



HAL
open science

Combining reference-star and angular differential imaging for high-contrast imaging of extended sources

S. Juillard, V. Christiaens, O. Absil, S. Stasevic, J. Milli

► To cite this version:

S. Juillard, V. Christiaens, O. Absil, S. Stasevic, J. Milli. Combining reference-star and angular differential imaging for high-contrast imaging of extended sources. *Astronomy & Astrophysics*, 2024, 688, 10.1051/0004-6361/202449747 . insu-04836831

HAL Id: insu-04836831

<https://insu.hal.science/insu-04836831v1>

Submitted on 15 Dec 2024

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.



Distributed under a Creative Commons Attribution 4.0 International License

Combining reference-star and angular differential imaging for high-contrast imaging of extended sources

S. Juillard^{1,*}, V. Christiaens^{1,**}, O. Absil^{1,***}, S. Stasevic^{2,3}, and J. Milli²

¹ STAR Institute, Université de Liège, Allée du Six Août 19c, 4000 Liège, Belgium
e-mail: sjuillard@uliege.be

² Univ. Grenoble Alpes, CNRS, IPAG, 38000 Grenoble, France

³ LESIA, Observatoire de Paris, Université PSL, Sorbonne, Université Paris Cité, CNRS, 5 place Jules Janssen, 92195 Meudon, France

Received 26 February 2024 / Accepted 3 June 2024

ABSTRACT

Context. High-contrast imaging (HCI) is a technique designed to observe faint signals near bright sources, such as exoplanets and circumstellar disks. The primary challenge in revealing the faint circumstellar signal near a star is the presence of quasi-static speckles, which can produce patterns on the science images that are as bright, or even brighter, than the signal of interest. Strategies such as angular differential imaging (ADI) or reference-star differential imaging (RDI) aim to provide a means of removing the quasi-static speckles in post-processing.

Aims. In this paper, we present and discuss the adaptation of state-of-the-art algorithms, initially designed for ADI, to jointly leverage angular and reference-star differential imaging (ARDI) for the direct HCI of circumstellar disks.

Methods. Using a collection of HCI datasets, we assessed the performance of ARDI in comparison to ADI and RDI based on iterative principal component analysis (IPCA). These diverse datasets were acquired under various observing conditions and include the injection of synthetic disk models at various contrast levels. We also considered reference stars with different levels of correlation with the science targets.

Results. Our results demonstrate that ARDI with IPCA improves the quality of recovered disk images and the sensitivity to planets embedded in disks, compared to ADI or RDI individually. This enhancement is particularly pronounced when dealing with extended sources exhibiting highly ambiguous structures that cannot be accurately retrieved using ADI alone, and when the quality of the reference frames is suboptimal, leading to an underperformance of RDI. We finally applied our method to a sample of real observations of protoplanetary disks taken in star-hopping mode, and propose to revisit the protoplanetary claims associated with these disks. Among eight proposed protoplanets claimed through velocity kinks or direct imaging, none of them were re-detected in our new processed images.

Key words. methods: data analysis – techniques: image processing – protoplanetary disks

1. Introduction

High-contrast imaging (HCI) is a technique that aims to observe faint signals near bright sources, such as exoplanets and circumstellar disks. After the majority of the starlight and aberrations caused by atmospheric turbulence have been removed by the coronagraph and the adaptive optics (AO), respectively, residual atmospheric speckles and quasi-static speckles not seen by the AO system remain (Racine et al. 1999). These speckles must be removed to reveal the faint circumstellar signal hiding in the vicinity of the star. Several observing strategies aim to provide a lever to distinguish the speckle field from the disk and planet signal, such as spectral differential imaging (SDI, Racine et al. 1999; Marois et al. 2000) and polarimetric differential imaging (PDI, Kuhn et al. 2001). In this paper, we focus on two strategies. Firstly, angular differential imaging (ADI, Marois et al. 2006) consists of capturing images over time while targeting a star and fixing the orientation of the telescope pupil, causing the field of view to rotate due to the Earth's rotation. The discriminant employed by the ADI strategy to differentiate the speckle from

actual signals is that, unlike the field of view that rotates according to the parallactic angle, the speckle field resulting from imperfections in the instrument remains mostly fixed. Angular diversity is typically leveraged by discarding the static component (which represents the speckle field) to retrieve the rotating circumstellar signal. Secondly, reference-star differential imaging (RDI, Mawet et al. 2009; Lafrenière et al. 2009; Ruane et al. 2019) involves using observations of reference star(s) that share similar characteristics with the scientific target of interest, either taken during data acquisition or selected from an archive library. The primary goal is to acquire reference images that exhibit a similar speckle pattern. References are leveraged as a model of the speckle field, which is to be subtracted from the observed data after being appropriately scaled in flux.

The choice of an appropriate observing strategy depends on the specific scientific goals, the nature of the target being observed, and the quality of available resources. Each of these high-contrast imaging strategies has its strengths and limitations. ADI is effective at distinguishing faint point-like sources, but it does require a sufficient parallactic angle rotation to function correctly and can cause deformations in extended sources due to rotational invariance (Milli et al. 2012). These deformations depend on the amount of rotation and the corresponding

* F.R.S.-FNRS PhD Research Fellow.

** F.R.S.-FNRS Postdoctoral Fellow.

*** F.R.S.-FNRS Senior Research Associate.

degree of rotational symmetry, considering the disk morphology (Juillard et al. 2023). This limitation not only hampers most face-on disks from being captured, but more generally applies to any extended source morphology, leading to challenging-to-correct deformations that can mislead the interpretation of the astronomer. Meanwhile, RDI is efficient for stable speckle patterns and can produce high-fidelity images of disks, but generally requires dedicated observations of an appropriate reference star in very similar observing conditions. It is less effective in the case of particularly challenging datasets, such as faint disks, time-varying speckle patterns, or when no well-suited reference star is available, which can be the case for archival ADI datasets.

To best utilize the capabilities of each strategy, employing a suitable post-processing algorithm is required. Principal component analysis (PCA, Soummer et al. 2012; Amara & Quanz 2012) is widely used and can be applied to both ADI and RDI, but it can suffer from overly aggressive PSF (point spread function) subtraction, resulting in deformations of the circumstellar signal (Milli et al. 2012). In the case of ADI, PCA and other PSF-subtraction algorithms such as median subtraction (c-ADI, Marois et al. 2006) and LOCI (locally optimized combination of images, Lafrenière et al. 2007) work well for point-like sources but generally struggle with extended structures such as disks. Despite many efforts to compensate for these limitations, more recent algorithms such as iterative PCA (IPCA Pairet et al. 2021; Stapper & Ginski 2022) or inverse problem (IP) approaches (MAYONNAISE, REXPACO, and MUSTARD; Pairet et al. 2021; Flasseur et al. 2021; Juillard et al. 2022) are still prone to deformations of extended structures, in particular when these structures include signals that appear static throughout the ADI image sequence. This means that angular diversity is not a suitable discriminator to distinguish this type of signal from the quasi-static speckle field. Regarding RDI, other algorithms such as data imputation with sequential non-negative matrix factorization (DI-sNMF, Ren et al. 2018) and the Karhunen-Loève transform with data imputation (DIKL, Ren 2023) have demonstrated efficiency in optimally extracting circumstellar signals and mitigating the risk of over-subtraction. While RDI can be highly effective, it necessitates that the speckle pattern remain relatively stable over time, and requires sufficiently correlated reference stars to produce high fidelity images.

Combining diverse observing strategies might assist in mitigating the distinct weaknesses of each approach, as exemplified by Lawson et al. (2022), who combined RDI with PDI, or by Wahhaj et al. (2013), Christiaens et al. (2019), and Flasseur et al. (2022), who merged ADI with SDI (ASDI) for extended sources. Previous studies in the scientific literature have explored the simultaneous use of RDI and ADI. These attempts typically involve using techniques such as median subtraction or PCA, and applying them to the concatenation of the reference and science cubes (e.g., Carter et al. 2023, Wallack et al. 2024). However, there has not been a dedicated effort to identify the most effective method of combining ADI and RDI, especially in the context of disk imaging. Previous research in this context was rather focused on comparing the performance of ADI with RDI (Ruane et al. 2019; Xie et al. 2022). In response to the described challenges and inspired by the many efforts made in the past few years to develop efficient novel algorithms to retrieve extended sources using ADI and RDI strategies separately, this paper aims to explore how one can adapt the IPCA algorithm to jointly leverage angular and reference-star differential imaging (ARDI) in a simultaneous approach. In Sect. 2, we explain the theoretical foundation and limitations of state-of-the-art IPCA algorithms employing ADI and RDI separately. In Sect. 3, we explain how

Table 1. Notations.

Notation	Definition
Constants	
m	Width of an image (i.e., the number of pixels on one side of the square images)
n	Number of frames in the ADI cube
r	Number of frames in the reference library
q	Number of dimensions in the PCA low-rank subspace (aka rank)
Matrices	
$Y \in \mathbb{R}^{n,m \times m}$	ADI dataset, i.e., a cube of vectorized images
$L_r \in \mathbb{R}^{r,m \times m}$	Library of r reference stars
$\bar{S} \in \mathbb{R}^{n,m \times m}$	Estimated speckle fields for each image in Y
$\bar{d} \in \mathbb{R}^{+,m \times m}$	Estimated circumstellar signal (forced to be positive)
$U \cdot \Sigma \cdot V$	Singular value decomposition of a matrix
Operators	
$\ _ \ $	Absolute value
$\mathcal{H}_q(Y)$	Operator ($\mathbb{R}^{n,m \times m} \mapsto \mathbb{R}^{n,m \times m}$) returning the estimated speckle fields via PCA using q principal components.
$\mathcal{H}_q^{L_r}(Y)$	Operator ($\mathbb{R}^{n,m \times m} \mapsto \mathbb{R}^{n,m \times m}$) returning the estimated speckle fields via PCA using the q principal components of a set of reference L_r
$\ _ \ _1$	l_1 -norm, sum of the absolute value of all elements in a matrix
\mathcal{Q}_Y	Rotation operator ($\mathbb{R}^{n,m \times m} \mapsto \mathbb{R}^{n,m \times m}$) creating a cube of n images where each frame contains the same image rotated according to the parallactic angles of the sequence of observations Y . This operator is used to inject simulated disks or planets into the data.
\mathcal{Q}_Y^{-1}	Median image of a cube after de-rotation ($\mathbb{R}^{n,m \times m} \mapsto \mathbb{R}^{m \times m}$)
$\[_ ; _]$	Concatenation of two collections of images

we adapted the IPCA algorithm to jointly leverage ADI and RDI. We also explore the question of parameter optimization. For Sect. 4, we conducted a comprehensive comparative evaluation of ARDI against RDI and ADI, all using an IPCA-based algorithm, and assessed their performance across diverse scenarios. In Sect. 5, we explain how we tested IPCA with ARDI on a sample of real observations of protoplanetary disks, and reevaluated the planet candidates claimed in these disks.

Notations. In the rest of the paper, we use the following notations, which are shared with a previous publication (Juillard et al. 2023): $Y \in \mathbb{R}^{n,m \times m}$ denotes the ADI sequence (also known as the ADI cube), with n being the number of frames, and where each frame is a vectorized square image of size $m \times m$ pixels. The ADI sequence is composed of the circumstellar signal d and the speckle field S . The estimate of a parameter is written with an over bar (e.g., \bar{S}). We use $\bar{S} \in \mathbb{R}^{n,m \times m}$ to denote the cube of the estimated speckle field, which is composed of n unique frames. We use $\bar{d} \in \mathbb{R}^{+,m \times m}$ to denote the estimation of the circumstellar signal, which is one single positive image, commonly shared by all frames of the cube. Finally, $L_r \in \mathbb{R}^{r,m \times m}$ denotes the library of reference stars containing r frames. Matrix products are written using a dot symbol (\cdot). A complete list of all the notation used in this paper is presented in Table 1.

2. IPCA algorithms for ADI and RDI processing

PCA is an unsupervised, model-less, statistical procedure that involves creating an orthogonal subspace to describe data through singular value decomposition (SVD). SVD generalizes the eigendecomposition of a square normal matrix with an orthonormal eigenbasis matrix. This decomposition is expressed as $Y = U \cdot \Sigma \cdot V$, where $Y \in \mathbb{R}^{n,m \times m}$ is a vectorized set of images, Σ is a diagonal rectangle matrix ($\mathbb{R}^{m,n}$) containing the singular values, and the two matrices U and V represent the left and right singular vectors. In our specific application, $U \in \mathbb{R}^{n,n}$ corresponds to the temporal principal components (PCs) of the data and $V \in \mathbb{R}^{m \times m, m \times m}$ corresponds to the spatial PCs. They are arranged such that the first PCs represent components that explain most of the variance in the data cube (e.g., the more static part, shared by most images), while the higher-rank PCs represent the more unique features of the dataset, such as the detector noise. To capture the quasi-static speckle field, we will use the orthogonal subspace formed by V . It is composed of principal components $[v_1, \dots, v_{m \times m}]$, each having the size of one image ($\mathbb{R}^{m \times m}$). By retaining low rank PCs and discarding higher-rank PCs, it is possible to create a subspace $V_q \in \mathbb{R}^{q, m \times m}$ of q images, forming a collection of PCs $[v_1, \dots, v_q]$ that capture the most spatially static and stable features of the data (i.e., the quasi-static speckles). Then, projecting the ADI images onto this subspace enables the extraction of the quasi-static speckle image at each frame such that $\bar{S} = Y \cdot V_q \cdot V_q^T$. We define the operator $\mathcal{H}_q(Y) = \bar{S}$ to encapsulate all operations described above to estimate the speckle field. Finally, subtracting \bar{S} from the original data creates a residual cube that hopefully still contains the circumstellar signals. The disk estimate \bar{d} is obtained from the median frame of the derotated residual cube.

Despite an appropriate choice for the number of PCs q used to build the speckle field estimate \bar{S} , PCA systematically leads to partial subtraction of extended signals (Esposito et al. 2014; Christiaens et al. 2019). The causes of these deformations were described by Pueyo (2016) as over-subtraction, which refers to the partial projection of circumstellar signal onto a low-rank subspace, and self-subtraction, which refers to the signal of interest being partially captured in the PCA low-rank subspace. While these effects can be attributed to the algorithms, they are linked to the validity of the underlying model used in post-processing. In particular, it has been observed through diverse algorithms that constraining the positivity of the circumstellar signal greatly prevents the risk of over-subtraction (Ren et al. 2018; Flassieur et al. 2021; Juillard et al. 2023). However, when using ADI-only for extended sources, the self-subtraction effect is partially caused by the inherent ambiguity of the model: the rotation-invariant flux. The deformations caused by this ambiguity cannot be corrected solely in post-processing if no additional information is provided (Juillard et al. 2023).

Iterative PCA has demonstrated proficiency in preserving disk signal when leveraging angular diversity (Pairet et al. 2021; Stapper & Ginski 2022; Juillard et al. 2023). The iterative aspect enables the usage of a model that constrains the positivity of the estimates to enhance the performance of PCA. In this section we reintroduce the IPCA method using either ADI or RDI. These approaches will enable us to establish, in Sect. 4, a consistent comparison with our proposed adaption of IPCA for ARDI, as all three strategies, or combinations of strategies, make use of an IPCA-based algorithm. Alongside the presentation of these algorithms, we propose an explanation of the concept of a fixed point, which is fundamental to IPCA algorithms. Understanding the concept of a fixed point is not only essential for

these algorithms, but also crucial for appreciating parameter optimization.

2.1. ADI with IPCA

IPCA consists of iteratively performing PCA on the science data cube Y while subtracting the previously estimated disk signal \bar{d}_i at each step, while imposing positivity, thereby preventing over-subtraction. The process starts with $\bar{d}_0 = 0$ (an image of zeros). For a given rank q , one iterative step is detailed as follows:

$$\begin{aligned} \bar{S}_{i+1} &= \mathcal{H}_q(Y - Q_Y(\bar{d}_i)), \\ \bar{d}_{i+1} &= \|Q_Y^{-1}(Y - \bar{S}_{i+1})\|, \end{aligned} \quad (1)$$

where $Q_Y : \mathbb{R}^{m \times m} \mapsto \mathbb{R}^{n, m \times m}$ is the rotation operator associated with the data cube Y , which creates a cube of n images where each frame contains the same image rotated according to the parallactic angles of the sequence of observations Y . The inverse operation $Q_Y^{-1} : \mathbb{R}^{n, m \times m} \mapsto \mathbb{R}^{m \times m}$ corresponds to the median image of the derotated input. The theoretical foundation of this process hinges on the concept of a fixed-point algorithm. A fixed point for a function f , is a point that remains unchanged by the function, such that $f(x_{\text{fix}}) = x_{\text{fix}}$. The stationary point can be found by iterating over the sequence $x_{i+1} = f(x_i)$. Optimization methods, such as the well-established Newton's method to find roots or minima, rely on this principle. In our specific application, the aim is to reach the following fix-point d_{fix} , such that $d_{\text{fix}} = \|Q_Y^{-1}(Y - H_q(Y - Q_Y(d_{\text{fix}})))\|$, while finding the optimal description of the quasi-static component via PCA using q components: $\bar{S} = \|H_q(Y - Q_Y(\bar{d}))\|$. The rank q determines the variability of the estimated speckle field. For rank $q = 1$, the speckle field is considered to be morphologically static and can only vary in intensity. Selecting a large value for q incorporates more high-variance components into the speckle field description, providing a more intricate representation, capable of capturing greater variability. However, it is crucial to maintain a low rank to avoid encapsulating astrophysical circumstellar contributions in the speckle field estimation. This constraint is commonly referred to as the low-rank approximation.

When performing an IPCA step, as described in Eq. (1), it is possible to either keep the same rank q , or to increase q . Subtracting the disk estimation from the original cube before performing the next PCA operation enables more disk signal to be preserved while increasing the rank. However, at each step, the disk estimate will also contain noise and artifacts that will then propagate through the iterations. Hence, wisely choosing the starting rank q , when to increase it through the iterations, and when to end the process, is required to obtain an optimal result. The choice of parameters will influence the amount of noise propagating in the early iterations, and depends on the contrast and variability of the speckle field. When the low-rank approximation is not suitable for capturing enough of the speckle field, as is the case when the amplitude of the variations in the speckle field is larger than the signal intensity, it is necessary to set the starting rank q to a higher value to avoid the propagation of bright artifacts that appear in the first iterations. For the same purpose, the number of iterations per rank should also be limited, especially at a low rank. Theoretically, multiple sets of parameters should lead to very similar results. In practice, this statement holds true for well-behaved datasets (e.g., bright disks observed under stable conditions), but it can be more sensitive when strong residuals appear in early iterations. The estimate reached after a few iterations has been observed to correct for

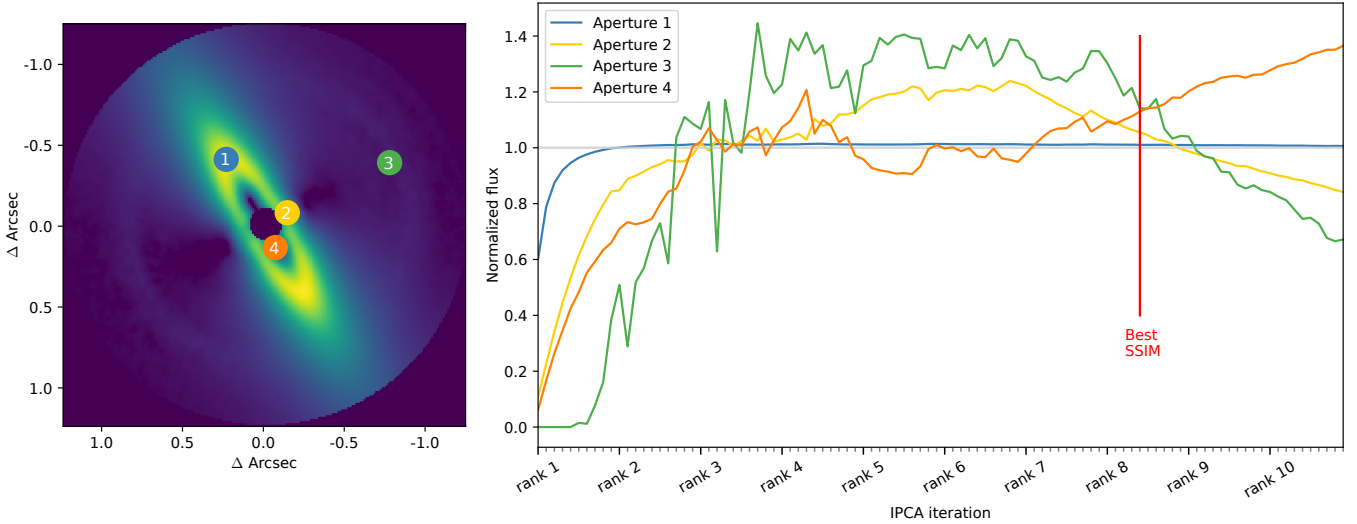


Fig. 1. Best disk image, according to the SSIM metric, obtained with the IPCA algorithm leveraging ADI only while processing a synthetic dataset (disk #B, cube #3 at a contrast of 10^{-3} ; see Sect. 4). The image on the left displays different apertures placed on (1) the brightest part of the disk, (2) a region of the disk that is rotation invariant, (3) a noise region, and (4) the region near the coronagraphic region where the variation of the speckle field is higher. On the right-hand side, we plot the evolution of the integrated flux computed in these different apertures through the IPCA iterations. The red vertical line indicates the position of the best image among all iterations according to SSIM. Each curve in the figure is normalized by its mean. The x -axis represents each iteration of the IPCA process, starting at rank 1, up to rank 10, with 10 iterations per rank. The gray horizontal line, placed at $y = 1$, indicates the position where the flux is equal to its mean.

over-subtraction and recover the signal with its correct intensity (Stapper & Ginski 2022; Juillard et al. 2023). However, it is important to note that when using ADI-only for imaging extended signals, the estimate will still suffer from deformation stemming from rotation-invariant components (Juillard et al. 2023). Indeed, algorithms like PCA are designed to remove all signals that appear static and quasi-static. Therefore, without a method to prevent ambiguous rotation-invariant signals from being assigned to the speckle field, the contribution from the disk will inevitably not be fully preserved.

In Fig. 1, we show the evolution of the disk estimate through IPCA iterations on an example synthetic dataset. The IPCA is parameterized such that it performs ten iterations per rank, while starting at rank $q = 1$ and increasing the rank up to $q = 10$. The image on the left provides the best estimation of the injected disk according to the structural similarity index measure (SSIM, Wang et al. 2004). This metric, which we also use in Sect. 4, is particularly adapted for assessing image similarity. It evaluates not only the pixel-wise differences between the images but also considers structural information such as contrast, luminance, and structure on various windows of the two images. Indeed, ranking disk image estimations is a multifactorial problem that often necessitates finding a satisfactory compromise between fewer geometrical biases and fewer noise residuals or errors in the overall restored flux.

We can observe that the flux in the aperture #1, placed in the brightest part of the disk, converges rapidly and stays stable through the iterations, which is the desired behavior of a reached fixed-point. In contrast, in other areas, such as those with only noise or rotation-invariant flux, no convergence is observed.

2.2. RDI with IPCA

The implementations of the RDI algorithm found in the scientific literature (Schneider et al. 2009; Soummer et al. 2012; Ren et al. 2018) involves the subtraction of an estimated speckle pattern

based on a scaled version of the reference images. We propose an adapted version of the IPCA algorithm that employs an iterative process to optimize the projection of the estimated reference. For a given set of references $L_r \in \mathbb{R}^{r,m \times m}$ of r reference images, we subsequently refer to $V_q^{\text{ref}} \in \mathbb{R}^{q,m \times m}$ as the collection of q PCs estimated by PCA on the reference library. We define the operator $\mathcal{H}_q^{L_r}(Y)$, which, similarly to $\mathcal{H}_q(Y)$, returns \bar{S}_{i+1} by projecting the science cube into the principal components, which are now computed on the reference library. We note that $\mathcal{H}_q(Y) = \mathcal{H}_q^{L_r}(Y)$. The subsequent iterative process can then be expressed as:

$$\begin{aligned} \bar{S}_{i+1} &= \mathcal{H}_q^{L_r}(Y - Q_Y(\bar{d}_i)), \\ \bar{d}_{i+1} &= \|Q_Y^{-1}(Y - \bar{S}_{i+1})\|, \end{aligned} \quad (2)$$

where \bar{S}_{i+1} is obtained by projecting the data cube Y onto the subspace created from the references, while subtracting the previous disk estimation at each step. Unlike in ADI-based IPCA, the PCs computed via the references are not updated throughout the process. This approach seeks the optimal projection of the reference star PCs on the dataset, while considering disk signal as positive.

We display in Fig. 2 the convergence process to find the optimal projection while using IPCA-RDI of rank $q = 1$. This example was produced using the same test dataset as in Fig. 1. In this example, we can observe that the flux exhibits a very clear convergence, which is achieved in only a few iterations. This iterative process enables the over-subtraction effect to be mitigated, and significantly improves the results compared to classical PCA. RDI however remains intrinsically sensitive to the stability of the speckle pattern over time, and necessitates the availability of sufficiently correlated reference stars.

Other algorithms, such as DI-sNMF (Ren et al. 2018), have also demonstrated efficiency in preserving the disk signal and preventing over-subtraction. However, for consistency in the test pipeline used for performance comparisons in Sect. 4, we use the

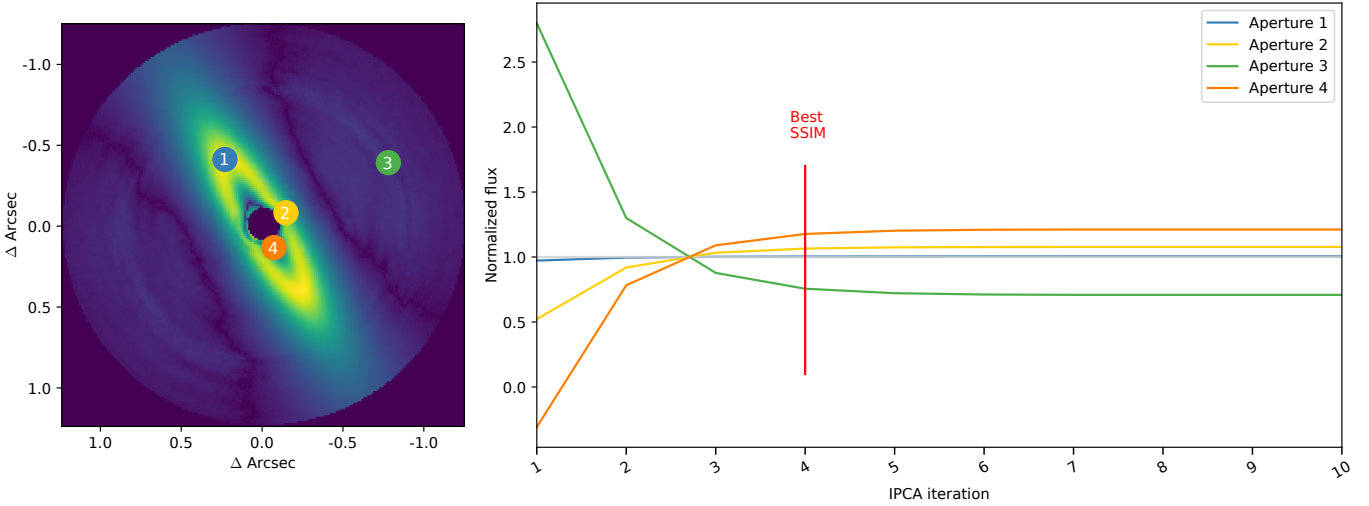


Fig. 2. Same as Fig. 1, but for RDI-only IPCA. The set of reference frames used here is presented in Sect. 4 and is referred to as “optimal”. The x -axis represents each iteration of the IPCA process, for ten iterations using a rank-1 PCA.

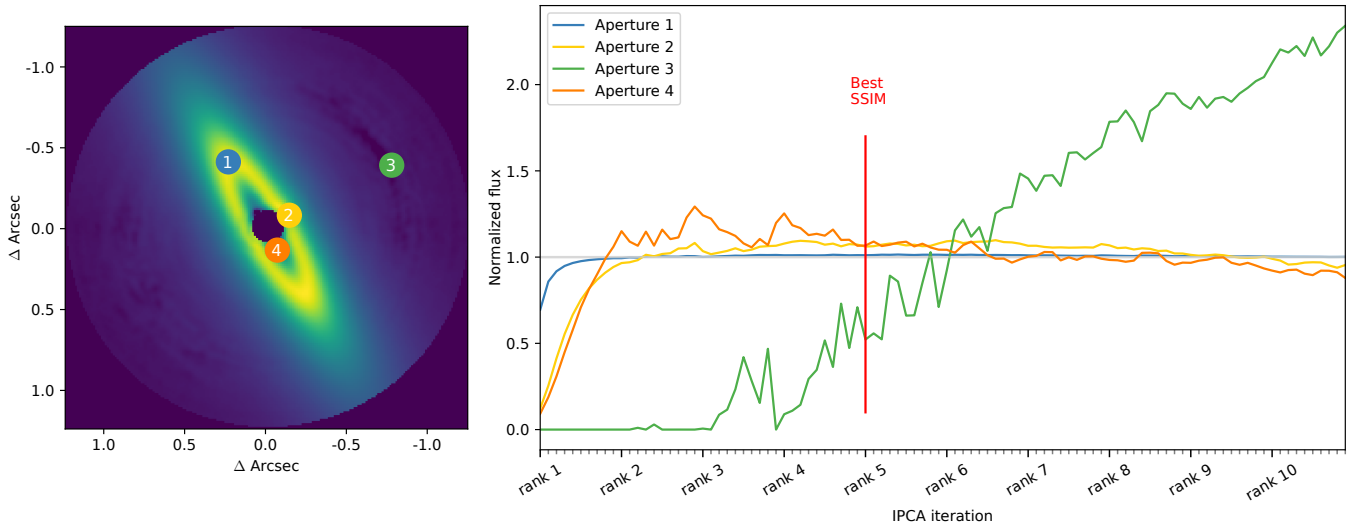


Fig. 3. Same as Fig. 1 but by leveraging the ARDI strategy with an IPCA algorithm, using the same reference frames as in Fig. 2.

IPCA implementation for RDI, similar to what we implemented for the ADI and ARDI strategies.

3. Combining angular and reference-star differential imaging (ARDI)

Simultaneously employing both RDI and ADI strategies is expected to help mitigate the individual limitations of each strategy. The combination of RDI with ADI holds the potential of offering several advantages in the field of high-contrast imaging for exoplanet and disk detection around stars. Here, our primary focus is to adapt the IPCA approach, initially designed for ADI, to leverage ARDI. Other advanced algorithms using both ADI and RDI, including IP-based approaches, could also be considered, but are beyond the scope of this paper. In this section we also explore the problem of parameter optimization for IPCA.

To combine RDI with ADI using IPCA, we propose injecting the reference star(s) into the ADI cube to compute the PCA low-rank subspace, by concatenating the most correlated reference frames to the ADI cube of images, so that both are used

simultaneously to build the speckle field estimate through PCA. The process can be expressed as follows:

$$\begin{aligned}
 L_{n+r,i} &= [Y - Q_Y(\bar{d}_i); L_r], \\
 \bar{S}_{i+1} &= \|\mathcal{H}_q^{L_{n+r,i}}(Y - Q_Y(\bar{d}_i))\|, \\
 \bar{d}_{i+1} &= \|\mathcal{Q}_Y^{-1}(Y - \bar{S}_{i+1})\|,
 \end{aligned} \tag{3}$$

where the operation $[_ ; _]$ represents the concatenation of two collections of images, creating a set of images $(L_{n+r,i})$ with a size of $\mathbb{R}^{n+r,m \times m}$ that contains both the data and reference images. The ratio between the size r of the reference library and the number n of frames in the ADI cube must be chosen carefully for optimal results. We refer to r/n as the reference frame ratio. The proposed algorithm can be classified as semi-supervised, as it simultaneously leverages reference frames for the speckle field (which can be regarded as labeled data), and the actual (unlabeled) data to provide an optimal estimation.

In Fig. 3, we monitor the evolution of the disk estimate through the IPCA iterations while using the ARDI strategy, similarly to the test presented in Fig. 1 with ADI only. We used

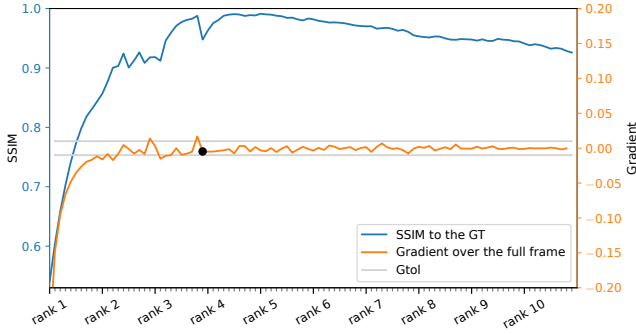


Fig. 4. Evolution of the IPCA with ARDI estimate through the iterations, using the same dataset as in Fig. 3. The blue curve relates to the left y -axis and displays the value of the SSIM for each estimate. The orange curve relates to the right y -axis and represents the gradient computed over the full frame, as described in Eq. (4). The g_{tol} has been set to 10^{-2} and is represented on the figure by the two gray lines. The black dot indicates the iteration where the gradient curve reaches the g_{tol} and does not exceed it in further steps. The x -axis represents each iteration, where the minor ticks represent a single iteration, and the major ticks represent a rank update.

identical parametrization (ten ranks, ten iterations per rank), and a reference frame ratio (r/n) of one. Conversely to what we observed with ADI only, three out of the four apertures seem to converge when using reference images. Results obtained with ARDI do not exhibit the deformation due to rotation invariant flux (located in aperture #2). However, the background noise in the final image increases throughout the iterations. Consequently, determining where to stop the iterative process involves choosing a good compromise between noise subtraction (non-propagation), filtering, and recovery of extended signals (convergence). A common criterion to assess convergence is setting a stopping criterion based on the gradient of the iterate. In our case, the criterion, based on the disk estimate, reads as follows:

$$\left\| \frac{|\bar{d}_i|_1 - |\bar{d}_{i+1}|_1}{|\bar{d}_i|_1} \right\| < g_{\text{tol}}, \quad (4)$$

where g_{tol} stands for the gradient tolerance and defines the precision threshold to be reached to assess convergence. It is a threshold to the gradient of the iterate, computed in our case by comparing disk estimates between two consecutive iterations, i and $i + 1$.

We show in Fig. 4 the evolution of the SSIM with respect to the ground truth (the injected disk), as well as of the gradient (Eq. (4)), as a function of the iterations. Additionally, we present in Appendix A the evolution of the gradient across patches of 15×15 pixels. Firstly, we can observe that multiple iterations achieve a similar SSIM, corresponding to different estimates of equal quality. Secondly, the iteration where convergence is reached according to the gradient criterion occurs close to the iteration where the SSIM is the highest (even though the criterion serendipitously falls into a local minimum in the SSIM curve). The region over which the gradient is computed can introduce bias, as it may not accurately represent the convergence behavior of the signal of interest. Indeed, if the field of view is too large or if the signal-to-noise ratio of the circumstellar signals is low, then the regions where there is no signal can become dominant in the computation of the gradient. Conversely, a few pixels of significantly higher intensity may dominate the gradient computation, potentially masking the convergence of the

full disk, including its fainter components. When computing the convergence criterion over patches (Fig. A.1), we observe that only patches containing signal achieve convergence, and that the brightest part of the disk converges faster than the fainter parts. Overall, we conclude that monitoring the convergence curves can provide a relevant indicator of the optimal iteration to be chosen. Nevertheless, assessing the convergence of a numerical method in computer science is a recurring problem for which proposed solutions are known to lack reliability in some cases (Nocedal et al. 2002). Moreover, it is particularly challenging in our case of application, given that the iterative process is not guaranteed to converge.

An alternative method for combining RDI and ADI is to utilize the PCs computed via RDI as the initialization of IPCA with the ADI cube. Hence, only the initialization differs and is expressed by the following equations:

$$\begin{aligned} \bar{S}_0 &= \mathcal{H}_q^{Lr}(Y), \\ \bar{d}_0 &= \|Q_Y^{-1}(Y - \bar{S}_0)\|. \end{aligned} \quad (5)$$

Then, the iterative steps are identical to the IPCA with ADI presented in Sect. 2.1. While we did not observe any significant nor systematic improvement when choosing this method, it remains relevant and can in some cases produce results consistent with those obtained with IPCA-ARDI.

4. Testing algorithms

4.1. Tests based on simulated data

In this section we compare the three different strategies: RDI, ADI, and ARDI, all utilizing the same core algorithm, IPCA. We consistently employed the same evaluation pipeline as presented in Juillard et al. (2023), which was previously used to compare three different algorithms for processing datasets using ADI alone. This test pipeline consists of a total of 60 test datasets composed of five different disk morphologies, injected at three different contrast levels (10^{-3} , 10^{-4} , and 10^{-5}), into four different observing ADI sequences of stars without any known circumstellar signal, reflecting different observing conditions. The test datasets are available on Zenodo¹.

The datasets, obtained through the High-Contrast Data Center (HCDC), were acquired using the Infrared Dual-Band Imager and Spectrograph (IRDIS, Dohlen et al. 2008; Vigan et al. 2010) camera of the Spectro-Polarimetric High-contrast Exoplanet Research coronagraphic system on the Very Large Telescope (VLT/SPHERE, Beuzit et al. 2019). The test datasets all consist of the $H2$ channel from the dual-band $H23$ set. They were chosen to exhibit a diverse range of characteristics, including low Strehl ratio with a 26° rotation (ID#1), wind-driven halo (ID#2), an unstable speckle field (ID#3), and good Strehl ratio with an 80° field rotation (ID#4). The raw data processed with the data handling software (Pavlov et al. 2008) of the HCDC (Delorme et al. 2017), which performs dark, flat, and bad pixel correction on a coronagraphic sequence. For future reference, we computed the mean and standard deviation of the Pearson correlation coefficients (PCC) between each unique pair of frames in the ADI cube. The mean PCC are as follows: Cube #1: $\mu = 0.99$; Cube #2: $\mu = 0.97$; Cube #3: $\mu = 0.93$; Cube #4: $\mu = 0.96$, with standard deviations below 0.001 for all the cubes.

The injected disks represent a range of scenarios for both debris and protoplanetary disks. As detailed in

¹ Datasets available at <https://zenodo.org/records/11442267>

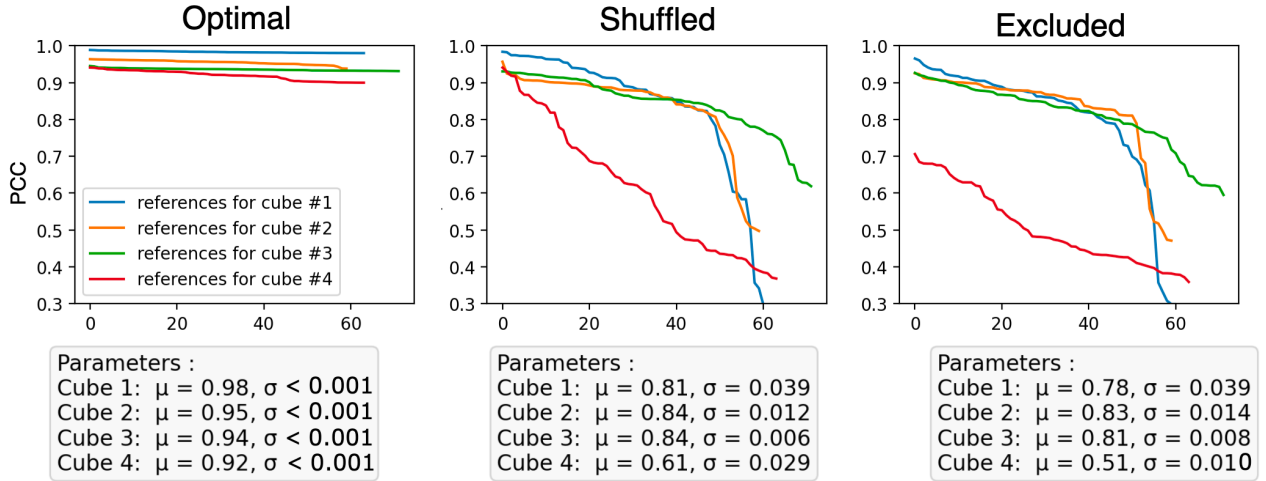


Fig. 5. Pearson correlation coefficient of the selected reference frames computed individually and averaged over each frame of the raw ADI data before the injection of simulated disks or planets, for three different selections of references: (left) optimal references selected; (middle) a random sample of each cube’s dedicated references, in an even proportion; and (right) a random sample of references excluding the optimal references. These coefficients were calculated within an annulus with an outer radius of 60 pixels and inner radius of 30 pixels, and are classified from the most correlated to the least correlated. The text box below each figure displays the mean PCC (μ) and the variance (σ) of the reference sample relative to its associated data cube.

Juillard et al. (2023), this selection consists of two 75° inclined disks with varying sharpness levels (A and B), a 45° inclined disk with two concentric rings (C), a nearly face-on disk with azimuthal flux variation (D), and a hydrodynamical simulation of a disk with embedded spiral structures and a companion (E). The contrast of the injected disks is determined by measuring the integrated flux within a full width at half-maximum (FWHM)-sized aperture, centered at the peak intensity of the disk, and then dividing this value by the integrated flux within an FWHM-sized aperture of the stellar point spread function. However, we made an exception for the synthetic disk #E, where we measured the flux at the companion location.

We performed two series of tests. In the first series, we compared the performance of ADI, RDI and ARDI using an IPCA-based algorithm on our 60 test datasets, while using an optimal choice of reference stars. The reference frames were selected from a set of archival IRDIS observations taken with the same filter, coronagraph, and exposure time as the test datasets. These reference targets were observed between 2014 December 11 and 2021 June 1, and the raw data were calibrated through the same process as the test datasets. For each test dataset, the PCC was calculated between the frames of the dataset and the reference targets, excluding any observations of the dataset star taken at different epochs. The PCC was calculated within a circular annulus between $0''.31$ and $0''.67$, which captures both the dominant speckle region and position of the waffle pattern, used for precise star centering of a coronagraphic sequence (Zurlo et al. 2014), if it was included in the observation. For each frame in the dataset, the 300 best correlated reference frames were identified, and those that appeared in this selection for more than 30% of the dataset frames were selected for the final reference library.

The left-hand side plot in Fig. 5 (“Optimal”) displays the PCC for each frame of the reference library. The PCC value for each reference frame corresponds to the average of the correlation coefficients computed individually with each frame of the ADI cube before disk injection. The reference frames are then sorted from the most correlated to the least correlated following this methodology.

For the second series of tests, we aimed to assess how ARDI could assist in cases of inadequate reference stars. We tested the ARDI-IPCA algorithm for three different qualities of reference frames. In the “Optimal” case, we used the most correlated frames, which is the same selection used in the first series, where we compare ADI, RDI, and ARDI. In the “Shuffled” case, we randomly selected a sample from the reference libraries dedicated to each of our four ADI sequences into one common reference library. In this test, the random selection contains 25% of frames from the “Optimal” reference library. In the “Excluded” case, we selected for each ADI cube a sample only from references dedicated to the three other ADI cubes, creating a selection that excludes optimal references. The PCC computed for each of the three selections of references is presented in Fig. 5. For the IPCA processing in ADI, RDI, and ARDI, we used our implementation of the GreEDS algorithm² (Pairet et al. 2021; Juillard et al. 2023). This Python implementation utilizes PyTorch, enabling the processing of a dataset from the test pipeline in less than a minute on a standard laptop. For all tests, we used an identical parameterization: ten ranks and a reference frame ratio (r/n) of one. The number of iterations increases by one at each rank (e.g., one iteration is performed at the first starting rank q , two at rank $q + 1$, and so on). This enables improved results compared to setting the number of iterations per rank to 10, as proposed in Juillard et al. (2023), especially for faint disks.

A table representing the best algorithm for each dataset and for the three different strategies (ADI, RDI, and ARDI), is shown in Fig. 6. The quality of the extracted disk images is assessed using the SSIM metric. SSIM is chosen for its ability to cover multiple aspects of image similarity and its proven capacity to approximate structural differences akin to human visual perception (Wang et al. 2004). In Appendix B, we display the same table including bar plots in each cell to indicate the SSIM value obtained by each algorithm for each test case. However, given that assessing the quality of disk image estimation involves a multifactorial challenge, often requiring a balanced

² <https://github.com/Sand-jrd/GreEDS>

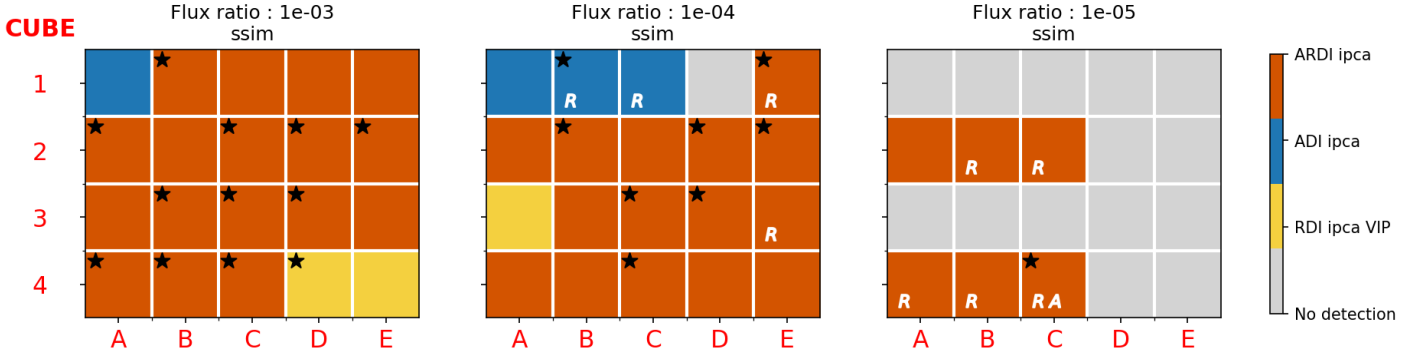


Fig. 6. Results of our systematic tests comparing the three strategies RDI, ADI, and ARDI using the IPCA-based algorithm. Each cell represents a different synthetic dataset, combining disk morphologies labeled along the x -axis with datasets before injection along the y -axis. Each figure contains three tables representing three different levels of contrast: 10^{-3} (left), 10^{-4} (middle), and 10^{-5} (right). The color of the cell indicates which algorithm performed best according to the SSIM metric. The white letters on a cell indicate which algorithm(s) did not detect the disk (R \rightarrow IPCA-RDI, A \rightarrow IPCA-ADI, A_R \rightarrow IPCA-ARDI). Gray cells mean that no algorithm detected the disk. Black stars indicate that the five metrics (SSIM, Spearman, Pearson, Euclidean, and SAD) selected the same winner.

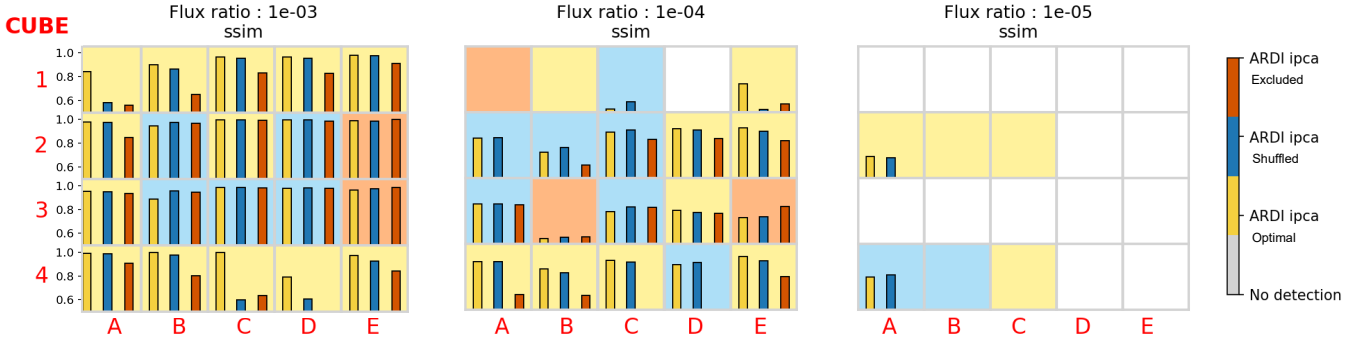


Fig. 7. Results of our systematic tests comparing different qualities of references while utilizing ARDI for IPCA. The organization of this table is identical to Fig. 6, except that the value of the SSIM of the disk estimation compared to the ground truth is also represented for each algorithm within each cell.

trade-off between minimizing geometric biases and reducing residual noise, we also consider four additional metrics – namely, Pearson, Spearman, sum of absolute differences (SAD), and Euclidean distance – in Appendix C.1. The cases where all these metrics unanimously favor a specific algorithm are highly likely to stand out as clearly superior and are marked with a black star in the upper left corner of the table. The field of view considered for calculating the metrics affects the results. We calculated the metrics using a $1''$ -radius aperture, excluding a 6-pixel radius inner circle located behind the coronagraphic mask. Complete results, including the estimated disks (d) for each algorithm and residual plots in Appendix D are available on Zenodo³. The results for the second series of tests, aiming to compare three different selections of reference frames, are displayed in Fig. 7, using a table with bar plots as in Appendix B.

4.2. Discussion of the test results

The results of the first series of tests show that using the combined ARDI strategy, using optimal reference frames, is most often selected as the best strategy according to the SSIM (as well as other metrics). A significant visual improvement is observed for disk #E, as shown in Figs. D.1E, D.2E, D.3E, and D.4E, particularly at a medium contrast of 10^{-4} . At this contrast, RDI

struggles to capture the planet, while ADI has difficulty capturing the disk structure. We show in Fig. 8 the injected and retrieved disk images in four cases where ARDI performed significantly better than using either ADI or RDI individually. The combination of both strategies proves to be the most relevant, as the weaknesses of each algorithm are effectively compensated for when used together. However, the degree of improvement provided by ARDI is not uniform across all datasets, including seven cases for which ADI or RDI alone perform better than ARDI according to SSIM. We propose to examine each of these specific cases:

- ADI outperforms ARDI in dataset #1 for disks A to C with a contrast of 10^{-4} , as shown in Figs. D.1A, D.1B, and D.1C. We observe that ADI achieves superior performance in scenarios where the dataset has a low Strehl ratio. Interestingly, this is also the same context where RDI shows lower performance. This observation suggests that ADI performs better when the information provided by the reference dataset is largely irrelevant. dataset #1A (Fig. D.1A) with a contrast of 10^{-3} is an exception, where ADI is selected as the best method based on the SSIM metric in Fig. 6. However, it is important to note that there is no consensus among the metrics. In fact, three of the five metrics (SAD, Euclidean and Pearson) favor ARDI, while one, Spearman correlation, indicates a preference for RDI. Upon visual inspection, it is obvious that ADI does not provide the most satisfactory result in this particular case.

³ Appendix D available at <https://zenodo.org/records/11442350>

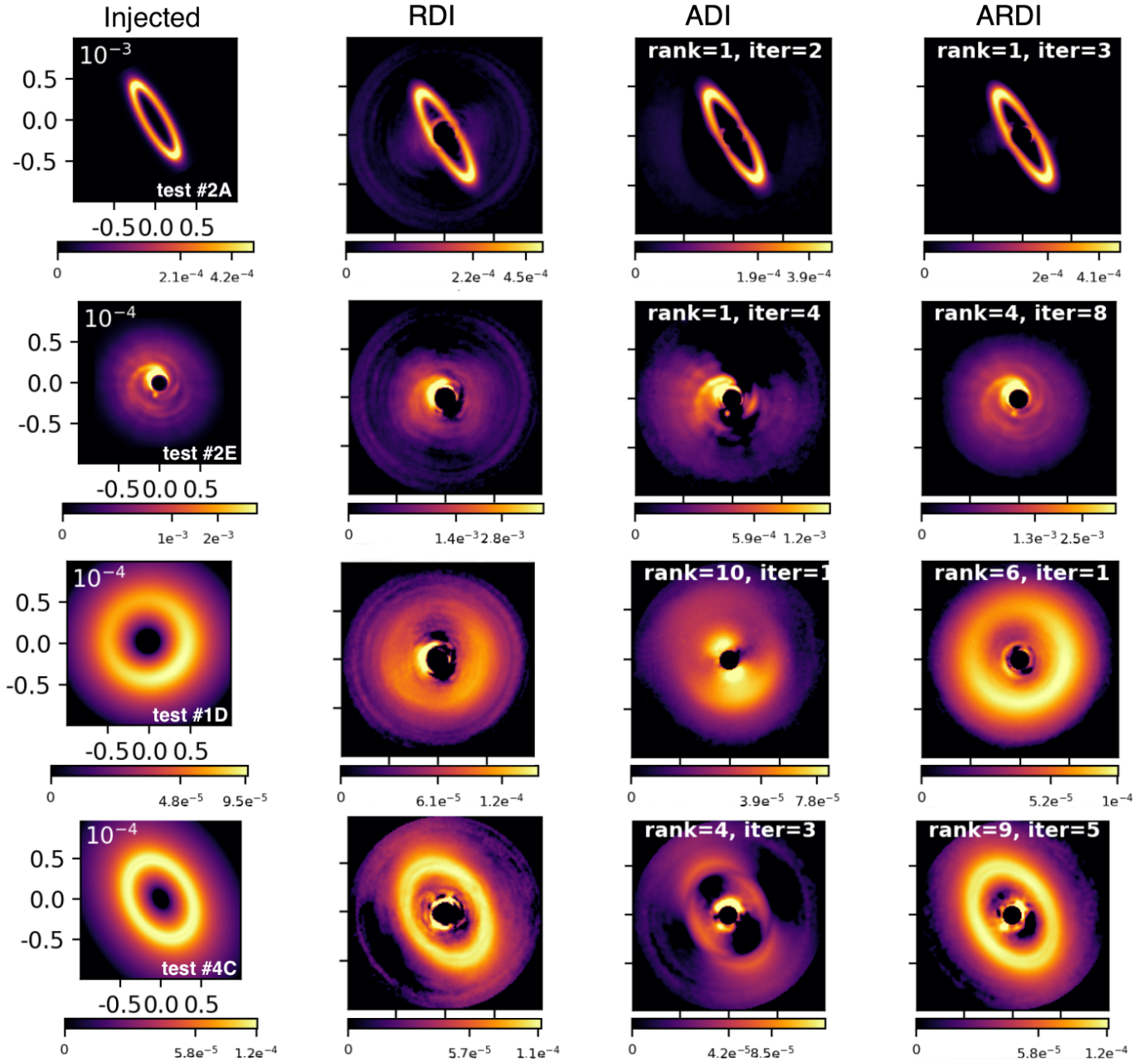


Fig. 8. Four examples of IPCA disk estimations obtained using RDI, ADI, and ARDI strategies, for a test dataset where ARDI performed significantly better than ADI and RDI. From top to bottom, the test dataset are cube #2 with disk A injected at a contrast of 10^{-3} , cube #2 with disk E at a contrast of 10^{-4} , cube #1 with disk D at a contrast of 10^{-4} , and cube #4 with disk C at a contrast of 10^{-4} . From left to right, the images display the injected disk and the results using RDI, ADI, and ARDI, utilizing the IPCA algorithm. These results were selected as they demonstrate ARDI’s capability to compensate for the individual limitations of each strategy compared to using each method individually, in a variety of scenarios in terms of observing conditions, disk morphologies, and contrasts.

- RDI seems to outperform ARDI in dataset #4D and #4E (Figs. D.4D, D.4E) with a contrast of 10^{-3} , as well as in dataset #3A (Fig. D.3A) with a contrast of 10^{-4} . We note that RDI excels primarily in the case of face-on disks and favorable observing conditions. For instance, datasets #4D and #4E are almost perfectly recovered by RDI, suggesting that the angular diversity does not provide significantly more relevant information in these cases, as the information contained in the reference frames is sufficient on its own.

In terms of the restoration of the planet in disk #E, we observe that at a contrast of 10^{-3} , the planet is detected by all three strategies. However, ADI underperforms compared

to the two other strategies, and the flux is not correctly retrieved as indicated by the error plots in Figs. D.1E, D.2E, D.3E, and D.4E. However, at a contrast of 10^{-4} , the planet is systematically and clearly detected with ADI, while with RDI no clear point-like feature stands out at the planet location. In fact, at a higher contrast, RDI struggles to capture the planet, while ADI has difficulty capturing the disk structure. ARDI captures the planet and the disk structure for cubes #2, #3, and #4 but fails to recover both the disk and the planet for cube #1. At lower contrast, neither ARDI nor ADI using IPCA seems to stand out for the task of planet detection only. Moreover, note that point-like artifacts can

be observed in other synthetic disk estimates where no planets were injected, as in cube #3 disk #B at contrast 10^{-4} (Fig. D.1E).

Working with a comprehensive battery of tests that require automation and standardization has the downside of potentially missing the best results that a particular algorithm could provide in specific cases. Indeed, while the behavior of these different strategies across a broader variety of scenarios remain relevant in most cases, a more specific optimization of the IPCA parameters might have significantly improved the quality of some of the disks, and potentially enabled more detections (e.g., by starting at a higher rank and reducing iteration per rank when processing the faintest disks of our test gallery).

In our second series of tests, we used the IPCA algorithm with three different selections of reference stars with different levels of quality according to PCC (see Fig. 5). The dataset #4 (very stable) seems more sensitive to the selection of reference frames, while datasets showing more speckle variability (#3, #2, #1) seem less sensitive. The mean PCC calculated for a non-optimal selection is significantly lower for cube #4 (mean PCC $\mu < 0.6$) compared to other cubes (mean PCC $\mu \sim 0.8-0.5$). As a result, the estimates produced when using these suboptimal references significantly under-performed, with SSIM hardly reaching 0.8 in most cases, while for other cubes, suboptimal references could reach similar SSIM as the optimal and semi-optimal selections (see Fig. 7). Most interestingly, the three different selections achieved very similar results in terms of SSIM in many cases. In more than 40% of cases, nonoptimal references achieved a higher SSIM than the optimal selection. The cases where the references contained 25% of optimal references are very close to an optimal selection, even for a contrast of 10^{-5} . This suggests that selecting high-quality references, as assessed via PCC, has limited impact on the results in the final post-processed images using IPCA with ARDI. Instead, the outcomes appear to be more sensitive to the stability of the dataset.

5. Demonstration on protoplanetary disk datasets

To test our algorithm on real observations, we used the pre-processed datasets on protoplanetary disks from Ren et al. (2023). The sample consists of 48 datasets including 29 unique different young stars hosting protoplanetary disks with known substructures previously observed in scattered light. Among them, four were reported to have exoplanet candidates (or claims thereof) from high-contrast imaging: HD 100546 (Quanz et al. 2013, 2015; Currie et al. 2014), HD 169142 (Biller et al. 2014; Reggiani et al. 2014; Gratton et al. 2019; Hammond et al. 2023), LkCa 15 (Kraus & Ireland 2012; Sallum et al. 2015), and MWC 758 (Reggiani et al. 2018; Wagner et al. 2019, 2023). Additionally, localized deviations from Keplerian rotation possibly associated with the presence of planets were reported in HD 163296 (Izquierdo et al. 2022; Pinte et al. 2020, 2018) and HD 97048 (Pinte et al. 2019). All observations were made in the broad-band K_s filter using the IRDIS camera of VLT/SPHERE. These disks were previously presented by Ren et al. (2023) in both polarized light using the IRDAP (van Holstein et al. 2020) pipeline for polarimetric data processing, and in total intensity using RDI through the DI-sNMF algorithm. The references were acquired via the “star-hopping” mode (Wahhaj et al. 2021), which means that the telescope is regularly pointed at a nearby reference star during data acquisition, efficiently alternating science and reference observations with minimal overhead. This observing strategy ensures maximal similarity between the respective PSFs.

We present our reprocessing of the Ren et al. (2023) data with IPCA using the ARDI strategy in Fig. 9. For most datasets, a reference frame ratio (r/n) of one was employed, utilizing the dedicated reference star observed in star-hopping mode. The default parametrization method involves ten iterations per rank, starting from rank one to rank ten. However, for some datasets, the parametrization was individually optimized. In cases where the signal was faint, the starting rank was set to a higher value, and the number of iterations per rank was smaller. In situations where the angular diversity was too small, we increased the reference frame ratio. When reference stars were unsatisfactory (e.g., for HD 163296), we employed a combination of the most correlated references available from the references acquired in star-hopping mode from this sample of protoplanetary disk observations. In these cases, we also utilized a higher reference frame ratio (r/n) and used a higher starting rank. We manually selected the best image from the estimates generated for these parameters. IPCA parameters and the selected frame for each dataset can be found in Appendix E. Table E.1 also provides the mean and standard deviation of the PCC between the science data and the references, following the same procedure detailed in Sect. 4.1.

Among the 48 datasets, Ren et al. (2023) showcased high-quality results for 23 datasets⁴ including 15 unique systems using DI-sNMF, selectively discarding instances of non-detections and very low quality images. In the context of this article, we present all results processed with IPCA-ARDI, explicitly including the non-high quality images. Among these images, 33 datasets including 19 unique protoplanetary disks are visually recovered, exhibiting a range of image qualities from excellent (e.g., V351 Ori, HD 36112, or V1366 Ori) to bare detection (e.g., GM Aur or HD 163296). This variability is influenced by factors such as the quality of the references, the brightness of the disk, and the degree of angular diversity. Comparing the results, it is noteworthy that for the 18 datasets analyzed with both DI-sNMF (RDI) and IPCA (ARDI) detections, similar disk structures are observed.

5.1. Inspecting protoplanet claims in individual systems

Among the systems presented above, four have unconfirmed protoplanet claims from direct imaging observations (including two such claims for MWC 758), and two exhibit kinematic signatures that could indicate the presence of planets (including two candidates for HD 163296). The different images showing the position of the candidate(s) within their protoplanetary disk are provided in Figs. 10 and 11. We propose to review the different candidates in the following paragraphs.

HD 100546. A point-like feature, interpreted as a potential gas giant planet at 52 au from its star, was observed in 2013 with the NACO instrument installed at the VLT in the L' and M' bands (Quanz et al. 2013, 2015; Currie et al. 2014, 2015, 2017). These observations were processed using the ADI strategy with the PCA and LOCI algorithms. The potential candidate is located within the broad-disk signal, which was not recovered in the processed images shown in these publications. Indeed, the disk was filtered out due to the now well-known self-subtraction and over-subtraction effects (Juillard et al. 2023, 2022; Milli et al. 2012), which worsen when using a large number of principal components. Considering the

⁴ Only 18 datasets processed with DI-sNMF are presented in Ren et al. (2023), but among the non-presented ones, five more can be considered as clear detections (priv. comm.).

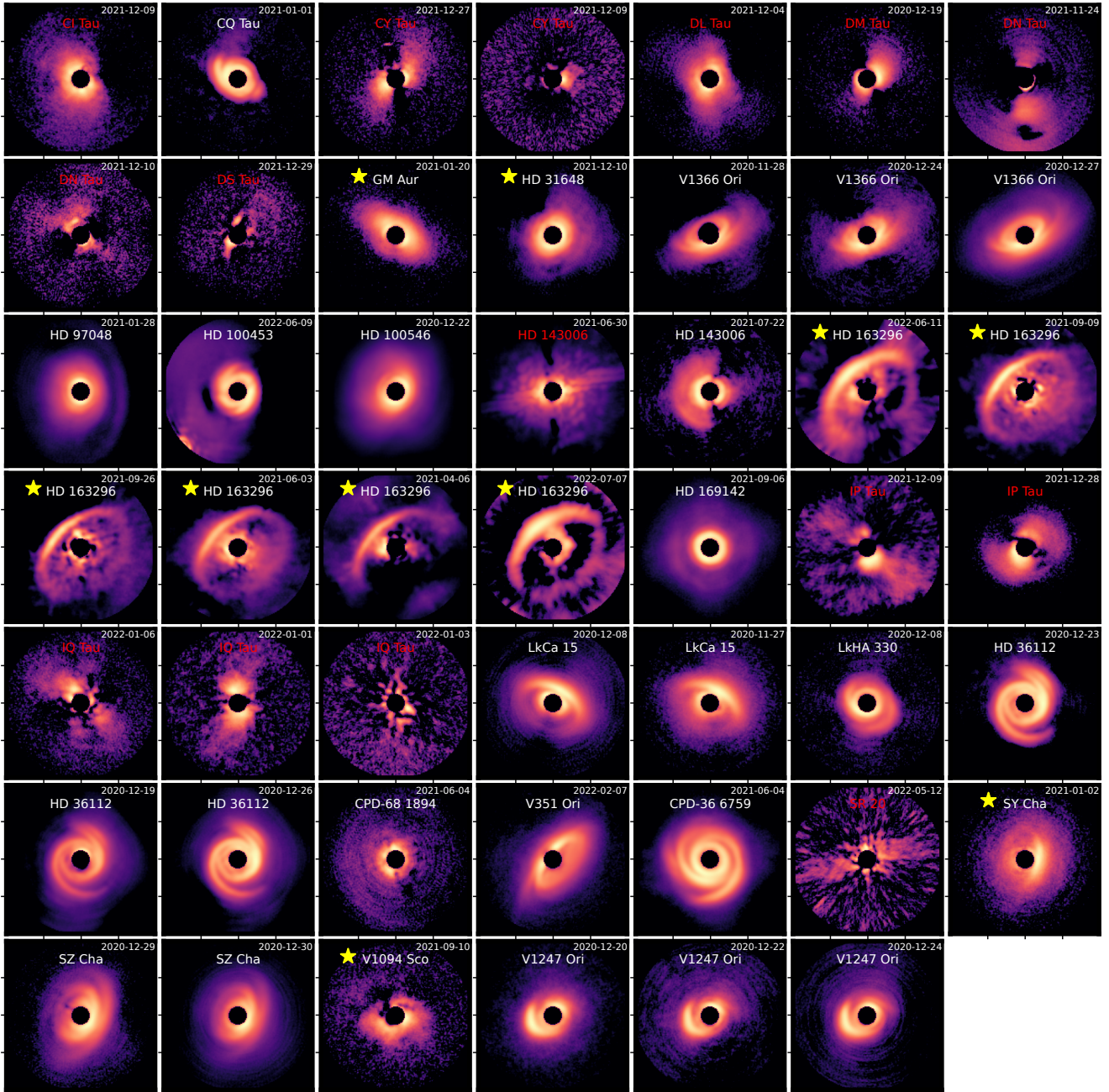


Fig. 9. Gallery of protoplanetary disks datasets, originally described by [Ren et al. \(2023\)](#), and re-processed here with IPCA-ARDI. The images have dimensions of 162×162 pixels ($\sim 2''$) with logarithmic scale color bars. Images marked by a yellow star are those that were not presented in [Ren et al. \(2023\)](#) due to the low quality of the produced disk estimation, but that we still consider as a detection in our IPCA-ARDI results. Images marked with a red name are those that are not considered clear detection with IPCA-ARDI.

expected Keplerian motion on a circular orbit between 2013 and 2020, the candidate would have moved by approximately 8° in position angle, which corresponds to a 6.5-pixel shift in the image (see [Fig. 10](#)). No companion is recovered in our images at those positions. In the ARDI images, as well as the DI-sNMF results, an extended spiral-like structure stands out rather than a point-like feature. Due to the presence of strong disk signal dominating in the area where the candidate is located, we cannot provide robust detection limits in the Ks filter for this candidate.

LkCa 15. Potential candidates were found using observations in L' and K bands taken with the near-infrared imager NIRC2 of the Keck II telescope in 2009 and 2010

([Kraus & Ireland 2012](#)), followed by several observations between 2014 and 2015 using the Large Binocular Telescope in L' , Ks , and $H\alpha$ filters ([Sallum et al. 2015](#)). The detections were interpreted as potential newly formed gas giant planets surrounded by dusty material. However, it was mentioned that the observed signal could represent a more complex structure rather than a single point source ([Kraus & Ireland 2012](#)). In fact, candidates buried within the disk pose a significant challenge for detection. Additionally, limitations arise due to the high variability of the speckle field in that specific image region and due to the small working angle. Recent observations of the disk support the hypothesis of a filtered disk signal, as there are no discernible point-like features (PLF) in the disk

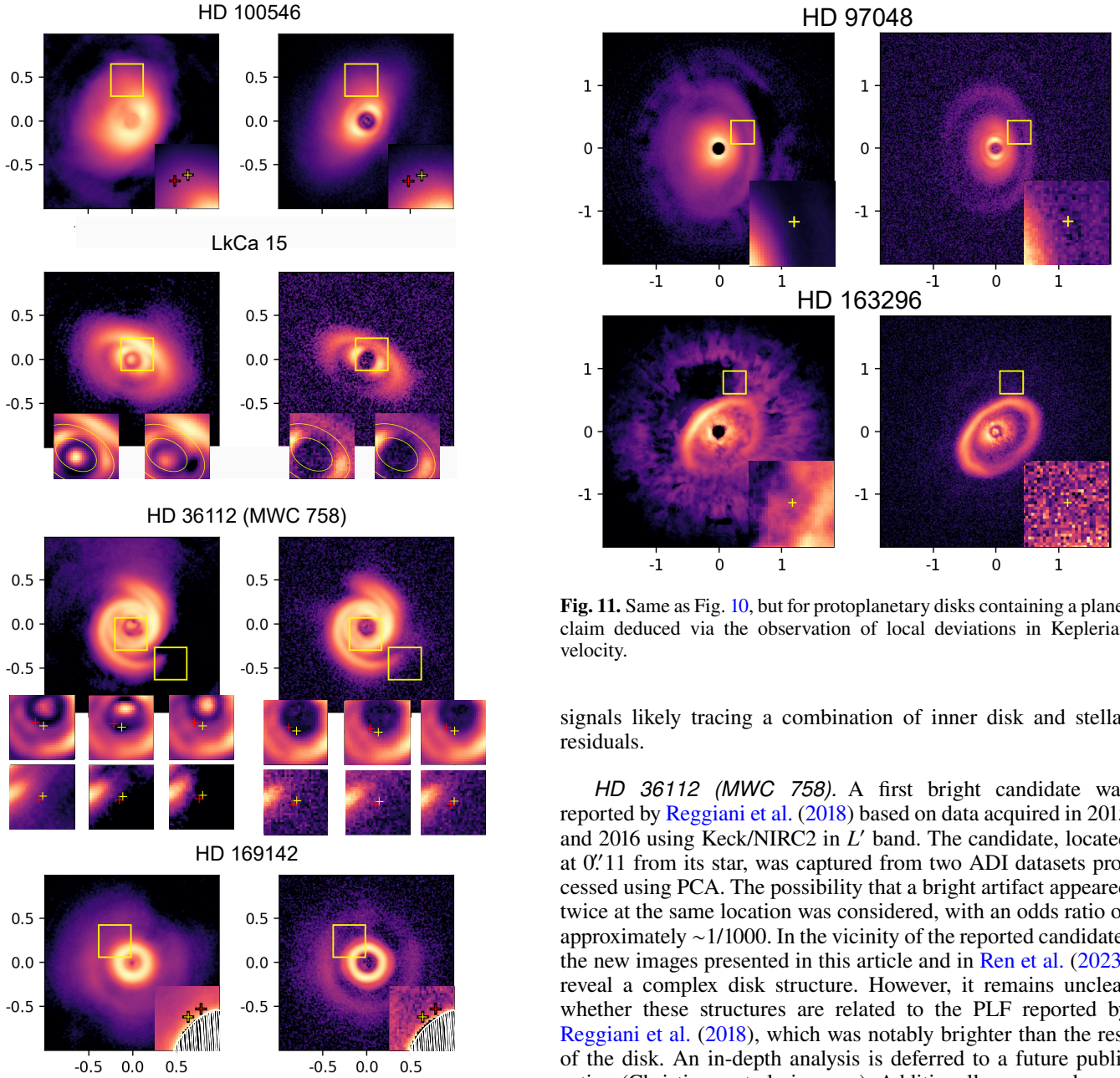


Fig. 10. ARDI images processed with IPCA (left) and polarized images processed with IRDAP (right) for all protoplanetary disks containing a planet claimed via direct imaging (in logarithmic scales). An approximate position of the candidate is marked with a yellow square on the full image, and a zoom on the candidate’s position is shown at the bottom (30×30 pixels, linear scales). In each sub-image, the yellow cross marks the candidate’s approximate position in the discovery paper, and the red cross the expected position on the date of acquisition of the SPHERE images presented in this work, considering a circular orbit. For LkCa 15, given the multiple claimed point-like features at varying positions, we only show two yellow orbits at 10 and 30 au defining the region where the candidate(s) were proposed.

(Currie et al. 2019). The images provided by both DI-sNMF and IPCA-ARDI corroborate this statement, since no PLF is observable. The second row in Fig. 10 presents the approximate positions where the diverse candidates should be located, displaying two bound orbits at 10 and 30 au from the star. Inferring detection limits for these candidates would be highly unreliable, as they are located within an area of strong residual

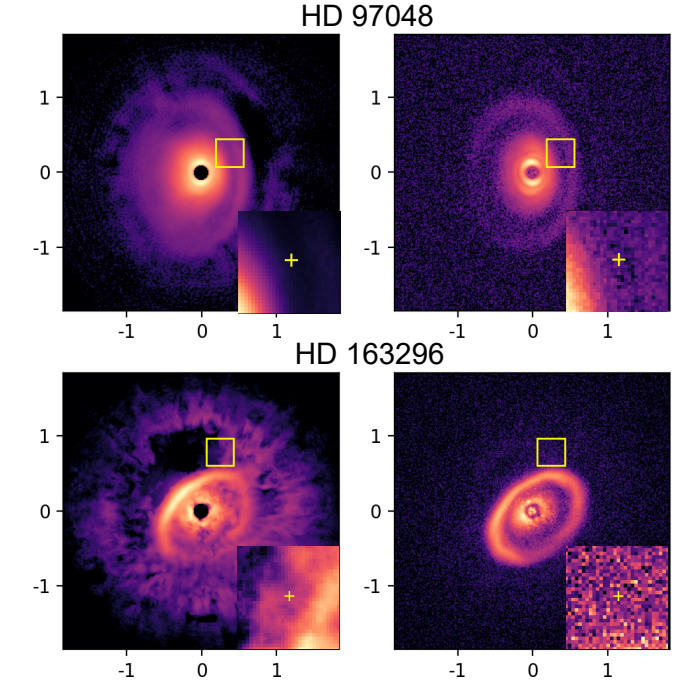


Fig. 11. Same as Fig. 10, but for protoplanetary disks containing a planet claim deduced via the observation of local deviations in Keplerian velocity.

signals likely tracing a combination of inner disk and stellar residuals.

HD 36112 (MWC 758). A first bright candidate was reported by Reggiani et al. (2018) based on data acquired in 2015 and 2016 using Keck/NIRC2 in L' band. The candidate, located at $0''.11$ from its star, was captured from two ADI datasets processed using PCA. The possibility that a bright artifact appeared twice at the same location was considered, with an odds ratio of approximately $\sim 1/1000$. In the vicinity of the reported candidate, the new images presented in this article and in Ren et al. (2023) reveal a complex disk structure. However, it remains unclear whether these structures are related to the PLF reported by Reggiani et al. (2018), which was notably brighter than the rest of the disk. An in-depth analysis is deferred to a future publication (Christiaens et al., in prep.). Additionally, a second, very red protoplanet candidate was claimed by Wagner et al. (2019, 2023) near the tail of the northern spiral using observations from 2016, 2017, 2018, and 2019 with LMIRCam mounted on the Large Binocular Telescope Interferometer. Further analysis of these two candidates, carried out by Boccaletti et al. (2021), did not re-detect the candidates proposed by either Reggiani et al. (2018) or Wagner et al. (2019), and did not find them consistent with being spiral-driving planets assuming linear density wave models. For these two candidates, the expected motion between the first observation and the SPHERE Ks-band observations (shown in Fig. 10), according to Keplerian motion on a circular orbit, would be respectively a 3.9-pixel shift and a 2.2-pixel shift. We do not re-detect them in our IPCA-ARDI images, nor did Ren et al. (2023) using DI-sNMF. We performed planet injection at the separation of the Wagner et al. (2019) candidate companion, and generated a contrast curve using PCA with ADI-only in Appendix F.1. Our analysis reveals a detection threshold of approximately 2.6×10^{-5} , which is consistent with the contrast limit established by Grady et al. (2013) at

this separation. According to the ATMO2020 model (Phillips et al. 2020) and assuming no extinction, our contrast threshold translates into a 1 Jupiter mass (M_{Jup}) sensitivity. This mass sensitivity does not match with the predicted mass of about 2–3 M_{Jup} inferred by Wagner et al. (2019). This tension could be explained by the usage of different planet evolution models, or possibly by extinction or by the presence of a circumplanetary disk (more details in Appendix F.1). Converting the contrast of 1×10^{-5} for the candidate in the L' band as presented in Wagner et al. (2019) into a mass using the ATMO2020 model for a consistent comparison, we find a mass of $0.5 M_{\text{Jup}}$, suggesting that our detection limit would not allow us to detect the candidate.

HD 169142. Reggiani et al. (2014) and Biller et al. (2014) independently observed a PLF in L' band using independent ADI datasets from VLT/NACO processed with PCA. Better images of the disk obtained with SPHERE suggested that this candidate at $0''.11$ separation may rather trace filtered signal from the inner ring of the disk (e.g., Ligi et al. 2019). Later, another candidate was proposed using observations in polarized light from 2015, along with two ASDI observations in 2017 and 2019 using the IRDIS and multiple IFS observations between 2015 and 2019 (Gratton et al. 2019; Hammond et al. 2023). However, detecting the candidate with ASDI required suppressing most of the disk signal. The observations were conducted in J and YJH bands and the position of the candidate between 2015 and 2019 appeared to follow a Keplerian motion. The signal was interpreted as a gap-clearing Jovian-mass protoplanet surrounded by a circumplanetary disk or envelope. Studies by Garg et al. (2021) and Law et al. (2023) using ALMA observations revealed several distinct chemical signatures that demonstrate a compelling link to ongoing giant planet formation. Furthermore, compact CO $J = 2-1$ and CO $J = 3-2$ emission counterparts were observed that coincide with the location of the candidate planet. Yet, no PLF appears in polarized light nor in total intensity for both the DI-sNMF and IPCA-ARDI of the K -band images presented here, at its predicted location at the epoch of the K_s observations considered in this work (2021) assuming Keplerian motion. In Appendix F.2, we attempt to infer the detectability of the candidate by injecting a fake planet at the same separation, using the estimated contrast in Hammond et al. (2023). In this test, we processed the dataset with IPCA-ARDI as well as plain PCA-ADI. In addition, we provide the contrast curve obtained via PCA-ADI in Fig. F.4. Our 5σ detection limit of 1×10^{-4} at the expected location of the companion translates into $2 M_{\text{Jup}}$ according to the ATMO2020 model, assuming no extinction. In comparison, Hammond et al. (2023) report a YJH contrast of $\sim 1.5 \times 10^{-5}$. Nevertheless, the detection of polarized intensity collocated with the protoplanet suggests the measured YJH light is dominated by scattered light, making the mass conversion using atmospheric models most probably irrelevant. Assuming similar expected contrast in K_s band as in YJH , our sensitivity limits do not allow us to confirm or refute the candidate.

HD 97048. A kinematic signature was detected in the gap of the disk around HD 97048 using high-spectral and spatial-resolution ALMA observations (Pinte et al. 2019). This signature could be explained by the presence of a superjovian planet in the disk, possibly at 130 au from the star. At this separation, the motion of the candidate between the date of the first detection and our dataset would have been less than a 1-pixel shift. Unfortunately, we were unable to directly observe any point-like feature at this location in total intensity. We injected fake planets at the same separation as the candidate to assess the contrast

sensitivity of our reduction at its separation, and found that a candidate below a contrast of 1×10^{-5} in K_s -band would not have been detected. Based on our contrast limit, the ATMO2020 models imply an upper mass limit of $0.5 M_J$ (see Appendix F.3). The discrepancy with the kinematically inferred mass may suggest the presence of extinction or the inadequacy of the ATMO2020 models for embedded protoplanets.

HD 163296. Using CO line observations, Pinte et al. (2018) initially reported a local deviation from Keplerian velocity, suggesting the presence of a massive body at approximately 260 au from the star. Subsequently, Izquierdo et al. (2022) reobserved the first candidate proposed by Pinte et al. (2018), along with a second localized deviation from Keplerian rotation, potentially associated with the presence of a giant planet at 94 au. The first candidate proposed by Pinte et al. (2018) and Izquierdo et al. (2022) at 260 au does not fit within the field of view of the images presented in Figs. 11 and 9. At these separations, the motion of the candidates between the date of the first detection and the new dataset is negligible. We show another post-processed image for the same dataset, with an uncropped field-of-view in Appendix F.4 as well as contrast curve with PCA-ADI. No PLF is observed at the locations of the two proposed sources in these direct imaging observations with both IPCA-ARDI and PCA-ADI, despite the estimated sensitivity reaching down to 3×10^{-4} in contrast (i.e., mass limit of $1 M_J$) for the candidate at 260 au, and 6×10^{-6} in contrast (i.e., mass limit of $4 M_J$) at 94 au (see details in Appendix F.4).

The direct observation of a protoplanet candidate within a protoplanetary disk in direct imaging can necessitate the use of aggressive post-processing techniques that lead to significant disk flux loss and alterations to the apparent morphology of the disk in order to reveal faint point-like sources. Many of the directly imaged protoplanet claims presented in this sample rely on poor-quality images, exhibiting signs of self- and over-subtraction. It remains unclear whether a PLF results from filtered disk signal, from artifacts inefficiently corrected by post-processing, or from an actual planet, especially for candidates buried within the disk signal. This ambiguity makes it challenging to both confirm and refute these candidates. Indeed, the absence of the proposed candidates cannot definitively discard the claims, as the wavelengths might be inappropriate, the contrast could be too high, or the candidate might be buried within the disk. Being able to capture high-quality images of the disk using techniques such as ARDI with IPCA is still relevant to prevent false candidates due to filtered disk signals.

6. Conclusion

We explored the possibility of combining the RDI and ADI observing strategies by adapting the IPCA algorithm originally designed for ADI datasets. Our main goal was to assess the extent to which combining information from ADI and RDI could enhance disk recovery compared to using each strategy independently, while employing IPCA-based methods. Our analysis encompassed systematic testing on a diverse set of 60 synthetic datasets, covering a range of observing conditions and disk morphologies. We tested our methods with both optimally chosen reference stars and with a random mix of reference stars to assess the impact of reference star quality on the estimated disk.

Our results revealed that ARDI consistently improves the recovery of extended signals compared to using these techniques individually. This improvement holds true across various

scenarios, encompassing different observing strategies and disk morphologies. A significant visual improvement from ARDI occurs when a planet is hosted within a disk. At lower contrast levels and for the more variable speckle pattern, RDI struggled to capture the planet, while ADI had difficulty capturing the disk structure, particularly due to the rotation-invariant flux, which is an inherent ambiguity of the observing strategy. The combination of these strategies emerges as a reliable approach to mitigate the limitations of individual observing strategies, addressing challenges that would otherwise be difficult to overcome.

We also explored the problem of parameters' optimization for IPCA algorithms. While a good choice of parameters is important, we found that multiple sets of parameters lead to similar results. Nevertheless, we identified different scenarios and matched them with specific parametrizations to provide optimal results in each case. For the brightest disks with a stable speckle field, it is better to start the iterative process at rank one and have multiple iterations per rank before increasing the rank (e.g., CPD-68 1894 or HD 36112 data presented in Sect. 5). This ensures that most of the self-subtraction effects have been mitigated before loosening the model of the speckle field by increasing the rank. Conversely, when the disk is faint relative to the speckle, it is necessary to start at a high rank and limit the number of iterations per rank to avoid the bright artifacts stemming from non-captured speckle field components from being propagated in the subsequent iterations (e.g., V1094 Sco or GM Aur datasets). Regarding the convergence of the iterative process, we acknowledge the difficulty in building a reliable stopping criterion, which is a common problem in computer science (Nocedal et al. 2002). It is particularly challenging in our case of application, as determining where to stop the iterative process involves choosing a good compromise between noise subtraction (non-propagation) as well as the filtering and recovery of extended signals (convergence), and also because different parts of the disk converge at different rates. Nevertheless, observing the convergence curves remains an interesting indicator of the optimal iteration to be reached.

We applied our IPCA algorithm with ARDI to reprocess a sample of 48 datasets that includes 29 different young stars surrounded by protoplanetary disks exhibiting known substructures, which were initially introduced by Ren et al. (2023) in both polarized light and total intensity, using RDI through the DI-sNMF algorithm. All disks from the 29 young protoplanetary disk gallery published by Ren et al. (2023) that were successfully recovered by DI-sNMF (RDI) were also successfully retrieved with IPCA-ARDI. Under optimal conditions with ideal reference stars, IPCA demonstrates capabilities comparable to those of DI-sNMF in preventing over-subtraction. Furthermore, the usage of ARDI enables the detection of fainter disks such as V1094 Sco.

While a few planet candidates were reported via direct imaging among the sampled young stars, we did not reobserve any of these candidates. This is at least partly due to limitations in the achieved sensitivity in the considered filter (K_s), which renders the detection of these companions challenging or unfeasible. While confirmed detections of protoplanets within disks remain rare, with only a select number being reported in the literature (e.g., Keppler et al. 2018; Haffert et al. 2019), studying young systems in their formative stages embedding protoplanetary disks housing nascent planets is crucial for understanding the intricate process of planetary system formation. The application of IPCA-ARDI holds promise in providing valuable insights by preserving both the structures of the disk and clear images of planets.

Acknowledgements. This project has received funding from the European Research Council (ERC) under the European Union's Horizon 2020 research and innovation program (grant agreement No. 819155), and from the Belgian Fonds de la Recherche Scientifique – FNRS. This work has made use of the SPHERE Data Centre, jointly operated by OSUG/IPAG (Grenoble), PYTHEAS/LAM/CESAM (Marseille), CA/Lagrange (Nice), Observatoire de Paris/LESIA (Paris), and Observatoire de Lyon (Galicher et al. 2018; Delorme et al. 2017). We are very grateful to Bin Ren and Myriam Benisty for sharing their data used in this work, and to Philippe Delorme, and Mariam Sabalbal for their help in selecting, providing, and pre-processing the SPHERE data cubes used for the tests. SS acknowledges funding from the European Research Council (ERC) under the European Union's Horizon 2020 research and innovation program (COBREX; grant agreement no. 885593).

References

- Amara, A., & Quanz, S. P. 2012, *MNRAS*, **427**, 948
- Baraffe, I., Chabrier, G., Barman, T. S., Allard, F., & Hauschildt, P. H. 2003, *A&A*, **402**, 701
- Baraffe, I., Homeier, D., Allard, F., & Chabrier, G. 2015, *A&A*, **577**, A42
- Beuzit, J. L., Vigan, A., Mouillet, D., et al. 2019, *A&A*, **631**, A155
- Biller, B. A., Males, J., Rodigas, T., et al. 2014, *ApJ*, **792**, L22
- Boccaletti, A., Pantin, E., Ménard, F., et al. 2021, *A&A*, **652**, L8
- Carter, A. L., Hinkley, S., Kammerer, J., et al. 2023, *ApJ*, **951**, L20
- Christiaens, V., Casassus, S., Absil, O., et al. 2019, *MNRAS*, **486**, 5819
- Christiaens, V., Gonzalez, C., Farkas, R., et al. 2023, *J. Open Source Softw.*, **8**, 4774
- Currie, T., Muto, T., Kudo, T., et al. 2014, *ApJ*, **796**, L30
- Currie, T., Cloutier, R., Brittain, S., et al. 2015, *ApJ*, **814**, L27
- Currie, T., Brittain, S., Grady, C. A., Kenyon, S. J., & Muto, T. 2017, *RNAAS*, **1**, 40
- Currie, T., Marois, C., Cieza, L., et al. 2019, *ApJ*, **877**, L3
- Delorme, P., Meunier, N., Albert, D., et al. 2017, in *SF2A-2017: Proceedings of the Annual meeting of the French Society of Astronomy and Astrophysics*, eds. C. Reylé, P. Di Matteo, F. Herpin, E. Lagarde, A. Lançon, Z. Meliani, & F. Royer, Di
- Dohlen, K., Langlois, M., Saisse, M., et al. 2008, *SPIE Conf. Ser.*, **7014**, 70143L
- Esposito, T. M., Fitzgerald, M. P., Graham, J. R., & Kalas, P. 2014, *ApJ*, **780**, 25
- Flasseur, O., Thé, S., Denis, L., Thiébaud, E., & Langlois, M. 2021, *A&A*, **651**, A62
- Flasseur, O., Thé, S., Denis, L., Thiébaud, E., & Langlois, M. 2022, *SPIE Conf. Ser.*, **12185**, 121853U
- Galicher, R., Boccaletti, A., Mesa, D., et al. 2018, *A&A*, **615**, A92
- Garg, H., Pinte, C., Christiaens, V., et al. 2021, *MNRAS*, **504**, 782
- Gomez Gonzalez, C. A., Wertz, O., Absil, O., et al. 2017, *AJ*, **154**, 7
- Grady, C. A., Devine, D., Woodgate, B., et al. 2000, *ApJ*, **544**, 895
- Grady, C. A., Muto, T., Hashimoto, J., et al. 2013, *ApJ*, **762**, 48
- Gratton, R., Ligi, R., Sissa, E., et al. 2019, *A&A*, **623**, A140
- Haffert, S. Y., Bohn, A. J., de Boer, J., et al. 2019, *Nat. Astron.*, **3**, 749
- Hammond, I., Christiaens, V., Price, D. J., et al. 2023, *MNRAS*, **522**, L51
- Izquierdo, A. F., Facchini, S., Rosotti, G. P., van Dishoeck, E. F., & Testi, L. 2022, *ApJ*, **928**, 2
- Juillard, S., Christiaens, V., & Absil, O. 2022, *A&A*, **668**, A125
- Juillard, S., Christiaens, V., & Absil, O. 2023, *A&A*, **679**, A52
- Keppler, M., Benisty, M., Müller, A., et al. 2018, *A&A*, **617**, A44
- Kraus, A. L., & Ireland, M. J. 2012, *ApJ*, **745**, 5
- Kuhn, J. R., Potter, D., & Parise, B. 2001, *ApJ*, **553**, L189
- Lafrenière, D., Marois, C., Doyon, R., Nadeau, D., & Artigau, É. 2007, *ApJ*, **660**, 770
- Lafrenière, D., Marois, C., Doyon, R., & Barman, T. 2009, *ApJ*, **694**, L148
- Law, C. J., Booth, A. S., & Öberg, K. I. 2023, *ApJ*, **952**, L19
- Lawson, K., Currie, T., Wisniewski, J. P., et al. 2022, *ApJ*, **935**, L25
- Ligi, R., Dorn, C., Crida, A., et al. 2019, *A&A*, **631**, A92
- Linder, E. F., Mordasini, C., Mollière, P., et al. 2019, *A&A*, **623**, A85
- Marois, C., Doyon, R., Racine, R., & Nadeau, D. 2000, *SPIE Conf. Ser.*, **4008**, 788
- Marois, C., Lafrenière, D., Doyon, R., Macintosh, B., & Nadeau, D. 2006, *ApJ*, **641**, 556
- Mawet, D., Serabyn, E., Liewer, K., et al. 2009, *Opt. Express*, **17**, 1902
- Milli, J., Mouillet, D., Lagrange, A. M., et al. 2012, *A&A*, **545**, A111
- Nocedal, J., Sartenaer, A., & Zhu, C. 2002, *Computat. Optim. Appl.*, **22**, 5
- Pairat, B., Cantalloube, F., & Jacques, L. 2021, *MNRAS*, **503**, 3724
- Pavlov, A., Feldt, M., & Henning, T. 2008, *Astronomical Data Analysis Software and Systems XVII*, eds. R. W. Argyle, P. S. Bunclark, & J. R. Lewis, *ASP Conf. Ser.*, **394**, 581
- Phillips, M. W., Tremblin, P., Baraffe, I., et al. 2020, *A&A*, **637**, A38

- Pinte, C., Price, D. J., Ménard, F., et al. 2018, [ApJ](#), **860**, L13
- Pinte, C., van der Plas, G., Ménard, F., et al. 2019, [Nat. Astron.](#), **3**, 1109
- Pinte, C., Price, D. J., Ménard, F., et al. 2020, [ApJ](#), **890**, L9
- Pueyo, L. 2016, [ApJ](#), **824**, 117
- Quanz, S. P., Amara, A., Meyer, M. R., et al. 2013, [ApJ](#), **766**, L1
- Quanz, S. P., Amara, A., Meyer, M. R., et al. 2015, [ApJ](#), **807**, 64
- Racine, R., Walker, G. A. H., Nadeau, D., Doyon, R., & Marois, C. 1999, [PASP](#), **111**, 587
- Reggiani, M., Quanz, S. P., Meyer, M. R., et al. 2014, [ApJ](#), **792**, L23
- Reggiani, M., Christiaens, V., Absil, O., et al. 2018, [A&A](#), **611**, A74
- Ren, B. B. 2023, [A&A](#), **679**, A18
- Ren, B., Pueyo, L., Zhu, G. B., Debes, J., & Duchêne, G. 2018, [ApJ](#), **852**, 104
- Ren, B. B., Benisty, M., Ginski, C., et al. 2023, [A&A](#), **680**, A114
- Ruane, G., Ngo, H., Mawet, D., et al. 2019, [AJ](#), **157**, 118
- Sallum, S., Follette, K. B., Eisner, J. A., et al. 2015, [Nature](#), **527**, 342
- Schneider, G., Weinberger, A. J., Becklin, E. E., Debes, J. H., & Smith, B. A. 2009, [AJ](#), **137**, 53
- Soummer, R., Pueyo, L., & Larkin, J. 2012, [ApJ](#), **755**, L28
- Stapper, L. M., & Ginski, C. 2022, [A&A](#), **668**, A50
- van Holstein, R. G., Girard, J. H., de Boer, J., et al. 2020, [A&A](#), **633**, A64
- Vigan, A., Moutou, C., Langlois, M., et al. 2010, [MNRAS](#), **407**, 71
- Wagner, K., Stone, J. M., Spalding, E., et al. 2019, [ApJ](#), **882**, 20
- Wagner, K., Stone, J., Skemer, A., et al. 2023, [Nat. Astron.](#), **7**, 1208
- Wahhaj, Z., Liu, M. C., Biller, B. A., et al. 2013, [ApJ](#), **779**, 80
- Wahhaj, Z., Milli, J., Romero, C., et al. 2021, [A&A](#), **648**, A26
- Wallack, N. L., Ruffio, J.-B., Ruane, G., Ren, B. B., Xuan, J. W. et al. 2024, [AJ](#) submitted [arXiv:2408.04048]
- Wang, Z., Bovik, A., Sheikh, H., & Simoncelli, E. 2004, [IEEE Trans. Image Process.](#), **13**, 600
- Xie, C., Choquet, E., Vigan, A., et al. 2022, [A&A](#), **666**, A32
- Zurlo, A., Vigan, A., Mesa, D., et al. 2014, [A&A](#), **572**, A85

Appendix A: Local convergence of IPCA-ARDI

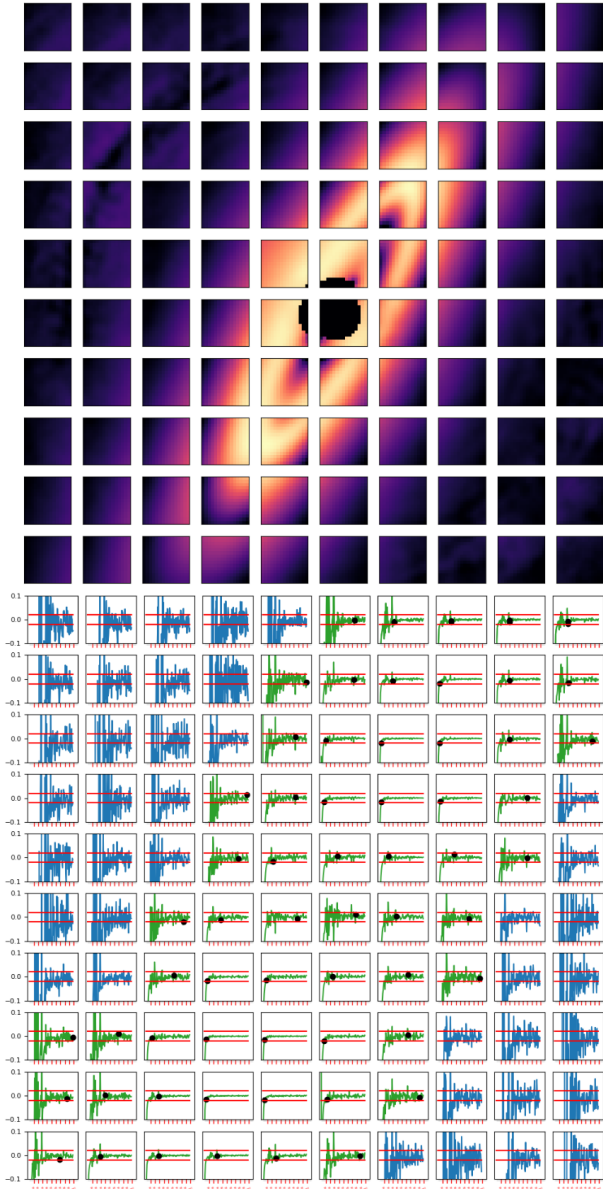


Fig. A.1. Evolution of the IPCA with ARDI estimate through the iterations. The dataset used is the same as in Fig. 3. The first grid of plots (top) displays the final estimation segmented into patches of size 15×15 pixels. The second grid of plots (bottom) corresponds to the gradient computed within these patches according to Eq. 4. The red lines indicate the g_{tol} threshold at $\pm 10^{-2}$. If the threshold is reached, the curves are displayed in green, and a black dot indicates at which iteration the signal has converged. Otherwise, the curves are displayed in blue.

In this appendix, we study the evolution of the integrated flux into different patches of the image through the IPCA iterative process when using ARDI. Figure A.1 shows that the brightest part of the disk converges faster than the rest of the image. We can also observe that the parts of the image where no signal is present do not seem to converge. It is important to note that the observed non-convergence on the outer parts of the image where there is no disk signal relates to the expression of the gradient. Indeed, from Eq. (4) it is clear that for an estimate d_i that tends toward zero, the quantity $\frac{d_i - d_{i+1}}{d_i}$ becomes more sensitive to minor variations.

Appendix B: Comparison of the SSIM score obtained by each algorithm

This appendix presents the results of the performance comparison for the three considered strategies (RDI, ADI, ARDI) using the IPCA algorithm, for each test dataset, using the structural similarity index measure (SSIM, Wang et al. 2004). For each of the three 4×5 -sized tables in Fig. B.1, each cell represents a different test dataset (see Sect. 4 and Juillard et al. (2023) for more details about the test datasets). In each cell, three bar plots represent the SSIM obtained between the ground truth and the final disk image obtained with each strategy using the IPCA algorithm. White cells indicate that no strategy detected the disk.

Appendix C: Comparison of results for different metrics

This appendix presents the performance of RDI, ADI, and ARDI, using the IPCA algorithm, for each test dataset according to four different metrics (see Fig. C.1: the Spearman rank correlation coefficient, the Pearson correlation coefficient, the Euclidean distance, and the sum of absolute differences (SAD)). A table using a fifth metric, the structural similarity index measure (SSIM, Wang et al. 2004), was already displayed in Fig. 6. In all these figures, the color of the cell indicates which algorithm performed the best according to the respective metrics. White cells mean that no strategy detected the disk. White letters on a cell indicate which strategy did not detect the disk (R \rightarrow RDI, A \rightarrow ADI, Ar \rightarrow ARDI). Black stars indicate that all five metrics (Spearman, SSIM, Euclidean, Spearman, Pearson, and SAD) agree on the strategy which achieved the best disk estimation. Overall, it can be observed that the various metrics yield relatively similar results. ARDI-IPCA is more frequently chosen as the best reconstruction by all metrics, with an average selection rate of 67% across all metrics (considering detection only). Among the tests where ARDI performed better, 45% of them are unambiguously better as all metrics agreed. ADI tends to perform better for contrast 10^{-4} and for the sharpest disks (A, B) where the deformations due to flux invariant to the rotation are minimal. RDI seems more appropriate for low contrast (10^{-3}), stable datasets with well correlated references (#3) and for the more face-on disks (C, D, E).

Appendix D: Results of disk estimations

This appendix presents the disk estimations obtained with RDI, ADI, and ARDI, using IPCA, for all 60 test datasets considered in this work (Figs. D.1A to D.4E). The figures of this appendix are available on Zenodo: <https://zenodo.org/records/11442350>.

Appendix E: IPCA parameters used to process the images presented in Fig. 9

In this appendix, Table E.1 provides a comprehensive summary of the parameters applied in the retrieval of the protoplanetary disk gallery using ARDI-IPCA, as shown in Sect. 5.

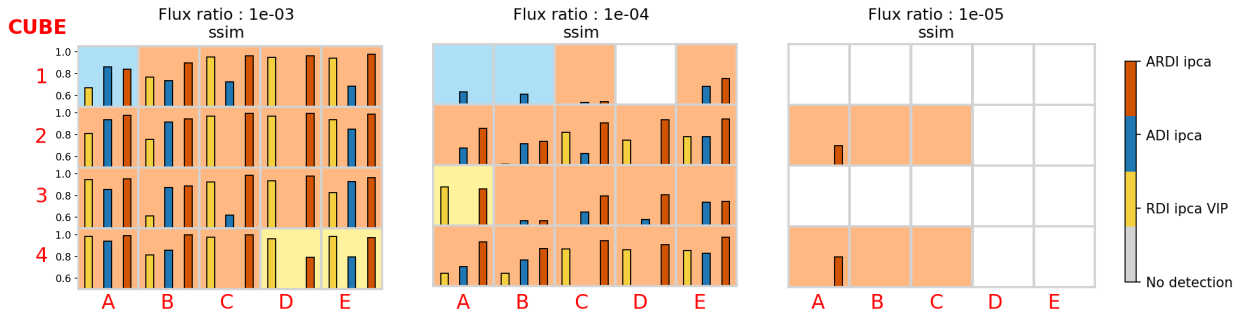


Fig. B.1. Results of our systematic tests comparing different strategies, namely RDI, ADI and ARDI with the IPCA algorithm. The organization of this table is identical to Fig. 6, except that the value of the SSIM of the estimation obtained by each algorithm compared to the ground truth is also represented within each cell by a colored histogram with RDI in yellow, ADI in blue, and ARDI in red.

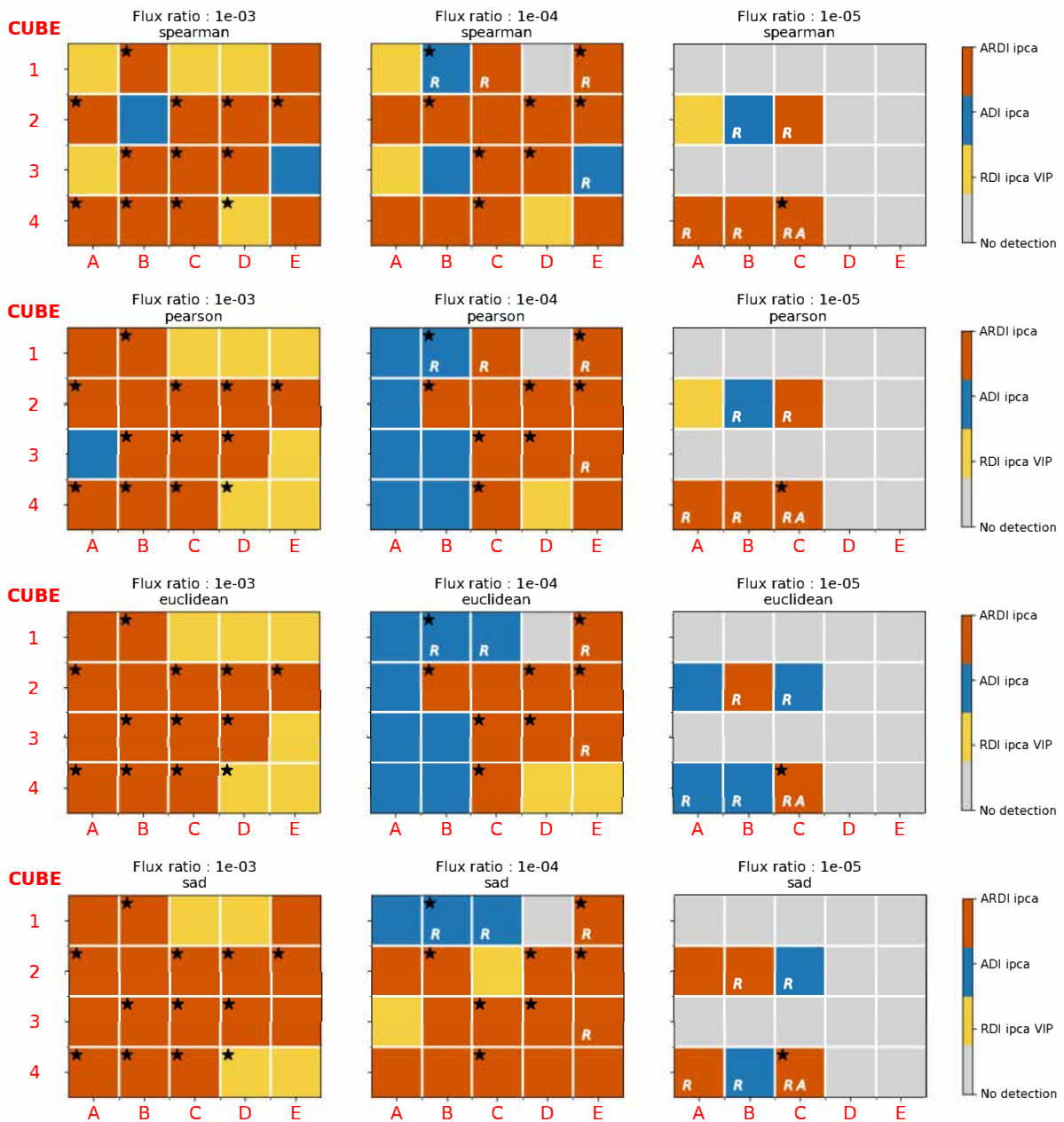


Fig. C.1. Same as Fig. 6 but for the four different metrics from top to bottom, respectively: Spearman's rank, Pearson correlation, Euclidean distance, and SAD.

Table E.1. IPCA parameters used to reprocess each dataset from Ren et al. (2023) in Sect. 5. For each dataset, this table provides the starting rank and the number of iterations per rank. Regarding the number of iterations per rank, “incr” stands for incremental and indicates that the number of iterations increases by one at each rank (e.g., one iteration is performed at the starting rank q , two at rank $q + 1$, and so on). Additionally, we also provide the number of iterations performed to obtain the final frame selected from visual assessment (“final iter.”), the rank that was reached at this iteration (“final rank”), and the mean and standard deviation of the PCC between science images and references.

Star name	Date	Start. rank	Iter. per rank	Final iter.	Final rank	Mean PCC	STD PCC
CI Tau	2021-12-09	10	incr	88	21	0.89	0.13
CQ Tau	2021-01-01	1	incr	85	8	0.83	0.02
CY Tau	2021-12-27	10	incr	10	13	0.93	0.04
CY Tau	2021-12-09	10	incr	10	13	0.94	0.02
DL Tau	2021-12-04	1	incr	10	1	0.96	0.01
DM Tau	2020-12-19	10	incr	105	23	0.80	0.05
DN Tau	2021-11-24	5	incr	6	7	0.79	0.25
DN Tau	2021-12-10	5	incr	6	7	0.88	0.05
DS Tau	2021-12-29	10	incr	6	12	0.85	0.07
GM Aur	2021-01-20	5	incr	8	9	0.80	0.04
HD 31648	2021-12-10	5	10	17	8	0.98	0.01
V1366 Ori	2020-11-28	1	10	41	4	0.70	0.36
V1366 Ori	2020-12-24	1	10	46	5	0.77	0.19
V1366 Ori	2020-12-27	1	10	41	4	0.74	0.18
HD 97048	2021-01-28	1	10	7	1	0.95	0.02
HD 100453	2022-06-09	1	10	85	8	0.45	0.06
HD 100546	2020-12-22	1	10	6	1	0.89	0.01
HD 143006	2021-06-30	1	10	85	8	0.75	0.11
HD 143006	2021-07-22	1	10	99	9	0.83	0.10
HD 163296	2022-06-11	15	10	145	28	0.87	0.06
HD 163296	2021-09-09	15	10	145	28	0.90	0.05
HD 163296	2021-09-26	15	10	145	28	0.92	0.03
HD 163296	2021-06-03	15	10	145	28	0.87	0.04
HD 163296	2021-04-06	15	10	145	28	0.88	0.03
HD 163296	2022-07-07	15	10	145	28	0.92	0.05
HD 169142	2021-09-06	1	10	31	3	0.79	0.26
IP Tau	2021-12-09	10	incr	105	23	0.94	0.02
IP Tau	2021-12-28	10	incr	105	23	0.91	0.03
IQ Tau	2022-01-06	10	incr	21	15	0.93	0.03
IQ Tau	2022-01-01	10	incr	91	22	0.94	0.03
IQ Tau	2022-01-03	10	incr	21	15	0.95	0.02
LkCa 15	2020-12-08	1	incr	12	2	0.76	0.03
LkCa 15	2020-11-27	1	incr	85	8	0.83	0.06
LkHA 330	2020-12-08	1	incr	54	10	0.89	0.01
HD 36112	2020-12-23	1	incr	22	2	0.93	0.01
HD 36112	2020-12-19	1	incr	23	3	0.93	0.01
HD 36112	2020-12-26	1	incr	51	5	0.92	0.01
CPD-68 1894	2021-06-04	1	incr	85	8	0.80	0.03
V351 Ori	2022-02-07	1	incr	85	8	0.78	0.02
CPD-36 6759	2021-06-04	1	incr	85	8	0.83	0.02
SR 20	2022-05-12	10	incr	22	15	0.98	0.01
SY Cha	2021-01-02	5	incr	151	19	0.72	0.06
SZ Cha	2020-12-29	1	incr	85	8	0.83	0.03
SZ Cha	2020-12-30	1	incr	12	2	0.83	0.03
V1094 Sco	2021-09-10	10	incr	65	20	0.80	0.05
V1247 Ori	2020-12-20	1	incr	85	8	0.75	0.04
V1247 Ori	2020-12-22	1	incr	11	1	0.75	0.06
V1247 Ori	2020-12-24	1	incr	18	2	0.72	0.06

Appendix F: Detectability of the protoplanet candidates

This appendix presents further analysis for the candidate protoplanets that are not mixed with coincident scattered light signal from the circumstellar disk, thus enabling the possibility of computing a reliable detection limit in terms of contrast, which can then be converted into a mass detection limit. The mass sensitivity is determined assuming no extinction, by converting our contrast limits using ATMO2020 evolutionary models (Phillips et al. 2020). The contrast curves are computed for a false positive probability equivalent to a 5σ detection in Gaussian statistics, including the t-student correction for small sample statistics, using routines from the Vortex Image Processing package⁵ (Gomez Gonzalez et al. 2017; Christiaens et al. 2023). The different subsections of this appendix detail our protoplanet detection limits and provide corresponding contrast curves and translation to masses where relevant.

F.1. HD 36112 (MWC 758)

This section presents a test of detectability for the protoplanet candidate proposed by Wagner et al. (2019) in the HD 36112 (MWC 758) system. The test involves injecting fake planets at the same separation as the candidate, using contrasts of 8×10^{-6} , 1×10^{-5} and 2×10^{-5} . Figure F.1 shows the results achieved with IPCA (left) on the dataset with the injected planet, in comparison to a result using a rank-3 PCA based on ADI only (right). Regarding the chosen rank for PCA, we processed the data for ranks 1 to 10 and identified the most effective rank for retrieving the fake companion through visual assessment. Additionally, in Fig. F.2, we present a 5σ contrast curve computed with rank-3 PCA using ADI only. We observe that only the simulated candidate at 2×10^{-5} is successfully recovered. The injected simulated planet is distinctly visible after post-processing the data with PCA and ADI-only, but this approach also introduces numerous point-like features attributable to the filtered disk. However, this same injected companion at 2×10^{-5} is undetectable in the image recovered with IPCA+ARDI considering the chosen rank and iteration. This is because the feature becomes blended with the brighter disk halo when the entire disk is recovered. The 5σ contrast curve with the rank-3 PCA corroborates our observation and shows a sensitivity reaching a contrast of $\sim 2.6 \times 10^{-5}$ at the separation of the candidate ($0''.6$), which is consistent with the 3×10^{-5} contrast limit established by Grady et al. (2013) at this separation. According to the ATMO2020 model and assuming no extinction, our contrast threshold translates into a $1M_{\text{Jup}}$ sensitivity. This mass sensitivity does not match with the predicted mass inferred by Wagner et al. (2019), which was $2 - 3M_{\text{Jup}}$, using the COND “hot-start” model (Baraffe et al. 2003, 2015). Indeed, the ATMO2020 model primarily focuses on the atmospheric properties of exoplanets, while the COND hot-start model focuses on the early stages of giant planet formation, particularly the rapid accretion of gas and dust from a protoplanetary disk and the subsequent generation of internal heat. Nevertheless, to enable a consistent comparison, we also convert the contrast of 1×10^{-5} for the candidate in the L' band as presented in Wagner et al. (2019) into a mass using the ATMO2020. We found a mass of $0.5M_{\text{Jup}}$, suggesting that our detection limit of $1M_{\text{Jup}}$ would not allow us to detect the candidate. All quoted mass conversions assume negligible extinction from circumplanetary and circumstellar material, which is likely

⁵ <https://github.com/vortex-exoplanet/VIP>

optimistic. Inclusion of extinction would inflate our mass limits, affecting more the mass inferred from the Ks-band contrast than the mass inferred from the L' band.

Regarding the protoplanet proposed by Reggiani et al. (2018), we can observe in the images that complex structures arise at the separation of the candidate, including point-like features, especially when using PCA with ADI-only. This could be explained by filtered disk signal.

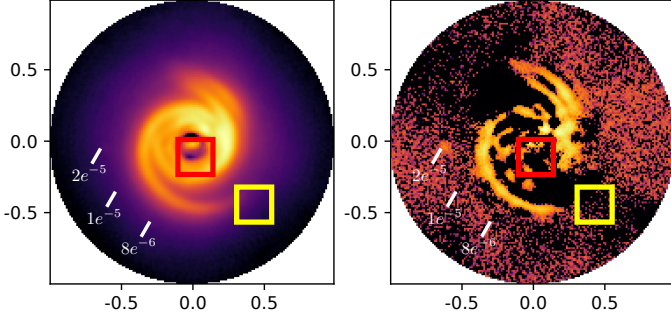


Fig. F.1. Results of IPCA (left) and 3-rank PCA with ADI only (right) for the VLT/SPHERE dataset of HD 36112 (Ren et al. 2023). Three simulated companions at a separation of $0''.6$ have been injected into the cube, with contrasts of 8×10^{-6} , 1×10^{-5} , and 2×10^{-5} . White lines indicate their positions. Two squares are displayed to mark the locations of the two claimed protoplanets: a red square at $0''.1$, indicating the candidate proposed by Reggiani et al. (2018), and a yellow square at $0''.6$, referring to the candidate proposed by Wagner et al. (2019).

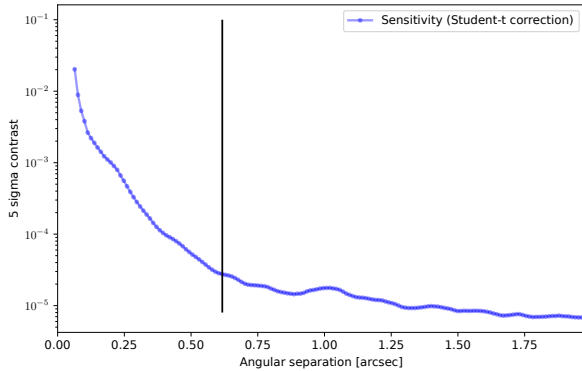


Fig. F.2. Contrast curve at 5σ confidence for the HD 36112 dataset, processed using rank-3 PCA with ADI only.

F.2. HD 169142

This section presents a test of detectability for the protoplanet candidate proposed by Gratton et al. (2019) and Hammond et al. (2023) in the HD 169142 system. The test involves injecting a fake planet candidate at the same separation of $0''.31$, using a contrast of 1.5×10^{-5} as estimated by Hammond et al. (2023). Similarly as in Appendix F.1, Fig. F.3 shows the results of IPCA and rank-3 PCA post-processing with ADI-only results for the dataset with the injected planet. Figure F.4 presents the 5σ contrast curve using rank-3 PCA with ADI-only. In Fig. F.3, the injected planet is barely visible, and only for the rank-3 PCA processing. Although it is at the noise level, it is still discernible due to its clear surrounding, likely a result of the over-subtraction effect. According to the contrast curve in Fig. F.4, the sensitivity at the separation of $0''.31$, where the candidate is located, is 1×10^{-4} . Moreover, PCA also reveals numerous point-like features due to the filtered bright inner disk. No such PLF appears

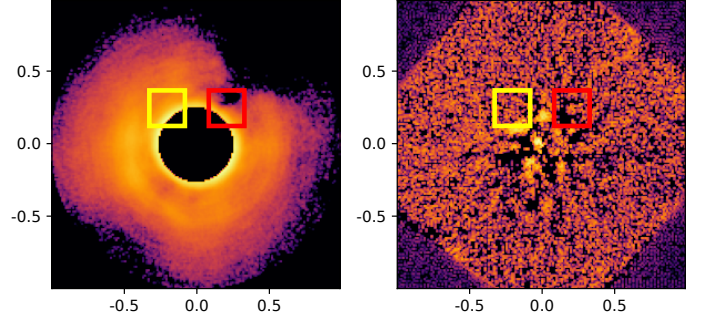


Fig. F.3. Results of IPCA with ARDI (left) and rank-3 PCA with ADI only (right), for an injected companion at a separation of $0''.31$ into the VLT/SPHERE dataset of HD 169142 (Ren et al. 2023), using a contrast of 1.5×10^{-5} . The position of the simulated candidate is indicated by a red square, while the location of the actual candidate proposed by Gratton et al. (2019); Hammond et al. (2023) is marked by a yellow square

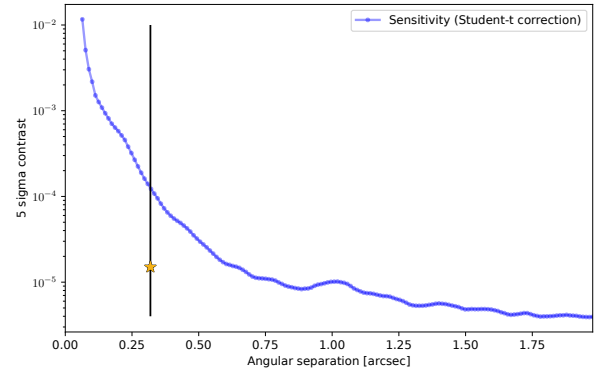


Fig. F.4. Contrast curve at 5σ confidence for the HD 169142 dataset, processed using PCA with ADI only. The separation of the candidate at $0''.31$ is indicated by a black vertical line. The yellow star indicate the expected contrast of the candidate, inferred by Hammond et al. (2023).

at the location of the actual protoplanet candidate. Our detection limit of 1×10^{-4} translates to $2M_{\text{Jup}}$, according to the ATMO2020 model, assuming no extinction. In comparison, the contrast in *YJH* bands reported by Hammond et al. (2023) corresponds to $1M_{\text{Jup}}$. Nevertheless, these results must be put into perspective regarding the presence of signal at the location of the protoplanet in polarized intensity images shown in Hammond et al. (2023), suggesting the presence of circumplanetary dust. Hence, the usage of the ATMO2020 model and the assumption of no extinction might not be appropriate. Therefore, the sensitivity limits in our data do not allow us to confirm or refute the candidate.

F.3. HD 97048

This section presents a test to infer the detection limit in the Ks-band for the protoplanet candidate deduced by Pinte et al. (2019) using local deviations in the Keplerian velocity in the disk surrounding HD 97048. The test involves injecting several fake planet candidates at the same separation as the candidate. For an injected planets at a contrast of 3×10^{-5} , we observe a marginal detection. Fake planets injected with a contrast below 1×10^{-5} are not recovered, indicating that we cannot detect planets smaller than one Jupiter mass with our images, according to the ATMO2020 model, assuming no extinction. Figure F.5 displays the image with the fake planet at contrast 3×10^{-5} ,

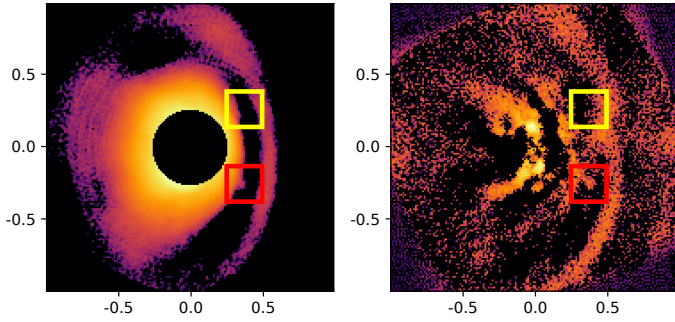


Fig. F.5. Results of IPCA (left), and rank-3 annular-PCA (right) with ADI only, for a fake companion at a separation of $0''.45$, using a contrast of 3×10^{-5} injected into the VLT/SPHERE dataset of HD 97048 (Ren et al. 2023). The location of the fake candidate is indicated by a red square, while the location of the actual protoplanet candidate inferred through the observation of kinks by Pinte et al. (2018) is indicated with a yellow square.

obtained with IPCA (left) and with 3-rank PCA with ADI-only (right). The injected simulated planet is visible with both PCA with ADI-only and IPCA with ARDI results. The PCA estimation also reveals numerous point-like features due to the filtered bright inner disk. Notably, no such point-like feature appears at the location of the actual protoplanet candidate. No contrast curve is provided for this dataset, as there is a disk signal at the same separation as the candidate.

F.4. HD 163296

This section presents the $6''$ -width uncropped observations of HD 163296 to check for any possible signal arising at the location of the detected kinks at 260 au ($2''.2$) from the star (Izquierdo et al. 2022; Pinte et al. 2018). We present in Fig. F.6 the results of the processed data using IPCA and PCA, both with ADI-only, and contrast curve is shown in Fig. F.7, indicating an achieved 5σ contrast of 6×10^{-6} at a separation of $2''.2$. In Fig. F.6, no signal is observed at the candidate location, suggesting that the companion might have a contrast below the detection limit of 6×10^{-6} . We injected multiple simulated companions at a separation of $2''.2$, with contrasts ranging from 1×10^{-3} to 3×10^{-6} . We found that for a contrast below 6×10^{-6} , the injected companion is hardly differentiable from remaining speckle patterns that have not been well subtracted. Additionally, the IPCA results reveal a halo spanning approximately from $2''.3$ to $4''.5$. In a previous study using data from the *Hubble* Space Telescope, Grady et al. (2000) reported extended signal between $2''.9$ and $3''.2$ in the southeast and northeast of HD 163296. However, our observation exhibits a distinctively circular halo, which does not match the disk inclination, suggesting that this signal may not originate from a circumstellar source and may not be related to the signal observed by Grady et al. (2000).

For the kink observed at a separation of $0''.85$, our detection limit at 5σ -contrast of 3×10^{-4} , translates to an apparent magnitude of $M_{Ks} \sim 13.6$. According to the ATMO2020 evolutionary models (Phillips et al. 2020), this magnitude corresponds to a mass sensitivity of approximately 4 Jupiter masses, while neglecting extinction. Hence, our observations may not reach sufficient depth to detect a Ks -band counterpart for this particular kink.

For the kink situated at $2''.2$, our detection limit at 5σ -contrast of 8×10^{-6} corresponds to an apparent magnitude of $M_{Ks} \sim 17.8$. This sensitivity level is lower than $1M_{\text{Jup}}$ (with $M_{Ks} < 16.3$)

based on the ATMO2020 evolutionary models and approximately 1 Jupiter mass ($M_{\text{J}} \sim 17.5$) according to the BEX models (Linder et al. 2019), assuming no extinction effects. Notably, based on the amplitude of the kink, Pinte et al. (2018) suggested a potential $2M_{\text{Jup}}$ planet, which would imply an expected apparent magnitude of $M_{Ks} \sim 15.6$ based on the BEX models. This discrepancy suggests a significant extinction effect in the Ks band, estimated to be $A_{Ks} \gtrsim 2.0$.

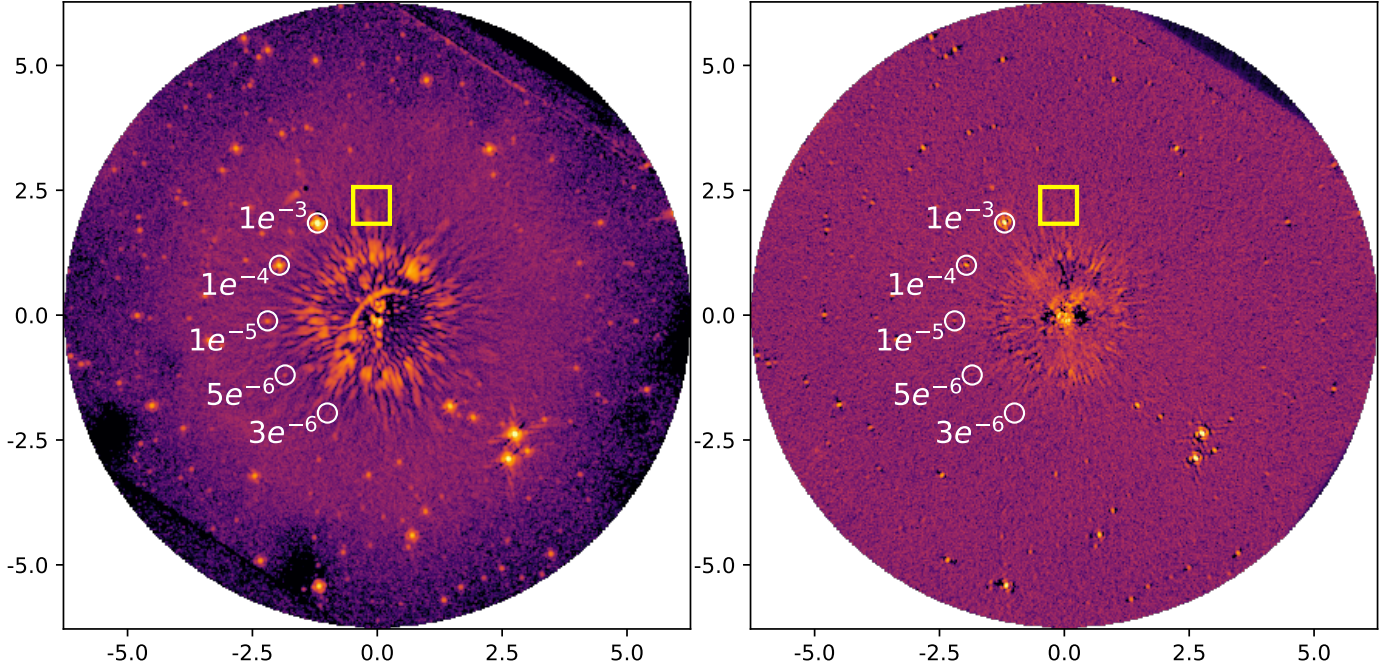


Fig. F.6. Results of IPCA (left) and rank-2 PCA (right) both with ADI only, for the uncropped observations of HD 163296. The position of the candidate, inferred by [Izquierdo et al. \(2022\)](#) and [Pinte et al. \(2018\)](#) from a local deviation in Keplerian velocity, is marked by the yellow square with dimensions 60×60 pixels. A zoom into the 60×60 pixels aperture is displayed in the bottom left of each image, with the planet's location indicated by an arrow on the smaller images. Five simulated companions at a separation of $2''.2$ have been injected into the cube, with contrasts ranging from 1×10^{-3} to 3×10^{-6} . White circle indicate their positions.

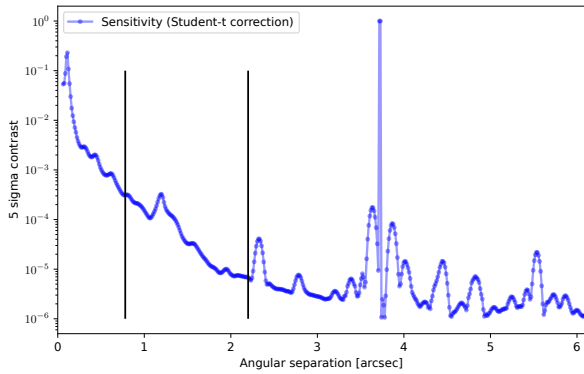


Fig. F.7. Contrast curve at 5σ confidence for the HD 163296 dataset, processed using PCA with ADI only. The separation of the candidate, inferred by [Izquierdo et al. \(2022\)](#) and [Pinte et al. \(2018\)](#) from a local deviation in Keplerian velocity at $2''.2$, is indicated by a black vertical line.