 \pm 13
NUV	–	–	–	–	0.63	45 \pm 6	51 \pm 6	69 \pm 6
FUV	–	–	–	–	0.68	51 \pm 6	51 \pm 6	75 \pm 6

Notes. ^aS4G classification from Buta et al. (2010); ^bfrom Hyperleda; ^cS4G classification from Herrera-Endoqui et al. (2015).

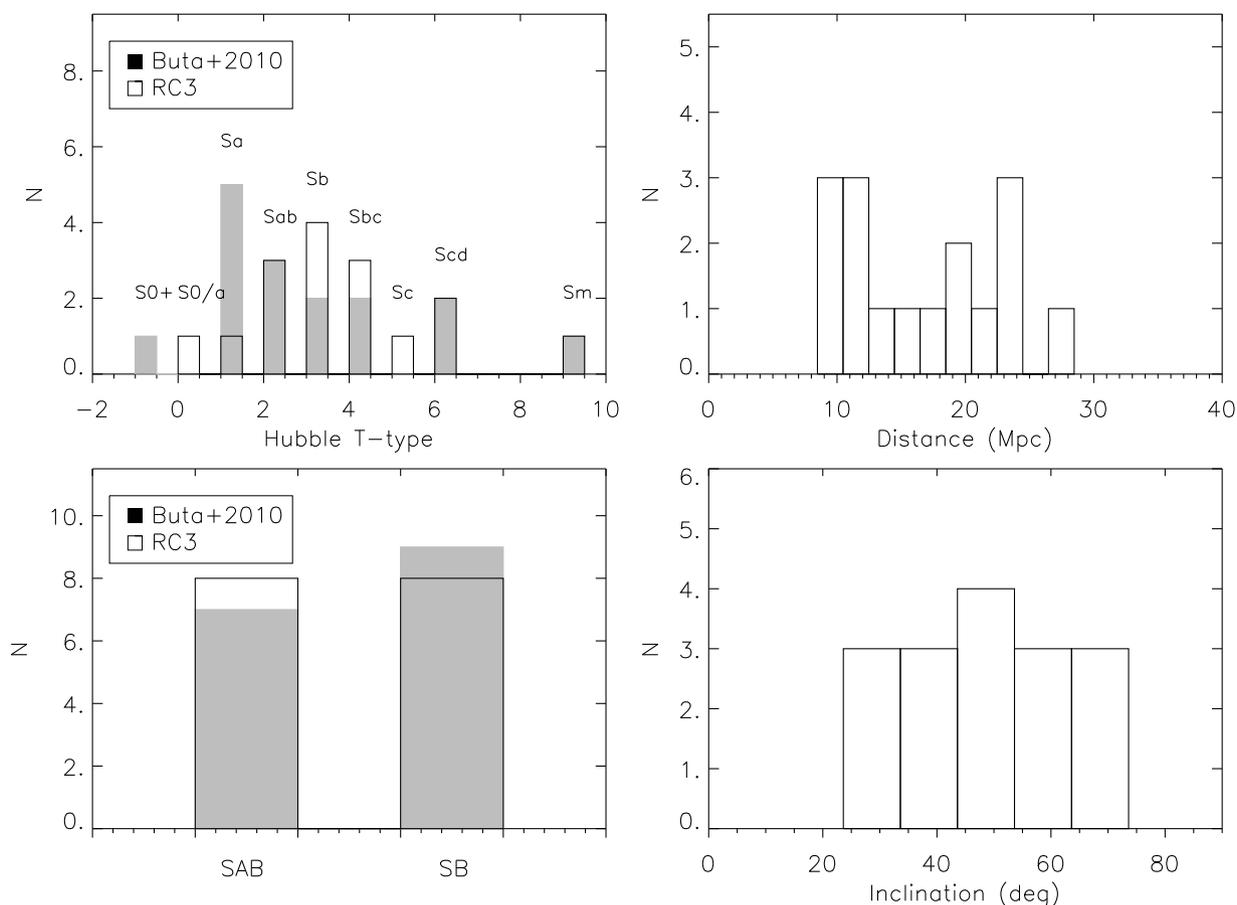


Figure 1. Distribution of properties in our galaxy sample, including (top left panel) Hubble T-type and (bottom left panel) bar type (SAB/SB) as given by RC3 and Buta et al. (2010); and (top right panel) distance and (bottom right panel) inclination. These figures show that our sample spans a wide range in galaxy properties.

We parametrize the shape of the bar in each galaxy of our sample by applying the widely used *ellipse-fitting* technique (Jedrzejewski 1987), which consists of tracing the projected light distribution with a series of elliptical isophotes across the full 2D galaxy image. For this we use the IRAF task ELLIPSE to fit elliptical isophotes of increasing semimajor axis (SMA), closely following the implementation presented by Menéndez-Delmestre et al. (2007). We used a constant SMA step equal to the resolution of the images (1.5 and 6 arcsec for the high- and low-resolution study, respectively), and allowed both the isophotes’ ellipticity (ϵ) and position angle (PA) to vary, while holding the centre of fitted ellipses fixed to the centre of the galaxy. The great advantage of this method is that while it is very efficient at picking out elliptical deviations in a galaxy’s photometry, it is fairly simple to apply to large samples and out to higher redshifts (e.g. Sheth et al. 2008).

The ellipse-fitting algorithm has its limitations: in the presence of bright star-forming knots, it may fail to converge onto a solution. This becomes a great concern at the bluest bands considered in our study, where compact star-forming regions may come to dominate the light distribution within the disc region. To circumvent this issue, we created custom-made masks for the images where the *ellipse-technique* would otherwise fail (Athanasoula et al. 1990). These masks were produced by hand, where the pixels occupied by either a foreground star or a compact star-forming region were replaced by the median value of the background adjacent to the masked region. In this manner we were able to

successfully apply the *ellipse-technique* to all the images in our study.

For each object the ensemble of isophote ellipticities (ϵ) and PA provides us with galaxy-wide profiles that allow us to recognize the signature of a bar (see Menéndez-Delmestre et al. 2007 for details): a monotonic increase in isophote ellipticity until the end of the bar is reached, at which point ellipticity drops sharply ($\Delta\epsilon \gtrsim 0.1$) to settle onto the ellipticity of the disc due to its inclination (Regan & Elmegreen 1997). Additionally, an abrupt variation in the PA profile ($\Delta\text{PA} \gtrsim 10$ deg) may accompany the signature of the end of a bar, depending on the global geometry of the system; if the disc PA is equal to that of the bar, no such PA change is observed. Example ellipticity and PA profiles for galaxies in our sample are shown in Fig. 2 for the *higher resolution optical-through-MIR* and *low-resolution UV-through-MIR* studies; a complete set of figures for our entire sample is shown in Appendix A and Appendix B.

We note that the ellipticity and PA profiles of certain galaxies display a deviation from the ideal bar signature just described. These deviations may be linked to: presence of dust lanes, leading to deviations from a monotonic increase in ellipticity within the bar region and particularly evident in the optical bands (e.g. NGC 1097, NGC 3351, NGC 3627, NGC 4321, NGC 4579, NGC 4725, NGC 7552); presence of nuclear bars, leading to a small increase in ellipticity followed by a drop close to the central regions of the galaxy (e.g. NGC 1097, NGC 1291, NGC 3198 with SMA \sim 15, 20, 10 arcsec, respectively); and the presence of prominent spiral arms,

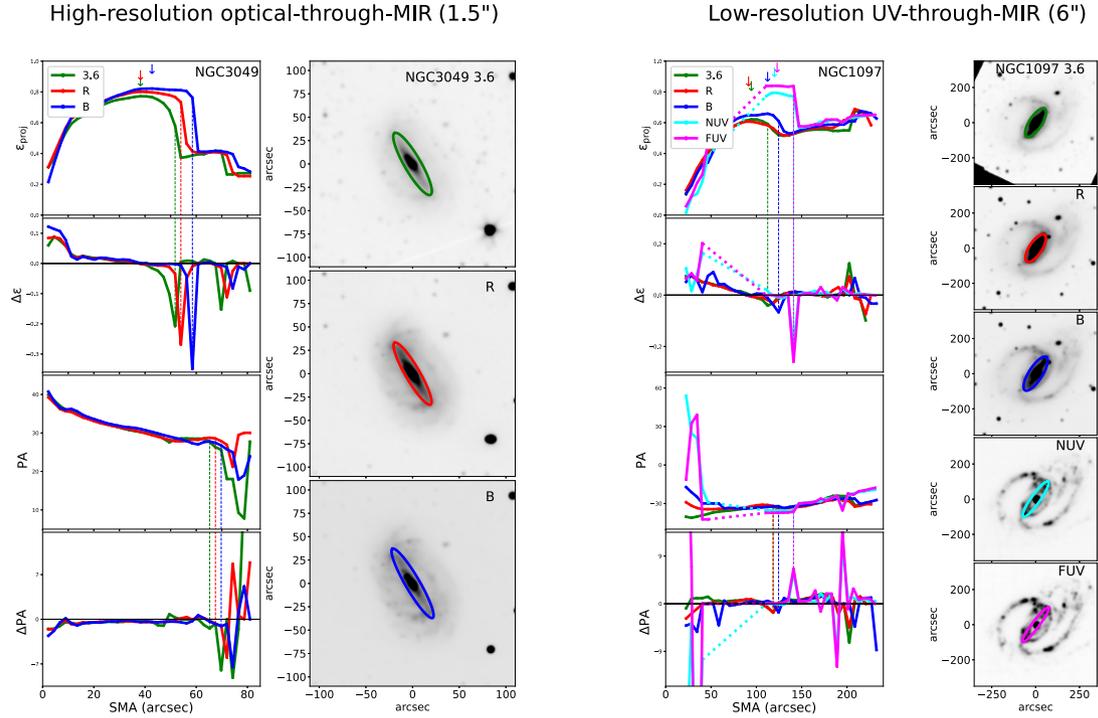


Figure 2. Ellipticity and position angle radial profiles for NGC 3049 (left) and NGC 1097 (right), the former exemplifying our methodology for the *higher resolution optical-through-MIR* study (including *B*, *R*, and $3.6\ \mu\text{m}$ bands) and the latter for the *low-resolution UV-through-MIR* study (including *FUV*, *NUV*, *B*, *R*, and $3.6\ \mu\text{m}$ bands). In both cases, we show radial profiles for ellipticity and position angle, as well as their respective variations at each isophotal fit. The vertical down-pointing arrows indicate the semimajor axis (SMA) of maximum ellipticity, i.e. $a_{\epsilon_{\text{max}}}$ for each band. The vertical dashed lines indicate the radial location of the maximum variation in ellipticity (top two panels) and PA (bottom two panels). The dashed portions of the radial profiles indicate regions where the *FUV/NUV* emission deviates significantly from a smooth light distribution, consequently impacting our isophotal ellipse-fitting approach (see Section 3). From top to bottom, postage stamp images in the $3.6\ \mu\text{m}$, *R*, and *B* bands for NGC 3049 and in the $3.6\ \mu\text{m}$, *R*, *B*, *NUV*, and *FUV* bands for NGC 1097 are shown with an overlaid ellipse showing the isophote of maximum ellipticity. Similar figures for all of our sample are shown as part of Appendix A for the *higher resolution optical-through-MIR* study and Appendix B for the *low-resolution UV-through-MIR* study.

delaying the ellipticity drop at the end of the bar and resulting in a bar signature with a broad, flat ellipticity peak (e.g. NGC 4559).

The maximum ellipticity in the bar signature corresponds to the *bar ellipticity* (ϵ_{bar}). Although the SMA at which the maximum bar ellipticity is reached (hereafter, $a_{\epsilon_{\text{max}}}$) is commonly adopted as a measure the bar length (Wozniak & Pierce 1991; Jungwiert, Combes & Axon 1997; Laine et al. 2002; Sheth et al. 2003; Menéndez-Delmestre et al. 2007), a number of alternative definitions of bar length – also based on the ellipticity and PA profiles of a galaxy – are found in the literature (e.g. Athanassoula & Misiriotis 2002; Erwin 2005; Gadotti et al. 2007).

In an effort to investigate the wavelength variation of a more complete set of commonly used measurements, we consider the following *three* alternative measures of bar length in our analysis: $a_{\epsilon_{\text{max}}}$, defined to be the SMA of the isophotal fit with the maximum ellipticity (e.g. Wozniak & Pierce 1991; Regan & Elmegreen 1997); $a_{\Delta\epsilon}$, the SMA of the isophote where the largest ellipticity drop takes place; and $a_{\Delta\text{PA}}$, the SMA of the isophote where the sharp change in PA occurs, when present (Erwin 2005). To identify the bar length based on all three definitions ($a_{\epsilon_{\text{max}}}$, $a_{\Delta\epsilon}$, and $a_{\Delta\text{PA}}$) we analyse not only the ellipticity and PA profiles, but also profiles of the *variation* in these parameters, shown in Fig. 2. Following Menéndez-Delmestre et al. (2007), to effectively identify a profile as a bar signature we required a bar to display a projected ellipticity, ϵ_{max} , greater than 0.2, and the end of the bar to be marked by a change in ellipticity, $\Delta\epsilon > 0.1$, with an optional accompanying change in the position angle $\Delta\text{PA} > 10^\circ$.

In this work, both $a_{\Delta\epsilon}$ and $a_{\Delta\text{PA}}$ are identified within a window of ~ 30 arcsec beyond where the maximum ellipticity is reached.

Based on a least-squares fitting approach, the *ellipse* task provides uncertainties for the geometric parameters of each isophotal fit. These errors, however, only reflect internal uncertainties to convey the amount of deviation from an ellipse in the light distribution at each SMA for the best-fitting ellipticity and PA combination. In an effort to include into our reported errors the uncertainty in the assumed SMA for the isophotal fit, we adopt the same approach as Menéndez-Delmestre et al. (2007) in the case of $a_{\epsilon_{\text{max}}}$. Considering that the presence of spiral arms may cause the highly elliptical isophotes that characterize the end of the bar to undergo a soft change in PA with close-to-constant ellipticity (see Fig. 2 for an example), the determination of exactly which SMA corresponds to the maximum ellipticity may become uncertain. Menéndez-Delmestre et al. (2007) quantify the error in the bar size measurement by considering the SMA range encompassing the tip of the ellipticity peak in the bar signature; that is, the region over which $\epsilon > (\epsilon_{\text{max}} - \delta\epsilon)$, where $\delta\epsilon = 0.01$, a value we adopt as the uncertainty on the ellipticity of the bar. We adopt this value as it corresponds to the maximum error found by the least-squares fit *ellipse* task for the ellipticity parameter. For the two alternative measures of bar length – $a_{\Delta\epsilon}$ and $a_{\Delta\text{PA}}$ – we adopt the following approach. Although $a_{\Delta\epsilon}$ is defined as the SMA of the isophote where the largest ellipticity drop takes place, the ellipticity change is not necessarily sharp, extending over a few arcsecs (see ellipticity profile of NGC 3049 as traced by the $3.6\ \mu\text{m}$ band in Fig. 2);

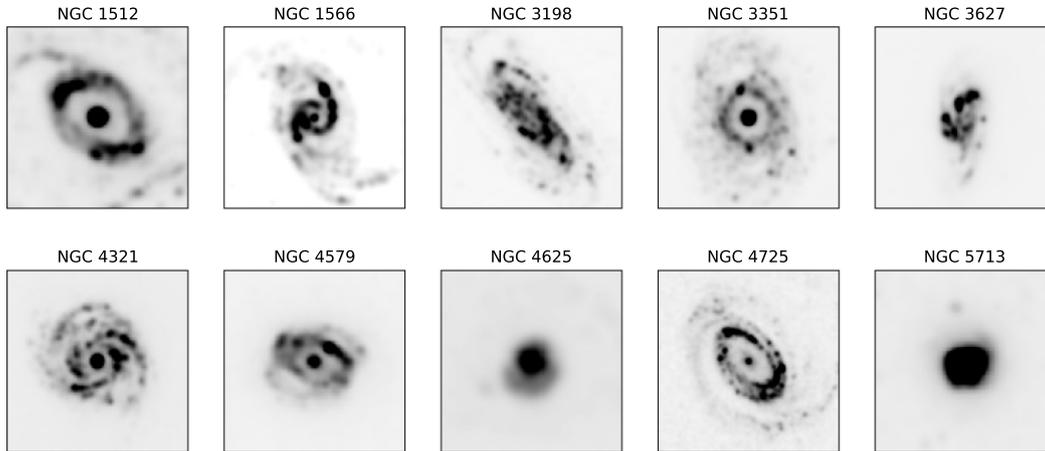


Figure 3. GALEX *NUV* images of the 10 galaxies in our sample whose bars become untraceable in the ultraviolet bands.

we consider the SMA range over which the ellipticity changes around the ellipticity drop at the end of the bar and associate an uncertainty to the bar length equivalent to the full width at half-maximum (FWHM) of a Gaussian fit to the ellipticity variation. In the case of extra-sharp drops in ellipticity, we establish a minimum uncertainty equivalent to one resolution element (1.5 arcsec in the *higher resolution optical-through-MIR* and 6 arcsec in the *low-resolution UV-through-MIR*). We define an equivalent approach to the uncertainties in bar length measures based on $a_{\Delta PA}$, considering the variation in the PA profile.

We note that in this study we are mainly concerned with the *relative* bar properties in the different bands at which each galaxy is studied. For this reason, we present the bar properties prior to any correction for inclination (i.e. deprojection). For the scope of this work, we compare the measured *projected* bar sizes and ellipticities for each galaxy from the UV to the MIR as part of the *low-resolution UV-through-MIR* study and from the optical to the MIR as part of the *higher resolution optical-through-MIR* study.

4 RESULTS

As part of the *low-resolution UV-through-MIR* study, we apply the above method to the entire array of *FUV*-, *NUV*-, *B*-, *R*-, and 3.6 μm -band images at the common resolution of ~ 6 arcsec for the 16 barred galaxies in our sample. We repeat this analysis for the same galaxies using the restricted set of *B*-, *R*-, and 3.6 μm -band images at the finer resolution of ~ 1.5 arcsec as part of the *higher resolution* study. The resulting collection of individual bar measurements for the *higher resolution optical-through-MIR* and the *low-resolution UV-through-MIR* studies are shown in Table 1.

4.1 Half of all Bars disappear in the UV

The first result of this analysis is that there are multiple cases where we lose the ability to identify bars in the *UV* bands. Fig. 3 shows the 9 out of the 16 galaxies in our sample with bars that disappear in the *NUV* images; an additional bar disappears when we go out to the *FUV* bands (NGC 4321). This represents 56 and 63 percent of the full sample for the *NUV* and *FUV* bands, respectively.

In many of these cases (e.g. NGC 1512, NGC 3351, NGC 4321, NGC 4579, NGC 4725), the ellipticity profiles in the UV bands trace what could be erroneously identified as a bar signature: a maximum in ellipticity, followed by a sharp drop. However, careful examination of the images shows that this ellipticity maximum does not spatially

coincide with the end of the bar as seen in the optical or MIR, but to nodes of star formation beyond the bar region, at times significantly offset from the bar PA (e.g. NGC 1512, NGC 3351; see Appendix B).

The images on the *FUV* and *NUV* bands often present great challenges for the implementation of our methodology. Compact star-forming regions dominate the brightness distribution, hindering the convergence of the *ellipse* routine, particularly within the central regions of the galaxy (e.g. NGC 1512, NGC 1566, NGC 3351, NGC 4321, NGC 4579, NGC 4725). Therefore, these ‘fake’ bar signatures, although displaying a maximum in ellipticity followed by a sharp drop are never preceded by the monotonic increase that is characteristic of the ellipticity profile in a real bar (see Appendix B).

4.2 Bar Properties as a function of waveband

In Figs 4–7, we present the measured bar length and ellipticity as a function of wavelength for both the *higher resolution optical-through-MIR* and the *low-resolution UV-through-MIR* studies. The former, based on higher resolution images, enables a finer evaluation of the variation in bar properties from optical-through-MIR, while the latter allows us to extend our bar characterization to the rest-frame UV. We choose to not elaborate on the optical-through-MIR behaviour within the lower resolution study, as this is done at a finer resolution within the higher resolution study.

The bar length measurements on Table 1 show the following systematics for the bulk of our sample: $a_{\epsilon_{\max}} \lesssim a_{\Delta\epsilon} \lesssim a_{\Delta PA}$. These trends are a direct consequence of our adopted bar length definitions. The location of maximum ellipticity in a galaxy’s ellipticity profile necessarily precedes (or at the closest, coincides) with the ellipticity drop, which in turn means that the bar length measured at the location of maximum ellipticity ($a_{\epsilon_{\max}}$) is equal to or lower than that measured at the ellipticity drop ($a_{\Delta\epsilon}$): $a_{\epsilon_{\max}} \lesssim a_{\Delta\epsilon}$. Since the PA change – when present – is typically found at larger SMA values beyond the ellipticity drop (see Fig. 2), the bar length measured at the location of PA variation ($a_{\Delta PA}$) is (in most cases) significantly larger than either the bar length at maximum ellipticity or at the ellipticity drop. For completeness, we note that our ellipse-based bar length measurements $a_{\epsilon_{\max}}$ are in good agreement with other visual bar length measurements (e.g. Herrera-Endoqui et al. 2015 based on 3.6 μm band S^4G images), which we have included for the reader’s appreciation in Table 1.

In Fig. 4, we compare our bar length measurements from the optical through the MIR, based on our three definitions. The bar

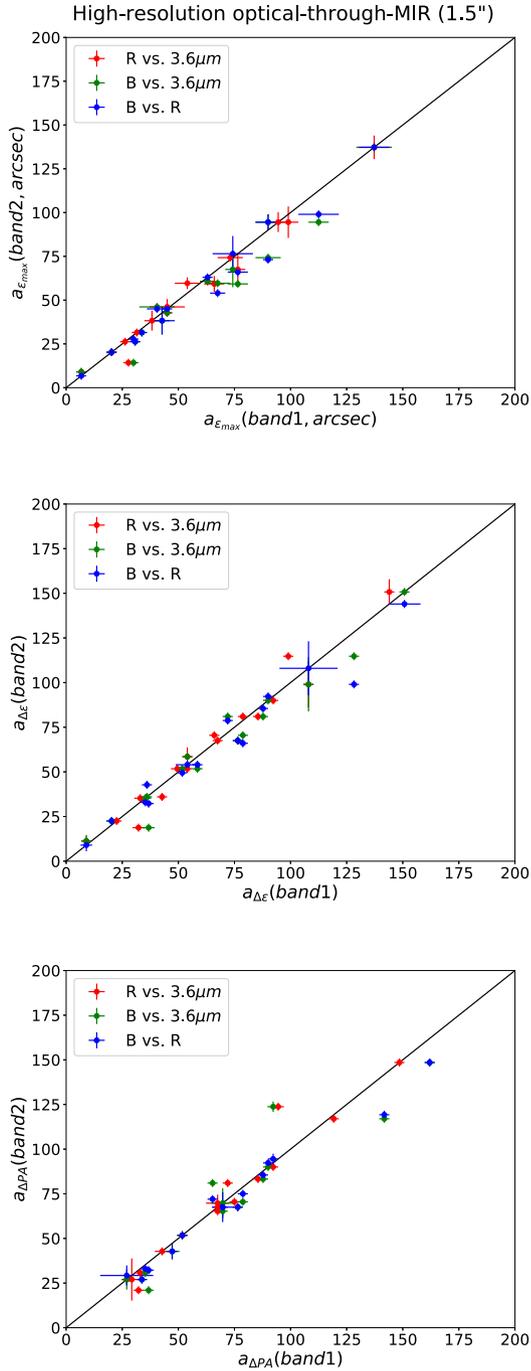


Figure 4. Measured bar sizes as a function of wavelength for the galaxies in our sample; each panel corresponds to one of the three bar size definitions adopted for the analysis (see Section 3 for details). Each galaxy is represented by three data points, where the value on the x -axis corresponds to that measured in the bluer band of the waveband pair being considered. We note that the bulk of individual bar size measurements in the B , R , and $3.6 \mu\text{m}$ bands is offset downwards from the identity line (solid line), indicating that bar sizes measured in bluer bands are larger. On average, the bar is measured to be ~ 9 per cent longer in the B band compared to that in the $3.6 \mu\text{m}$.

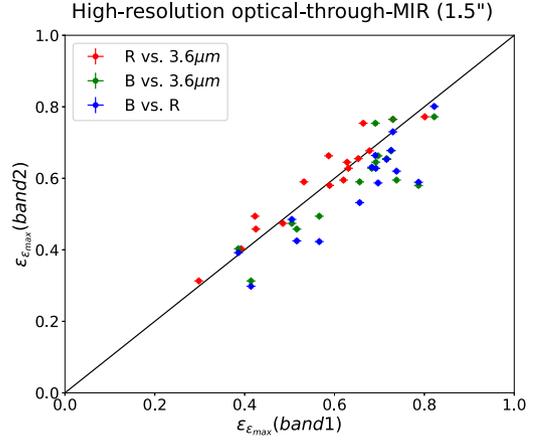


Figure 5. Following a similar format to that of Fig. 4, distribution of measured bar ellipticities in our sample. The individual measurements of ellipticities are generally offset downwards from the unity line (solid line), indicating that bar ellipticities measured in the B band are for the most part higher than those measured in the R and $3.6 \mu\text{m}$ bands.

length measurements in different bands are set against each other, with the bar length measured in the bluer band displayed in the x -axis. Most of our galaxies lie below the line of equality, pointing to bars being longer at bluer bands. The only exception to this trend is when comparing the two reddest bands (R and $3.6 \mu\text{m}$ bands) where the measured lengths are roughly consistent with each other, suggesting that similar stellar populations are traced at these bands. Based on median values measured as part of our *higher resolution optical-through-MIR* study, a bar measured following the ellipticity-maximum method is measured to be 3.6 per cent longer in the B -band compared to in the R -band and 8.6 per cent longer in the B band compared to that measured in the $3.6 \mu\text{m}$ (see Table 2).

In Fig. 5, we compare the measured bar ellipticities for each galaxy in the B , R , and $3.6 \mu\text{m}$ bands and find that the bar ellipticity is measured to be consistently higher in the B - than in the R - and $3.6 \mu\text{m}$ bands. We find that ~ 90 per cent (~ 80 per cent) of the sample have measured ellipticity ratios $\epsilon_B/\epsilon_R > 1$ ($\epsilon_B/\epsilon_{3.6\mu\text{m}} > 1$). With such skewed distributions towards high ellipticity ratios, this indicates that bar ellipticity is measured to be higher in the B band than in the redder R and $3.6 \mu\text{m}$ bands, or, equivalently, that bars appear thinner in the B band. We find that the ellipticity of a bar is typically measured to be ~ 8 – 10 per cent higher in the B band than in the $3.6 \mu\text{m}$ and R bands. However, in ~ 30 per cent of our sample this difference increases to $\gtrsim 20$ – 35 per cent.

We summarize our results for the *higher resolution optical-through-MIR* study – both for bar length and bar strength multiband measurements – in Table 2, where we also display the results of statistical tests used to evaluate the significance of these results. We perform two different statistical tests: a paired difference test and a Wilcoxon signed-rank test. The two serve a similar purpose, although the latter does not assume the data to be normally distributed. With these tests we essentially seek to verify whether bar length and bar strength measurements in different bands yield distinct means, i.e. if a set of galaxies is measured to have bars with statistically significant longer semimajor axis lengths (or higher ellipticity) in a given band with respect to another. The resulting p -values convey the probability of the null-hypothesis that the samples have identical means. We consider a threshold of $p < 0.05$ to exclude this hypothesis. In Table 2, we quantify (in terms of per cent) the impact that looking at different bands has on the bar length ($a_{\epsilon_{\text{max}}}$, $a_{\Delta\epsilon}$) and the bar ellipticity for all

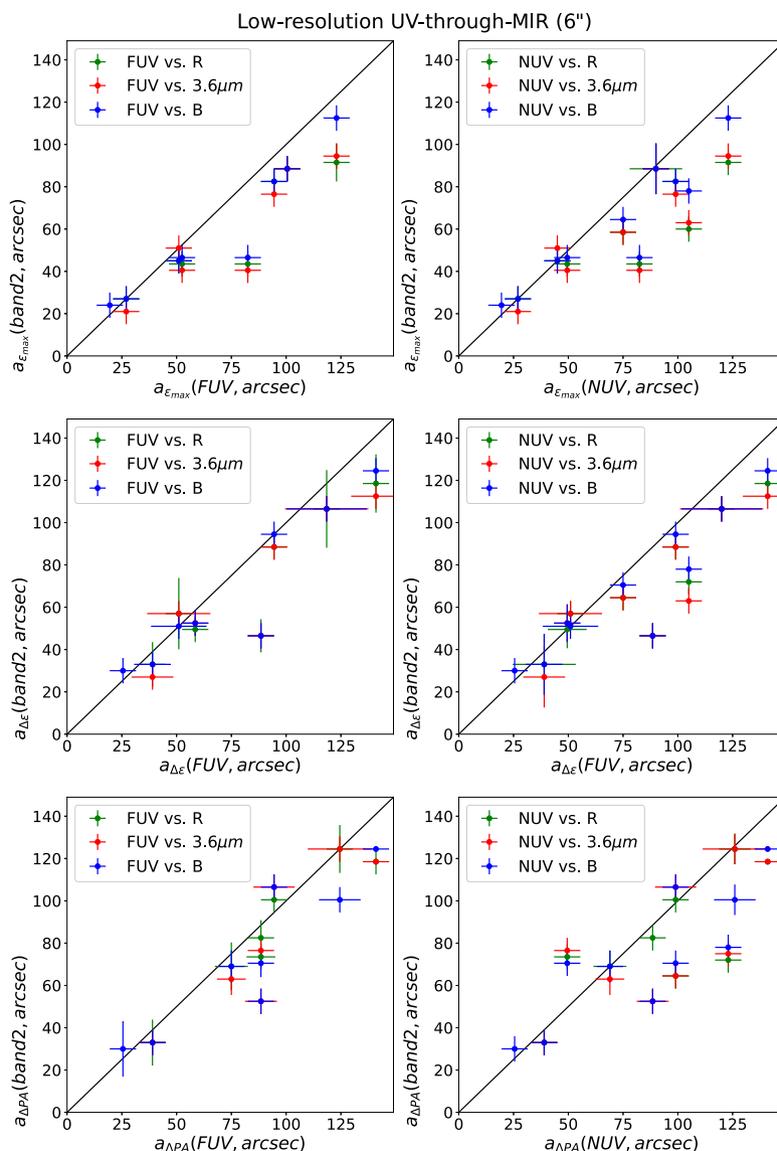


Figure 6. Following a similar format to that of Fig. 4, distribution of measured bar lengths in our sample based on the *low-resolution UV-through-MIR* study. Bar lengths measured in the *B*, *R*, and $3.6\ \mu\text{m}$ bands are shown relative to those measured in the *NUV* (left) and *FUV* (right) bands, for the 7(6) galaxies in our sample where bars are still identifiable in the *NUV* (*FUV*). These plots show that the results found for the *higher resolution optical-through-MIR* study extend out to *UV* bands.

band pairs, as well as the p -values for the Wilcoxon statistical test. We complement the results shown in Table 2 for both statistical tests in Appendix C. Based on this analysis, Table 2 shows that band-shifting effects are statistically significant when comparing bar length and ellipticity (strength) measurements in the *B* band and in the $3.6\ \mu\text{m}$ band. To a first-order approximation, we may put forward a correction for the band-shifting effects on the measurement of bar properties when considering the rest frame *B* and $3.6\ \mu\text{m}$ bands of 9 per cent and 8 per cent for the bar length and bar ellipticity, respectively.

We explored the potential difference in band-shifting effects for stronger versus weaker bars, as well as for longer versus smaller bars. However, we find no correlation between either of these bar properties (length or ellipticity) and the shifts caused by making these measurements at different bands.

To evaluate whether these results extend to the rest-frame UV, we consider the bar measurements based on the *low-resolution UV-through-MIR* study. An immediate consequence of degrading the

resolution is that, although bar lengths are not significantly affected (variation typically within ~ 5 per cent; see Table 1) the smaller bars in our sample are washed out (e.g. NGC 4625; see Appendix B). Furthermore, the measured ellipticities drop by ~ 10 per cent at the coarser resolution, compared to the measurements made within the same band for the finer resolution of the *optical-through-MIR* study.

Following the same format of Figs 4 and 5, we show in Figs 6 and 7 the bar length and ellipticity measurements for the *low-resolution UV-through-MIR* study. The great majority of bars are measured to be longer in the *FUV* and *NUV*, compared to measurements in the optical and MIR, to a degree significantly higher than that in the *higher resolution optical-through-MIR* study. We measure bars to be ~ 8 per cent longer in the *NUV* than in the *R* band (see Table 3); however, this increases to ~ 25 per cent when comparing to that measured in the $3.6\ \mu\text{m}$ band. In terms of ellipticity, we observe that the bar ellipticity in the *NUV* band is typically ~ 20 per cent higher when comparing to values determined in the *R* and $3.6\ \mu\text{m}$ bands.

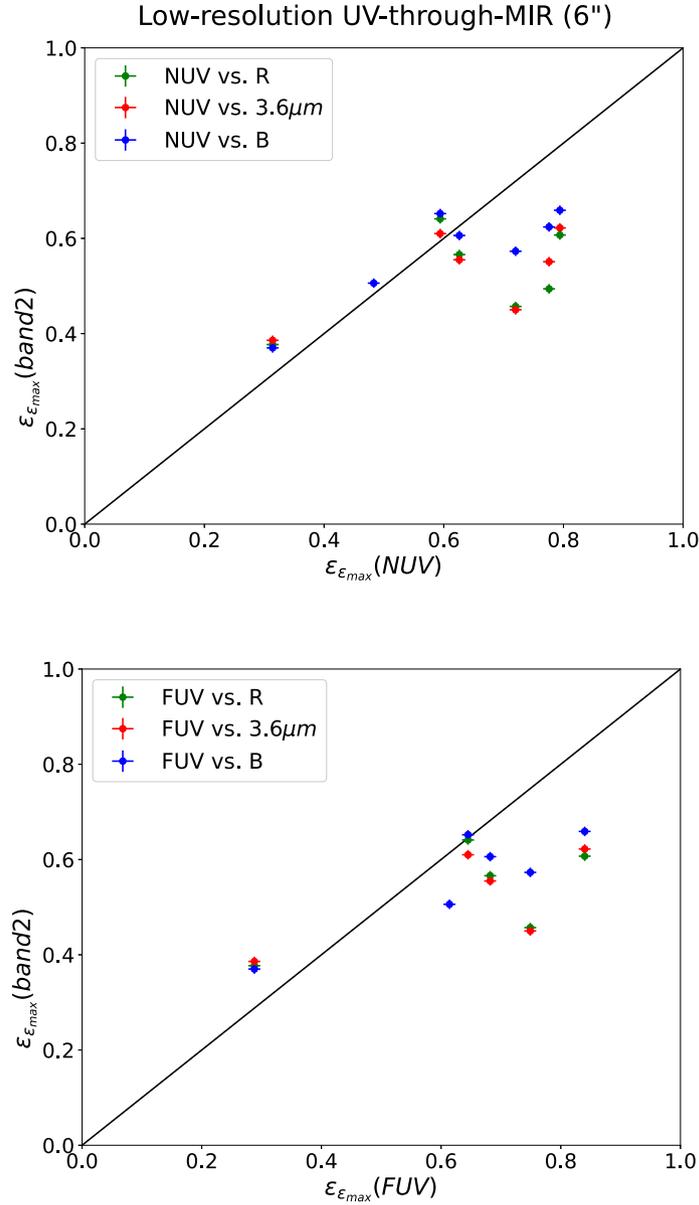


Figure 7. Following a similar format to that of Fig. 6, distribution of measured bar ellipticities in our sample based on the *low-resolution UV-through-MIR* study.

Table 2. Impact of measuring bar length and ellipticity in different bands for the *higher resolution optical-through-MIR* study – sample average of relative bar measurements and statistical significance.

Band	<i>R</i>	3.6
<i>B</i>	$a_{\epsilon_{\max}}$: 3.6 per cent (0.073)	$a_{\epsilon_{\max}}$: 8.6 per cent (0.023)
	$a_{\Delta\epsilon}$: 3.6 per cent (0.130)	$a_{\Delta\epsilon}$: 0.0 per cent (0.090)
	ϵ_{\max} : 9.9 per cent (0.001)	ϵ_{\max} : 8.0 per cent (0.005)
<i>R</i>		$a_{\epsilon_{\max}}$: 0.0 per cent (0.167)
		$a_{\Delta\epsilon}$: -1.4 per cent (0.899)
		ϵ_{\max} : -2.7 per cent (0.065)

Note. Percentages associated with the mean ratios between measurements for blue and red bands, where the bluer band is in the numerator (e.g. $a_{\epsilon_{\max}}$ is measured to be 3.6 per cent longer in the *B* band compared to in the *R* band and 8.6 per cent longer in the *B*-band compared to that measured in the 3.6 μm . Numbers in parentheses indicate the p value for a Wilcoxon test of two related samples, investigating the difference between both measurements. Smaller values reinforce the probability that one sample has a larger values than the other. See Section 4.2 for details.

The p -values do not exclude the possibility of these difference being just a result of a random selection of the same parent population; however, we note that this trend is merely based on 7 (6) galaxies with bars still identifiable at *NUV* (*FUV*) bands, placing this result in the realm of small-number statistics. We provide context and discuss the impact of these results in the next section.

We note that although there is a tendency for bars to appear slightly longer and thinner in the *FUV* band – compared to the *NUV* band, these differences fall within the uncertainties imposed by the 6 arcsec resolution we adopted for the *low-resolution UV-through-MIR* study.

5 DISCUSSION

5.1 Bars in the UV

Half of all bars in our sample disappear when viewed in the rest-frame UV. At $\lambda_{FUV, NUV} = 1516 \text{ \AA}, 2267 \text{ \AA}$, the light traced by GALEX *FUV*, *NUV* bands is dominated by young OB star complexes along spiral

Table 3. Impact of measuring bar length and ellipticity in different bands for the *low resolution UV-through-MIR study* – sample average of relative bar measurements and statistical significance.

Band	<i>NUV</i>	<i>B</i>	<i>R</i>	3.6
<i>FUV</i>	$a_{\epsilon\max}$: 3.0 per cent (0.109)	$a_{\epsilon\max}$: 11.1 per cent (0.078)	$a_{\epsilon\max}$: 13.6 per cent (0.068)	$a_{\epsilon\max}$: 28.6 per cent (0.066)
	$a_{\Delta\epsilon}$: 0.0 per cent (0.655)	$a_{\Delta\epsilon}$: 11.3 per cent (0.078)	$a_{\Delta\epsilon}$: 18.2 per cent (0.188)	$a_{\Delta\epsilon}$: 11.4 per cent (0.188)
	ϵ_{\max} : 7.2 per cent (0.062)	ϵ_{\max} : 16.9 per cent (0.219)	ϵ_{\max} : 20.5 per cent (0.188)	ϵ_{\max} : 22.9 per cent (0.188)
<i>NUV</i>		$a_{\epsilon\max}$: 1.7 per cent (0.225)	$a_{\epsilon\max}$: 7.7 per cent (0.068)	$a_{\epsilon\max}$: 25.4 per cent (0.156)
		$a_{\Delta\epsilon}$: 12.7 per cent (0.116)	$a_{\Delta\epsilon}$: 15.4 per cent (0.104)	$a_{\Delta\epsilon}$: 19.0 per cent (0.156)
		ϵ_{\max} : 3.3 per cent (0.469)	ϵ_{\max} : 20.7 per cent (0.219)	ϵ_{\max} : 20.2 per cent (0.219)
<i>B</i>			$a_{\epsilon\max}$: 0.0 per cent (0.027)	$a_{\epsilon\max}$: 8.2 per cent (0.014)
			$a_{\Delta\epsilon}$: 2.5 per cent (0.048)	$a_{\Delta\epsilon}$: 3.4 per cent (0.031)
			ϵ_{\max} : 6.7 per cent (0.000)	ϵ_{\max} : 6.8 per cent (0.049)
<i>R</i>				$a_{\epsilon\max}$: 0.0 per cent (0.465)
				$a_{\Delta\epsilon}$: 0.0 per cent (0.141)
				ϵ_{\max} : -1.6 per cent (0.194)

Note. Same as Table 2 for the low-resolution images.

arms, inner and outer star-forming rings and nuclear emission from an AGN or nuclear starburst (see Fig. 3). Although various works have pointed to star formation found associated with the bar structure (e.g. Díaz-García et al. 2020; Fraser-McKelvie et al. 2020), bar stellar orbits are mainly occupied by old stellar populations that contribute to the emission redwards of the Balmer discontinuity (3646 Å) and 4000 Å break. Hence, our ability to recognize the presence of a bar in a galaxy is compromised when our analysis is based on rest-frame UV imaging, even considering the relatively deep coverage provided for the galaxies in our sample by the GALEX NGS (Gil de Paz et al. 2004; Bianchi et al. 2003a, b), down to a surface brightness of $\mu_{FUV,NUV} = 27.25$ mag arcsec⁻², 27.35 mag arcsec⁻², respectively (Thilker et al. 2007). According to Erwin, Pohlen & Beckman (2008), typical surface brightness values at the outer extent of galaxy bars are approximately $\mu_R \sim 21$ –22 mag arcsec⁻². Considering that *NUV* – *r* colours of galaxies can vary between 1 and 2 mag for bluer systems dominated by young stellar populations to 5–6 mag for early type galaxies dominated by older stellar populations (Wyder et al. 2007), bars hosting young stellar populations should be easily detectable in the UV at our detection limit of $\mu_{NUV} = 27.4$ mag arcsec⁻². However, this would be close to the noise level for older stellar populations. We conclude that a significant fraction of bars are bound to disappear, since they are stellar structures traditionally hosted by older stellar populations. Indeed, Díaz-García et al. (2020) have shown, for similar detection limits, that UV emission is detected in bars which host star-forming regions. The UV emission along these structures will not be detected in passive bars unless surface brightness limits are considerably deeper (i.e. 5 mag) than the optical data at hand. In order to identify the stars that typically support the bar structure, we must probe emission redwards of the Balmer/4000 Å breaks.

The impaired ability to trace bars in the UV had been previously emphasized by Sheth et al. (2003), where the authors identify an acute drop in the local bar fraction based on the visual identification of bars in 139 local spirals using SDSS *u* band ($\lambda_{SDSS\ u} = 3453$ Å) imaging. Considering that the evolution of the bar fraction with redshift has been the focus of numerous studies (Abraham et al. 1999; Sheth et al. 2003; Elmegreen, Elmegreen & Hirst 2004; Jogee et al. 2004; Sheth et al. 2008; Melvin et al. 2014), the artificial paucity of bars in the UV acquires critical importance. At a fixed photometric band the morphology of a hypothetical barred galaxy at increasingly higher redshifts is imaged at incrementally bluer rest-frame wavebands. Therefore, it is imperative that we restrict the analysis to redshifts where we probe emission redwards of the Balmer/4000 Å breaks.

5.2 Bars appear thinner and longer in bluer bands

Our results show that at bluer wavebands both the measured bar length and the bar ellipticity increase. Although we find that 50–60 per cent of the bars disappear in the UV, the results on bar ellipticity and length extend to those cases in which the bar is still visible in the UV. The increase in maximum ellipticity and the location of the bar end further away from the galaxy centre can be clearly seen in the ellipticity profiles for NGC 1097 in Fig. 2.

At face value, we attribute the increase in bar length towards bluer bands to the prominent star-forming knots that are frequently present at the end of bars. While the contribution from bar stellar populations diminishes as we probe bluer emission, these star-forming regions at both ends of the bar become significantly brighter. The fitted isophotes stretch further out into the disc, resulting in the ‘artificial lengthening’ of the bar signature. The trend found for the bar ellipticity may be driven by the apparent bulge size at the different wavebands, since the semiminor axis of the measured bar isophotes are limited by the size of the central bulge. In this manner, a larger, more prominent bulge reduces the bar ellipticity (see also Díaz-García et al. 2016). Bulges are composed primarily of older, redder stars and are thus significantly more prominent – with a larger apparent size – at longer wavelengths (e.g. Möllenhoff 2004). For a given barred galaxy, at shorter wavelengths the bulge appears smaller than at longer wavelengths and thus the bar isophotes may reach higher ellipticities in the bluer bands. Interestingly, Speltincx, Laurikainen & Salo (2008) find an increase of ~ 25 per cent in bar strength from *H* band to *B* band for 152 galaxies from the Ohio State University Bright Spiral Galaxy Survey, based on the gravitational bar torque method (Q_b , the maximum tangential force normalized by the radial force). These authors also attribute this increase to the diminished prominence of the bulge in optical bands, as the radial forces that it introduces may be underestimated, leading to an increase in the measured bar strength. The impact of an increased contribution from the bulge, with radially oriented stellar orbits, has also been the centre of discussion when considering the reduced ellipticities associated with bars in lenticular galaxies (e.g. Speltincx, Laurikainen & Salo 2008; Buta et al. 2010). Work by Hilmi et al. (2020), based on simulated Milky Way (MW)-like galaxies formed in a cosmological context, also present findings pointing to longer and stronger bars when the bar ends are connected to the spiral structure.

We note that boxy/peanut bulges (see Athanassoula 2016 for a review), which have been shown to reside in the majority of massive disc galaxies (Erwin & Debattista 2017), have been identified

in NGC 1097, NGC 1291, NGC 1512, NGC 3627, and NGC 4725 (Erwin & Debattista 2013, 2017; Laurikainen et al. 2014). These structures, identified based on the isophotes in moderately inclined galaxies (Erwin & Debattista 2013), may in principle affect bar ellipticity measurements, due to the projection of a vertically extended bar structure. Although the expectation would be that these boxy/peanut bulges – which typically host older stellar populations – present rounder isophotes, with lower ellipticities, the galaxies in our sample with boxy/peanut bulges do not appear to present a particular trend within our general results. We consider that both the presence of boxy/peanut bulges, as well as the variation in bulge prominence associated with a diverse range of morphological Hubble types, may also present interesting trends. Such an analysis, however, calls for a more thorough analysis based on a broader sample and goes beyond the scope of this paper.

Results from recent works by Neumann et al. (2020) and Bittner et al. (2021) show that the stellar populations in bars are on average older and more metal-rich than the surrounding regions of the galaxy – both of which contribute to redder colours. They also show that along the major axis of the bar one finds, on average, younger stellar populations compared to the sides/edges of the bar, due to kinematic fractionation (Athanasoula, Rodionov & Prantzos 2017; Fragkoudi et al. 2017), where bars can separate stellar populations that differ by their initial disc kinematics (Debattista et al. 2017). This leads to the ends of the bar showing younger stars because they will have, on average, more elongated orbits than older stars. So even in the absence of recent star formation we may expect to see longer and more eccentric bars in the bluer wavelengths. These findings suggest that the higher values of bar length and ellipticity we find in bluer bands may not necessarily be ‘artificial’, as the young stars may indeed be in bar orbits.

We note that the increase in bar ellipticity (and bar strength) at bluer bands is to some extent related with the *ellipse-fit* method (and gravitational bar torque method) used. A full-2D decomposition of the bar and bulge components would allow for the bar’s ellipticity and strength to be measured with no contamination from the bulge. This has been in fact performed by Gadotti (2008) using the BUDDA decomposition code (de Souza, Gadotti & dos Anjos 2004) on *V*- and *R*-band images for a similarly sized sample of 17 barred galaxies, where the authors conclude that the bar ellipticity is typically underestimated by ~ 20 per cent when based on the peak in the ellipticity profile. This supports the use of complex methods to identify and characterize bars that are based on modelling and breakdown of galaxy components to obtain enriching results for individual galaxies and small samples; e.g. BUDDA (de Souza et al. 2004), GALFIT (Peng et al. 2002), GASP2D (Méndez-Abreu et al. 2008, 2014). Although these complex 2D decompositions of barred galaxies are hard to implement for large samples, several works have analysed significant samples: e.g. ~ 3500 galaxies in Kruk et al. (2018), ~ 800 galaxies in S^4G (Salo et al. 2015), 291 galaxies in Gadotti (2009), 162 galaxies with GASP2D (Méndez-Abreu et al. 2017). However, these decomposition approaches are not necessarily practical nor efficient means of identifying and characterizing bars at high redshift, where samples are plagued with low signal-to-noise imaging and surface-brightness dimming. Precisely the relative simplicity of the ellipse-fit methodology – the analysis of one-dimensional (1D) profiles of ellipticity and position angle adjustment based on the elliptical isophotes in 2D brightness distribution of barred galaxies – allows it to be extended efficiently to large samples out to the distant Universe. For this very reason, studies have in the past years extended the use of the ellipse-fit method towards large samples (e.g. Regan & Elmegreen 1997; Zheng et al. 2005; Knapen, Shlosman & Peletier

2000; Laine et al. 2002; Laurikainen & Salo 2002; Sheth et al. 2002, 2003; Marinova & Jogee 2007; Menéndez-Delmestre et al. 2007; Aguerri, Méndez-Abreu & Corsini 2009; Cheung et al. 2013) and at increasingly earlier cosmological times (e.g. Sheth et al. 2008; Barazza et al. 2009). We must therefore consider with great care the caveats that the widely applied ellipse-fit method entail.

Our results show that measuring bar properties at different bands introduces potentially artificial effects on the measured ellipticity and length. In going from the 3.6 μm - to the *B* band, bars appear longer and thinner by ~ 9 per cent and ~ 8 per cent, respectively; extending the waveband range further to the blue, bars increase their median length and ellipticity at *NUV/FUV* bands by ~ 25 and ~ 20 per cent with respect to their canonical ‘true values’ at 3.6 μm . These variations are significant in the light of the reported differences in bar properties with respect to: (1) Hubble type – with early-type bars displaying higher ellipticities than late types by ~ 10 – 20 per cent (e.g. Menéndez-Delmestre et al. 2007; Aguerri, Méndez-Abreu & Corsini 2009; Hoyle et al. 2011); (2) environment – e.g. Barazza et al. (2009) find that bars in cluster environments are longer than in the field; and (3) active versus normal galaxy type – where an increase in bar ellipticity of ~ 30 per cent and bar strength Q_b of ~ 10 per cent has been measured in non-active galaxies, with respect to active ones (Seyferts, starbursts) by Chapelon, Contini & Davoust (1999) and Laurikainen, Salo & Buta (2004), respectively. The fact that measuring bar properties in different bands may lead to similar variations that are found when analysing the impact that other galaxy properties may have on bar properties, alerts us on the cautions that must be taken when discussing these intrinsic changes in bar structural parameters at different cosmological times, free of band-shifting effects.

The bar length measurement is also important in the context of the bar pattern speed, where the latter is often parametrized by the bar rotation rate, defined as the ratio between the corotation resonance (CR) radius R_{CR} and the bar semimajor axis R_{bar} , R_{CR}/R_{bar} (Debattista & Sellwood 2000). Considering that the bar interacts with other galaxy components, exchanging energy, mass and angular momentum, *N*-body simulations show that as bars grow, they also slow down (e.g. Weinberg 1985; Debattista & Sellwood 2000; Martínez-Valpuesta, Shlosman & Heller 2006; Athanasoula, Machado & Rodionov 2013). Although this result seemed observationally supported at first (e.g. Aguerri, Beckman & Prieto 1998; Rautiainen, Salo & Laurikainen 2008; Garma-Oehmichen et al. 2020), more recent work has pointed to bars staying fast through their lifetimes (e.g. Aguerri et al. 2015; Guo et al. 2019), bringing tension into the picture with cosmological simulations (e.g. Algorry et al. 2017 based on EAGLE and Peschken & Łokas 2019 based on Illustris). Recent work by Fragkoudi et al. (2021) adds to this context, suggesting that simulations reproduce the observed tendency for fast bars all the way to $z \sim 0$ when considering less prominent dark matter haloes that in turn act less as an angular momentum sink, limiting significantly the bar slow-down predicted by earlier simulations. Within this context, although the 9 per cent difference in bar length that we report here, when measuring in the *B* and 3.6 μm bands, is unlikely to make a difference between studies reporting different populations of slow and fast bars, our results suggests that future studies calculating bar pattern speeds should do so based on consistent measurement of bar lengths in the same band.

5.3 Corrections for high-redshift bar studies

Theoretical work shows that a bar grows more elliptical and longer with time (e.g. Sellwood 1981; Athanasoula 2003), which in turn

reflects changes within the underlying disc dynamical structure. Hence, an evolution in the median structural parameters of bars with cosmological time may be interpreted as an indication of a transformation in the dynamical maturity of discs. Sheth et al. (2008) shows that bars in the most massive galaxies ($>10^{10.5}M_{\text{Sun}}$) are already in place at $z \sim 0.8$, while active bar growth (via an increase of the bar fraction with decreasing redshift) is still today taking place in the less massive systems ($<10^{10.5}M_{\text{Sun}}$). Within the context of bars being long-lived structures (e.g. Kraljic, Bournaud & Martig 2012), this suggests that the median bar grows – both in size and strength – towards lower redshifts.

Considering that at a fixed photometric band the morphology of a hypothetical barred galaxy at increasingly higher redshifts is imaged at incrementally bluer rest-frame wavebands, our results indicate that the bar size and the bar ellipticity will suffer a systematic shift towards higher values. That is, unless we make careful corrections for band-shifting effects, we may perceive a shift in the bar length and bar ellipticity distributions towards *longer* and *thinner* – thus *stronger* – bars in the high-redshift population. These potentially ‘artificial’ lengthening and thinning of the bar structure with redshift could mask any intrinsic evolution of the bar structural parameters. Barazza et al. (2009), based on the ellipse-fitting technique applied to HST/ACS *i*-band imaging of 63 barred galaxies, undertook the pioneering investigation of bar properties out to intermediate redshifts; they found no variation in the median bar size and ellipticity in three redshift bins: $0.4 < z < 0.55$, $0.55 < z < 0.7$, $0.7 < z < 0.8$. Although within this redshift range the *i*-band merely shifts from rest frame *V* to *B* band, the potential impact of band-shifting effects in studies comparing bar properties in different redshift bins such as this one has not yet been well discussed. The results of our study caution high-redshift bar studies not merely against relying on rest-frame UV imaging when identifying the presence of a bar, but also against tracing the evolution in bar properties – even in the optical and near-IR – at face value, with no consideration of band-shifting effects. This is particularly important in the light of the current availability of large data sets based on deep, high-resolution optical imaging.

The Wide-Field Camera 3 (WFC3) on the *HST* provides near-IR imaging with the F160W filter that allows us to probe rest-frame *B* band out to $z \lesssim 2.5$. Our results indicate that high-redshift studies will be able to put interesting constraints on any detected evolution in the bar size distribution. With a spatial angular resolution of $0.13 \text{ arcsec pixel}^{-1}$ for the WFC3 IR channel corresponding to $\sim 1 \text{ kpc}$ at $0.8 \lesssim z \lesssim 2.5$, WFC3 allows us to trace evolution in bar properties for what corresponds to a typical local bar, with $a \sim 4.2 \text{ kpc}$ (Menéndez-Delmestre et al. 2007). Recent results by Kim et al. (2021) are pioneering the exploration of bar properties in large samples out to $z \sim 0.8$, exploiting rest-frame optical data based on *HST*. With the advent of the *JWST*, the significantly finer angular resolution of the MIRI and NIRcam imagers (MIRI: $0.11 \text{ arcsec pixel}^{-1}$ at $5.6 \mu\text{m}$; NIRcam: $0.03, 0.6 \text{ arcsec pixel}^{-1}$ at 2 and $4 \mu\text{m}$, respectively) and their wavelength coverage out to MIR wavelengths ($\lambda \sim 5\text{--}25 \mu\text{m}$) will allow investigations of bar properties down to even subtypical bar size values at higher redshifts (Costantin et al. 2023; Guo et al. 2023).

We emphasize that the ability of reliably quantifying any change in the bar size distribution in light of band-shifting effects does not preclude the bias against smaller bars imposed by a decrease in spatial resolution. As studies probe larger distances, these become increasingly biased against the smaller bars that fall below the instrument’s physical spatial resolution. A number of studies have clearly established the need for careful consideration of bar size limit in sample selection when investigating the presence of bars at higher

redshifts (Sheth et al. 2003; Erwin 2005; Menéndez-Delmestre et al. 2007; Erwin 2018). Cosmological surface brightness dimming also irretrievably impacts our ability to study structures in the distant Universe, as has been discussed by previous works (e.g. Sheth et al. 2008; Melvin et al. 2014; Kruk et al. 2019). Sheth et al. (2008) explored in detail the impact that surface brightness dimming due to cosmologically significant distances had in their ability to detect underlying discs and showed that their measured bar fractions did not have a dependence on the surface brightness of their galaxies, especially at the low surface brightness end of their sample. Kruk et al. (2019), with a focus on the redshift evolution of boxy/peanut-shaped bulges, explicitly explore the impact of cosmological surface brightness dimming by artificially redshifting their SDSS galaxy images out the redshifts covered by their COSMOS barred population at redshifts up to $z \sim 1$. Taking these works into consideration and our bar characterization approach (based on the ellipse fit methodology, which inherently demands the detection of an underlying disc), we note that the ellipticity signature on which our bar length and ellipticity measures are based on will continue to be an effective approach to measure bar properties at higher redshifts. Our work suggests that, once bar size limit considerations and band-shifting effects are taken into account, high-redshift studies are in a position to reliably trace intrinsic changes in the distribution of bar sizes as a function of redshift.

Interestingly, the increase in measured bar ellipticity as bluer rest-frame wavebands are probed due to band-shifting effects may have an unexpected consequence that it may facilitate the identification of weak bars at high redshift. Although intrinsically weaker bars – and their shallow ellipticity signature – are more difficult to pick out in general, the ‘artificial’ enhancement of the bar ellipticity due to band-shifting effects may make them increasingly easier to recognize at higher redshifts. In this way, studies are in the advantageous position of being able to probe a wider range in bar strengths at increasingly higher redshifts.

6 CONCLUSIONS

We present a detailed study of bar properties as a function of wavelength for a sample of 16 large nearby barred galaxies, spanning the wavelength range from the ultraviolet (GALEX FUV/NUV) through the optical (SINGS B, R) out to the MIR (*Spitzer*/IRAC/ $3.6 \mu\text{m}$). While the MIR provides the optimal window to probe the details of stellar structure in galaxies, including bars, we choose to include the UV bands in our study to extend the validity of our calibration to redshifts beyond $z \sim 0.8$, when *HST*/optical bands start tracing rest-frame UV.

Based on the ellipticity and position angle profiles resulting from fitting elliptical isophotes to the full 2D light distribution of each galaxy, we determine the bar length and bar ellipticity for each galaxy, in each band. Our main results are the following:

(1) We find that at bluer wavebands both the bar length and the bar ellipticity increase. We attribute the increase in bar length to the frequent presence of star forming knots at the end of bars: these regions become more prominent in bluer bands, resulting in an apparent lengthening of the bar. Recent studies have also pointed to distinct stellar populations in bars, where kinematic fractionation may lead to younger stars pertaining to an otherwise older bar; this suggests that the bluer bands may perhaps be allowing us to trace these younger stars in bar orbits. The increase in bar ellipticity, on the other hand, we interpret as a result of the fact that the bulge, composed primarily of old and red stars, is less prominent at bluer bands, allowing for thinner ellipses to be fit within the bar region. Part

of the ellipticity increase, however, could be due to the differences between the properties of the bar as outlined by the young stars, compared to the older ones, in good agreement with what is seen in simulations. The resulting effect is that bars appear longer and thinner when traced in bluer wavebands.

(2) We find that, to first order, bars are measured to have a size ~ 9 per cent longer in the *B* band, compared to their ‘canonical’ size at $3.6\ \mu\text{m}$ and that their ellipticity increases by ~ 8 per cent. This result mainly places constraints on the impact that band-shifting effects may have for bar studies at high redshift.

(3) Although we find that $\gtrsim 50$ per cent of the bars disappear in the GALEX *NUV/FUV* bands, the results on bar ellipticity and length extend to those cases in which the bar is still visible in the UV. Our results can be used as a reference to implement band-shifting corrections to reliably gauge any intrinsic redshift evolution of bar properties beyond $z \sim 0.8$, when optical filters start tracing rest-frame UV bands.

These results are pioneering in providing a reference for band-shifting corrections in order to reliably gauge any intrinsic redshift evolution of bar properties beyond $z \sim 0.8$. This opens the door for current and future studies to exploit current high-resolution imaging surveys and extend detailed morphological studies of bars out to the high-redshift Universe.

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DATA AVAILABILITY

The data were derived from images in the public domain: <http://irsa.ipac.caltech.edu/data/SPITZER/S4G>. The measurements produced as part of our work are available in the article. Additional access will be shared on reasonable request to the corresponding author.

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APPENDIX A: RADIAL PROFILES FOR THE HIGH-RESOLUTION STUDY IN THE *B*-, *R*-, 3.6- μ M BANDS

We show in Fig. A1 the ellipse-fit results for all the galaxies in our sample. Following the same format as in Fig. 2, for each galaxy we show profiles for ellipticity and its variation, as well as position angle and its variation in the *B*, *R*, and 3.6 μ m band. The vertical down-pointing arrows indicate the semimajor axis (SMA) of maximum ellipticity, i.e. $a_{\epsilon_{\max}}$ for each band. Postage stamp images in the *B*, *R*, and 3.6 μ m bands, from top to bottom, are shown for each galaxy with an overlaid ellipse showing the isophote of maximum ellipticity.

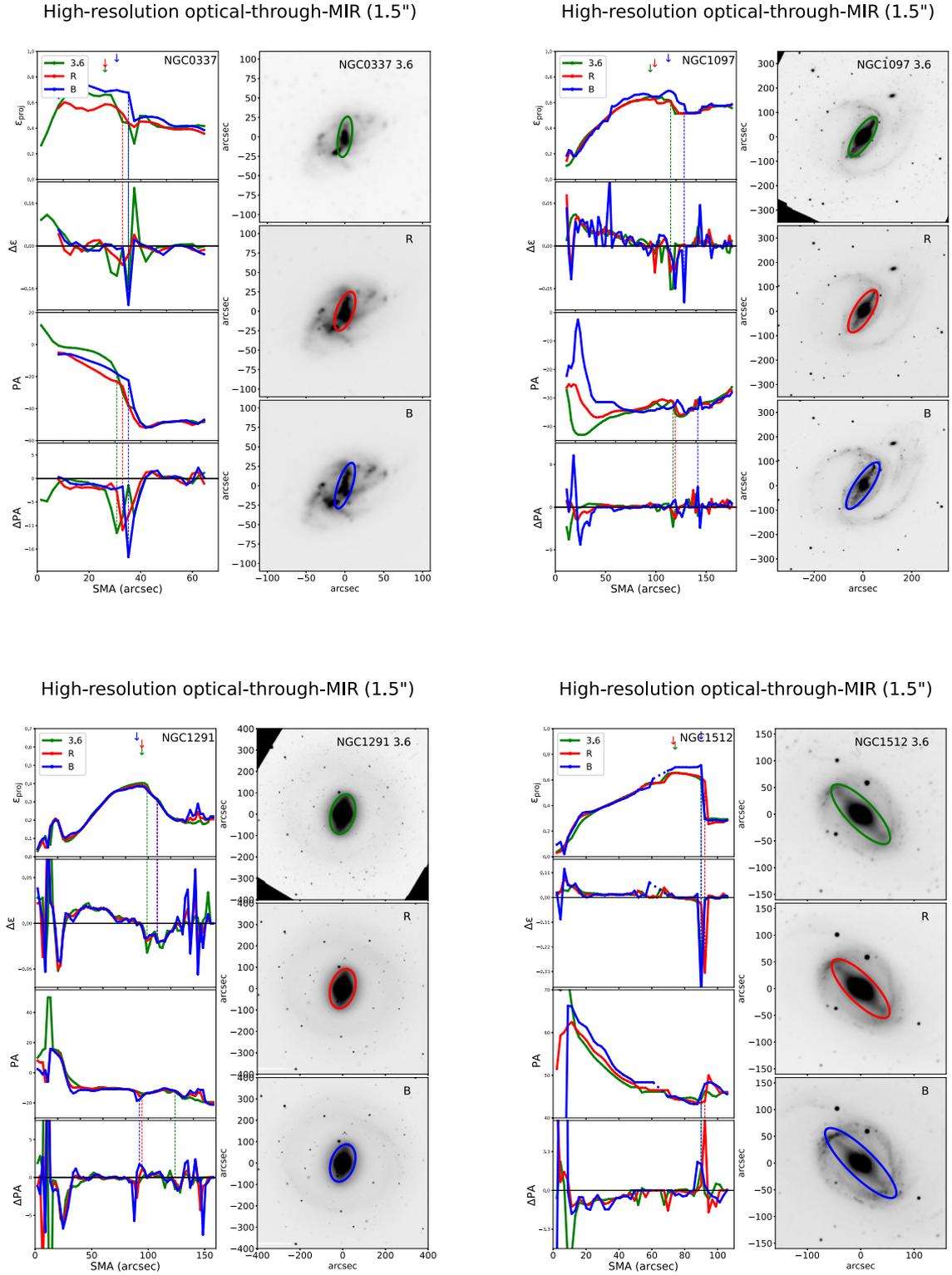


Figure A1. Ellipticity and PA profiles for all galaxies in our sample as part of the *higher resolution opt-through-MIR* study, following the format of Fig. 2.

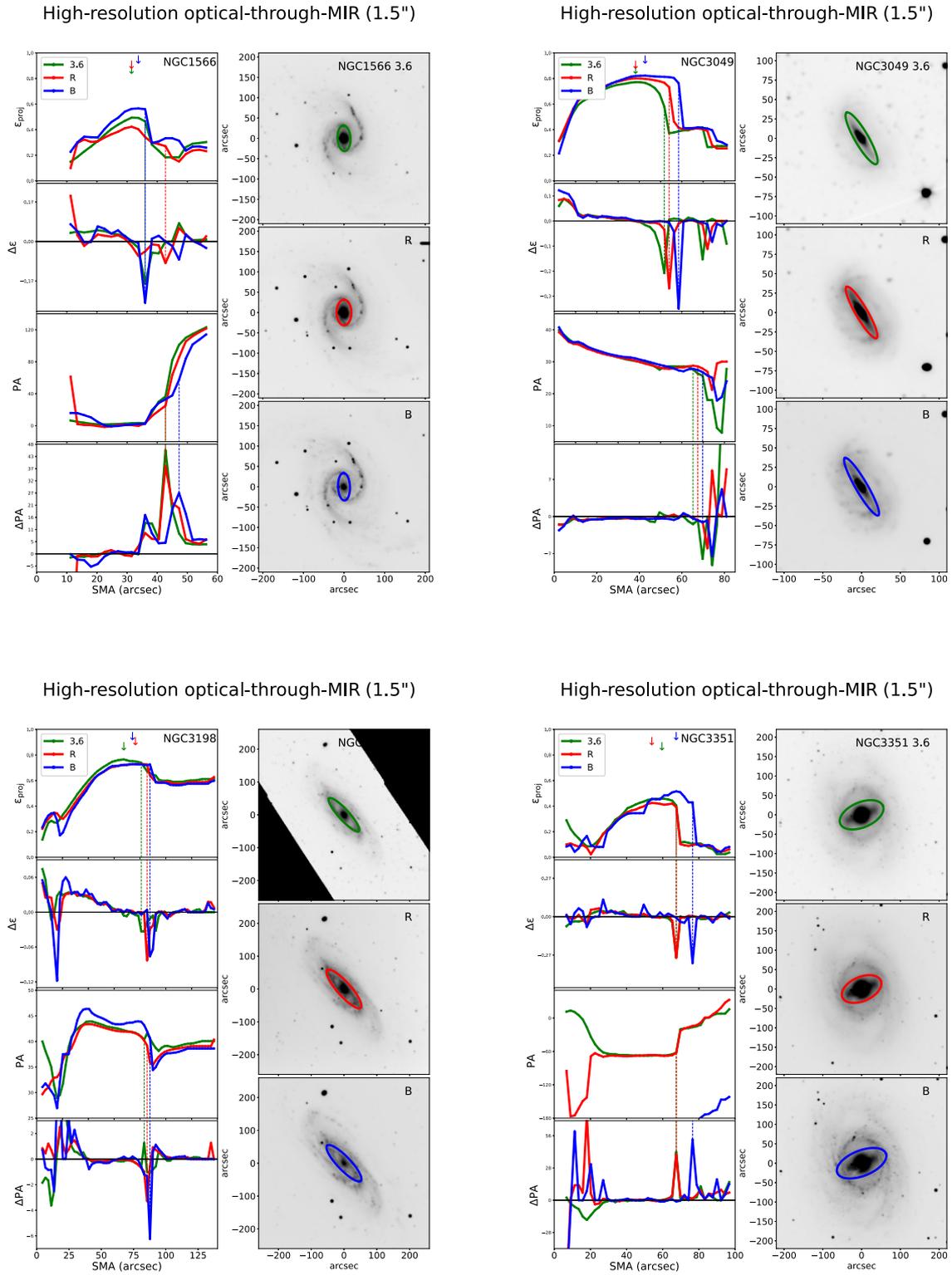


Figure A1 – continued

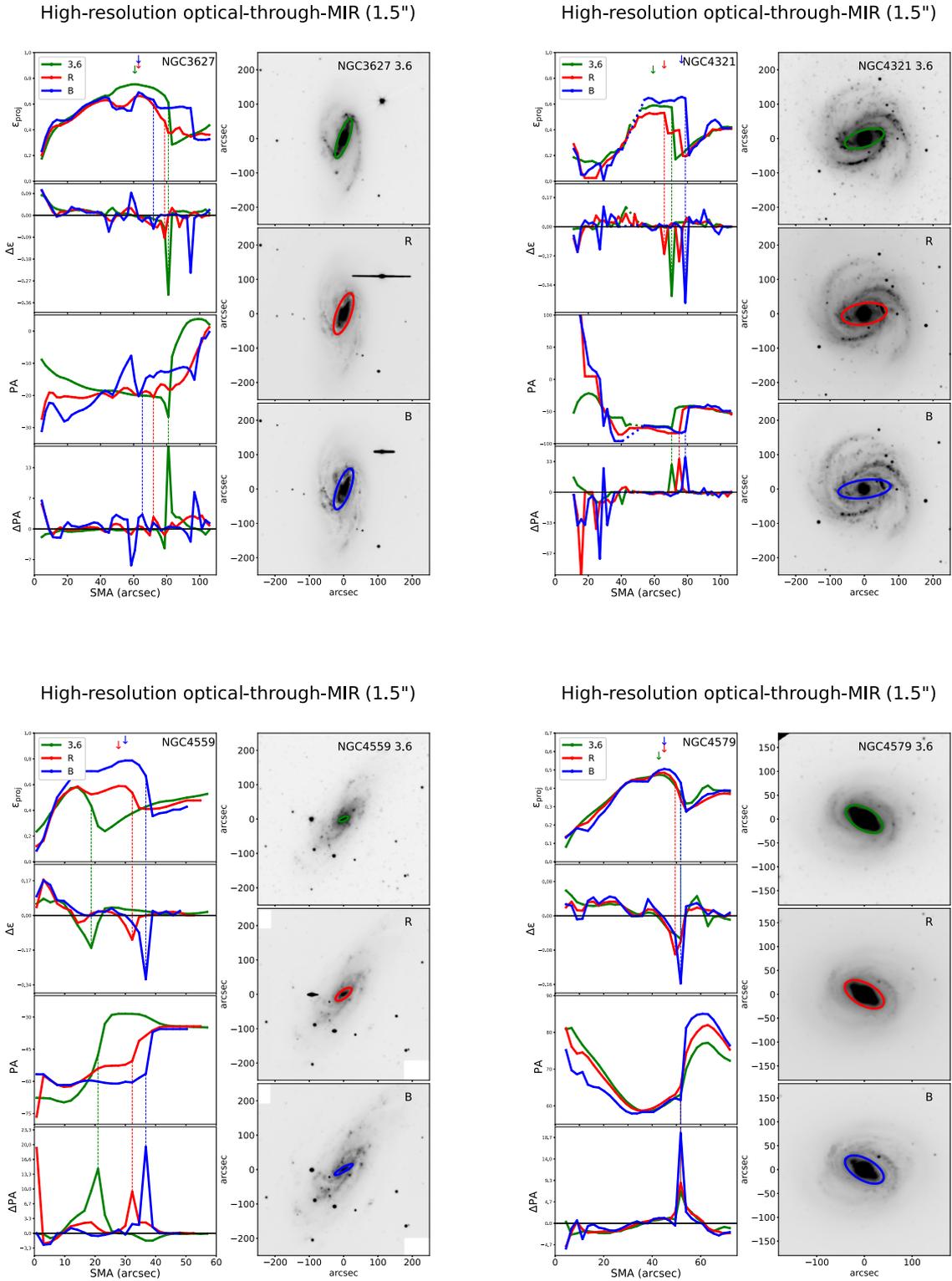


Figure A1 – continued

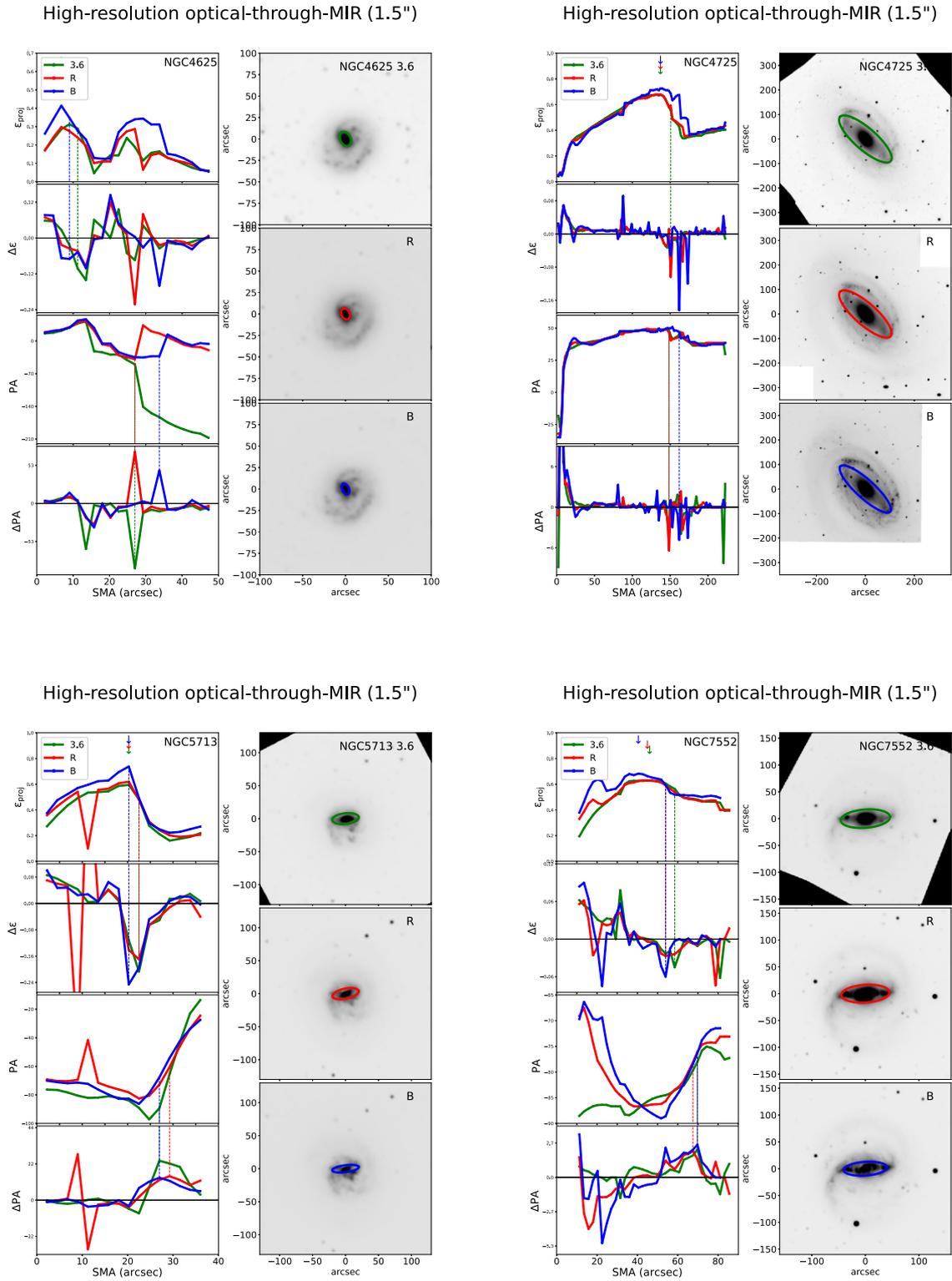


Figure A1 – continued

APPENDIX B: RADIAL PROFILES FOR THE LOW-RESOLUTION STUDY IN THE FUV-, NUV-, B-, R-, 3.6 μ M BANDS

Fig. B1 shows the ellipticity and PA profiles for all galaxies in our sample as part of the *low resolution UV-through-MIR* study, following the format of Fig. 2.

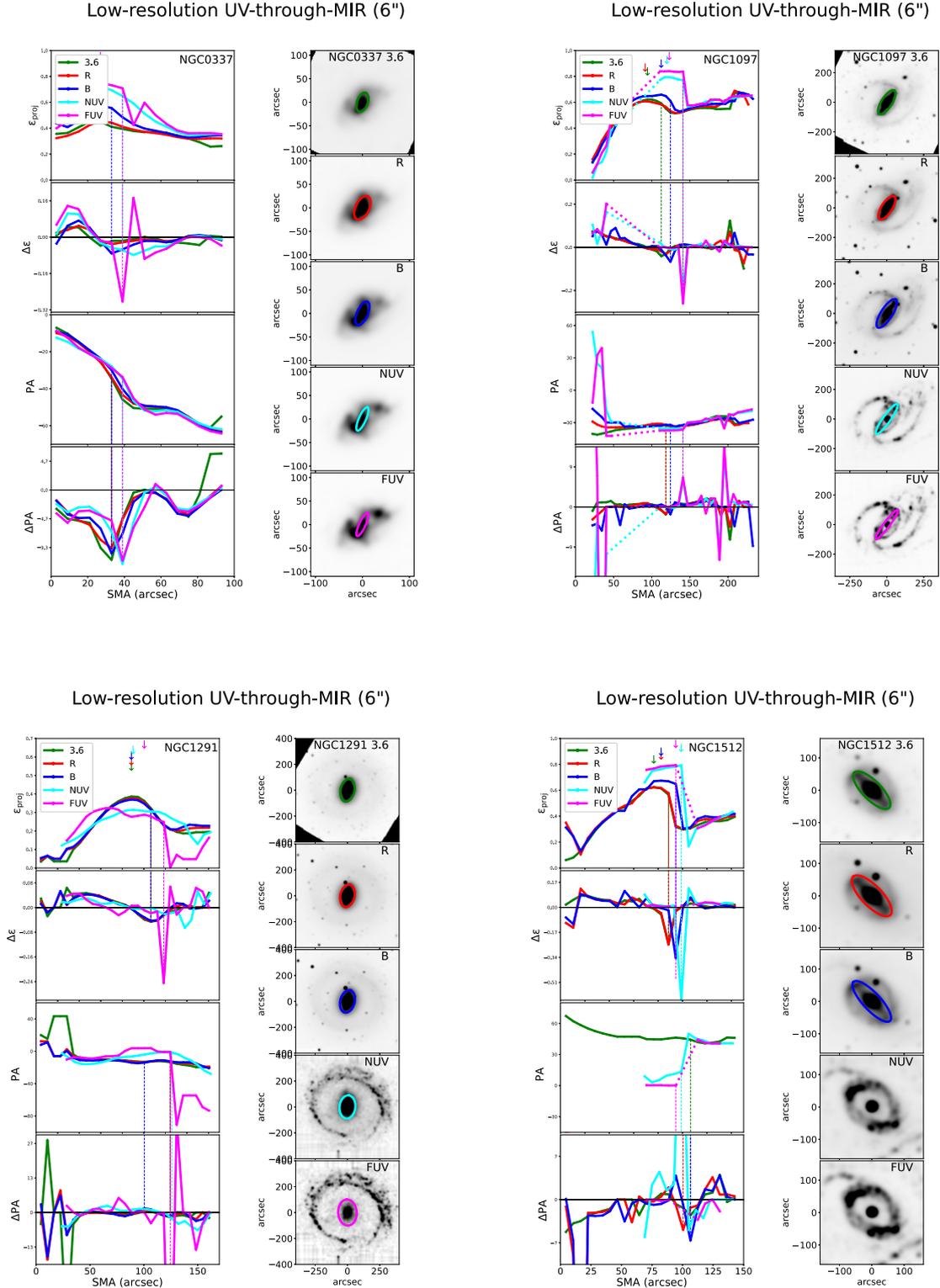


Figure B1. Ellipse-fit results for the UV-through-MIR analysis.

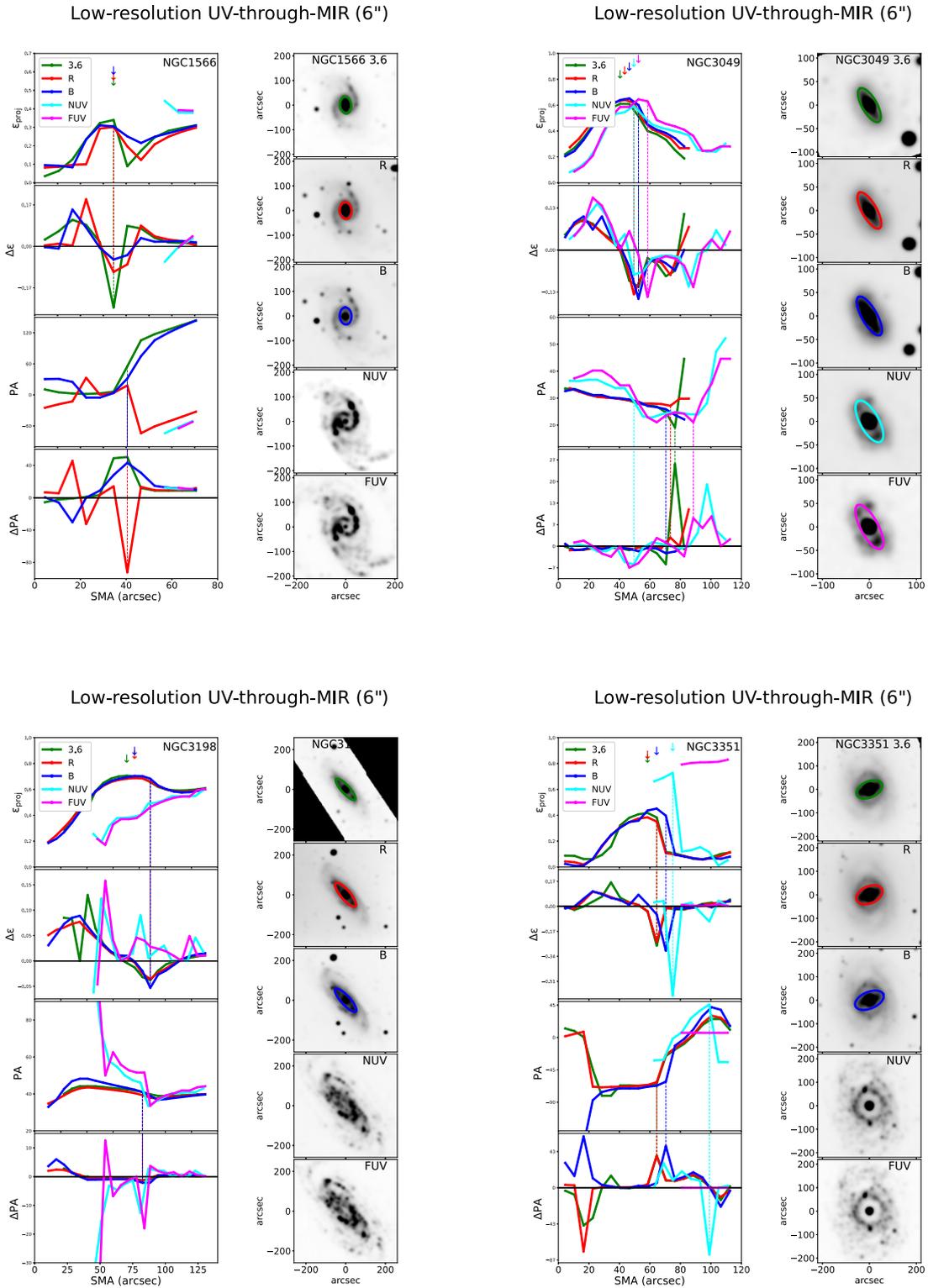


Figure B1 – continued

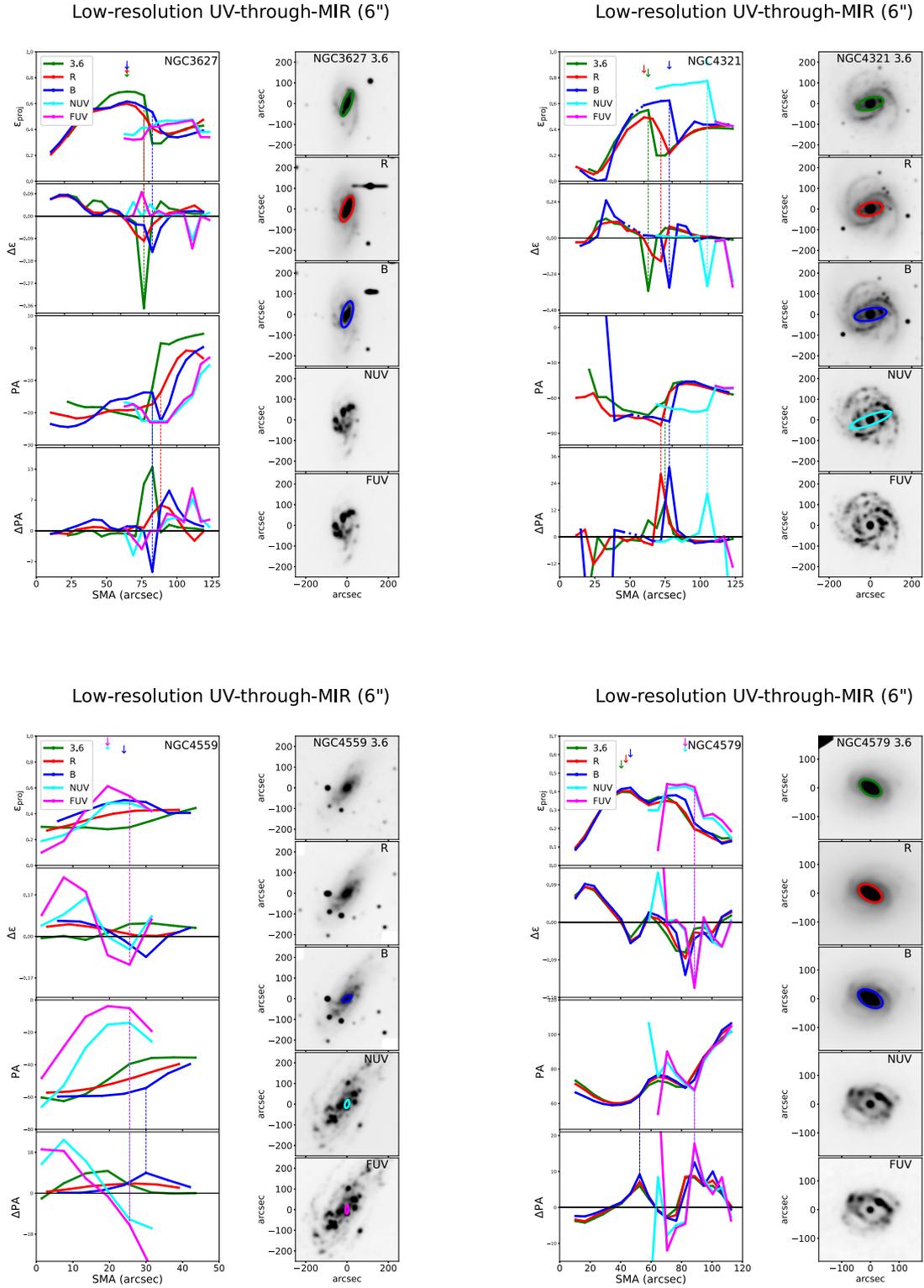


Figure B1 – continued

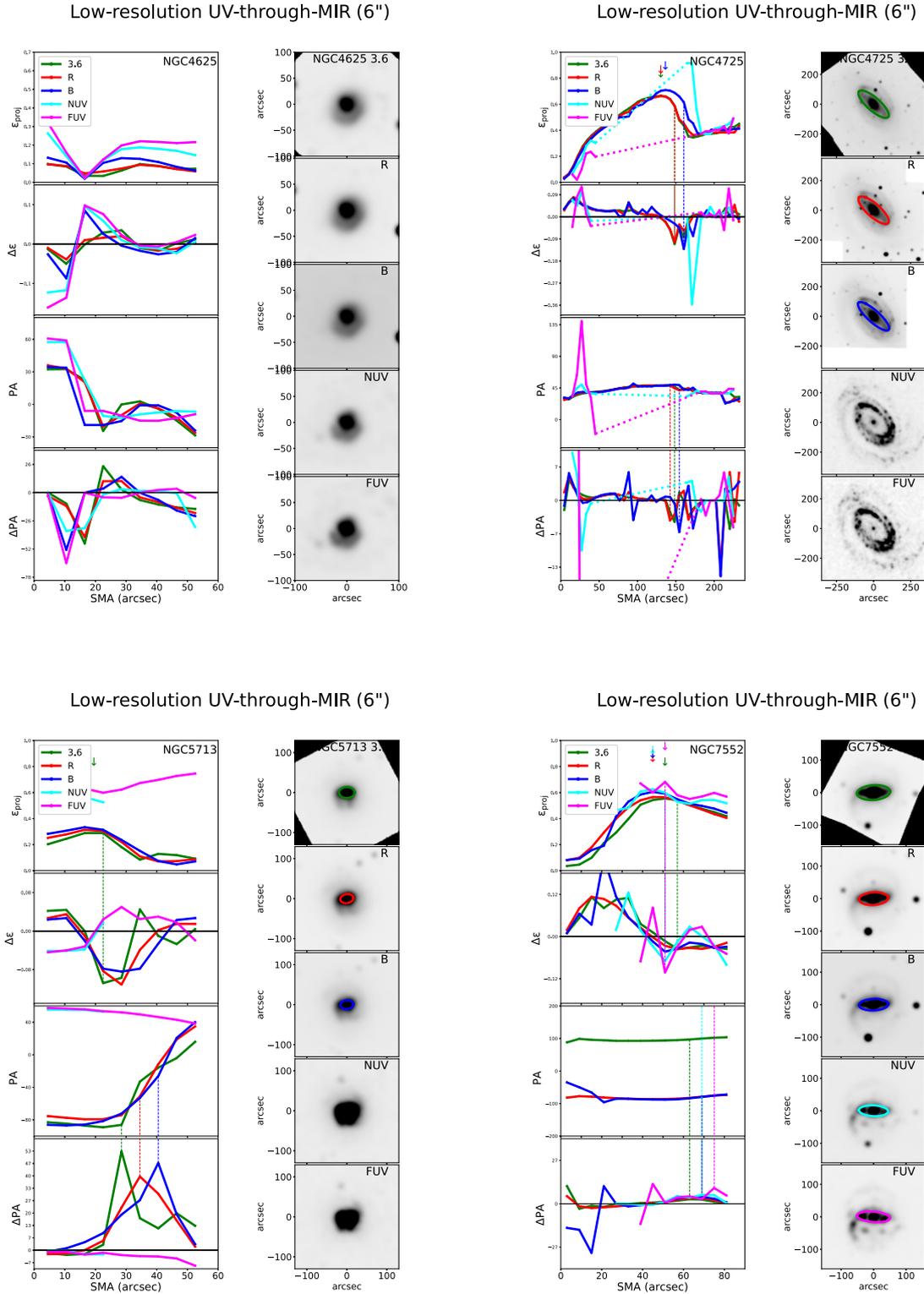


Figure B1 – continued

APPENDIX C: RESULTS ON THE STATISTICAL ANALYSIS

We summarize in Fig. C1 our bar length and bar strength multiband comparison, with the statistical significance of these results indicated by the paired t-test and Wilcoxon test for both the *higher resolution*

optical-through-MIR and the *low resolution UV-through-MIR* study. The results for the Wilcoxon test with bar lengths measured by the location of the bar ellipticity maximum ($a_{e\text{max}}$) are presented on Tables 2 and 3. We here show the two alternative methods for bar length measurement that we explored: $a_{\Delta E}$ and $a_{\Delta PA}$.

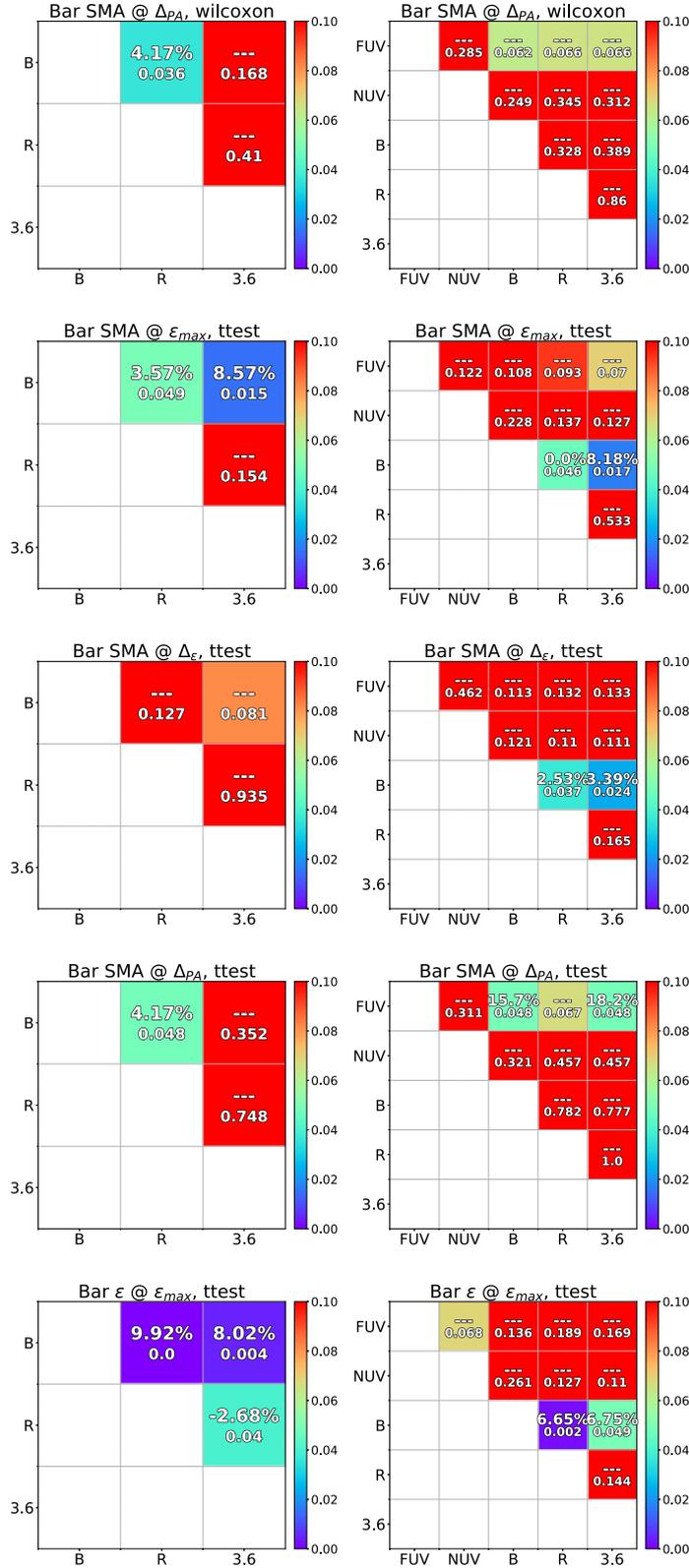


Figure C1. Comparison of bar ellipticity and bar length measurements in different bands, with significance of each comparison tested with the paired t -test and the Wilcoxon test, for both the high-resolution and the low-resolution study.

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