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PHANGS-JWST: Data-processing Pipeline and First Full Public Data Release

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 ³⁷ Department of Physics, University of Surrey, Guildford GU2 7XH, UK ³⁸ Department of Astrophysical Sciences, Princeton University, 4 Ivy Lane, Princeton, NJ 08544, USA ³⁹ Jodrell Bank Centre for Astrophysics, Department of Physics and Astronomy, University of Manchester, Oxford Road, Manchester M13 9PL, UK Received 2024 January 25; revised 2024 April 19; accepted 2024 May 8; published 2024 July 9 Abstract The exquisite angular resolution and sensitivity of JWST are opening a new window for our understanding of the

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Universe. In nearby galaxies, JWST observations are revolutionizing our understanding of the first phases of star formation and the dusty interstellar medium. Nineteen local galaxies spanning a range of properties and

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⁴³ https://coolresearch.io/

morphologies across the star-forming main sequence have been observed as part of the PHANGS-JWST Cycle 1 Treasury program at spatial scales of \sim 5–50 pc. Here, we describe pjpipe, an image-processing pipeline developed for the PHANGS-JWST program that wraps around and extends the official JWST pipeline. We release this pipeline to the community as it contains a number of tools generally useful for JWST NIRCam and MIRI observations. Particularly for extended sources, pjpipe products provide significant improvements over mosaics from the MAST archive in terms of removing instrumental noise in NIRCam data, background flux matching, and calibration of relative and absolute astrometry. We show that slightly smoothing F2100W MIRI data to 0...9 (degrading the resolution by about 30%) reduces the noise by a factor of \approx 3. We also present the first public release (DR1.1.0) of the pjpipe processed eight-band 2–21 μ m imaging for all 19 galaxies in the PHANGS-JWST Cycle 2 Treasury program.

Unified Astronomy Thesaurus concepts: Star formation (1569); Spiral galaxies (1560); Surveys (1671); Astronomy data reduction (1861); Young star clusters (1833); Interstellar medium (847); Interstellar dust (836)

Materials only available in the online version of record: figure sets

1. Introduction

Near- and mid-infrared (N/MIR) wavelengths provide vital windows for studying star formation and the interstellar medium (ISM) in galaxies. Stars are born enshrouded deep within dusty molecular clouds and due to this, infrared observations are much more effective at detecting these earliest phases than optical (e.g., Allen et al. 2004; Robitaille et al. 2008). Emission from polycyclic aromatic hydrocarbons (PAHs) is also present at NIR and MIR wavelengths (e.g., Tielens 2008; Hensley & Draine 2023). This PAH emission gives crucial information on the size and charge distributions of the PAHs, and characterizing the evolution of PAHs through the ISM is central to understanding their life cycle (e.g., Engelbracht et al. 2005; Chastenet et al. 2019; Wolfire et al. 2022). However, localizing these young stars and the clouds in which they are born *requires* high spatial resolution observations, tens of parsecs or better. For previous MIR telescopes including Spitzer, such resolution could only be achieved in the nearest, $D \lesssim 5$ Mpc, galaxies. JWST is finally allowing us to study galaxies at these spatial resolutions at distances up to ~ 20 Mpc. This will enable us to determine both the galaxy-wide and local processes that affect how stars are born, exert feedback on the surrounding ISM, and then die, enriching the ISM for future stellar generations (the baryon cycle).

Studying this baryon cycle that surrounds star formation and feedback has been the goal of the Physics at High Angular resolution in Nearby GalaxieS (PHANGS; Leroy et al. 2021)⁴⁴ collaboration through a series of multiobservatory surveys. The latest of these is PHANGS-JWST (Lee et al. 2023), a Cycle 1 Treasury Program to observe 19 galaxies in eight JWST filters from 2 μ m to 21 μ m. These observations reveal a complex web of dust and gas as well as dust-embedded stellar clusters that are hidden from similar resolution Hubble Space Telescope (HST) observations. The goals of this program are wideranging, and a number of results have already been published on, for example, PAH properties (Chastenet et al. 2023a, 2023b; Dale et al. 2023; Egorov et al. 2023; Sandstrom et al. 2023), feedback-driven bubbles (Barnes et al. 2023; Watkins et al. 2023), and embedded stellar clusters (Hoyer et al. 2023; Rodríguez et al. 2023).⁴⁵

Most of the PHANGS-JWST first results used data from the first four galaxies observed (NGC 628, NGC 7793, NGC 1365, and IC 5332). Observations of all 19 targets are now complete, and studies are beginning to exploit the full sample (Belfiore et al. 2023; Pathak et al. 2024). This work presents a full public data release (DR1.1.0⁴⁶) for the PHANGS-JWST Cycle 1 program and the associated pipeline (pjpipe⁴⁷; Williams et al. 2024). It also details the lessons we have learned during data acquisition for this Treasury program. The survey and much of the technical information have already been described in Lee et al. (2023), and we present here the significantly upgraded view of the reduction approach, emphasizing the (occasionally dramatic) applied changes.

We have undertaken the task of building a pipeline to robustly deal with shortcomings we have discovered using the official JWST pipeline that are general to imaging observations, or in some cases more specific to observations where emission fills the field of view. These are 1/f noise in NIRCam data, astrometric alignment, background matching between tiles in a mosaic,⁴⁸ and calibrating flux levels in the absence of a true background. The pipeline also addresses the reality of widely varying resolution across the JWST bands, as well as the opportunity for a small smoothing operation to dramatically improve the signal to noise at the long-wavelength MIRI bands. pjpipe offers robust implementations to address these issues, and the main focus of this paper is to describe our algorithms, illustrate their application to PHANGS-JWST data, and explain how these link with the official JWST pipeline. We note that most of the issues addressed by pjpipe are not specific to PHANGS-JWST. Indeed, other programs have already made use of this pipeline to improve their data products (Peltonen et al. 2023; J. Chastenet 2024, in preparation).

The layout of this paper is as follows. We briefly introduce the PHANGS-JWST survey in Section 2, before detailing our data-processing steps in Section 3, noting differences between this and our early science efforts (Lee et al. 2023). We provide details of our quality assurance checks in Section 4,

⁴⁴ http://phangs.org/

⁴⁵ For the full set of these "PHANGS-JWST first results," we direct interested readers to <u>https://iopscience.iop.org/collections/2041-8205_PHANGS-JWST-First-Results.</u>

 ⁴⁶ Some issues were found with the astrometry in NIRCam observations of NGC 3627 and have been updated to a DR1.1.1. Other data remain unchanged.
 ⁴⁷ https://pjpipe.readthedocs.io/

⁴⁸ We will refer to these throughout as "mosaic tiles," as opposed to the individual single exposure "tiles" (in the case of the NIRCam short images, these are the individual detector images, rather than the combination of the multiple detectors). These mosaic tiles are formed by the combination of all (four) of the dithers in a single pointing (for the NIRCam short images, this combines the four individual detectors as well, for a total of 16 input exposures to a mosaic tile).

highlighting the remaining known issues in Section 4.3. We compare our early science mosaics and those obtained from MAST in Section 5. A description of the output files for the data release is given in Section 6. We summarize in Section 7 and present an atlas of our 19 galaxies in Figure 9. Our data (both combined mosaics and individual tiles for point-spread function (PSF) photometry) are available publicly at https:// archive.stsci.edu/hlsp/phangs/phangs-jwst.

2. Data Overview

PHANGS-JWST is a JWST Cycle 1 Treasury Program, building upon existing PHANGS efforts to map nearby (D < 20 Mpc) galaxies with Atacama Large Millimeter/ submillimeter Array (Leroy et al. 2021) CO observations, Multi Unit Spectroscopic Explorer (MUSE; Emsellem et al. 2022) optical integral field spectroscopy, and high-resolution HST (Lee et al. 2022) optical imaging. Specifically, this Cycle 1 Treasury (program 2107, PI: Lee) has observed the 19 mainsequence galaxies that form the PHANGS-MUSE subset of the full sample, which have been observed with JWST between 2022 June 6 and 2024 February 6.49

For full details of the survey, we refer the reader to Lee et al. (2023) and we briefly summarize details relevant to this paper here. The observations cover eight bands, spanning $2-21 \,\mu\text{m}$, using both NIRCam (F200W, F300M, F335M, F360M) and MIRI (F770W, F1000W, F1130W, F2100W). Our observations are primarily mosaics, with a maximum 2×2 tiling pattern. Because our targets are large in the sky, even these mosaics are often not sufficient to cover the full galaxy, though the coverage always captures most of the infrared emission and star formation activity from the galaxy (Figure 1), and we place the JWST footprint to overlap existing data. Details of the observations are given in Table 1. We use a minimum of four dithers per integration to fully sample the PSF, up to a maximum of 16 overlapping dithers in the case of F200W. To maximize the observing efficiency, we observed off positions for the MIRI data in parallel with our NIRCam science observations. A collage of our mosaics is shown in Figure 1. A rich tapestry of structure is clearly visible, from extended dust emission to pointlike, embedded stellar clusters.

3. Data Processing

In Lee et al. (2023), we presented our processing methodology for the four galaxies observed during the first month and a half of our program. Since then, we have seen significant improvements in a number of directions. First, the reference files (for this release, we use reference files jwst_1201.pmap⁵⁰) have improved as the in-flight performance of JWST and its instruments have been better understood. Second, improvements in the official JWST pipeline⁵¹ (Bushouse et al. 2023) have fixed a number of bugs and improved many of the underlying reduction steps (for our DR1.1.0 reprocessing, we use pipeline version 1.13.4). Finally, we have identified and understood several issues associated with observations of extended, bright sources that fill the field of view and developed robust strategies for many of these. This paper will focus on this last point, but we note our latest data

processing represents a large improvement both over our original data release (e.g., improved flux calibration, background matching, and astrometry), as well as the mosaics that are currently available in the MAST archive (see Section 5).

Our pipeline (pipipe) wraps around the official JWST pipeline (Bushouse et al. 2023), adding a number of bespoke steps not present in the official pipeline, which we show schematically in Figure 2. We make pipipe publicly available at https://github.com/phangsTeam/pjpipe, which also includes the configuration required to reproduce our reduction. We show the main output of pipipe, the mosaics, in Figure 9. For users who wish to use pipipe for their science, documentation and examples of configuration files are available at https://pjpipe.readthedocs.io/en/. As an opensource project, we also encourage users to contribute bug fixes and new features. By making our code fully available, we hope that some of the general techniques we have developed for our image processing will also feed back into improvements in the official JWST pipeline.⁵² A current list of known issues with the JWST pipeline is maintained at https://jwst-docs.stsci. edu/known-issues-with-jwst-data (some of which are dealt with by pjpipe).

3.1. Level One Processing

The Level One stage of the pipeline transforms uncalibrated images to count-rate images via a number of steps (e.g., flagging bad pixels, subtracting reference pixels, and fitting slopes to ramps). The details for each step are given at https://jwst-pipeline. readthedocs.io/en/latest/jwst/pipeline/calwebb_detector1.html.53 We essentially keep this step at the observatory-recommended defaults, although we process individual exposures (tiles) from each dither group in observation sequence rather than simultaneously. Processing the observations sequentially allows us to use the output files from the previous observation to identify persistence, i.e., faint residual images caused by previous exposure to bright sources appearing in subsequent integrations. Then data affected by persistence can be flagged in the subsequent persistence step. In our observations, we have not noticed persistence occurring from previous observations, but this could be a potential issue for observations immediately following, for example, observations of solar system planets. We also attempt to mitigate saturation by setting suppress one group to False in the ramp fitting step so that flux can be obtained for any pixels where the first group is not saturated.

3.2. Level Two Processing

Level Two processing applies calibrations and corrections to count-rate images on a per-exposure level. This stage performs background subtraction (for MIRI data), initial world coordinate assignment, flat-fielding, and photometric calibration. The details for each step here are given at https://jwst-pipeline.readthedocs. io/en/latest/jwst/pipeline/calwebb image2.html. As with Level One, we keep most of these at the observatory-recommended defaults.

⁴⁹ An up-to-date observing log can be found at https://www.stsci.edu/cgi-bin/ get-visit-status?id=2107&markupFormat=html&observatory=JWST&pi=1 ⁵⁰ https://jwst-crds.stsci.edu/

⁵¹ https://github.com/spacetelescope/jwst

 $[\]frac{52}{10}$ Many of the pjpipe algorithms have been presented at the "Improving JWST Data Workshop," and the materials and discussions from that workshop are available at https://outerspace.stsci.edu/display/JEA/Improving+JWST +Data+Products+Workshop.

 $^{^{53}}$ For the specific documentation for this version, see Bushouse et al. (2023).

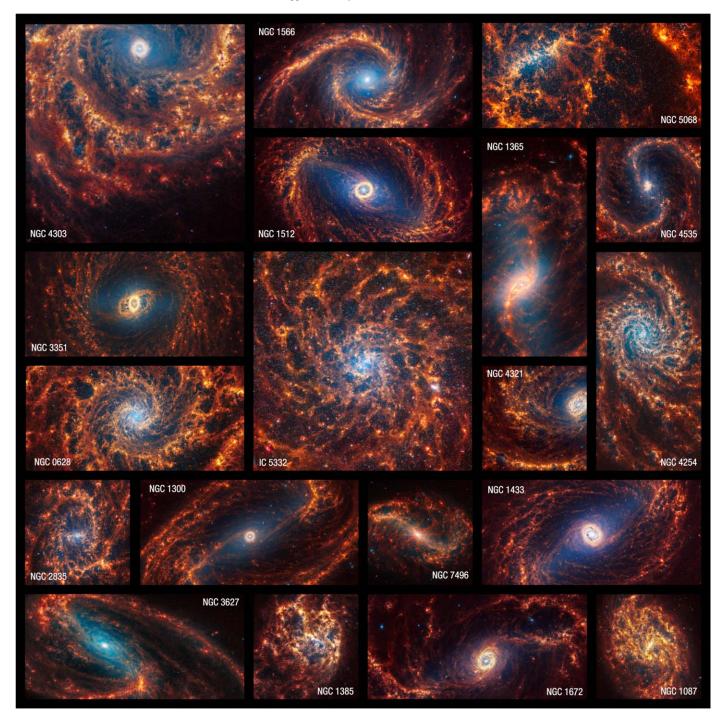


Figure 1. Color mosaics showing the JWST images for our full 19 galaxy sample developed to support a 2024 January press release, https://webbtelescope.org/ contents/news-releases/2024/news-2024-105. Red is a combination of the MIRI filters (F770W, F1000W, F1130W, and F2100W), green is a combination of NIRCam and MIRI (F360M and F770W), and blue is purely NIRCam (F300M and F335M). F200W imaging has also been obtained but is not included in the color composites. Note that in many cases, these mosaics are zoomed, and so do not show the full extent of coverage of the galaxy. For the full coverage of each observation, see Figure 9.

In this stage, we combine into a single master background frame the MIRI background observations in image space. A background estimate is calculated per pixel after applying sigma clipping and combining all dithers executed during the observation. In the case of PHANGS-JWST, the MIRI background observations are obtained in parallel during the NIRCam science observations and as a result, the dither used is also the one executed to obtain the NIRCam images. When constructing the background image, we use a relatively strict sigma clip of 1.5σ . Our background fields occasionally contain clear foreground stars or background galaxies and this strict tolerance limits the number of negative artifacts seen in our calibrated images, although in a few cases, unavoidable artifacts persist. These are extended background galaxies that the dither pattern fails to remove and appear as negative spots in our background images. Work is ongoing to mitigate this in this step (e.g., by masking sources and interpolating over missing data), but our current data occasionally contain these THE ASTROPHYSICAL JOURNAL SUPPLEMENT SERIES, 273:13 (20pp), 2024 July

 Table 1

 Overview of the Observational Setup for Our JWST Filters

Filter	Total Expo- sure Time ^a	Avg. Tile Overlap	5σ Point- source Sensitivity	1σ Surface Brightness Sensitivity
	(s)	1	 (μJy)	(MJy sr ⁻¹)
		NIRCar	n	
F200W	1202.5	16	0.039	0.071
F300M	386.5	4	0.084	0.052
F335M	386.5	4	0.092	0.050
F360M	429.5	8	0.107	0.054
		MIRI		
F770W	88.8	4	1.028	0.13
F1000W	122.1	4	1.253	0.13
F1130W	310.8	4	2.637	0.24
F2100W	321.9	4	5.152	0.27

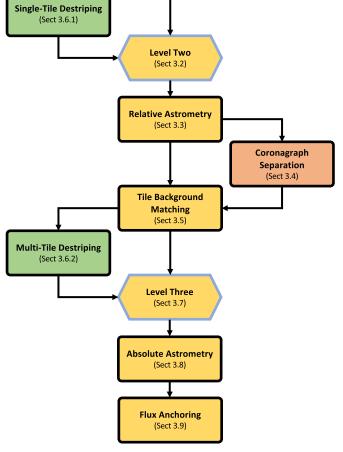
Note. We use a four-point dither pattern in all cases, so the average tile overlap is an indication of how many repeats we perform in the same area. Exposure times and overlaps are from Lee et al. (2023), but sensitivities have been updated using the same method presented in Lee et al. (2023), reflecting the new data processing.

^a This is the total exposure time per mosaic tile.

artifacts. For our larger mosaics, we expect the overall noise level to be slightly lower because we acquire background observations for every NIRCam position, meaning we generate a background from more observations.

This background estimate is then subtracted from the science image. Following this step, we obtain fully calibrated individual tiles (in units of MJy sr^{-1}), to which we perform a number of bespoke steps before mosaicking. We note that for this data release, the flux calibration does not take the decrease in sensitivity at long MIRI wavelengths into account. This will be updated in a future version of the data.

When constructing MIRI background images, the proximity of the background to the on-source image in *time* appears to be a critical factor. For PHANGS-JWST, the initial NIRCam observations, and hence the MIRI background images, failed for IC 5332, and the successful NIRCam science and MIRI background observations were only obtained months later. Another PHANGS-JWST target, NGC 7496 was observed immediately prior to IC 5332. Figure 3 compares the results subtracting backgrounds obtained close in position but not in time (i.e., using the later IC 5332 for the background) and close in time but not in position (i.e., using NGC 7496). We perform a sigma-clipped rms measurement on these three images, using a source mask created from the "close in time, not in position" backgrounds. We obtain $rms = [0.80, 0.46, 0.36] MJy sr^{-1}$ in these three cases. This shows that using the MIRI backgrounds obtained months after the MIRI science observations led to large-scale artifacts in the final mosaics, highlighting the temporal variability of the MIRI backgrounds. In the current release, we use the NGC 7496 "off" observations that ended 14 minutes before the IC 5332 science observations started to construct the background for IC 5332. For pointings where the background observations fail, we would therefore recommend reobserving both science and background given the clear time variability of the MIRI background.



All Instruments

Level One (Sect 3.1)

NIRCam-Specific

Figure 2. Schematic overview of our data reduction. While most stages are instrument-agnostic, there are some specific to either NIRCam or MIRI. Steps that are present in the official JWST pipeline are highlighted as blue-outlined hexagons.

3.3. Relative Astrometric Alignment

As noted in Rigby et al. (2023), the relative pointing accuracy after guide star acquisition is excellent, to within a few milliarcseconds. However, given that many of our observations are mosaics that use multiple tiles, we found that the absolute alignment between these guide stars is insufficient for our purposes. This could be due to a number of reasons the star tracker locking onto an incorrect target, that target being extended at the JWST resolution, or outstanding uncertainties in the location of the chips relative to the guide star tracker, for instance. We expect that as the guide star catalogs improve through the life cycle of JWST and as the instruments are better understood, these issues will become less common and severe, but for these observations, we require correction to get a good common astrometric alignment between the various mosaic tiles.

Using tweakreg, we align the mosaic tiles using a single shift for each set of dithers. tweakreg uses an iterative rejection algorithm to match sources between point-source catalogs, in order to align a set of input images. For each instrument and individual mosaic tile, the tile is observed in

MIRI-Specific

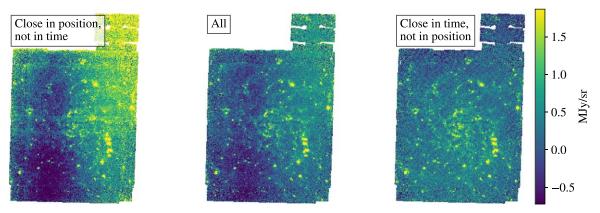


Figure 3. Illustration of how the temporal variability of MIRI backgrounds affects the final science images, using F2100W IC 5332 observations. Left: MIRI background using later NIRCam parallel observation, as the initial observation failed. This background is therefore taken close in position to IC 5332, but not close in time. Middle: combination of IC 5332 background (close in position) and NGC 7496 background (close in time) observations. Right: using only the NGC 7496 (close in time) background observation. Clearly, using backgrounds taken closer in time leads to a better, and overall flatter background.

filter sequence so that all observations of a specific tile with an individual instrument use the same guide star. Therefore, because of the excellent relative pointing accuracy, initial testing found that the astrometric correction is almost identical for the four NIRCam images, while the four MIRI filters also share a single correction that is different from the NIRCam correction. We therefore base our astrometric solutions on two bands, F300M for the NIRCam observations and F1000W for the MIRI observations. The F1000W band is chosen instead of the shorter-wavelength F770W band because it is not expected to be as heavily dominated by PAH emission, making stellar point sources within the galaxies clearer. It also has improved angular resolution compared to the F1130W band, allowing for better centroiding of point sources. For NIRCam, we work with F300M because we found the sheer density of point sources in F200W to be a detriment to solving for a good astrometric correction so using F200W as the reference band often led to a poor initial alignment.

For each visit (i.e., each mosaic tile), we calculate a single x, y shift and apply it to all individual exposures targeting that mosaic tile. The offsets derived during this step are typically small, on the order of 0."1. We experimented with including rotations but found this degraded the astrometric solution, without measuring any significant rotation. Note that we apply the same shift to the NIRCam short and long filters, although given some uncertainties in the exact chip position we will allow for a further shift during our Level Three processing (Section 3.7).

3.4. MIRI Coronagraph Separation

The MIRI data contain Lyot coronagraph observations. The coronagraph is separate from the primary imager but can still be used to obtain science-quality data because it shares the same optical path as the primary MIRI imager. Our testing showed that the background levels for the coronagraph observations were often significantly different compared to those on the primary image. To account for these different background levels, we separate the coronagraph from the primary imager observations by editing the data quality array in the Level Two images. We produce two images, one image flags the primary imager as "nonscience." This means we can perform the background matching (Section 3.5) separately for these two

parts of the CCD and include the coronagraph in our final mosaics after accounting for their different background levels.

3.5. Background Matching

Particularly for MIRI data, the background levels can vary substantially, even after subtracting the background calculated in the Level Two processing. This is mainly driven by instrumental backgrounds (primarily the telescope thermal background at longer wavelengths, and scattered light at shorter wavelengths), which we observe to vary at the $\sim 0.1 \text{ MJy sr}^{-1}$ level between dithers even within a single mosaic tile. These must be corrected for in order to provide consistent level matching between individual observations, and tiles in a mosaic.

The JWST pipeline uses the skymatch algorithm to provide consistent levels between tiles. It does this by calculating a single sigma-clipped average within overlap regions, and then minimizing the difference among the background values for the various overlapping observations (i.e., minimizing the difference of the medians). However, we found no combination of parameters that allowed this algorithm to produce good results across our full data set. Instead, "jumps" between adjacent tiles always remained visible. This likely reflects the fact that our overlap regions are full of bright extended emission, with no real "background" to sigma-clip to.

We therefore take a different approach. For each pixel in the overlap region between a pair of images, including those that contain signals from the source, we calculate the difference between values at that pixel in the different images. Then we take the median of the pixel-by-pixel differences to calculate the average offset between these two images. To do this, we first reproject all images to a common pixel grid and then calculate the median difference of all overlapping pixels. When calculating these median differences between pairs of tiles, we use sigma clipping to reject outlying differences. We then adjust the background level of the images being considered, in order to minimize the differences among all image pairs being considered. This is done using least-squares minimization, weighting each image pair by the rms spread of the difference histograms for each tile overlap. To avoid including noisy and untrustworthy overlaps between images, we do not include any image pairs with a total pixel overlap of less than 0.2% of the average number of pixels in an overlap region. We also do not

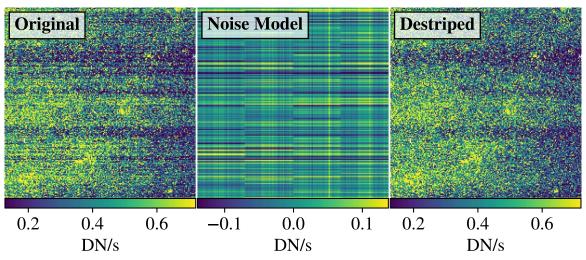


Figure 4. Single-tile destriped model for one F200W observation of NGC 628. The left panel shows the original data, the middle panel is our stripe model, and the right panel shows the destriped data. Although effective at removing much of the striping, some low-level, large-scale stripes remain.

include image pairs that do not overlap by at least 25% along at least one axis of the image or image pairs that have rms values that are more than 2σ greater than the average rms value of all the other image pair differences (i.e., very noisy difference histograms, which often occur when there are few overlapping pixels). These cuts typically remove the small diagonal overlaps that occur in the centers of the 2×2 mosaics or similar (or the centers of the four NIRCam short chips). These choices mirror those used by Montage (Jacob et al. 2010) for their background matching step.

In practice, we perform this background matching in two stages. First, we match the level between dithers within a visit group. Then we stack these and match the tiles in the overall mosaic. We find that this process produces excellent background matching both for NIRCam and MIRI tiles, significantly better than what we could achieve with the skymatch step in the JWST pipeline.

3.6. NIRCam Destriping

NIRCam data are well documented to suffer from 1/f (pink) noise (Schlawin et al. 2020). For example, this is clearly visible in the left panel of Figure 4 as horizontal striping is pervasive across the entire image. This noise is due to the way NIRCam data are read out and has power at a number of spatial scales. This 1/f noise is particularly challenging to remove from our data because diffuse emission also pervades our observations with power on overlapping spatial scales.

Reference pixels are present as a four-pixel border around the NIRCam chips, and the Level One processing provides a firstorder correction to remove this noise, but the striping is still persistent even following this refpix step. Fortunately, algorithms have been developed for mitigating this (e.g., remstriping; Bagley et al. 2023). Unfortunately, these have mostly been designed for blank fields and high redshift observations.

To mitigate this noise, we adopt a two-stage approach. First, we employ a single-tile destriping technique that uses a single exposure. Second, we use temporal information from overlapping exposures to remove remaining large-scale stripes that the single-tile destriping does not correct. Because the 1/f noise is instrumental we apply this destriping to the data *before* our Level Two processing and flat fields are applied.

Note that for our initial reduction of the first four galaxies (presented in Lee et al. (2023), and associated papers), we used a robust principal component analysis (PCA) method. Further testing has revealed that this method is extremely sensitive to source masking, and using a median filter yielded more robust results. In particular, our testing with the PCA method often showed imprints of the source in the noise model, indicating that we are removing real emission. The same is not true in the case of the median filter. We do believe that PCA may ultimately represent the optimal strategy when there is limited extended emission in a tile, but our data contain too much emission for the PCA to work successfully.

3.6.1. Single-tile Destriping

First, following Bagley et al. (2023), we measure 1/f-noiseinduced stripes on flat-fielded images but subtract the stripes before the flats are applied. We start by filtering our data using a Butterworth (1930) filter, which removes large-scale structures (both emission and large-scale 1/f noise). We opt to use this filter because its smooth frequency response suppresses the Fourier "ringing" that is often seen around discontinuous sources in many high-pass filters.

Having filtered the data, we then removed the 1/f noise using a median filter at a number of scales. The 1/f noise is present in both the x and y directions, and we subtract striping from these independently. To avoid oversubtraction, we create a source mask and exclude any remaining bright sources after the filtering. Then, we calculate and subtract a row median in the y direction (collapsing along columns), before calculating and subtracting a median in the x direction (i.e., collapsing along the rows). In this step, we split the data per amplifier because the noise properties are different in each amplifier. For rows or columns where more than 80% of the data are masked, we fall back to the full row or column median, to avoid spurious values.

We show an example of this single-tile destriping in Figure 4. We find the amplitude of the noise removed in this step to be typically around the 0.05 DN s⁻¹ level. This step effectively removes the 1/f noise on small scales while retaining the large-scale structure present in the images. However, large-scale "ripples" remain, which we mitigate with our second destriping technique (Section 3.6.2).

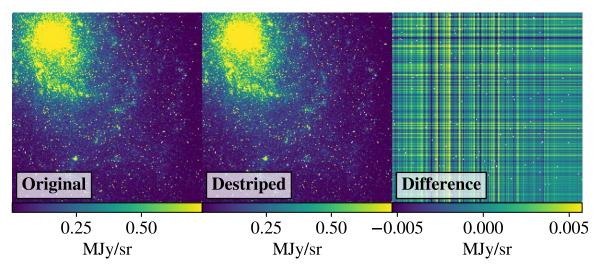


Figure 5. Applying the single-tile destriping to Spitzer observations of M33. The left and middle panels show the original and destriped data, respectively, and are locked at the same color scale. The right panel shows the difference image, on the order of $0.001 \text{ MJy sr}^{-1}$, essentially below the noise level of the original image. The white pixels in the difference image are due to pixels with no data; these are not present in the JWST noise models but are present here due to differences in the specifics of the data between JWST and Spitzer.

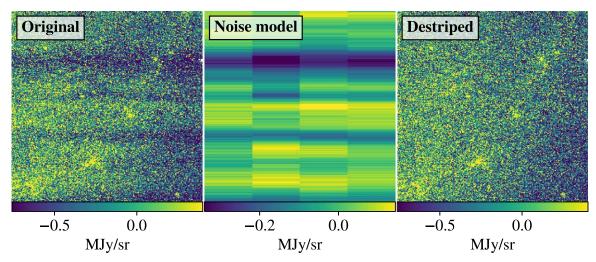


Figure 6. Multitile destriping for one F200W observation of NGC 628. Left: uncorrected data. Middle: noise model. Right: corrected data. This step successfully removes the majority of the remaining large-scale striping in the data, leaving a mostly flat background.

We can see in the middle panel of Figure 4 that there is no clear imprint of the galaxy in the noise model, indicating that this method does not remove flux in the image. To test this further, and ensure that this destriping procedure does not alter the true flux levels to a significant extent, we take Spitzer observations of M33 at a similar physical resolution. These data do not suffer from 1/f noise, and so this test is to ensure that our destriping algorithm essentially does not subtract anything in the absence of 1/f noise. We apply the destriping to this image after ensuring it looks like a calibrated JWST image to the pipeline, with the exception of the lack of 1/f noise. The difference is shown in Figure 5, with the difference (i.e., the "stripes" measured here) being on the order of 0.001 MJy sr⁻¹. We are therefore confident that our destriping does not add or remove flux to any great degree.

3.6.2. Multitile Destriping

Although our previous step effectively removes small-scale 1/f noise, large-scale noise is still present because it survives the Butterworth filtering. This is particularly visible in the short

NIRCam bands but is prevalent throughout the NIRCam observations. The stripes are typically at around the 1σ level compared to the error map but can be larger (particularly for the NIRCam short wavelengths). We take advantage of the fact that this noise is not correlated between different exposures to remove striping from each individual exposure while leaving the real galactic structure intact.

In detail, for each exposure, we create a sigma-clipped median image of all overlapping exposures. For PHANGS-JWST, there are at least four overlapping exposures for all bands and up to 12 for F200W. When creating this image, we make sure that the overall flux level matches between tiles by performing this after our background matching step (Section 3.5). For each individual exposure, we then create a difference image by comparing that exposure to this median image. Because it is present in both the median and the image, real structure will not be present in this difference image but it will include any large-scale ripples. Using this difference image to effectively remove any large-scale stripes still left in the data.

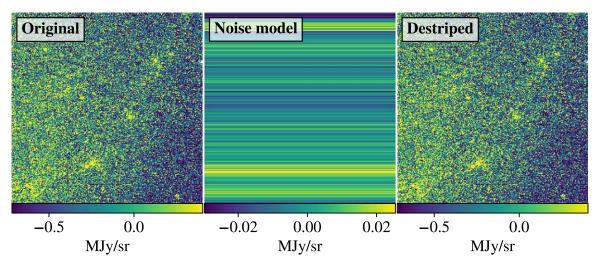


Figure 7. Multitile destriping for one F200W observation of NGC 628, having smoothed the stacked image perpendicular to the stripe axis. Left: uncorrected data. Middle: noise model. Right: corrected data. This step successfully removes the remaining striping, so having completed this step our backgrounds are flattened.

An example of this procedure is shown in Figure 6. Typically, the amplitude of the noise removed here is at around the 0.05 MJy sr^{-1} level and is worst for the F200W band. We also attempted to perform destriping using only this method, but found that the small-scale stripes remained; hence, both the single- and multitile destriping are necessary to remove the 1/f noise on all scales.

Particularly in the short NIRCam observations, there remains some very low-level striping in the data following this first pass using overlapping tiles. This reflects the fact that while the 1/fnoise is essentially random, given its prevalence it can interfere constructively when creating an average image. We therefore carry out a second pass where we create a smoothed, stacked image, performing a large-scale (a tenth the size of the full array) median filter along the axis perpendicular to the direction of the stripes in the stacked image. This creates a new median image as in the previous step but with any remaining largescale noise filtered out. We then subtract the large-scale median filtered image from the individual tiles and then take a row-byrow median to construct a last noise estimate. We calculate the median across the full tile rather than per amplifier, as the median filter redistributes flux, potentially leading to a bias along rows completely full of emission. We show an example of this final step in Figure 7 and note that these last remaining stripes have significantly smaller amplitude than those subtracted in the previous destriping steps, at around the 0.01 MJy sr^{-1} level. While we have not performed detailed tests on how much our destriping lowers the noise in our imaging, initial tests in blank fields appear to lower the rms noise level by around 10%, consistent with other algorithms (e.g., Bagley et al. 2023). We leave a full exploration of optimal destriping methods and noise improvements to future work. Having carried out these destriping steps, the images appear stripe-free.

3.7. Level Three Processing

Level Three processing combines individually calibrated exposures into a final mosaic via drizzling (Fruchter & Hook 2002), along with astrometric alignment, background matching, and cosmic ray rejection. The details for each step here are given at https://jwst-pipeline.readthedocs.io/en/latest/jwst/pipeline/calwebb_image3.html.

We disable astrometric alignment for the MIRI images, having already performed this step manually (Section 3.3); we find that keeping this on tends to degrade the astrometric solutions, given the small number of point sources in each MIRI exposure. We do, however, leave this on with a strict tolerance for the NIRCam images. The derived shifts are significantly less than one pixel, highlighting that our initial alignment step is performing well.

We also disable level matching for all bands, as we have implemented our bespoke solution. We do, however, subtract a constant value as the lowest sigma-clipped value in each group of tiles. This would get us to a "true" background level in the case of blank fields in our observations, although as we will see in Section 3.9, this is not generally the case.

We perform outlier flagging with default parameters, although to speed this step up we keep all images in memory (in_memory=True; if this setting is False then this step takes on the order of hours, rather than minutes). F200W is the most memory-intensive band here and can use up to 250 GB of RAM. Following Level Three processing, we have final mosaics in all of our eight bands, which are aligned north-up. These mosaics are produced at the native pixel scale of the data (i.e., 0."031 for NIRCam short, 0."063 for NIRCam long, 0."11 for MIRI). We use the default pixfrac of 1.0 and note that as we perform another round of astrometric alignment following the mosaicking (Section 3.8), the images are not aligned to the same pixel grid.

3.8. Absolute Astrometric Alignment

Following Level Three processing, our mosaics have good internal alignment but are not astrometrically locked to an external reference frame. Given the relatively small fields of view for our PHANGS-JWST mosaics, they often include few Gaia sources. The Gaia sources that are in the field are often extended H II regions or stellar clusters within the galaxy, and not suitable point sources on which to centroid a final astrometric solution. Therefore, we rely on an external catalog to supply our absolute astrometric reference.

For PHANGS-JWST, these are luminous red stellar sources, a mix of AGB stars and red supergiants, that are bright in our JWST bands. These stars are identified from DOLPHOT

(Dolphin 2000, 2016) catalogs produced from the HST imaging, and their astrometric calibration descends from those data (Lee et al. 2022). To select these sources, we use a combination of detection quality and color-magnitude cuts. To obtain bright, uncrowded, well-PSF-fitted pointlike sources, we use the following DOLPHOT parameter criteria: $S/N_{F814W} \ge 10$, $crowd_{F814W} \leq 0.01$, PhotQualFlag_{F814W} = 0, OBJTYPE = 1, $-0.025 < \text{sharp}_{F814W} \le 0.01$ and $-0.15 < \text{round}_{F814W} \le 0.4$. The crowding parameter describes how much brighter a star would be (in magnitudes) if the flux from neighboring stars was not removed. By definition, the value is zero for an isolated star. The object type is a classification for the morphology of the sources, where 1 corresponds to a well-detected star. The quality flag provides information on artifacts, such as bad or saturated pixels, and whether the extent of the source is not fully captured because it is at the edge of the chip. A flag of zero indicates no such issues. Sharpness is a morphological parameter that is computed relative to a point source: if a source is perfectly fit by the PSF model, then the value is zero. Positive values indicate sources are more centrally peaked than the PSF (e.g., a cosmic ray), while negative values indicate those that are more extended. Finally, roundness is a symmetry parameter, with a value of zero for perfectly circular sources. For further discussion, we refer the reader to Thilker et al. (2022) and the DOLPHOT (Dolphin 2000, 2016) documentation⁵⁴ sharpness and roundness limits are slightly adjusted if necessary per galaxy. For the color-magnitude cuts, we require F814W in the range of 3-5 mag brighter than the foreground extinction corrected tip of the red giant branch and take only sources with *V-I* color in the range of $2 \leq F555W - F814W < 4$ mag.

tweakreg is used to perform the alignment, but because our astrometric solutions can often be different from the internal astrometry by an arcsecond or more, we take a twostep approach. First, we use a relatively wide search area with a loose tolerance to find matches, using a simple x/y shift (i.e., shift in tweakreg). Then we use this initial guess with significantly stricter tolerances to improve the initial solution. In this second step, we also allow for rotation (rshift). We generally find small rotations, although occasionally up to ~0°.01, leading to subpixel shifts across the full mosaics.

We perform this procedure separately for each NIRCam band. For MIRI, there is a lack of true point sources at many of the longer wavelength bands (particularly $21 \ \mu$ m). For this, we calculate a solution from the F1000W band (as the highest resolution MIRI band with many clear point sources), and apply this solution to all the MIRI bands as they have consistent internal astrometry. At this stage, we also backpropagate these astrometric solutions to the individual tiles, as these are used as the inputs to the stellar photometry package DOLPHOT (Dolphin 2000, 2016).

3.9. Anchoring the Background Level

After the tile background matching step above, our mosaic tiles have a consistent internal background (Section 3.5), with little or no jump in intensity from tile to tile within a mosaic. However, there is still uncertainty in the overall intensity of the background level. That is, at this stage the intensity of empty sky near the galaxy is not necessarily zero, and our empirical tests suggest background-level variations by several times ~ 0.1 MJy sr⁻¹ from galaxy to galaxy and band to band to be

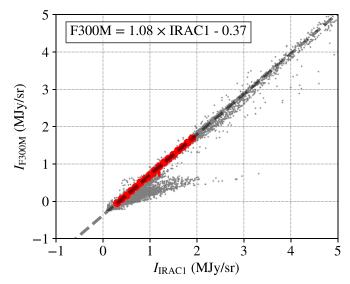


Figure 8. External anchoring for NGC 628 at F300M to IRAC1 (both convolved to 4'' resolution). The gray dots indicate individual pixels, and the red dots are binned medians for the region over which we fit to anchor. The best-fit line is shown in dashed gray, with the fit parameters given in the top left.

common, even up to $\sim \pm 1 \text{ MJy sr}^{-1}$ in some cases. Given the small field of view of many PHANGS-JWST mosaics compared to the angular extent of the galaxy, it is often not trivial to identify a region of true galaxy-free sky to establish the background level.

In order to establish a consistent, correct background level across all bands we adopt a two-step "anchoring" procedure that builds on the approach described in Leroy et al. (2023). First, we solve for an internally consistent background level across all NIRCam bands and all MIRI bands. Then, we set the overall background level by comparing one MIRI and one NIRCam band to wide-area Spitzer or Wide-field Infrared Survey Explorer (WISE) images as these include plenty of galaxy-free sky and therefore have well-established sky background levels.

All stages of this "anchoring" procedure assume that near-IR or mid-IR intensities of the same region at different bands exhibit a linear relation, on average, at low to moderate intensities. We solve for the slope and intercept of this relation. The slope represents the average band ratio, which may reflect underlying dust or stellar physics but does not matter to this exercise. The intercept represents the intensity offset between the current sky background level in the two images. Empirical testing (e.g., Figure 8 and see Appendix B of Leroy et al. 2023) shows that tight linear relations hold among every near-IR and mid-IR band pair that we have examined. That is, the bands appear to approach the empty sky in approximate lock-step.

We apply this procedure to solve for four sets of offsets for each source:

- 1. The sky background offset between the NIRCam F300M image and a wide-field reference Spitzer IRAC1 3.6 μ m image (when available) or a wide-field WISE1 3.4 μ m image (when IRAC1 data are not available).
- 2. Internal sky background offsets among NIRCam images.
- 3. The sky background offset between the MIRI F770W image and a wide-field reference IRAC4 8 μ m image (when available) or a wide-field WISE3 12 μ m image (when IRAC4 data are not available).

^{4.} Internal sky background offsets among the MIRI images.

Then we use these to place the NIRCam and MIRI data on a common sky background level that is "anchored" to share the sky background level of the wide-field reference image.

An example of the procedure is shown in Figure 8. In detail, we conduct the fit by constructing medians within bins defined to span a range of intensities close to zero (which we adjust somewhat from source to source and band to band to capture the behavior near the background level). The median binning approach downweights band ratio variations in individual regions; e.g., see the outlying spur at low intensities in Figure 8. The focus on low intensities avoids, e.g., curvature in the band-to-band intensity relations due to the influence of changing stellar and dust colors in HII regions or galaxy centers. Note that in this approach the solved slope is a nuisance term. Though the band ratios are of physical interest elsewhere, here we only care that we can solve for a linear relation and use this to adjust the background level of the image to a consistent, externally validated scale.

The procedure appears to work empirically. The bands show tight, clear linear relations with one another, almost always similar to Figure 8. The ratio images among bands in low-intensity parts of galaxies, which tend to be sensitive to background-level discrepancies, also appear mostly free of background-related artifacts. Most directly, in a few cases we do have some regions that contain mostly empty sky (which we will also use to study the noise properties as a function of resolution in Section 4.1). After this anchoring procedure, the intensity histograms appear centered at intensities within a few times ~0.01 MJy sr⁻¹ of zero. We caution that the anchoring to a true, correct sky

We caution that the anchoring to a true, correct sky background will only be as good as the reference image. Our IRAC1 images come mostly from S⁴G (Sheth et al. 2010), our IRAC4 images from SINGS (Kennicutt et al. 2003) or the Local Volume Legacy (Dale et al. 2009), and our WISE1 and WISE3 images from z0MGS (Leroy et al. 2019). In each case, we conducted an additional round of by-hand background fitting to the reference image, but we do note that this remains an area where additional work (as well as improvements to the algorithm) may further improve the images at the few times 0.01 MJy sr⁻¹ level.

3.10. PSF Matching

Much of the PHANGS science relies on multiwavelength photometry, for which performing an analysis at matched resolution is important. We therefore produce images convolved to a common resolution as a standard part of the pipeline. We use the method outlined in Aniano et al. (2011) to produce convolution kernels that can be convolved with the data to produce a common, Gaussian-resolution image. Our code, called jwst_kernels, is available at https://github. com/francbelf/jwst_kernels. We assume the JWST PSFs computed using the WebbPSF simulation tool (Perrin et al. 2014) version 1.2.1 taking the detector effects into account.

Based on the W_{-} metric from Aniano et al. (2011; their Equation (21), which quantifies the level of negative values in the kernel, and is related to the necessity of moving power from large to smaller scales), the minimum "very safe" common resolution across our JWST data set is determined by the F2100W image and is 0."86. In testing, we find that this resolution occasionally produces ringing around bright point sources. This is likely due to the highly nonaxisymmetric and non-Gaussian nature of the JWST PSFs. We therefore produce beam-matched images at 0.185 (close to the "very safe" value), 0.93 and 1, with the latter two being very conservative choices, which may nonetheless be preferable for comparison with low-resolution data.

Users interested in mitigating the impact of the JWST PSF may wish to convolve the data with a suitable kernel to obtain the smallest "safe" Gaussian PSF for a specific band, different from our default F2100W. We report in Table 2 the Gaussian PSFs corresponding to "very safe," "safe," and "aggressive" kernels according to Aniano et al. (2011). Suitable kernels for all these bands can be generated using our jwst kernels code with the make_jwst_kernel_to_Gauss function. For matched PSF work, we recommend experimenting with the kernels presented here; if resolution is particularly important, then opting for an "aggressive" kernel may be the optimal choice. However, these more aggressive convolutions are more prone to artifacts, especially around very bright sources. We therefore strongly suggest using visual inspection to ensure any convolved images are artifactfree to the level the science requires.

We also carry the convolution through to the error maps, by following standard error propagation (i.e., convolving the square of the error map with the square of the kernel, and taking the square root of each image). This assumes the noise is uncorrelated at the pixel level. This appears to be true to first order based on our tests of the scale dependence of the noise (Section 4.1) but is unlikely to hold in all cases given that the PSF is often significantly oversampled. As the community continues to better understand the noise properties, including the contribution of sky noise and interpixel correlations of noise within the JWST instruments, we expect these error maps to become more sophisticated.

4. Data Quality

In this section, we detail the quality of the data, with particular attention to how the noise properties scale as a function of resolution. We also describe our quality assurance process and note outstanding issues with the data. For an example of the quality of our data, we show the mosaics for NGC 628 in Figure 9, with equivalent plots for the other galaxies in the figure set.

4.1. Noise Scaling with Resolution

Of the bands used in this survey, the F2100W band has the largest native PSF (FWHM of 0."674), much larger than the 0."11 size of the pixels. The F2100W band also has the highest noise level of any band in the PHANGS-JWST data set. As a result, any signal-to-noise cuts performed in an analysis that includes F2100W data will be dominated by the large number of pixels cut across these maps. The F2100W noise is instrumental in nature and correlated at the pixel level. The pixel size is much smaller than the PSF (the F2100W PSF is oversampled by ~10 × in area relative to Nyquist sampling). Therefore the noise in these images can be dramatically decreased by smoothing the images to a slightly worse resolution. Smaller gains can also be realized for the other bands.

Figure 10 demonstrates this effect, showing the empirically measured noise in several regions with no source emission in the MIRI bands. The figure shows the empirically estimated

FWHM for Very S	Safe, Safe, and Aggressiv	ve Convolution Kernels to a G	aussian PSF for Our JWST Filter	s, Calculated Using the Metr	ics of Aniano et al. (2011)
Instrument	Filter	Source PSF	Very Safe Gauss.	Safe Gauss.	Aggressive Gauss.

Table 2

liisuument	The	FWHM (arcsec)	FWHM (arcsec)	FWHM (arcsec)	FWHM (arcsec)
NIRCam	F200W	0.064	0.089	0.077	0.069
NIRCam	F300M	0.104	0.141	0.125	0.115
NIRCam	F335M	0.115	0.161	0.139	0.122
NIRCam	F360M	0.122	0.170	0.147	0.130
MIRI	F770W	0.238	0.320	0.286	0.258
MIRI	F1000W	0.314	0.415	0.375	0.339
MIRI	F1130W	0.360	0.478	0.436	0.400
MIRI	F2100W	0.651	0.863	0.777	0.699

4.2. Noise Estimates Compared to Error Maps

Using the background regions from the previous subsection, we compare the error measured at the native resolution to the pipeline-calculated error map. We take the standard deviation of the pixels within the background region in the science image for the four MIRI bands (we do not use NIRCam as due to the source density, we are unable to find any background regions within our images, although given the good agreement between our calculated sensitivities and the exposure time calculator estimates, these error maps are likely also reasonable), and compare these to the mean of the pixels within the same region of the error map image. We note that this error map does not account for pixel-correlated noise. We have not included this effect at this point but will revisit this at some point in the future. For an example of how this can be accounted for, we refer readers to Yang et al. (2023). This comparison is shown in Figure 11—there is a good agreement here, with the error map value typically being higher by around 5%-10%. We therefore deem these error maps suitable for providing estimates of the noise for observations where no background regions are present.

4.3. Quality Assurance and Outstanding Issues

Having created our final mosaics, we perform a round of internal quality assurance (QA). We have examined each individual image for issues such as poor astrometric alignment (which manifests as "blurring" between tiles), stepping between tiles which indicates poor level matching, and other defects in the data quality (e.g., remaining 1/f noise). We find that each mosaic is free of these issues.

We then performed a round of interband quality checks, producing three-color images (of F200W, F300M, and

F1000W) using SAOImageDS9 (Joye & Mandel 2003) to assess the NIRCam short, NIRCam long, and MIRI overall alignment. Here, we are essentially checking that the astrometry between these mosaics is good, as large misalignments will lead to clear blurring between the color channels. We find that the NIRCam-to-NIRCam agreement is very good (with any misalignment significantly less than a NIRCam pixel), as is the NIRCam-to-MIRI agreement in most cases (with any misalignment less than a MIRI pixel), although there are some exceptions. The mosaics themselves generally look excellent and are ready for a wide range of science cases, but there are some outstanding issues that remain unresolved in this data release. We expect that these may improve as the characterization of the instrumental features of JWST improves, and substantial improvements will be followed by new version releases. In particular, our QA has noted the following issues:

4.3.1. PSF Spikes and Saturation in Bright Galaxy Centers

A few bright sources lead to large PSF spikes that can cover a significant portion of the observations. These spikes primarily affect MIRI images and they mostly originate from bright point sources, both star clusters and in active galactic nuclei, in galaxy centers. We have experimented with using WebbPSF to subtract these and so recover the affected extended emission. However, we find that WebbPSF models matched to the observed sources tend to oversubtract the spikes, particularly near the bright sources themselves. This could be due to the "brighter-fatter" effect (Argyriou et al. 2023) whereby charge can leak into neighboring pixels, leading to an increase in the PSF size of 10%-25% in bright sources. Additionally, limitations in the current theoretical model for the PSF are likely prevalent in the outskirts of the PSF wings. Whatever the root cause, we currently do not have a robust method for removing these PSF spikes. Lacking the ability to correct for these, many of our analyses identify and mask regions affected by PSF spikes before analyzing the data.

Additionally, in the brightest galaxy centers, a small number of pixels are saturated by the first readout and so are unrecoverable. It may be possible to estimate the flux here via a curve-of-growth analysis (e.g., Liu et al. 2023) or PSF fitting. However, we have not yet developed a robust, automated solution for this and so in some cases, a number of central pixels do not have measured flux values. This tends to be more prevalent in the MIRI bands, although some NIRCam wavelengths are also affected. These bright centers also occasionally cause negative horizontal banding. As of now, we have not corrected for this and note that this may

⁵⁵ Assuming the F2100W beam is Gaussian, this is smoothing with a Gaussian of FWHM 0.62″.

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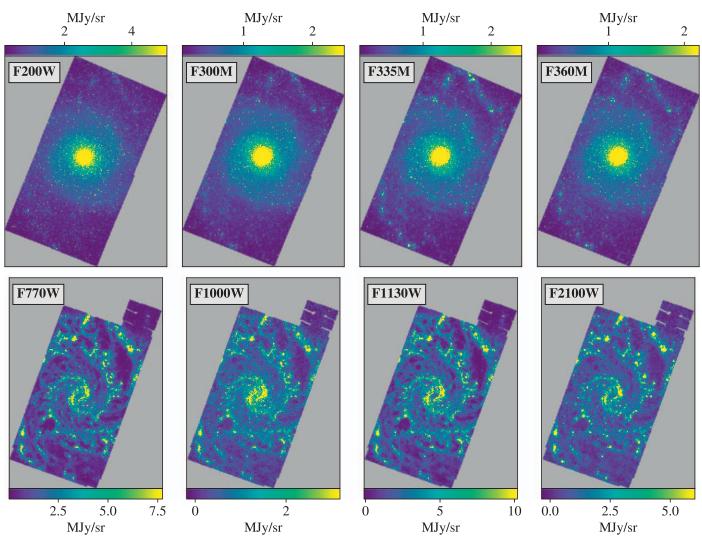


Figure 9. Final mosaics for each band for NGC 628. NIRCam filters are shown in the top row, and MIRI on the bottom. Each image is on a linear stretch between the 2nd and 98th percentiles of the pixel distribution, to highlight the fainter structure in the images. The complete figure set showing per-filter mosaics for our full galaxy sample is available online.

(The complete figure set (19 images) is available in the online article.)

affect the level matching between tiles, as the centers typically fall in the overlap regions of our mosaics.

4.3.2. Astrometry Issues

In Figures 12 and 13, we show the quality of the alignment using sources confidently matched by eye between HST, NIRCam, and MIRI observations. While we can see that the NIRCam-to-HST alignment is generally very good-the median offset across all NIRCam bands with respect to HST is 0."012 (around 0.5 NIRCam short or 0.2 NIRCam long pixels) with a median rms scatter of 0."006. However, the agreement is less good between MIRI and NIRCam-the median offset across all MIRI bands with respect to F300M is 0."05 (around 0.5 MIRI pixels) with a median rms scatter of 0."03. This is mainly driven by the much lower density of point sources in the MIRI data, and the source extraction not being optimized for the JWST PSF shape. As well as this, in some cases, a point source at the MIRI resolution (particularly F2100W) can be an extended source at the higher NIRCam resolution (often a background galaxy), making this comparison more complicated. The shifts are typically less than a pixel

but can be more severe in some cases; our worst case is NGC3351, where the shift between NIRCam and MIRI is around one MIRI pixel. In a few cases, there are multiple peaks in the distribution of residual astrometric offsets, indicating some persistent unresolved issues with the tile-to-tile matching.

We expect the astrometry to improve in future releases as we improve our alignment algorithms—fitting bespoke PSF models to sources may yield better centroiding⁵⁶—and as the overall astrometric properties of the JWST instruments and guide star catalog improve. There is also ongoing work to optimize the matching algorithm, which may improve the astrometric solution (e.g., JHAT; Rest et al. 2023). Our tests here hint that a more hands-on approach may be required; however, for this release, we prefer a repeatable, automated approach. We intend to revisit this topic in the future. We deem our current absolute astrometry to be acceptable for most science cases in our data set, although for high-precision, multiband aperture photometry some manual correction may still be required.

⁵⁶ This is currently in development by the JWST-FEAST team, with their algorithm at http://github.com/Vb2341/One-Pass-Fitting.

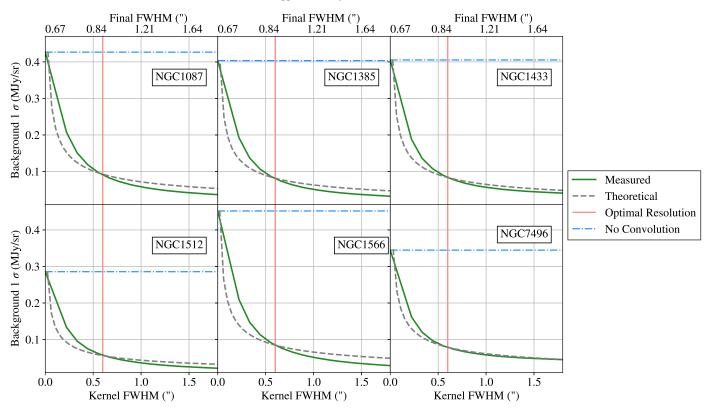


Figure 10. Variations in the standard deviation of an empty sky region after convolving the F2100W maps with a 2D Gaussian with the FWHM value displayed on the lower x-axis. The approximate final resolution is listed on the upper x-axis. Green lines show the measured standard deviations. Vertical red lines represent the Gaussian kernel that will approximately produce our chosen 0."90 final resolution, a kernel FWHM of 0."60. Horizontal blue dash-dot lines represent the background 1σ of the regions at the native resolution, before any convolution. The gray dashed line represents the theoretical prediction for uncorrelated noise, normalized to the measured curve at our chosen final resolution. Each panel represents a different map that includes an empty sky region in all of the MIRI bands.

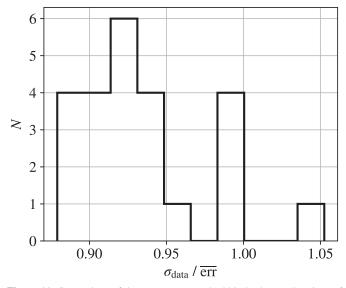


Figure 11. Comparison of the scatter measured within background regions of MIRI data to the mean of the pipeline-generated error map value within the same region (one region per galaxy).

5. Existing Mosaic Comparison

In this section, we show the level of changes in our data processing relative to both our early science data release and the mosaics obtainable through the MAST service. Overall, our latest version (DR1.1.0) is a significant improvement over both of these products and so we suggest users prefer our latest versions over these other available versions.

5.1. Comparison to Early Science Release

Our first point of comparison is to the data we released as part of our early science release. This comprised four galaxies (IC 5332, NGC 628, NGC 1365, and NGC 7496) that were observed in the first few months of JWST science operations. A comparison between our latest version and the early science NGC 628 images is shown in Figure 14, and analogous figures for the other three galaxies are shown in the Figure 14 figure set. The differences are stark in the short NIRCam band, which are caused both by our differences in destriping (PCA tended to subtract too much flux around bright sources that extended over much of the field of view) as well as the background matching (steps between tiles are clearly visible). These issues are significantly improved in our latest version. For the longwavelength NIRCam bands, the differences are minor, a few tenths of a MJy sr⁻¹ at most.

The MIRI images also see significant improvements, both in astrometry and background matching. The astrometry differences are most clear at F2100W, but indeed are present across the four MIRI filters. The background is also much more consistent between tiles in our latest version, highlighted as steps between mosaic tiles (most clearly seen in F1130W). Calibration updates have also substantially changed since our early science, by as much as 20% in the F770W, though less in other bands.⁵⁷

We will be revisiting many of our early science results with the full Cycle 1 sample in due course, as the increased number

⁵⁷ https://www.stsci.edu/contents/news/jwst/2023/updates-to-the-miriimager-flux-calibration-reference-files

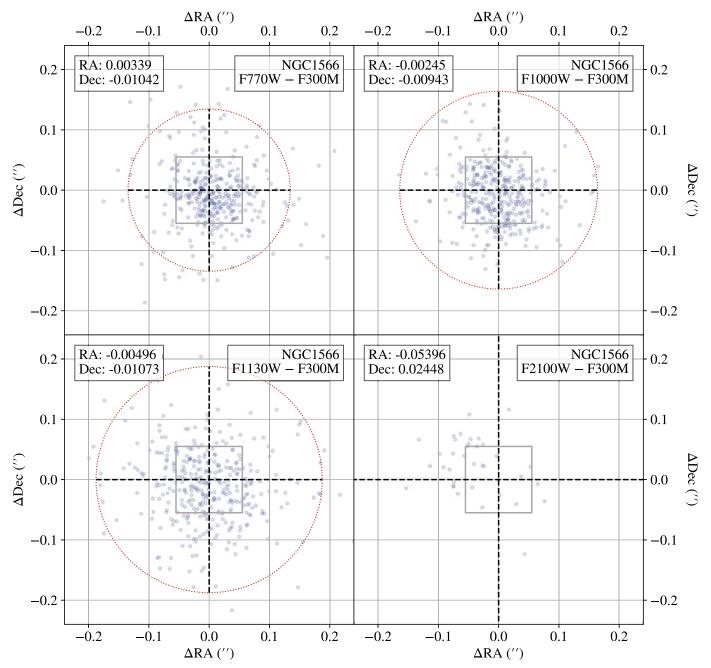


Figure 12. Diagnostic diagrams showing the MIRI-to-NIRCam relative astrometric alignment quality for example galaxy NGC 1566. Each of the panels displays the resulting R.A./decl. "scatterball" for a different MIRI band using our final MIRI-measured sky coordinate minus NIRCam F300M R.A./decl. The plotted points are plotted with transparency in order to allow assessment of peak location and distribution shape for each scatterball. The FWHM extent of the MIRI PSF in each band is drawn with a red dotted circle with a black dashed cross, and the pixel size on the MIRI detector is marked as a thick gray square. With text in each panel, we provide the sigma-clipped mean residual R.A./decl. offset from NIRCam F300M in arcseconds. This galaxy is in the top third of galaxies in ranked quality of MIRI-to-NIRCam alignment.

statistics and various changes to the data may affect results. For the most part, our First Results work took a conservative approach and many focused on source detection rather than detailed photometry. These results are unlikely to be affected by our improvements to the data processing. However, those that used flux measurements are likely to vary at the 10%–20%level, especially those using long-wavelength MIRI data, and any new publications will supersede the First Results (e.g., Sutter et al. 2024). Initial testing has also shown that the improvements to astrometry allow for better measurements of the concentration index (see Rodríguez et al. 2023). This means that in the F200W band, we can better discriminate star clusters from individual stars (M. J. Rodríguez et al. 2024, in preparation).

5.2. Comparison to MAST Mosaics

Here, we make a comparison to the Level Three mosaics that are available from MAST⁵⁸. These images are automatically processed using observatory-recommended default settings per

⁵⁸ As downloaded on 2024 April 2: http://mast.stsci.edu/.

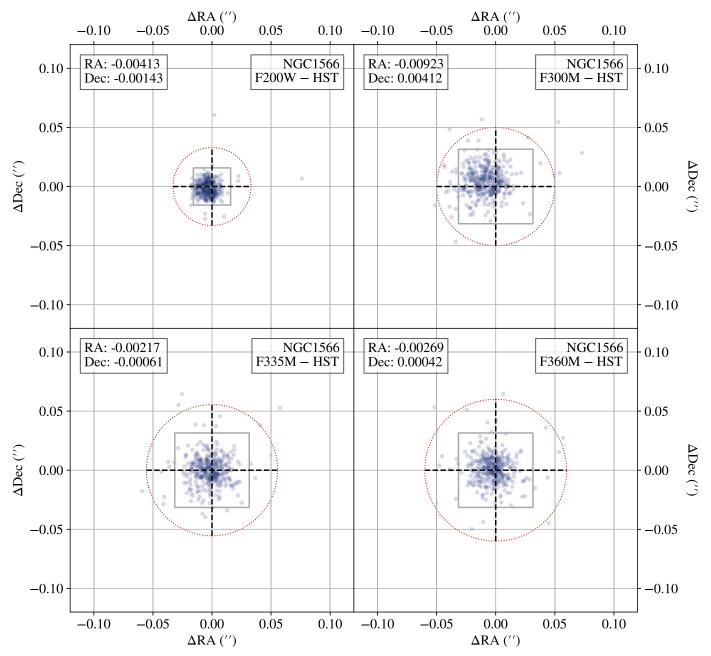


Figure 13. Diagnostic diagrams showing the NIRCam-to-HST astrometric alignment quality for example galaxy NGC 1566. Each of the panels displays the resulting R.A./decl. "scatterball" for a different NIRCam band using our final NIRCam-measured sky coordinate minus HST R.A./decl. The points are plotted with transparency in order to allow assessment of peak location and distribution shape for each scatterball. The FWHM extent of the NIRCam PSF in each band is drawn as a red dotted circle with a black dashed cross, and the pixel size on either the short (F200W) or long (F300M, F335M, F360M) NIRCam detectors is marked as a thick gray square. With text in each panel, we provide the sigma-clipped mean residual R.A./decl. offset from HST in arcseconds. This galaxy is typical in quality of NIRCam alignment, though it does not show evidence for tile-to-tile intramosaic alignment residuals as some galaxies do.

band, and as such are designed to produce reasonable quality mosaics across the full breadth of science that JWST will cover.

We show an example comparison in Figure 15 for NGC 628 for all eight bands and for the other 18 galaxies in the figure set. The differences are generally subtle in the NIRCam bands: we see the striping that is accounted for in our data processing showing up, as well as astrometric offsets between the two different versions. However, differences are small, generally on the order of a few tenths of a MJy sr^{-1} .

The difference is more substantial for the MIRI bands. We see the pronounced effects of our improved tile matching

(Section 3.5) as clear steps between the tiles. Due to how we set up our observations, MAST does not associate the background observations with the science automatically. This leads to clear nonlinear gradients across the tiles that have not been accounted for. For nonparallel observations where the backgrounds are directly linked to the science observations in the observing metadata, this will not be an issue. Astrometric offsets also appear in the difference images as galactic structures that have not been subtracted. These trends are similar across all of our targets (see Figure 15). Particularly for the MIRI, we strongly recommend our processing over the THE ASTROPHYSICAL JOURNAL SUPPLEMENT SERIES, 273:13 (20pp), 2024 July

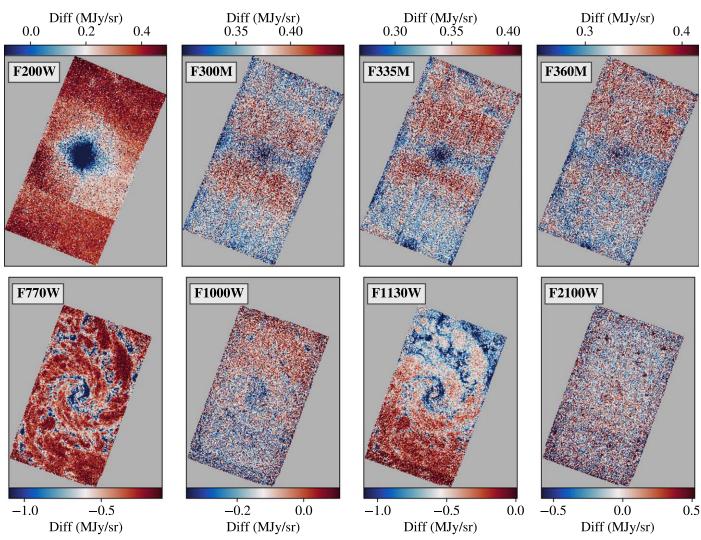


Figure 14. Comparison between our latest reprocessing and our early science mosaic for NGC 628. The difference here is the early science version subtracted from our latest reprocessed version, after reprojecting to the same world coordinate system (WCS). There are significant differences in the short NIRCam band, caused by both background-level matching and our new destriping method. At long NIRCam wavelengths, the differences are relatively minor. For MIRI, the improved calibration shows up as a galaxy imprint (most clear at F770W). Differences in astrometry show neighboring red/blueshifts and the background matching algorithm shows clear steps between mosaic tiles. The complete figure set showing equivalent comparisons for the other three galaxies that comprised our early science release is available online.

(The complete figure set (four images) is available in the online article.)

MAST-obtainable mosaics. For NIRCam, the differences are more subtle but our reprocessing efforts still represent an improvement over the automated JWST pipeline.

6. Released Data Products

Our data products are hosted at the MAST high-level science products web pages, at https://archive.stsci.edu/hlsp/phangs/ phangs-jwst, as well as on CANFAR at doi:10.11570/24.0090. Our primary data products are mosaics (img.fits), although we also provide compressed directories containing individually corrected tiles (tweakback.tar.gz), which are fully processed up to the final mosaic stage, and have absolute astrometric corrections back-propagated to them. The files are presented as multiextension .fits files, where the extensions are standard outputs from the JWST pipeline. These are detailed in Table 3. We note that pjpipe does not currently propagate the effects of any convolution to any arrays except the science and error arrays, so these other extensions will be invalid for PSF-matched images.

7. Summary

In this work, we have presented the overview of the steps we have taken in producing science-ready data for the PHANGS-JWST survey, implemented in our custom pipeline pjpipe (our DR1.1.0 was produced using pjpipeversion 1.1.0). This pipeline provides:

- 1. Wrappers around the official pipeline for Level One, Two, and Three processing (Sections 3.1, 3.2, and 3.7);
- 2. Optimized methods for both relative astrometry between mosaic tiles (Section 3.3) and absolute astrometry to external references (Section 3.8);
- 3. A number of algorithms to remove 1/*f* noise in NIRCam data (Section 3.6);

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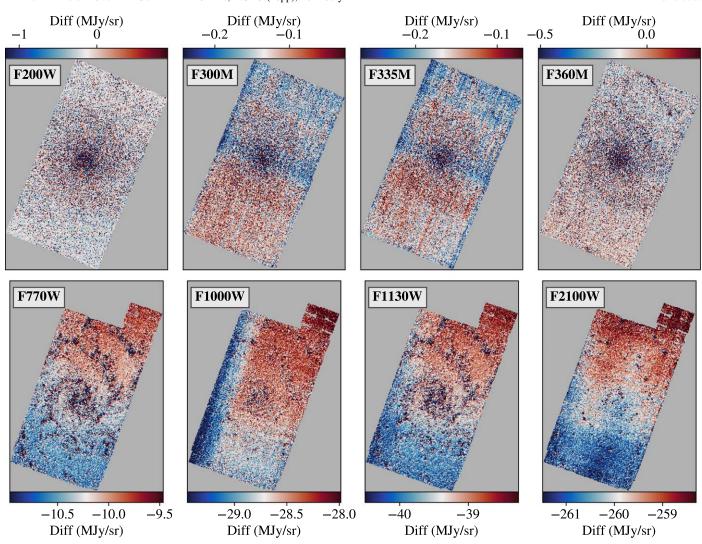


Figure 15. Comparison between our reprocessing and Level Three MAST-obtained mosaics for NGC 628. The difference here is the MAST version subtracted from our reprocessed version, after reprojecting to the same WCS. As can be seen, the astrometry is somewhat different in NIRCam, and the striping in the MAST mosaic is clear. The background levels are substantially different in the MIRI bands, due to a combination of our improved background matching algorithm and the fact that the MAST mosaics do not include our dedicated background observations. The complete figure set showing equivalent comparisons for the 18 other galaxies in our program is available online. (The complete figure set (19 images) is available in the online article.)

Table 3	
Description of Fits Extensions for Our Released Data Pro	ducts

Extension Name	Description	
SCI	2D data array containing the pixel values, in units of surface brightness (MJy sr ⁻¹)	
ERR	2D data array containing uncertainty estimates for each pixel. These values are based on the combined VAR_POISSON and VAR_RNOISE data	
	given as the standard deviation	
DQ	2D data array containing data quality flags for each pixel	
VAR_POISSON	2D data array containing the variance estimate for each pixel, based on Poisson noise only	
VAR_RNOISE	2D data array containing the variance estimate for each pixel, based on read noise only	
VAR_FLAT	2D data array containing the variance estimate for each pixel, based on uncertainty in the flat field	
CON ^a	3D context image, which encodes information about which input images contribute to a specific output pixel	
WHT ^a	2D weight image giving the relative weight of the output pixels	

Note.

^a Only present in img.fits files.

- 4. Improved background matching between mosaic tiles (Section 3.5), and the ability to separate the MIRI coronagraph from the main science observations to improve this process (Section 3.4);
- 5. Flux anchoring to external images, to provide absolute flux levels in images with no "true" background (Section 3.9);
- 6. Convolution algorithms to produce images at a matched angular resolution (Section 3.10).

While these techniques have been tailored specifically for our survey, they have flexible implementations and will be of general use in other observations of nearby galaxies.

We have extensively tested our processing methodology for 19 galaxies observed in four NIRCam and four MIRI bands, and our approach produces excellent mosaics across our sample. Further testing on the 55 galaxies that make up our Cycle 2 Treasury (program 3707, PI: Leroy) will help refine the data processing further. The automated pipeline enables rapid processing of these data so that we can make products available in a timely manner. We have developed a number of bespoke algorithms that will be of general use to the community, and hence our pipeline, pjpipe, is made public at https://github. com/phangsTeam/pjpipe and is pip-installable (pip install pjpipe). We have designed pjpipe to be flexible and easy to interface with, to maximize its usefulness. We make our first full public data release (DR1.1.0) mosaics available at https://archive.stsci.edu/hlsp/phangs/phangsjwst and stress that these represent a significant improvement on those we used in our early science papers, as well as those obtainable from the MAST database.

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Facility: JWST

Software: astropy (Astropy Collaboration et al. 2013, 2018, 2022), jwst (Bushouse et al. 2023), matplotlib (Hunter 2007), numpy (Harris et al. 2020), photutils (Bradley et al. 2022), SAOImageDS9 (Joye & Mandel 2003), scipy (Virtanen et al. 2020).

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References

- Allen, L. E., Calvet, N., D'Alessio, P., et al. 2004, ApJS, 154, 363
- Aniano, G., Draine, B. T., Gordon, K. D., & Sandstrom, K. 2011, PASP, 123, 1218 Argyriou, I., Lage, C., Rieke, G. H., et al. 2023, A&A, 680, A96
- Astropy Collaboration, Price-Whelan, A. M., Lim, P. L., et al. 2022, ApJ, 935, 167 Astropy Collaboration, Price-Whelan, A. M., Sipőcz, B. M., et al. 2018, AJ, 156, 123
- Astropy Collaboration, Robitaille, T. P., & Tollerud, E. J. 2013, A&A, 558, A33
- Bagley, M. B., Finkelstein, S. L., Koekemoer, A. M., et al. 2023, ApJL, 946, L12
- Barnes, A. T., Watkins, E. J., Meidt, S. E., et al. 2023, ApJL, 944, L22
- Belfiore, F., Leroy, A. K., Williams, T. G., et al. 2023, A&A, 678, A129
- Bradley, L., Sipócz, B., Robitaille, T., et al. 2022, astropy/photutils: v1.5.0, Zenodo, doi:10.5281/zenodo.6825092
- Bushouse, H., Eisenhamer, J., Dencheva, N., et al. 2023, JWST Calibration Pipeline, v1.11.3, Zenodo, doi:10.5281/zenodo.8157276
- Butterworth, S. 1930, Exp. Wirel. Wirel. Eng., 7, 536
- Chastenet, J., Sandstrom, K., Chiang, I.-D., et al. 2019, ApJ, 876, 62
- Chastenet, J., Sutter, J., Sandstrom, K., et al. 2023a, ApJL, 944, L11
- Chastenet, J., Sutter, J., Sandstrom, K., et al. 2023b, ApJL, 944, L12
- Dale, D. A., Boquien, M., Barnes, A. T., et al. 2023, ApJL, 944, L23
- Dale, D. A., Cohen, S. A., Johnson, L. C., et al. 2009, ApJ, 703, 517
- Dolphin, A. 2016, DOLPHOT: Stellar photometry, Astrophysics Source Code Library, ascl:1608.013
- Dolphin, A. E. 2000, PASP, 112, 1383

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- Emsellem, E., Schinnerer, E., Santoro, F., et al. 2022, A&A, 659, A191
- Engelbracht, C. W., Gordon, K. D., Rieke, G. H., et al. 2005, ApJL, 628, L29
- Fruchter, A. S., & Hook, R. N. 2002, PASP, 114, 144
- Harris, C. R., Millman, K. J., van der Walt, S. J., et al. 2020, Natur, 585, 357
- Hensley, B. S., & Draine, B. T. 2023, ApJ, 948, 55
- Hoyer, N., Pinna, F., Kamlah, A. W. H., et al. 2023, ApJL, 944, L25
- Hunter, J. D. 2007, CSE, 9, 90
- Jacob, J. C., Katz, D. S., Berriman, G. B., et al. 2010, arXiv:1005.4454
- Joye, W. A., & Mandel, E. 2003, in ASP Conf. Ser. 295, Astronomical Data Analysis Software and Systems XII, ed. H. E. Payne, R. I. Jedrzejewski, & R. N. Hook (San Francisco, CA: ASP), 489
- Kennicutt, R. C. J., Armus, L., Bendo, G., et al. 2003, PASP, 115, 928
- Lee, J. C., Sandstrom, K. M., Leroy, A. K., et al. 2023, ApJL, 944, L17
- Lee, J. C., Whitmore, B. C., Thilker, D. A., et al. 2022, ApJS, 258, 10
- Leroy, A. K., Sandstrom, K., Rosolowsky, E., et al. 2023, ApJL, 944, L9
- Leroy, A. K., Sandstrom, K. M., Lang, D., et al. 2019, ApJS, 244, 24
- Leroy, A. K., Schinnerer, E., Hughes, A., et al. 2021, ApJS, 257, 43
- Liu, D., Schinnerer, E., Cao, Y., et al. 2023, ApJL, 944, L19
- Pathak, D., Leroy, A. K., Thompson, T. A., et al. 2024, AJ, 167, 39
- Peltonen, J., Rosolowsky, E., Williams, T. G., et al. 2023, MNRAS, 527, 10668
- Perrin, M. D., Sivaramakrishnan, A., Lajoie, C.-P., et al. 2014, Proc. SPIE, 9143, 91433X
- Rest, A., Pierel, J., Correnti, M., et al. 2023, The JWST HST Alignment Tool (JHAT), v2, Zenodo, doi:10.5281/zenodo.7892935
- Rigby, J., Perrin, M., McElwain, M., et al. 2023, PASP, 135, 048001
- Robitaille, T. P., Meade, M. R., Babler, B. L., et al. 2008, AJ, 136, 2413
- Rodríguez, M. J., Lee, J. C., Whitmore, B. C., et al. 2023, ApJL, 944, L26
- Sandstrom, K. M., Koch, E. W., Leroy, A. K., et al. 2023, ApJL, 944, L8
- Schlawin, E., Leisenring, J., Misselt, K., et al. 2020, AJ, 160, 231
- Sheth, K., Regan, M., Hinz, J. L., et al. 2010, PASP, 122, 1397
- Sutter, J., Sandstrom, K., Chastenet, J., et al. 2024, ApJ, in press (arXiv:2405. 15102)
- Thilker, D. A., Whitmore, B. C., Lee, J. C., et al. 2022, MNRAS, 509, 4094
- Tielens, A. G. G. M. 2008, ARA&A, 46, 289
- Virtanen, P., Gommers, R., Oliphant, T. E., et al. 2020, NatMe, 17, 261
- Watkins, E. J., Barnes, A. T., Henny, K., et al. 2023, ApJL, 944, L24
- Williams, T., Egorov, O., Rosolowsky, E., et al. 2024, PJPipe, v1.1.0, Zenodo, doi:10.5281/zenodo.10458746
- Wolfire, M. G., Vallini, L., & Chevance, M. 2022, ARA&A, 60, 247
- Yang, G., Papovich, C., Bagley, M. B., et al. 2023, ApJL, 956, L12