



**HAL**  
open science

## **Sulphur storage in cold molecular clouds: the case of the NH<sub>4</sub>+SH<sup>-</sup> salt on interstellar dust grains**

J. Vitorino, J. -C. Loison, V. Wakelam, E. Congiu, F. Dulieu

### ► **To cite this version:**

J. Vitorino, J. -C. Loison, V. Wakelam, E. Congiu, F. Dulieu. Sulphur storage in cold molecular clouds: the case of the NH<sub>4</sub>+SH<sup>-</sup> salt on interstellar dust grains. *Monthly Notices of the Royal Astronomical Society*, 2024, 533, pp.52-62. <10.1093/mnras/stae1747>. <insu-04822435>

**HAL Id: insu-04822435**

**<https://insu.hal.science/insu-04822435v1>**

Submitted on 6 Dec 2024

**HAL** is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.



Distributed under a Creative Commons CC BY 4.0 - Attribution - International License

# Sulphur storage in cold molecular clouds: the case of the $\text{NH}_4^+\text{SH}^-$ salt on interstellar dust grains

J. Vitorino <sup>1</sup>★, J.-C. Loison <sup>2</sup>, V. Wakelam <sup>3</sup>, E. Congiu <sup>1</sup> and F. Dulieu <sup>1</sup>★

<sup>1</sup>CY Cergy Paris Université, Observatoire de Paris, PSL Research University, CNRS, LERMA, F-95000 Cergy, France

<sup>2</sup>Université de Bordeaux, Institut des Sciences Moléculaires, UMR 5255, F-33400 Talence, France

<sup>3</sup>Laboratoire d'Astrophysique de Bordeaux, Université de Bordeaux, CNRS, B18N, allée Geoffroy Saint-Hilaire, F-33615 Pessac, France

Accepted 2024 July 11. Received 2024 July 2; in original form 2024 May 27

## ABSTRACT

In comets and in the cold phase of the interstellar medium (ISM), ammonium salts are key molecular species due to their role in the retention of volatile compounds on cold surfaces. In the case of sulphur, the  $\text{H}_2\text{S}/\text{OCS}$  ratio observed in protostars could be explained by the presence of ammonium hydrosulphide ( $\text{NH}_4\text{SH}$ ) salts. However, laboratory data on the properties of  $\text{NH}_4\text{SH}$  in ISM cold relevant conditions are rather scarce, as they usually focus on the atmosphere of Jupiter. We propose to consolidate the laboratory data regarding  $\text{NH}_4\text{SH}$  on grains, by performing temperature programmed desorption experiments and Fourier transform infrared reflection spectroscopy. The salt was also exposed to H atoms to mimic the ISM conditions.  $\text{NH}_4\text{SH}$  was found to form *in situ* at 10 K, from a mixture of ammonia ( $\text{NH}_3$ ) and hydrogen sulphide ( $\text{H}_2\text{S}$ ). The  $\text{NH}_4^+$  infrared feature ( $1485\text{ cm}^{-1}$ ) is the most prominent one at 80 K. As pure species,  $\text{H}_2\text{S}$  and  $\text{NH}_3$  desorb at 76 and 90 K, respectively, whereas they are released into the gas phase at 153 K when adsorbed in the form of salt. The presence of water delays the desorption of the salt until the very end of the water desorption, but does not affect the desorption kinetics. During H-exposure, the salt is dissociated and no new product was detected. As a comparative study, salts have been included in the Nautilus gas–grain model. The results show a good correlation with the observations of IRAS 16293–2422B, as opposed to when  $\text{NH}_4\text{SH}$  is not included in the model.

**Key words:** astrochemistry – molecular processes – methods: laboratory: solid state – ISM: clouds – ISM: molecules.

## 1 INTRODUCTION

Laboratory data about ammonium hydrosulphide ( $\text{NH}_4^+\text{SH}^-$ , hereafter  $\text{NH}_4\text{SH}$ ) characteristics in interstellar medium (ISM) relevant conditions are rather scarce. Indeed, in terrestrial conditions, ammonium hydrosulphide is known to be formed via the mixing of gaseous hydrogen sulfide ( $\text{H}_2\text{S}$ ) and ammonia ( $\text{NH}_3$ ), which are corrosive chemical compounds with high toxicity for humans. Additionally, their use asks for a very precise leak control, especially for  $\text{H}_2\text{S}$ , given the intense odour detectable by smell well before the recommended exposure limit concentration.<sup>1</sup>  $\text{NH}_4\text{SH}$  is one of the most interesting salts potentially present in the ISM and in the upper layer of the coma in comets such as 67P/Churyumov–Gerasimenko (Altwegg et al. 2020). The detection of sulphur and ammonia clouds in overlapping layers by the Galileo probe on Jupiter, with a uniform mixture at 10 bars, also indicates that  $\text{NH}_4\text{SH}$  is a major compound of gas-giant planets atmosphere (Niemann et al. 1998; Pater et al. 2001; Wong et al. 2004; Li et al. 2020). This has been supported by theoretical models (Weidenschilling & Lewis 1973; Atreya et al. 1999; Friedson 2005), as well as spacecraft observations (Roman, Banfield & Gierasch 2013). Consequently, most laboratory

studies investigate the radiolysis and photodissociation of ammonium hydrosulfide. For example, the infrared and radiolytic data produced by Loeffler et al. (2015), allow them to predict that  $\text{NH}_4\text{SH}$  would be detectable mostly in crystalline form at 200 K in gas-giant clouds, but that ionizing radiations would destroy the crystalline structure, producing an amorphous mixture instead. The first documented laboratory experiments of ultraviolet (UV)-photolysis of mixed  $\text{H}_2\text{S}$ ,  $\text{NH}_3$ , and  $\text{H}_2\text{O}$  ices and pure  $\text{NH}_4\text{SH}$  were done by Lebofsky & Fegley (1976), in order to compare the spectra obtained with the observations of Jupiter and Saturn satellites, as well as Saturn rings. More recently, combined studies of UV–vis reflectance spectra of photolyzed  $\text{NH}_4\text{SH}$ , together with visible light imaging of these mixtures, have proven their usefulness for remote sensing campaigns on Jupiter, showing a colour-dependence on the irradiation dose and temperature (Loeffler et al. 2016; Loeffler & Hudson 2018).

Salts are ionic compounds, often formed by proton transfer in an acid–base reaction. The latter can take place at cryogenic temperatures due to a very low-activation energy (Theulé et al. 2013), while the salt created is usually much more refractory than any of the individual species involved in the ionic bond. For instance,  $\text{NH}_3$  and  $\text{HCOOH}$  or  $\text{CH}_3\text{COOH}$ , when interacting in form of ammonium formate ( $\text{NH}_4^+\text{COO}^-$ ) or ammonium acetate ( $\text{NH}_4^+\text{CH}_3\text{COO}^-$ ), can remain on dust grains far beyond their respective desorption temperature (Bergner et al. 2016; Kruczkiewicz et al. 2021). The same goes for reactions such as  $\text{NH}_3 + \text{HCN} \rightarrow \text{NH}_4^+\text{CN}^-$  (ammonium cyanide, Clutter & Thompson 1969; Gerakines, Moore &

\* E-mail: [julie.vitorino@cyu.fr](mailto:julie.vitorino@cyu.fr) (JV); [francois.dulieu@cyu.fr](mailto:francois.dulieu@cyu.fr) (FD)

<sup>1</sup>Hydrogen sulfide; SDS No. 295 442 [Print]; Merck Life Science S.A.S.: Saint-Quentin-Fallavier, France, 2024 Jan 3. <https://www.sigmaaldrich.com/FR/fr/product/aldrich/295442> (accessed 2024 Mar 4).

Hudson 2004; Noble et al. 2013; Gerakines, Yarnall & Hudson 2024), or  $\text{NH}_3 + \text{HNCO} \rightarrow \text{NH}_4^+ \text{OCN}^-$  (ammonium isocyanate, Demyk et al. 1998; Raunier et al. 2003; Van Broekhuizen, Keane & Schutte 2004). In each of these cases, the reactants are relatively volatile compounds, likely to sublimate from dust grains at comparatively low temperatures (i.e.  $<100\text{ K}$  for the smallest ones,  $<200\text{ K}$  for the largest ones). On the other hand, the salts formed sublimate at higher temperatures, allowing the snowline of both volatile species involved in the mixtures to be extended towards warmer regions. Therefore, salts are considered good candidates to be involved in the complexation of early chemistry.

In addition, salts would provide an explanation for some of the most significant depletions in the Universe. For example, a nitrogen depletion is observed in several environments including comets (Filacchione et al. 2019) and ISM (Savage & Sembach 1996), compared to the Solar System ratios (Lodders 2010). Likewise, sulphur is reported missing especially in dense cores and protostellar environments (Tieftrunk et al. 1994; Goicoechea et al. 2006). The abundance of S-bearing species is then approximately 1000 times lower than that found in diffuse clouds (Kushwahaa et al. 2023), which is also not in tone with their abundance in comets (Calmonte et al. 2016; López-Gallifa et al. 2024). So far, there is not one clear explanation to these depletions. Regarding sulphur, it is believed to be trapped in the dust grains cores (Perrero et al. 2024), or in the icy mantles in the form of salts. The latter hypothesis would explain the release of S-containing compounds in comets and enable their detection in larger amounts.

Moreover, the H<sub>2</sub>S presence in the early stages of pre-stellar cores development is a determining factor when it comes to understanding the carbonyl sulfide (OCS) as well as the S-bearing molecules content, especially in low-mass stars. The H<sub>2</sub>S on interstellar dust grains would indeed be converted to OCS throughout the collapse phase, as long as the core environment is cold enough (typically  $<20\text{ K}$ ) to maintain a high quantity of CO ice (Drozdovskaya et al. 2018). The transformation would mostly happen via grain-surface processes (Ferrante et al. 2008; El Akel et al. 2022). More elaborately, the rate of UV radiations in the pre-stellar core is also likely to determine the fluency of the conversion of H<sub>2</sub>S to OCS, as shown by laboratory experiments where a high UV irradiation involves a high conversion rate (Chen et al. 2015), and observations of IRAS 16293–2422 B, where such a rate would explain the amount and diversity of S-bearing molecules detected compared to the Solar system ratios (Ruffle et al. 1999; Jiménez-Escobar, Caro & Chen 2014; Martín-Doménech et al. 2016; Drozdovskaya et al. 2018; Phuong et al. 2018).

It is thus also necessary to take ammonium salts into account in gas-grain models of protostars, as they would significantly influence the proportion of sulphur available to create S-bearing molecules, despite the apparent depletion.

In this paper, we aim to investigate the nature of the products upon the desorption of NH<sub>4</sub>SH. We propose to densify the laboratory data on the formation, infrared (IR) signatures, and characteristics of the salt on a cold surface representing an interstellar dust grain. We performed experiments and analyses with varying ratios of incoming NH<sub>3</sub> and H<sub>2</sub>S in pure form, on a bare gold-plated surface, and in the case of a water ice substrate. The reactivity and destruction of the salt under H atoms exposure were also studied. Hereafter, Section 2 describes the experimental methods and setup reproducing the interstellar conditions. In Section 3, we present our infrared spectroscopy and mass spectrometry data on ices of NH<sub>3</sub> and H<sub>2</sub>S, on NH<sub>4</sub>SH salt, on NH<sub>4</sub>SH exposed to H atoms, and finally in the case of mixtures containing water molecules in different proportions.

In Section 4, we discuss the astrophysical implications and show that the inclusion of NH<sub>4</sub>SH in a gas grain model can better explain the OCS/H<sub>2</sub>S ratio observed towards protostars.

## 2 EXPERIMENTAL METHODS

All of the experiments have been performed with the VErS de NoUvelles Synthèses (VENUS) experimental setup, located in LERMA-CY and described in detail elsewhere (Congiu et al. 2020). It consists of a gold-coated circular surface of 9 mm diameter, mounted on a closed-cycle He cryostat and held inside an Ultra-High Vacuum stainless steel chamber. The base pressure is approximately of  $10^{-10}$  mbar. Chemical species under a gaseous, liquid, or solid state at room temperature are kept or converted into gas state in the molecular or atomic beamlines, then sent towards the surface by means of a differential pumping through two intermediate chambers. The beams meet on the surface and the gas-phase interactions above the cold sample are negligible. Indeed, the partial pressure does not exceed  $10^{-8}$  mbar in the volume of the beams, and the main chamber is maintained at  $10^{-10}$  mbar even when species are being deposited. The only exception is for H<sub>2</sub> diffusion when the atomic hydrogen beam is being used (the partial pressure of the chamber therefore going up to  $10^{-8}$  mbar). Overall, surface-surface and possibly gas-surface interactions are largely dominant compared to gas-gas interactions. Representing an interstellar dust grain, the surface temperature can vary between 6 and 350 K and is computer-controlled by means of a resistive heater clamped behind the sample holder.

### 2.1 RAIRS and TPD

A typical experiment consists of two phases. The species are deposited at a constant temperature for a sufficiently long time to allow their detection, forming an ice film on the surface (phase 1, deposition), before being sublimated and released in gaseous form during a linear increase in the temperature of the surface (phase 2, temperature programmed desorption, TPD). During both phases, we can monitor the ice evolution with a Vertex-70 Fourier Transform infrared (FT-IR) spectrometer, performing reflection absorption infrared spectroscopy (RAIRS), a non-destructive analysis technique. The surface is continuously scanned with a grazing infrared beam at an 83° angle, over a spectral range of 800–4500 cm<sup>-1</sup> (12.5–2.22 μm). The light coming from the sample with the desired angle is collected in a N<sub>2</sub>-cooled mercury cadmium telluride (MCT) detector. Every two minutes (corresponding to 256 scans), a mean spectrum is produced.

Once the deposition at low temperature is achieved, the temperature of the surface is linearly increased at a rate  $\beta = 0.2\text{ K s}^{-1}$ . The desorbing species are detected with a Hiden 51/3F vertically translatable quadrupole mass spectrometer (QMS), placed 5 mm in front of the surface during the course of the TPD. Every molecule and fragment can then be probed according to their respective mass, chosen by the user and measured in atomic mass units (amu). The ionization energy of the QMS was set at 30 eV, and the number of species entering the QMS head is monitored as a function of temperature. From this, numerous parameters can be derived, including the binding energies between the molecules and the gold surface.

### 2.2 Atomic, water, and molecular beams

In this work we have carried out a series of experiments, as shown in Table 1, for which we provide TPD and infrared data. They can be grouped into several categories: individual species (i, ii),

**Table 1.** List of the conducted experiments and their associated parameters (composition of the ice, coverage, deposition temperature, and H flux when used).

Exp.	Reactive system	Coverage (monolayers)	$T_{\text{dep}}$ (K)	H flux (atoms $\text{cm}^{-2} \text{s}^{-1}$ )
i	{H <sub>2</sub> S}	0.1; 0.15; 0.35; 0.65	10	–
ii	{NH <sub>3</sub> }	0.35; 0.4; 0.5; 1	10	–
iii	{H <sub>2</sub> S + H}	0.65 : 10	10	$8 \times 10^{12}$
iv	{H <sub>2</sub> S + NH <sub>3</sub> }	1.8 : 4.65	10	–
v	{H <sub>2</sub> S + NH <sub>3</sub> }	1.8 : 1.9	10	–
vi	{H <sub>2</sub> S + NH <sub>3</sub> } + {H}	1.8 : 1.9; 50	80; 10	$8 \times 10^{12}$
vii	{H <sub>2</sub> S + NH <sub>3</sub> + H}	1.8 : 1.9; 30	10	$8 \times 10^{12}$
viii	{H <sub>2</sub> S + NH <sub>3</sub> + H <sub>2</sub> O}	1.9 : 1.9; 12	10	–
ix	{H <sub>2</sub> S + NH <sub>3</sub> + H <sub>2</sub> O}	0.8 : 1.5; 30	10	–
x	{NH <sub>3</sub> + H <sub>2</sub> O}	1 : 6	10	–
xi	{H <sub>2</sub> S + NH <sub>3</sub> + H <sub>2</sub> O + H}	1.9 : 1.9 : 12 : 30	10	$8 \times 10^{12}$

hydrogenation of H<sub>2</sub>S (iii), salt formation (iv, v), salt with H atoms exposure (vi, vii), salt in a water ice matrix (viii, ix), NH<sub>3</sub> with water (x), and salt in a water ice matrix and under H-exposure (xi). When two or more species are deposited on the surface, we denote the codeposition by placing all the species in a pair of curly brackets. For instance, in the figures of this paper, we use the notation  $\{A + B\}_{\text{doseA:doseB}}^T$  to refer to a codeposition of species *A* and *B* on the surface, at temperature *T*, with estimated or calculated doses for each species. Conversely, a sequential deposition is indicated by a series of brackets, that is  $\{A + B\}_{\text{doseA:doseB}}^{T_{AB}} + \{C\}_{\text{doseC}}^{T_C}$ .

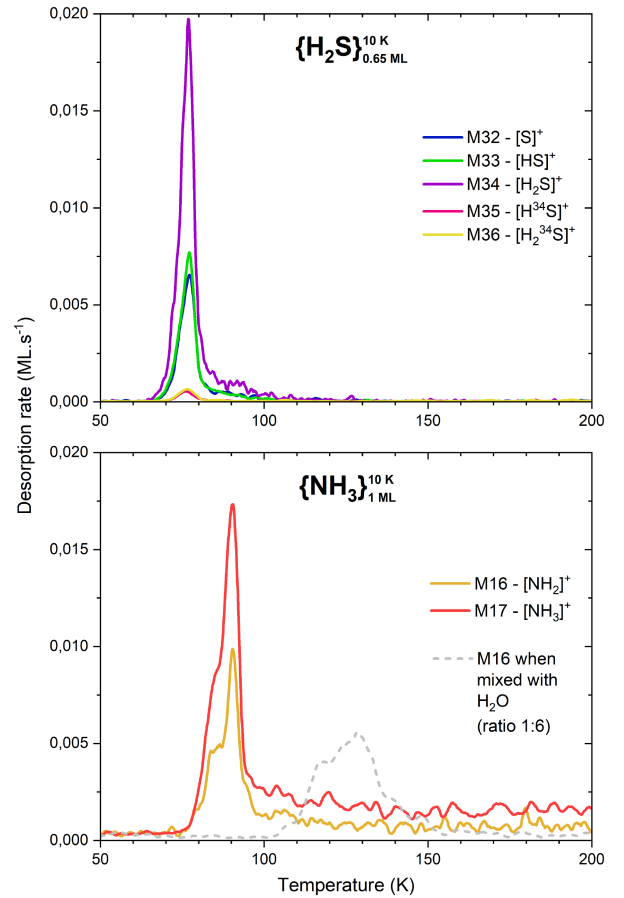
The molecular NH<sub>3</sub> beam was calibrated as described in Noble et al. (2011). A monolayer (ML), which defines the theoretical filling of all available adsorption sites on the sample and corresponds to  $\sim 10^{15}$  molecules  $\text{cm}^{-2}$ , was reached after 10 and 31 min of deposition, for NH<sub>3</sub> and H<sub>2</sub>S, respectively. This gives corresponding condensation rates – or fluxes – of  $\phi_{\text{NH}_3} = 1.7 \times 10^{12}$  molecules  $\text{cm}^{-2} \text{s}^{-1}$ , and  $\phi_{\text{H}_2\text{S}} = 5.4 \times 10^{11}$  molecules  $\text{cm}^{-2} \text{s}^{-1}$ . On the other hand, in the case of H<sub>2</sub>O, water vapour was leaked into the main chamber via a delivery manifold with known characteristics (a monolayer grown in approximately 5 min, giving  $\phi_{\text{H}_2\text{O}} = 3.3 \times 10^{12}$  molecules  $\text{cm}^{-2} \text{s}^{-1}$ ). For water, the thickness of an ice layer of 1 ML is 3 Å, which is also a good estimate for other species of comparable size.

In the case of hydrogen, only H<sub>2</sub> molecules can be stored in a gas bottle. The atoms were obtained by dissociating the molecules and creating a plasma in a quartz tube surrounded by a microwave cavity (Surfatron). The H atoms have a short residence time on one adsorption site, due to their propensity to diffuse and recombine rapidly to form H<sub>2</sub>. Therefore, their quantity in ML can only be estimated with the maximum number of reactions that can occur in a barrierless reaction such as NO + H (Congiu et al. 2012). One ‘monolayer’ of H atoms is then considered to be achieved in about 2 min when the flux is of  $\phi_{\text{H}} = 8 \times 10^{12}$  atoms  $\text{cm}^{-2} \text{s}^{-1}$ . The H<sub>2</sub> dissociation efficiency was measured before each experiment with H atoms, giving an average of  $\sim 62$  per cent, with a possible variation of at most 20 per cent between two different experiments.

### 3 RESULTS

#### 3.1 Single-component ices: NH<sub>3</sub> and H<sub>2</sub>S depositions

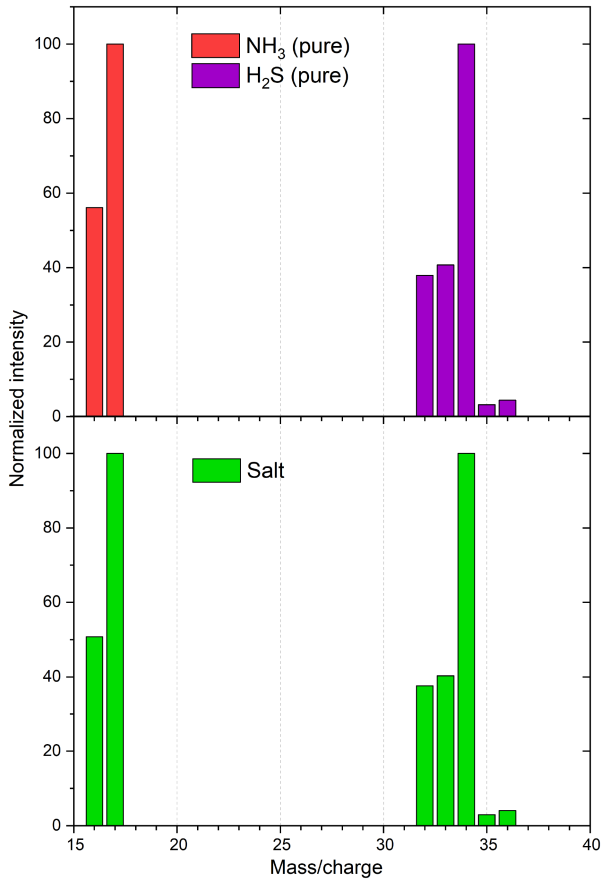
A characterization of the individual species was first carried out in order to determine their properties under our experimental conditions. Fig. 1 displays the TPD spectra of 0.6 monolayers of H<sub>2</sub>S, and 1 monolayer of NH<sub>3</sub>. Both species were deposited at 10 K and desorb from the surface at 76 and 90 K respectively, following a first order desorption behaviour. Several fragments can be identified for each



**Figure 1.** TPD spectra of 0.65 monolayers of H<sub>2</sub>S (top panel), and of 1 monolayer of NH<sub>3</sub> (bottom panel). Both species were deposited at 10 K. Each curve represents a fragment ionized by the QMS. The dotted line stands for  $m/z = 16$  when NH<sub>3</sub> is mixed with H<sub>2</sub>O.

species, as shown with each curve, giving a unique cracking pattern for each molecule (see upper panel of Fig. 2). The amounts of each fragment depend mainly on the ionization energy of the QMS, which in this work was set at 30 eV. The H<sub>2</sub>S and NH<sub>3</sub> major fragments correspond to the ionized unimpair molecules:  $m/z = 34$  for H<sub>2</sub>S<sup>+</sup> and  $m/z = 17$  for NH<sub>3</sub><sup>+</sup>.

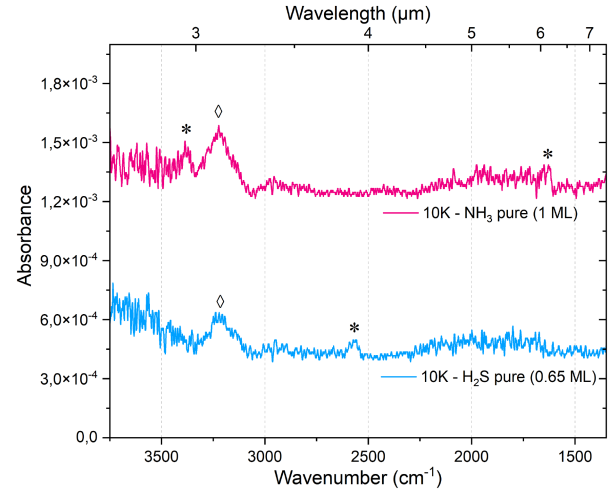
In the case of H<sub>2</sub>S, we detect between 3 and 4 per cent of <sup>34</sup>S, an amount that is not considered to significantly affect the nature and kinetics of the reactions with H and NH<sub>3</sub>. Even when less than 1



**Figure 2.** Mass spectra of NH<sub>3</sub> (top left), H<sub>2</sub>S (top right), and NH<sub>4</sub>SH (bottom). The intensities have been normalized to the most abundant fragment for each of the pure species ( $m/z = 17$  for NH<sub>3</sub>,  $m/z = 34$  for H<sub>2</sub>S). For NH<sub>4</sub>SH, the ammonium fragments are normalized to mass 17, while the hydrogen sulfide fragments are normalized to mass 34.

ML of H<sub>2</sub>S is deposited on the sample, the molecules quickly exhibit a multilayer behaviour under our conditions, making it difficult to determine the time after which a monolayer is reached. Moreover, the areas under the main fragment curves of both H<sub>2</sub>S and NH<sub>3</sub> are not directly comparable to each other due to their different ionization cross-sections at 30 eV, of respectively  $\sigma_{\text{H}_2\text{S}}^{\text{tot}} = 3.18 \text{ \AA}^2$  (Vinodkumar et al. 2011) and  $\sigma_{\text{NH}_3}^{\text{tot}} = 1.57 \text{ \AA}^2$  (Kumar & Kumar 2020), giving a ratio of  $\frac{\sigma_{\text{H}_2\text{S}}}{\sigma_{\text{NH}_3}} \sim 2.02$ . Consequently, the determination of the monolayer completion for H<sub>2</sub>S (31 min) was based on the value obtained for NH<sub>3</sub> (10 min). Measurements in the salt (cf. Section 3.2) further confirm a calculated ratio of  $\sim 1.9$  between the two species, which is in good agreement with the total cross-sections found in the literature.

The infrared bands of H<sub>2</sub>S and NH<sub>3</sub> are shown in Fig. 3 with asterisks (\*), and the values are listed in Table 2. When very little product is present on the surface, the only relevant and usable signature of H<sub>2</sub>S is a small absorption band at 2570 cm<sup>-1</sup>, testifying to the HS<sup>-</sup> stretching mode ( $\nu_3$ ), in agreement with other studies (Jiménez-Escobar et al. 2014; Loeffler et al. 2015; Oba et al. 2018). The detection of H<sub>2</sub>S with a grazing infrared beam is rather difficult, due to the very broad nature of its main signature and to the associated low-band strength, determined directly from experimental measurements at  $A(\text{S-H}) = 1.12 \times 10^{-17} \text{ cm molecule}^{-1}$  (Hudson & Gerakines 2018). Ammonia infrared signatures are present at 3380



**Figure 3.** Infrared spectra of 0.65 monolayers of H<sub>2</sub>S (bottom curve) and of 1 monolayer of NH<sub>3</sub> (top curve). The species signatures are shown with \*, while the diamonds (◇) indicate residual water present on the detector window, that has not been deposited on the surface.

and 1642 cm<sup>-1</sup> ( $\nu_3$  and  $\nu_4$  NH bending modes of NH<sub>3</sub>). They also correspond to the known features documented in the literature (Gálvez et al. 2010; Kruczkiewicz et al. 2021). On both spectra, as well as in the other spectra presented in this work, the broad, sharp absorption band at 3205 cm<sup>-1</sup>, indicated by diamonds (◇), is due to residual water condensed on the IR detector window.

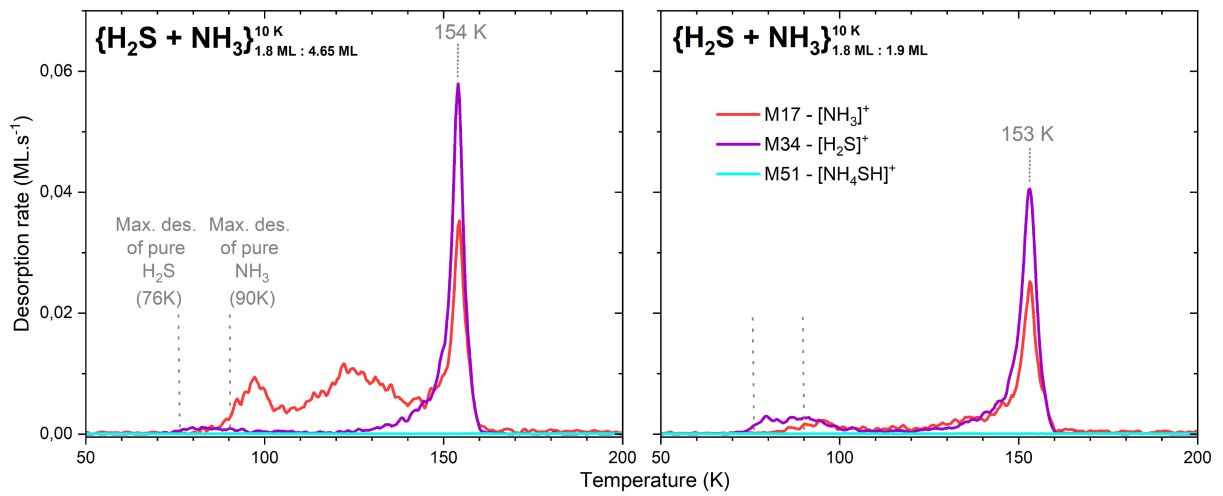
### 3.2 Salt formation: H<sub>2</sub>S:NH<sub>3</sub> mixtures

When mixed together, ammonia and hydrogen sulfide interact to form ammonium hydrosulfide. Fig. 4 shows the TPD results of experiment (iv) on the left panel, and of experiment (v) on the right panel. In both cases, NH<sub>3</sub> and H<sub>2</sub>S gases were codeposited on the surface held at 10 K, via two different molecular beams. In (iv), NH<sub>3</sub> was in excess (4.65 ML) compared to H<sub>2</sub>S (1.8 ML), while in (v) both formed an  $\sim 1:1$  mixture, (1.9 ML of NH<sub>3</sub> with 1.8 ML of H<sub>2</sub>S). Only the desorption curves of the major fragments resulting from the temperature increase are shown. The absence of signal for the  $m/z = 51$  confirms that, being a salt, NH<sub>4</sub>SH is an electrovalent compound and does not contain any covalent bond forming a molecule, but rather two ions HS<sup>-</sup> and NH<sub>4</sub><sup>+</sup> (Bragin et al. 1977; Loeffler et al. 2015). This is the case for many salts (Raunier et al. 2003; Noble et al. 2013; Kruczkiewicz et al. 2021; Gerakines et al. 2024). Therefore, we observe  $m/z = 34$  and  $m/z = 17$  as the main components of the salt, desorbing at 153 and 154 K in the two experiments. With very little desorption of pure molecular H<sub>2</sub>S and NH<sub>3</sub>, we note that 6 per cent and 18 per cent of H<sub>2</sub>S remain pure in the exp. (iv) and (v) respectively, while 17 per cent and 14 per cent of NH<sub>3</sub> remain pure in exp. (iv) and (v), respectively. These results, together with those shown before in Fig. 2, confirm the protection of both NH<sub>3</sub> and H<sub>2</sub>S against the temperature increase. They are stored in salt form at temperatures at which the volatile species would already have desorbed, only to be released in their original form during the sublimation. This behaviour is however not unique to NH<sub>4</sub>SH; most salts tend to dissociate into two more volatile species, when facing a sufficient temperature increase (Noble et al. 2013; Kruczkiewicz et al. 2021; Gerakines et al. 2024).

By calculating the ratio between two representative masses of a chemical species, it is possible to evaluate its conservation in the

**Table 2.** List of infrared features observed in our experiments and found in the literature. <sup>(a)</sup>(Smith 1991), <sup>(b)</sup>(Oba et al. 2018), <sup>(c)</sup>(Kruczkiewicz et al. 2021), <sup>(d)</sup>(Gálvez et al. 2010), <sup>(e)</sup>(Ferraro, Sill & Fink 1980), <sup>(f)</sup>(Bragin et al. 1977), <sup>(g)</sup>(Hudson & Gerakines 2018), <sup>(h)</sup>(Fathe et al. 2006), <sup>(i)</sup>(Loeffler et al. 2015), <sup>(j)</sup>(d’Hedencourt & Allamandola 1986), <sup>(k)</sup>(Bouilloud et al. 2015), <sup>(l)</sup>(Hudgins et al. 1993).

Experiment	Exp. no.	H <sub>2</sub> S		NH <sub>3</sub>		NH <sub>4</sub> SH		H <sub>2</sub> O		
		$\nu_1^{(a)}$	$\nu_3^{(a,b,g,h)}$	$\nu_3^{(c,d)}$	$\nu_4^{(c,d)}$	$\nu_4^{(c,d,i)}$	$\nu_3^{(e)}$	$\nu_4^{(e,f)}$	$\nu_1^{(l)}$	$\nu_2^{(j,k)}$
{H <sub>2</sub> S}	i	–	2568	–	–	–	–	–	–	–
{NH <sub>3</sub> }	ii	–	–	3383	1642	–	–	–	–	–
{H <sub>2</sub> S + NH <sub>3</sub> }	iv	–	–	3364	1631	1488	3007	1814	–	–
	v	–	2567	3371	1631	1483	–	–	–	–
{H <sub>2</sub> S + NH <sub>3</sub> } + {H}	vi	–	–	3371	–	1475	3007	–	–	–
{H <sub>2</sub> S + NH <sub>3</sub> + H}	vii	–	2572	3366	1628	1480	3003	–	–	–
{H <sub>2</sub> S + NH <sub>3</sub> + H <sub>2</sub> O}	viii	–	2554	3390	1641	1478	–	–	3312	1684
	ix	2527	–	3385	1644	1487	–	–	3336	1681
{NH <sub>3</sub> + H <sub>2</sub> O}	x	–	–	3398	1643	–	–	–	3341	1690
{H <sub>2</sub> S + NH <sub>3</sub> + H <sub>2</sub> O + H}	xi	–	2551	3388	1643	1478	3009	–	3310	1681



**Figure 4.** TPD spectra of H<sub>2</sub>S:NH<sub>3</sub> mixtures. The left panel shows results of the experiment (iv), where NH<sub>3</sub> was more abundant than H<sub>2</sub>S on the surface. The right panel shows the results of the experiment (v), where NH<sub>3</sub> and H<sub>2</sub>S formed a homogeneous mixture on the surface. Only the major fragments are shown. The grey dotted lines indicate the temperatures at which each product desorbs in its pure form.

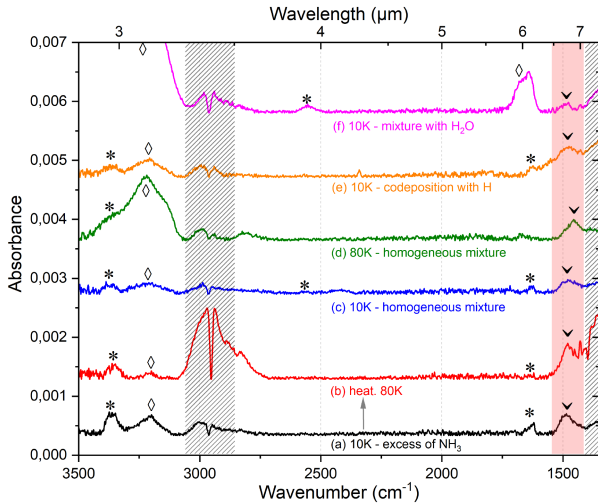
final detection by the QMS. In fact, with a fixed cracking pattern and non-varying experimental conditions, the ratio between each of the fragments is supposed to remain identical. However, if a reaction involving a different species produces a new product, the latter will have different binding, cohesive, and interaction energies, thus altering the proportions between the masses. To determine whether the salt desorbs as an independent entity (namely a hypothetical molecular NH<sub>4</sub>SH) or breaks down into two pure species, we calculated the ratios between  $m/z = 33$  and  $34$  for H<sub>2</sub>S, and  $m/z = 16$  and  $17$  for NH<sub>3</sub>. The values are reported in Table 3. We find that the 33/34 ratio is not different in pure H<sub>2</sub>S or in the salt, being constant at  $\sim 0.4$ . The 16/17 ratio is slightly different when calculated from a TPD of purely NH<sub>3</sub> or of salt ( $\sim 0.55$  and  $\sim 0.5$ , respectively). However, the total ratio remains in the same range, indicating the dissociation of the salt into two different species.

In the experiment with an excess of NH<sub>3</sub> (iv), there is also a peculiar desorption of  $m/z = 17$  between 90 and 142 K, with a maximum at 117 K. It represents 46 per cent of the total ammonia sent on the surface during this experiment. Ammonia ice alone shows a peak at 90 K, but if it is likely to form H-bonds with other components on the surface, especially with water, its desorption is delayed towards higher temperatures, as shown by the dashed grey line of Fig. 1

**Table 3.** Proportions between different  $m/z$  ratios in individual species and in the salt.

Experiment	33/34	16/17
{H <sub>2</sub> S} (i)	0.41	–
{NH <sub>3</sub> } (ii)	–	0.56
{H <sub>2</sub> S + H} (iii)	0.41	–
{H <sub>2</sub> S + NH <sub>3</sub> } – excess of NH <sub>3</sub> (iv)	0.41	0.51
{H <sub>2</sub> S + NH <sub>3</sub> } – homog. mixture (v)	0.40	0.48
{H <sub>2</sub> S + NH <sub>3</sub> } + {H} (vi)	0.41	0.57
{H <sub>2</sub> S + NH <sub>3</sub> + H <sub>2</sub> O} (viii)	0.40	–
{H <sub>2</sub> S + NH <sub>3</sub> + H <sub>2</sub> O} (ix)	0.39	–

(lower panel) (Suresh et al. 2024). In the case where ammonia is mixed with H<sub>2</sub>S, we believe that the same mechanism occurs between NH<sub>3</sub> and the salt. The cracking pattern of ammonia is preserved, suggesting that it is not chemically bonded to any compound, but is only slightly interacting with NH<sub>4</sub>SH. We refer to these molecules as ‘interacting NH<sub>3</sub>’, as opposed to ‘non-interacting NH<sub>3</sub>’, which characterizes molecules that do not interact with the salt. We also note that some of the ‘non-interacting NH<sub>3</sub>’ could be related to the



**Figure 5.** Infrared spectra for the main experiments. *a, c, d, e, f* correspond to experiments iv, v, vi, vii, viii. The spectrum *b* was recorded at 80 K, during the heating phase after the deposition of {H<sub>2</sub>S + NH<sub>3</sub>} (iv). The spectrum *d* was recorded at 80 K, before the H bombardment phase of the salt (experiment vi). The arrows ∇ indicate the salt signature and the symbol \* denotes the signatures of pure H<sub>2</sub>S and NH<sub>3</sub>. The symbol ◇ indicates the bands of H<sub>2</sub>O, whether caused by ice condensed on the MCT detector window (*a, b, c, d*), or ice intentionally deposited on the surface (*f*), or condensed on the surface due to residual pressure of H<sub>2</sub>O in the main chamber (*d*).

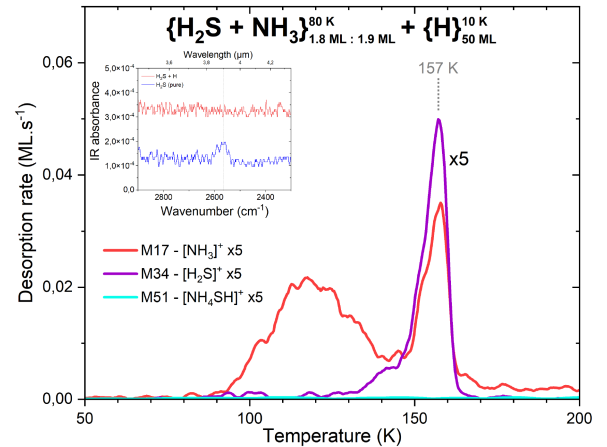
uneven coverage between the two beams used to mix NH<sub>3</sub> and H<sub>2</sub>S. However, this amount should not exceed 10 per cent.

In the mid-IR region, NH<sub>4</sub>SH has visible spectral signatures. The spectra associated with simple H<sub>2</sub>S:NH<sub>3</sub> mixtures (experiments iv, v, and first phase of vi), correspond to the curves *a, c, and d* in Fig. 5. The absorption bands and corresponding wavenumbers are listed in Table 2. Already at 10 K, the most prominent feature observed in all NH<sub>4</sub>SH experiments performed in this work is the ν<sub>4</sub> NH bending mode of NH<sub>4</sub><sup>+</sup>, which appears around 1485 cm<sup>-1</sup>. As more NH<sub>4</sub>SH is formed when more NH<sub>3</sub> molecules are available on the sample (exp. iv), the absorbance increases and the band becomes more prominent. Similarly, we detect larger NH<sub>3</sub> signatures (ν<sub>3</sub> near 3370 cm<sup>-1</sup> and ν<sub>4</sub> near 1635 cm<sup>-1</sup>) due to its excess, in the spectrum *a*. Despite a high noise level, mainly caused by the vibrations of the experimental setup and the initial distortion of the baseline, delimited by hatching in Fig. 5, we attribute the 2900–3000 cm<sup>-1</sup> absorptions to the NH<sub>4</sub><sup>+</sup> ion (Ferraro et al. 1980).

As the sample was gradually heated in experiment (iv), we monitored the evolution of the ice with the FT-IR. The main signature of NH<sub>4</sub>SH at 1485 cm<sup>-1</sup> becomes the sharpest when the temperature of the sample reaches 80 K. This agrees well with the spectrum obtained during the deposition at a constant temperature of 80 K (profile *d*, first phase of exp. vi). Whereas Bragin et al. (1977) find that the salt crystallizes around 100 K, Loeffler et al. (2015) observed amorphous NH<sub>4</sub>SH up to 120 K, followed by its crystallization between 130 and 160 K. On the other hand, the crystallization phase of H<sub>2</sub>S occurs between 70 and 80 K (Fathe et al. 2006), which corresponds well to the sharp maximum observed under our experimental conditions.

### 3.3 Salt behaviour under H atoms exposure

To the best of our knowledge, the behaviour of NH<sub>4</sub>SH exposed to H atoms has not been studied in the laboratory under conditions relevant to the ISM. Studies on NH<sub>4</sub>SH have mainly focused on



**Figure 6.** TPD spectra resulting from the {NH<sub>4</sub>SH} + {H} experiment (vi). Only the major fragments are shown. Each of them has been multiplied by 5 for readability and scale consistency purposes. *Inset*: infrared spectra of H<sub>2</sub>S (bottom curve) and of H<sub>2</sub>S + H codeposited (top curve). The grey dashed line shows the absorption band reference (2570 cm<sup>-1</sup>).

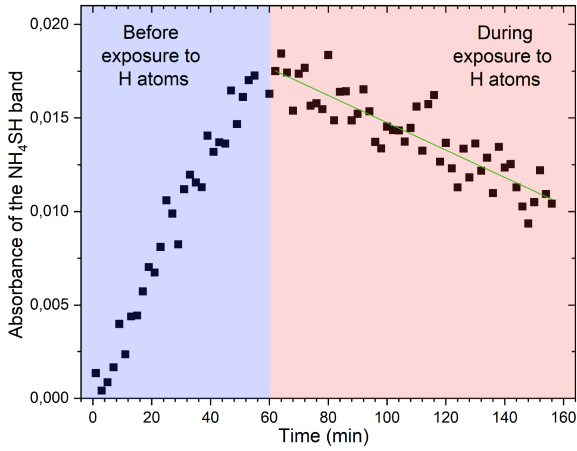
its behaviour in mixed ices (Altwegg et al. 2020), under photolysis (Lebofsky & Fegley 1976), with temperature changes, under proton irradiation (Loeffler et al. 2015, 2016; Loeffler & Hudson 2018), and under electron irradiation (Mahjoub et al. 2016). In contrast, the H-exposure of H<sub>2</sub>S has been the subject of some experimental works (Oba et al. 2018; El Akel et al. 2022; Santos, Linnartz & Chuang 2023) and computational studies (Garrod & Herbst 2007; Lamberts & Kästner 2017). In this study, three types of H<sub>2</sub>S and NH<sub>4</sub>SH H-exposure experiments were performed (Table 1, iii, vi, vii). In all cases, the H atoms flux remained constant at φ<sub>H</sub> = 8 × 10<sup>12</sup> atoms cm<sup>-2</sup> s<sup>-1</sup>.

We found the exposure of H<sub>2</sub>S to H atoms to be very effective. With a continuous flow of hydrogen atoms during the deposition, the final amount of product detected in {H<sub>2</sub>S + H} corresponds to about 38 per cent of that of a signal of unaltered {H<sub>2</sub>S}. However, its infrared signature becomes practically invisible (see inset of Fig. 6), giving us information about the signal-to-noise limit for the detection of H<sub>2</sub>S with our RAIRS instrument, which is about 0.25 ML (2.5 × 10<sup>14</sup> molecules cm<sup>-2</sup>). The loss of product can be explained by the chemical desorption that takes place during the reaction processes



According to equation (1) and equation (2), once H<sub>2</sub>S undergoes H abstraction to form an HS radical, it is likely to either recombine with another H to reform H<sub>2</sub>S or undergo chemical desorption, that is to be ejected by ejected from the surface and return to the gas phase due to excess energy that is too large to dissipate (Dulieu et al. 2013). The competition between the different surface processes, together with the H + H addition, has previously been studied quantitatively (Lamberts & Kästner 2017). The cracking patterns of the residual H<sub>2</sub>S remain identical, except for the *m/z* = 35 and 36 which are significantly reduced, possibly indicating a preferential chemical desorption of H<sub>2</sub><sup>34</sup>S and H<sup>34</sup>S over H<sub>2</sub><sup>32</sup>S.

Exposure of the salt to hydrogen atoms is also effective in reducing the amount of product remaining on the surface, although not as fast as the {H<sub>2</sub>S + H} reaction. Fig. 7 shows the time evolution of the infrared band regions of the main NH<sub>4</sub>SH signature (at 1485 cm<sup>-1</sup>) during the experiment (vi). The salt was grown by



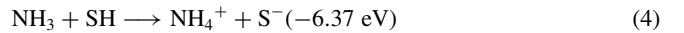
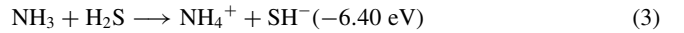
**Figure 7.** Scatter plot of the infrared band areas of the main signature of  $\text{NH}_4\text{SH}$  ( $1485\text{ cm}^{-1}$ ) as a function of time, in the experiment vi. The salt was grown for 60 min at 80 K and then exposed to H atoms for 100 min at 10 K with our atomic beam. The decreasing solid line is a linear fit of the decrease in the IR band area.

one hour  $\{\text{H}_2\text{S} + \text{NH}_3\}$  codeposition on the surface, which was maintained at 80 K for infrared data readability purposes. We were careful to maintain stable conditions to ensure the lowest possible presence of pure, non-interacting  $\text{NH}_3$  and  $\text{H}_2\text{S}$  on the sample. The surface was then cooled to 10 K before being exposed to hydrogen atoms for 100 min, corresponding to approximately 43 ML of H. The two phases (deposition and H-exposure) are highlighted in Fig. 7 with a blue and red background colour, respectively. We estimate that the salt is destroyed with a half-life of 120 min by sending  $\phi_{\text{H}} = 8 \times 10^{12}$  atoms  $\text{cm}^{-2} \text{ s}^{-1}$ , which is about 1/3 of the destruction obtained after sending  $\sim 80$  times more hydrogen atoms than the amount of  $\text{NH}_4\text{SH}$  present on the surface. It should be noted, however, that this rate of destruction may depend on the intrinsic mobility of H on the surface, which may be slightly different depending on the experimental conditions. In fact, the reformation of  $\text{H}_2$  consumes most of the H atoms sent to the surface. Here, the destruction of the salt can be fitted with a linear rather than an exponential curve, meaning that the reactivity of the salt can be complete or end with a small plateau. We therefore conclude that, while the salt form does not prevent all hydrogenation or abstraction reactions from occurring, it greatly reduces the reactivity rate, compared to when only molecular  $\text{H}_2\text{S}$  is exposed to H.

The TPD traces obtained as a result of this two-phase experiment are shown in the Fig. 6. Each fragment curve has been multiplied by 5, in order to keep the same scale as before. Since there is less  $\text{NH}_4\text{SH}$  on the surface, the shift of its desorption temperature from 153 to 157 K is not surprising. As mentioned above, the deposition conditions at 80 K were such that the amounts of pure  $\text{NH}_3$  and  $\text{H}_2\text{S}$  were as low as possible. Despite this, we detect 1.1 ML of excess  $\text{NH}_3$  as its ‘interacting’ form after H-exposure, in addition to 0.6 ML of  $\text{NH}_4^+$  constituting the salt, which gives a total amount of 1.7 ML of residual ammonia on the surface (i.e.  $\sim 90$  per cent of the total amount of ammonia sent by the beams). An equivalent quantity of 0.55 ML of  $\text{H}_2\text{S}$  constituting the salt remains on the sample after sending H atoms (i.e.  $\sim 31$  percent of the initial  $\text{H}_2\text{S}$ ). We also conducted a triple codeposition experiment (vii), in which  $\text{NH}_3$  molecules,  $\text{H}_2\text{S}$  molecules, and hydrogen atoms were simultaneously sent to the surface at 10 K, with fluxes identical to experiment (vi). In contrast to the results of experiment (vi), in experiment (vii) small amounts of pure non-interacting  $\text{NH}_3$  and  $\text{H}_2\text{S}$  remained on the surface (0.2 ML

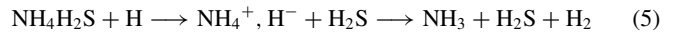
and 0.1 ML, respectively), in addition to a continuous desorption of  $\text{NH}_3$  (90–150 K) due to H-bonding (0.85 ML). In this case, between 0.7 and 0.8 ML of each species is incorporated into the salt. Again, it appears that  $\sim 95$  per cent of  $\text{NH}_3$  remains on the surface after H-exposure, whereas for  $\text{H}_2\text{S}$  it is  $\sim 46$  per cent. The infrared spectrum associated with this experiment is shown in Fig. 5(e). Apart from larger  $\text{NH}_3$  signatures, no other special features were observed. The small sharp peak at  $2342\text{ cm}^{-1}$  corresponds to  $\text{CO}_2$  due to a minor leak.

Overall, our results show that the consumption of  $\text{H}_2\text{S}$  when facing a hydrogen atoms exposure, even in salt form, does occur, but at a reduced rate. The reaction of  $\text{NH}_4\text{SH}$  with hydrogen atoms may disrupt the ionic crystal structure of the salt to form  $\text{NH}_3 + \text{SH} + \text{H}_2$ . The process is globally exothermic and is equivalent to the reaction described by reaction 1. In this case, the SH formed could react with the excess of hydrogen atoms to reform  $\text{H}_2\text{S}$ , some of which would pass into the gas phase by chemical desorption. This could explain the loss of  $\text{NH}_4\text{SH}$  and  $\text{H}_2\text{S}$  observed experimentally. However, this reaction does not seem possible because the first step should be  $\text{NH}_4\text{SH} + \text{H} \rightarrow \text{NH}_3\text{SH} + \text{H}_2$ , which is endothermic and should rather form  $\text{NH}_4\text{S}$ . In fact, the ionic crystal energy required to form  $\text{NH}_4\text{S}$  is similar to that required to form the salt  $\text{NH}_4\text{SH}$ :

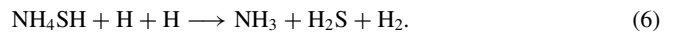


In reaction 4, if there is an excess of hydrogen atoms,  $\text{NH}_4\text{SH}$  should be formed, which then acts as a catalyst for the hydrogen recombination reaction ( $\text{H} + \text{H} \rightarrow \text{H}_2$ ), without loss of  $\text{H}_2\text{S}$  by chemical desorption.

Another possibility is the formation of the salt  $\text{NH}_4\text{H}_2\text{S}$  ( $\text{NH}_4^+, \text{H}_2\text{S}, \text{e}^-$ ). With an excess of H, we would then have the following endothermic reaction:



The general reaction would be:

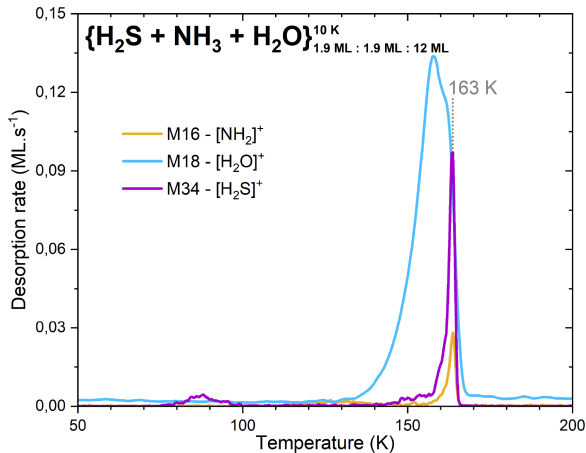


Some of the  $\text{H}_2\text{S}$  formed will give back  $\text{NH}_4\text{SH}$ , but some may also react with hydrogen atoms to form  $\text{H}_2 + \text{SH}$  and lead to a loss of  $\text{H}_2\text{S}$  by chemical desorption during its reformation. It is clear from our results that no pure  $\text{H}_2\text{S}$  remains on the surface after 100 min of exposure to H. The fact that  $\text{NH}_4\text{SH}$  is less sensitive to hydrogenation than  $\text{H}_2\text{S}$  can be explained by the fact that a part of the salt is reformed after reaction with H, while the other part is consumed indirectly by the chemical desorption of  $\text{H}_2\text{S}$  according to equation (1), and by the fact that the whole mechanism requires several hydrogen atoms.

Looking at the infrared spectra obtained for this experiment, we observe a progressive destruction of the salt signature at  $1485\text{ cm}^{-1}$  during the H-exposure phase. No particular new features were observed, and no quantification of the  $\text{NH}_3$  reformation could be calculated due to a severe baseline deformation.

### 3.4 Salt behaviour in water ice mixtures

The behaviour of  $\text{NH}_4\text{SH}$  in the environment of other compounds was investigated by Altwegg et al. (2020), who carried out TPD experiments where  $\text{CH}_3\text{OH}$ ,  $\text{NH}_3$ ,  $\text{H}_2\text{S}$ , and  $\text{H}_2\text{O}$  were simultaneously deposited on a gold-coated surface held at 50 K. They observed the formation of ammonium hydrosulfide in the ice mixture, as in previous works where the use of  $^{15}\text{N}$  confirmed the  $\text{NH}_3$  chemistry



**Figure 8.** TPD spectra of a H<sub>2</sub>S:NH<sub>3</sub>:H<sub>2</sub>O mixture, where all species have been codeposited at 10 K. The fragment  $m/z = 16$  accounts for [NH<sub>2</sub>]<sup>+</sup>, to avoid the confusion between [NH<sub>3</sub>]<sup>+</sup> and [OH]<sup>+</sup> that  $m/z = 17$  would result in.

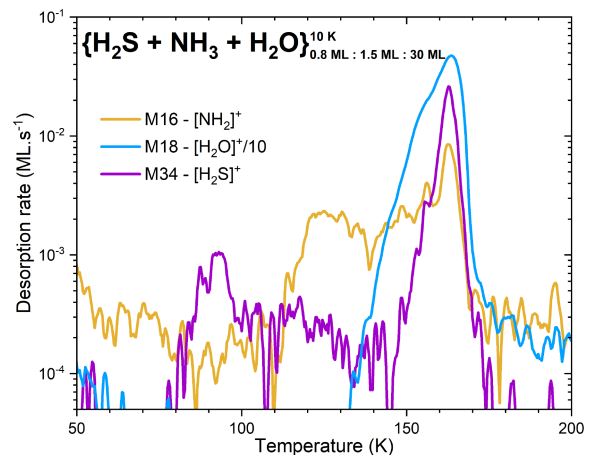
in a similar mixture before its irradiation by electrons and thermal processing (Mahjoub et al. 2016).

Other works have studied more generally the behaviour of ammonium salts in water ice mantles. For example, for salts that desorb at temperatures higher than water (NH<sub>4</sub><sup>+</sup>HCOO<sup>-</sup>, NH<sub>4</sub><sup>+</sup>CH<sub>3</sub>COO<sup>-</sup>, NH<sub>4</sub><sup>+</sup>CN<sup>-</sup>, among others), no or little influence of H<sub>2</sub>O on the TPD profiles of the salts and on the IR spectra has been registered (Noble et al. 2013; Kruczkiewicz et al. 2021). Conversely, for a salt that sublimates in the water desorption region like NH<sub>4</sub>SH, we found a shift of 10 K in the desorption temperature, as it can be seen in Fig. 8 which shows the results of experiment (viii). The kinetic order of the desorption does not change, but its maximum appears at the very end of the H<sub>2</sub>O desorption, towards 163 K. The NH<sub>4</sub>SH binding energy would therefore be greater for H<sub>2</sub>O than for gold, implying that even in a codeposition situation, NH<sub>4</sub>SH remains bound to water until there is not enough left to hold it on the now too hot surface. We also note that this happens after the water has crystallized.

Given the double assignment of  $m/z = 17$  to ionized OH<sup>+</sup> and NH<sub>3</sub><sup>+</sup>,  $m/z = 16$  has been chosen here to represent NH<sub>3</sub>. In this experiment, there is no non-interacting NH<sub>3</sub> on the surface; only 0.6 ML of ammonia interacts with either water or NH<sub>4</sub>SH. In comparison, 70 percent and 86 percent of the total quantity of ammonia and hydrogen sulfide are constituting the NH<sub>4</sub>SH salt.

Regarding the infrared results, the major changes undoubtedly concern the water band, located in the absorption regions of NH<sub>3</sub> (3400 cm<sup>-1</sup> for the –OH stretching mode and 1650 cm<sup>-1</sup> for the H<sub>2</sub>O bending mode). Additionally, a slight blue-shift occurs for the H<sub>2</sub>S absorption, measured at 2554 cm<sup>-1</sup> instead of ~2570 cm<sup>-1</sup> (see Table 2).

In order to verify the influence of a water ice matrix on the desorption of the salt under more astrophysically relevant conditions, we performed experiment (ix). The TPD results are displayed in Fig. 9. The IR spectrum is not shown, due to the lack of signal other than H<sub>2</sub>O. Approximately 0.8 ML of H<sub>2</sub>S and 1.5 ML of NH<sub>3</sub> have been diluted in 30 ML of H<sub>2</sub>O at 10 K, as a result of regular valve openings throughout the deposition. The readability of the results is therefore improved by using a logarithmic scale and dividing the H<sub>2</sub>O component by a factor of 10. As in experiment (viii), very small returns from NH<sub>4</sub>SH to H<sub>2</sub>S and NH<sub>3</sub> are visible (0.1 ML of each, approximately). The desorption is also shifted to 163 K, following the same pattern as the end of the water desorption. Our results are



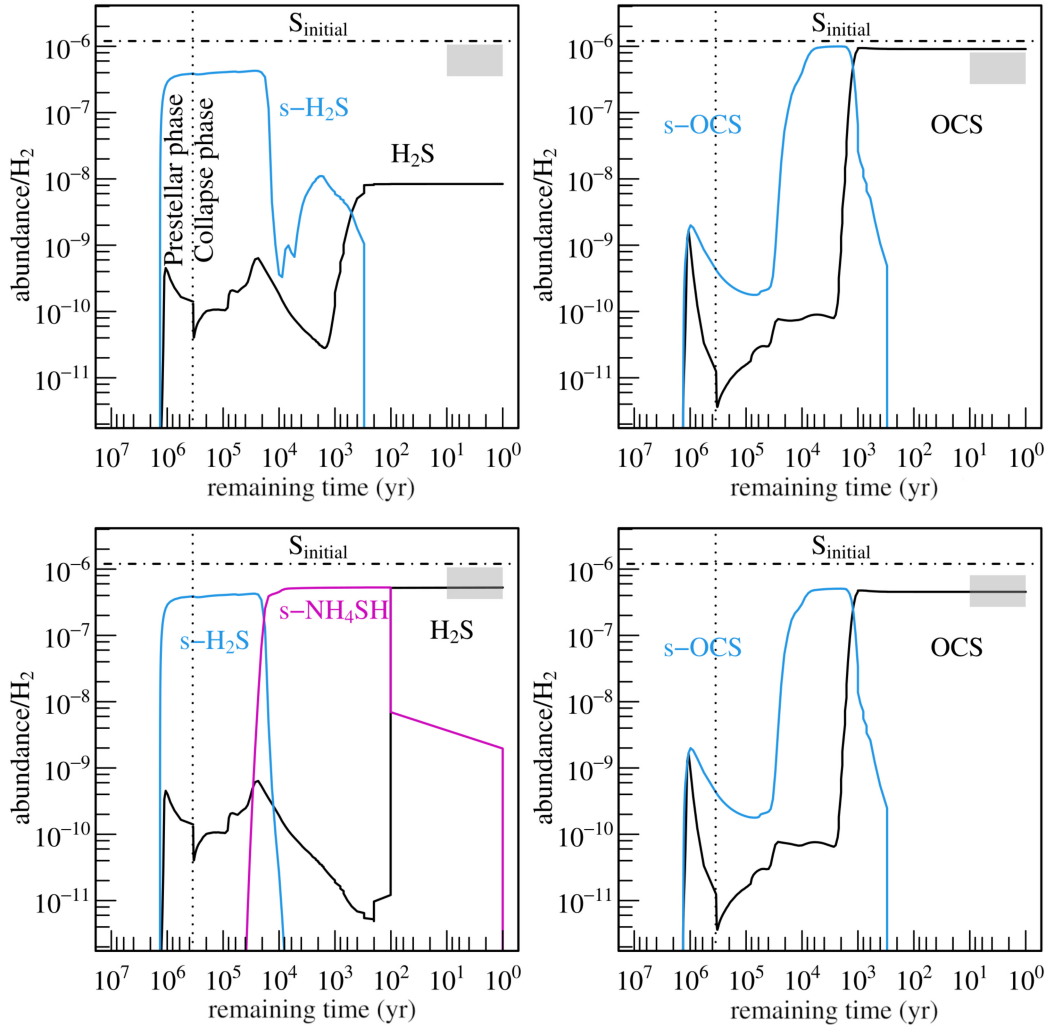
**Figure 9.** TPD spectra of a H<sub>2</sub>S:NH<sub>3</sub>:H<sub>2</sub>O mixture in more astrophysically relevant conditions. About 1.5 ML of NH<sub>3</sub> and 0.8 ML of H<sub>2</sub>S have been deposited bit by bit in a water ice matrix of ~30 ML. A logarithmic scale is used for readability purposes.

consistent with the single NH<sub>4</sub>SH TPD spectrum published to date, where the desorption peak occurs at 168 K in a mixed ice containing H<sub>2</sub>O, CH<sub>3</sub>OH, H<sub>2</sub>S, and NH<sub>3</sub> (Altwegg et al. 2020).

Overall, water seems to affect the desorption and decomposition of NH<sub>4</sub>SH only by delaying it, shifting its temperature by 10 K under our experimental conditions, regardless of the mixture ratios between different species.

#### 4 ASTROPHYSICAL IMPLICATIONS

H<sub>2</sub>S and OCS are the main sulphur species observed in the gas phase in protostars (Kushwahaa et al. 2023), with variable ratios. It is therefore important to understand the evolution of their ratios through modelling, as well as the potential role of salts that form when *s*-H<sub>2</sub>S (*s*- meaning solid, on surfaces) is mobile on ice, during a temperature increase. To model the abundances in protostars, we used the Nautilus gas–grain model in its three-phase form (Ruaud, Wakelam & Hersant 2016). The model was used to simulate the abundances of atoms and molecules in neutral and ionic form as a function of time, employing kida.uva.2014 (Wakelam et al. 2015). The basic chemical network has recently been updated to better describe COMs on grains and in the gas phase (Manigand et al. 2021; Loison et al. 2022; Agúndez et al. 2023), but also to describe the formation of salts. 800 species are identified in the network and are involved in 9000 individual reactions. In the model that we used, the elements are initially in their atomic or ionic form (the ones with an ionization potential below 13.6 eV are considered to be fully ionized), and the C/O elemental ratio is equal to 0.71. The initial sulphur abundance is set to  $6.0 \times 10^{-7}$  with respect to H, corresponding to a depletion by a factor of 20 compared to the typical cosmic abundance in dense molecular clouds (Fuente et al. 2016). The grain surface and the mantle are both chemically active in these simulations, while the accretion and desorption processes are allowed only between the surface and the gas phase, except for the sputtering by cosmic rays. The dust-to-gas ratio is 0.01 by mass. A sticking probability of 1 is assumed for all neutral species, and desorption can occur via thermal and non-thermal processes, including chemical desorption and sputtering of ices by cosmic ray collisions (Wakelam et al. 2021). The formalism of the surface reactions and a more



**Figure 10.** Fractional abundances ( $n/n_{\text{H}_2}$ ) of  $\text{H}_2\text{S}$ , and  $\text{NH}_4\text{SH}$  (left panels) and  $\text{OCS}$  (right panels) in gas phase (black lines) and on the grain surface (blue and purple lines) as a function of time during the collapse phase predicted by the Nautilus astrochemical model. The time axis is reversed to better visualize the abundance evolution. Upper panels: without  $\text{NH}_4\text{SH}$  formation pathway. Lower panels: with inclusion of  $\text{NH}_4\text{SH}$  formation pathway. The abundance of sulphur initially in the gas phase is fixed at  $1.2 \times 10^{-6}$ , that is a depletion factor of 20 compared to the solar abundance (Asplund et al. 2009). The grey rectangle at the top right of each panel represents the observed abundances towards IRAS 16293–2422B (Kushwahaa et al. 2023), assuming an uncertainty of  $3 \times 10^{-7}$  in the y-axis.

detailed description of the simulations can be found in Ruaud et al. (2016).

To study the chemistry in protostellar envelopes, we have used the same approach as in Manigand et al. (2021). We perform two steps: a prestellar and a collapse phase. In the first step, conditions typical of the cold core are assumed for  $10^6$  yr (gas and dust temperature of 10 K, a proton density of  $2 \times 10^4 \text{ cm}^{-3}$ , a visual extinction of 4.5 mag, a cosmic ray ionization rate of  $1.3 \times 10^{-17} \text{ s}^{-1}$ , and a standard external UV field of 1 G0). The final composition of the cold core is then used as the initial condition for the collapse phase, which is the same as in Aikawa et al. (2008, 2012), Wakelam et al. (2014), and Manigand et al. (2021). The selected parcel of gas begins to collapse at a radius of about 10 000 au and ends-up at a radius of 15 au, where the gas temperature reaches  $\sim 265$  K and a density of  $4 \times 10^8 \text{ cm}^{-3}$ .

In the version of the NAUTILUS code used in this work, we have introduced two types of salts: salts stable in the gas phase and ammonium salts, that would form by the Langmuir–Hinshelwood

mechanism. The former, such as table salt  $\text{Na}^+\text{Cl}^-$ , can be considered as classical molecules, even if they dissociate as ions in ice (Acharyya, Woon & Herbst 2023). On the other hand, ammonium salts do not exist in the gas phase and are only stable when the ions resulting from a proton transfer from an acid to the base (here, ammonia) are sufficiently stabilized by interactions with water molecules, in order to compensate for the exothermicity of the proton transfer. Calculating the interaction energy between ions and the structure of amorphous ices is therefore very complicated, as it is necessary to estimate for which systems these energies are sufficient to ensure the formation of ion pairs, that is salts. To do this, we have used a method similar to that of Woon (2012), which consists in calculating the most stable structures for acid-base- $(\text{H}_2\text{O})_n$  systems by varying the size of the water cluster. If a proton transfer occurs for  $n \leq 8$ , we consider the salt formation, barrierless when exothermic (Woon 2012). The ammonium salts thus introduced are  $\text{NH}_4^+\text{CN}^-$ ,  $\text{NH}_4^+\text{Cl}^-$ ,  $\text{NH}_4^+\text{NCO}^-$ ,  $\text{NH}_4^+\text{NCS}^-$ ,  $\text{NH}_4^+\text{HCOO}^-$ , and  $\text{NH}_4^+\text{SH}^-$ . This would explain, among other things, the ice

absorption band usually observed at 2165 cm<sup>-1</sup> (4.62 μm) and attributed to OCN<sup>-</sup>. However, since the bands attributed to OCN<sup>-</sup> and NH<sub>4</sub><sup>+</sup> are seen in dense molecular cloud ices (McClure et al. 2023), it also implies that either NH<sub>3</sub>, HNCO, or HOCN are mobile at low temperatures, so that the diffusion energy is much lower than the usual value of 0.4 times the adsorption energy.

In this first version of the NAUTILUS code with salts taken into account, and despite our experimental results showing a significant reactivity, even reduced, of salts with hydrogen atoms, we do not include salt destruction reactions except for desorption. Due to the instability of ammonium salts in the gas phase, unlike usual molecules, we cannot calculate their desorption temperature with the usual models, including the proxy approach developed in Wakelam et al. (2017). We have therefore chosen to set the desorption energies for the destructive process  $s\text{-NH}_4\text{SH} \rightarrow \text{NH}_3 + \text{H}_2\text{S}$  so as to reproduce the experimental sublimation temperature of about 150 K for  $s\text{-NH}_4\text{SH}$ , corresponding to a binding energy of about 6000 K, slightly larger than that of water.

In models of ice formation in interstellar clouds, the gas-phase reactions of S<sup>+</sup> and S are not efficient enough to consume significant amounts of sulphur, which will eventually stick to the ice. The evolution of solid sulphur at low temperatures in dense clouds (around 10 K) is then dominated by reactions with atomic hydrogen, with the production of  $s\text{-H}_2\text{S}$ . As the H–SH bond of H<sub>2</sub>S is weaker than the H–H bond of H<sub>2</sub>,  $s\text{-H}_2\text{S}$  tends to react with  $s\text{-H}$  by tunneling effect, giving  $s\text{-H}_2 + s\text{-SH}$ . The  $s\text{-SH}$  radical can also react with  $s\text{-H}$  to give mainly  $s\text{-H}_2\text{S}$  (and probably  $s\text{-S}$ ), creating a loop and thus an equilibrium between  $s\text{-H}_2\text{S}$  and  $s\text{-SH}$  (and  $s\text{-S}$ ) (Lamberts & Kästner 2017; Vidal et al. 2017; Oba et al. 2018). Since both  $s\text{-SH}$  and  $s\text{-H}_2\text{S}$  have similar abundances at 10 K in our model and are the main reservoir of sulphur for evolved dense clouds, this equilibrium is controlled by the abundance and mobility of H on the grains.

The resulting model predictions are shown in Fig. 10. Inclusion of the experimental results in the NAUTILUS astrochemical gas–grain model corroborates the importance of solid-state NH<sub>4</sub>SH to better constrain the H<sub>2</sub>S and OCS abundances (Fig. 10, lower panels). When the dense cloud gives rise to the formation of protostars, the temperature in the collapse phase increases, leading to a significant increase in the mobilities of the various species present on the grains, especially radicals. Among the radicals we count  $s\text{-SH}$  and  $s\text{-S}$ , but also  $s\text{-HCO}$ , i.e. the most abundant radical on the grains due to that of  $s\text{-CO}$ , together with notable abundances of  $s\text{-CH}_2\text{OH}$  and  $s\text{-CH}_3\text{O}$ . In our model, the mobilities of  $s\text{-SH}$ ,  $s\text{-S}$ , and  $s\text{-HCO}$  are comparable and inferred from the adsorption energies of Wakelam et al. (2017). The reaction of  $s\text{-S}$  with  $s\text{-HCO}$  gives  $s\text{-SH} + s\text{-CO}$  as well as  $s\text{-H} + s\text{-OCS}$ , while the reaction of  $s\text{-SH}$  and  $s\text{-HCO}$  can give  $s\text{-HCOSH}$ ,  $s\text{-H}_2\text{S} + s\text{-CO}$ , and  $s\text{-H}_2 + s\text{-OCS}$ . The solid  $s\text{-HCOSH}$  species reacts quite easily with H by tunneling effect and are converted fairly rapidly to  $s\text{-OCS}$ . HCOSH was detected towards quiescent clouds in the Galactic Centre (Rodríguez-Almeida et al. 2021; García de la Concepción et al. 2022), which is in good agreement with our model, where the abundance of  $s\text{-HCOSH}$  in protostars is low.

## 5 CONCLUSIONS

In our experimental study, we took a closer look at the formation, reactivity, and desorption of ammonium hydrosulfide salts. These salts are in fact chemical intermediates with specific properties that give them an important role in molecular evolution during star formation processes, but also in the sublimation of ices in general, as in comets. In particular, we have shown that:

(i) NH<sub>4</sub>SH salts are formed at temperatures as low as 10 K from a mixture of H<sub>2</sub>S and NH<sub>3</sub>, and appear to crystallize around 80 K under our experimental conditions.

(ii) Even diluted in water, H<sub>2</sub>S, and NH<sub>3</sub> mixtures lead to the formation of ammonium hydrosulfide salts, and their desorption is shifted to 163 K, that is 10 K higher than in the H<sub>2</sub>S:NH<sub>3</sub> case, regardless of the mixture ratios in the H<sub>2</sub>O:H<sub>2</sub>S:NH<sub>3</sub> system.

(iii) In astrophysical contexts (i.e. protostars or protostellar discs), salts are likely to desorb with water, and in any case at higher temperatures than the individual pure species NH<sub>3</sub> and H<sub>2</sub>S.

(iv) Adopting an effective binding energy of ~6000 K for NH<sub>4</sub>SH and the release of the two components NH<sub>3</sub> and H<sub>2</sub>S in the gas phase seems a reasonable description of the desorption process.

(v) Including salts in present chemical networks helps to find a better agreement with the observed OCS/H<sub>2</sub>S ratio in protostars.

(vi) S-bearing ionic compounds could explain the non-detection of H<sub>2</sub>S in ices to date, as some of it would be in the form of salts.

(vii) Salts do not provide total protection of H<sub>2</sub>S against exposure to hydrogen atoms, but they do slow down the H-abstraction process.

## ACKNOWLEDGEMENTS

This work was funded by CY Initiative of Excellence (grant ‘Investissements d’Avenir’ ANR-16-IDEX-0008), Agence Nationale de la recherche (ANR) SIRC project (Grant ANR-SPV2024482020–2024), by the Programme National ‘Planétologie’ and ‘Physique et Chimie du Milieu Interstellaire’ (PCMI) of CNRS/INSU with INC/INP co-funded by CEA and CNES.

## DATA AVAILABILITY

The experimental data and conditions have been saved on our own laboratory server, under the form of metadata files associated with each experiment. They will be released publicly after a minimum of 3 yr after their production. For the time being, the data are available from the corresponding authors upon reasonable request.

## REFERENCES

- Acharyya K., Woon D. E., Herbst E., 2023, *MNRAS*, 527, 1722  
 Agúndez M. et al., 2023, *A&A*, 673, A34  
 Aikawa Y., Wakelam V., Garrod R. T., Herbst E., 2008, *ApJ*, 674, 984  
 Aikawa Y., Wakelam V., Hersant F., Garrod R. T., Herbst E., 2012, *ApJ*, 760, 40  
 Altwegg K. et al., 2020, *Nat. Astron.*, 4, 533  
 Asplund M., Grevesse N., Sauval A. J., Scott P., 2009, *ARA&A*, 47, 481  
 Atreya S. K., Wong M. H., Owen T. C., Mahaffy P. R., Niemann H. B., Pater I. D., Drossart P., Encrenaz T., 1999, *Planet. Space Sci.*, 47, 1243  
 Bergner J. B., Öberg K. I., Rajappan M., Fayolle E. C., 2016, *ApJ*, 829, 85  
 Bouilloud M., Fray N., Bénilan Y., Cottin H., Gazeau M.-C., Jolly A., 2015, *MNRAS*, 451, 2145  
 Bragin J., Diem M., Guthals D., Chang S., 1977, *J. Chem. Phys.*, 67, 1247  
 Van Broekhuizen F. A., Keane J. V., Schutte W. A., 2004, *A&A*, 415, 425  
 Calmonte U. et al., 2016, *MNRAS*, 462, S253  
 Chen Y. J. et al., 2015, *ApJ*, 798, 80  
 Clutter D. R., Thompson W. E., 1969, *J. Chem. Phys.*, 51, 153  
 Congiu E. et al., 2012, *ApJ*, 750, L12  
 Congiu E., Sow A., Nguyen T., Baouche S., Dulieu F., 2020, *Rev. Sci. Instrum.*, 91, 124504  
 Demyk K., Dartois E., D’Hendecourt L., Muizon M. J. D., Heras A. M., Breitfellner M., 1998, *A&A*, 339, 553  
 Drozdovskaya M. N. et al., 2018, *MNRAS*, 476, 4949  
 Dulieu F., Congiu E., Noble J., Baouche S., Chaabouni H., Moudens A., Minissale M., Cazaux S., 2013, *Sci. Rep.*, 3, 1338

- El Akel M., Kristensen L. E., Gal R. L., Van Der Walt S. J., Pitts R. L., Dulieu F., 2022, *A&A*, 659, A100
- Fathe K., Holt J. S., Oxley S. P., Pursell C. J., 2006, *J. Phys. Chem. A*, 110, 10793
- Ferrante R. H., Moore M. H., Spiliotis M. M., Hudson R. L., 2008, *ApJ*, 684, 1210
- Ferraro J. R., Sill G., Fink U., 1980, *Appl. Spectrosc.*, 34, 525
- Filacchione G. et al., 2019, *Space Sci. Rev.*, 215, 19
- Friedson A. J., 2005, *Icarus*, 177, 1
- Fuente A. et al., 2016, *A&A*, 593, A94
- Gálvez O., Maté B., Herrero V. J., Escribano R., 2010, *ApJ*, 724, 539
- García de la Concepción J. et al., 2022, *A&A*, 658, A150
- Garrod R. T., Herbst E., 2007, *A&A*, 467, 1103
- Gerakines P. A., Moore M. H., Hudson R. L., 2004, *Icarus*, 170, 202
- Gerakines P. A., Yarnall Y. Y., Hudson R. L., 2024, *Icarus*, 413, 116007
- Goicoechea J. R., Pety J., Gerin M., Teyssier D., Roueff E., Hily-Blant P., Baek S., 2006, *A&A*, 456, 565
- d'Hedencourt L. B., Allamandola L. J., 1986, *A&AS*, 64, 453
- Hudgins D. M., Sandford S. A., Allamandola L. J., Tielens A. G. G. M., 1993, *ApJS*, 86, 713
- Hudson R. L., Gerakines P. A., 2018, *ApJ*, 867, 138
- Jiménez-Escobar A., Caro G. M. M., Chen Y. J., 2014, *MNRAS*, 443, 343
- Kruczkiewicz F., Vitorino J., Congiu E., Theulé P., Dulieu F., 2021, *A&A*, 686, A236
- Kumar Y., Kumar M., 2020, *Chem. Phys. Lett.*, 740, 137071
- Kushwahaa T., Drozdovskaya M. N., Tychoniec L., Tabone B., 2023, *A&A*, 672, A122
- Lamberts T., Kästner J., 2017, *J. Phys. Chem. A*, 121, 9736
- Lebofsky L., Fegley M., 1976, *Icarus*, 28, 379
- Li C. et al., 2020, *Nat. Astron.*, 4, 609
- Lodders K., 2010, *Solar System Abundances of the Elements, Astrophysics and Space Science Proceedings*. Springer-Verlag, Berlin
- Loeffler M. J., Hudson R. L., 2018, *Icarus*, 302, 418
- Loeffler M. J., Hudson R. L., Chanover N. J., Simon A. A., 2015, *Icarus*, 258, 181
- Loeffler M. J., Hudson R. L., Chanover N. J., Simon A. A., 2016, *Icarus*, 271, 265
- Loison A. C. J.-C., Boulanger A., Caux E., Müller H. S. P., Wakelam V., Manigand S., Jørgensen J. K., 2022, *A&A*, 660, L6
- López-Gallifa Á. et al., 2024, *MNRAS*, 529, 3244
- Mahjoub A. et al., 2016, *ApJ*, 820, 141
- Manigand S. et al., 2021, *A&A*, 645, A53
- Martín-Doménech R., Jiménez-Serra I., Caro G. M. M., Müller H. S., Occhiogrosso A., Testi L., Woods P. M., Viti S., 2016, *A&A*, 585, A112
- McClure M. K. et al., 2023, *Nat. Astron.*, 7, 431
- Niemann H. B. et al., 1998, *J. Geophys. Res.: Planets*, 103, 22831
- Noble J. A., Dulieu F., Congiu E., Fraser H. J., 2011, *ApJ*, 735, 121
- Noble J. A., Theule P., Borget F., Danger G., Chomat M., Duvernay F., Mispelaer F., Chiavassa T., 2013, *MNRAS*, 428, 3262
- Oba Y., Tomaru T., Lamberts T., Kouchi A., Watanabe N., 2018, *Nat. Astron.*, 2, 228
- Pater I. D., Dunn D., Romani P., Zahnle K., 2001, *Icarus*, 149, 66
- Perrero J., Beitia-Antero L., Fuente A., Ugliengo P., Rimola A., 2024, *MNRAS*, 527, 10697
- Puong N. et al., 2018, *A&A*, 616, L5
- Raunier S., Chiavassa T., Marinelli F., Allouche A., Aycard J. P., 2003, *Chem. Phys. Lett.*, 368, 594
- Rodríguez-Almeida L. F. et al., 2021, *ApJ*, 912, L11
- Roman M. T., Banfield D., Gierasch P. J., 2013, *Icarus*, 225, 93
- Ruaud M., Wakelam V., Hersant F., 2016, *MNRAS*, 459, 3756
- Ruffle D. P., Hartquist T. W., Caselli P., Williams D. A., 1999, *MNRAS*, 306, 691
- Santos J. C., Linnartz H., Chuang K.-J., 2023, *A&A*, 678, A112
- Savage B. D., Sembach K. R., 1996, *ApJ*, 470, 893
- Smith R. G., 1991, *MNRAS*, 249, 172
- Suresh S. K., Dulieu F., Vitorino J., Caselli P., 2024, *A&A*, 682, A163
- Theulé P., Duvernay F., Danger G., Borget F., Bossa J. B., Vinogradoff V., Mispelaer F., Chiavassa T., 2013, *Adv. Space Res.*, 52, 1567
- Tiefrunk A., des Forets G. P., Schilke P., Walmsley C. M., 1994, *A&A*, 289, 579
- Vidal T. H. G., Loison J.-C., Jaziri A. Y., Ruaud M., Gratier P., Wakelam V., 2017, *MNRAS*, 469, 435
- Vinodkumar M., Bhutadia H., Limbachiya C., Joshipura K. N., 2011, *Int. J. Mass Spectrom.*, 308, 35
- Wakelam V., Vastel C., Aikawa Y., Coutens A., Bottinelli S., Caux E., 2014, *MNRAS*, 445, 2854
- Wakelam V. et al., 2015, *ApJS*, 217, 20
- Wakelam V., Loison J. C., Mereau R., Ruaud M., 2017, *Mol. Astrophys.*, 6, 22
- Wakelam V., Dartois E., Chabot M., Spezzano S., Navarro-Almida D., Loison J.-C., Fuente A., 2021, *A&A*, 652, A63
- Weidenschilling S. J., Lewis J. S., 1973, *Icarus*, 20, 465
- Wong M. H., Mahaffy P. R., Atreya S. K., Niemann H. B., Owen T. C., 2004, *Icarus*, 171, 153
- Woon D. E., 2012, *Comput. Theor. Chem.*, 984, 108

This paper has been typeset from a  $\text{\TeX}/\text{\LaTeX}$  file prepared by the author.